

Readme

The present scripts are made to help with the numerical analysis of data from radio nuclei transport experiments as demonstrated in *simulation_example.m*. The relevant functions are documented at the beginning of their file.

The used underlying physical model describes the concentration profile $c(t, x)$ of unabsorbed radioactive nuclei in the ground where t denotes time and x the depth below the surface. The temporal evolution of c is expected to obey the transport equation

$$\frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2} + v \frac{\partial c}{\partial x} - \lambda c \quad (1)$$

where D and v are diffusion and drift velocity coefficient respectively and λ is the decay constant of the nuclei. In the case of constant D and v , eq. 1 has an analytical solution. For the more complex case of spatially varying diffusion and drift velocity, a numerical scheme for solving the transport equation, as used in *transport_radio_nuclei.m*, is written down in the following.

We may discretize the spacial coordiante into segments $i = 1, \dots, N$ of thickness dx . We then discretize eq. 1 at each segment as

$$\frac{\partial c_i}{\partial t} = \frac{J_{in,i} - J_{out,i}}{dx} - \lambda c_i \quad (2)$$

where $J_{in,i}$ and $J_{out,i}$ are the conserved incoming and outgoing fluxes of each segment which we discretize as

$$J_{in,i} = \frac{D_{i-1} + D_i}{2} \frac{c_{i-1} - c_i}{dx} + \frac{v_{i-1} + v_i}{2} \frac{c_{i-1} + c_i}{2}, \quad J_{out,i} = J_{in,i+1} \quad (3)$$

thus leading to

$$\begin{aligned} \frac{\partial c_i}{\partial t} = & \frac{1}{2dx^2} [(D_{i-1} + D_i) c_{i-1} - (D_{i-1} + 2D_i + D_{i+1}) c_i + (D_i + D_{i+1}) c_{i+1}] + \\ & + \frac{1}{4dx} [(v_{i-1} + v_i) c_{i-1} + (v_{i-1} + v_{i+1}) c_i + (v_i + v_{i+1}) c_{i+1}] - \lambda c_i. \end{aligned} \quad (4)$$

This scheme then can be time iterated with conventional solvers. The boundary conditions read

$$J_{in,1} = v_1 c_0(t), \quad J_{out,N} = \left(\frac{v_{N-1} - v_N}{2} \right) (1.5c_N - 0.5c_{N-1}). \quad (5)$$

Here, the time course of the concentration at the surface $c_0(t)$ can be chosen freely. For the case of constant diffusion and drift velocity coefficientt this scheme simplifies to form in [Frissel et al. (1974)].