

Next generation platform for implementing fast gates in ion trap quantum computation

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Scalable trapped-ion quantum computation relies on the development of high-fidelity fast entangling gates in a many ion crystal. Conventional geometric phase gates either suffer from scattering errors or off-resonant carrier excitations. A potential route to achieve fast entanglement is creating a standing wave which can suppress the unwanted carrier coupling.

Non-Adiabatic Mølmer Sørenson Gates

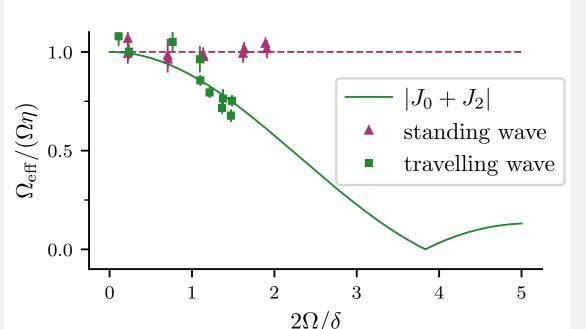
- Two qubit gates are implemented by coupling spin with shared motion of the ions.
- Fast entangling gates enable deeper quantum computational circuit depths for given level of incoherent error.
- But going fast excites multiple motional modes ("spectator modes") which can introduces errors.
- Mølmer Sørenson (MS) interaction is common two-qubit gate which requires a bichromatic field incident on the ions. Using travelling waves gives the Hamiltonian

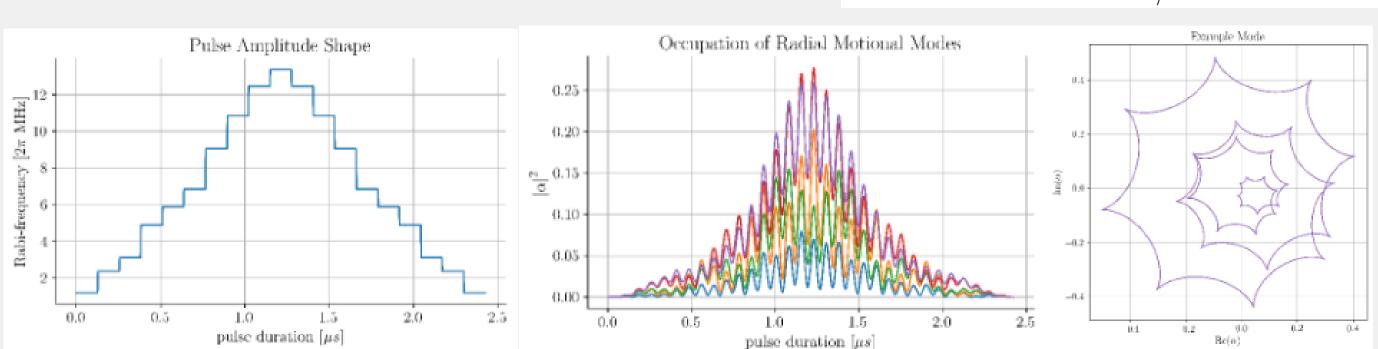
 $\hat{H}_{MS-TW} = \hbar \Omega \hat{S}_{\phi-\pi/2} \cos{(\delta t)} + \hbar \Omega \eta \hat{S}_{\phi} \cos{(\delta t)} (\hat{a} e^{-i\omega_z t} + \hat{a}^{\dagger} e^{i\omega_z t})$ with the first term being the carrier whilst the second is the desired coupling.

• As these terms do not commute, in the ??interaction picture?? this Hamiltonian may be expressed as [Xref Canzz?X]:

$$\hat{H}_{MS-TW} = \hbar \eta \Omega (J_0(2\Omega/\delta) + J_2(2\Omega/\delta)) \cdot \cos(\delta t) \hat{S}_{\phi}(\hat{a}e^{-i\omega_z t} + \hat{a}^{\dagger}e^{i\omega_z t})$$

- Carrier term: causes the spin dependent force coupling to be modulated by (JO+J2).
- Spectator excitation: Amplitude shaped pulses [ref vera] effectively remove "spectator" error by closing phase loops of all excited modes.





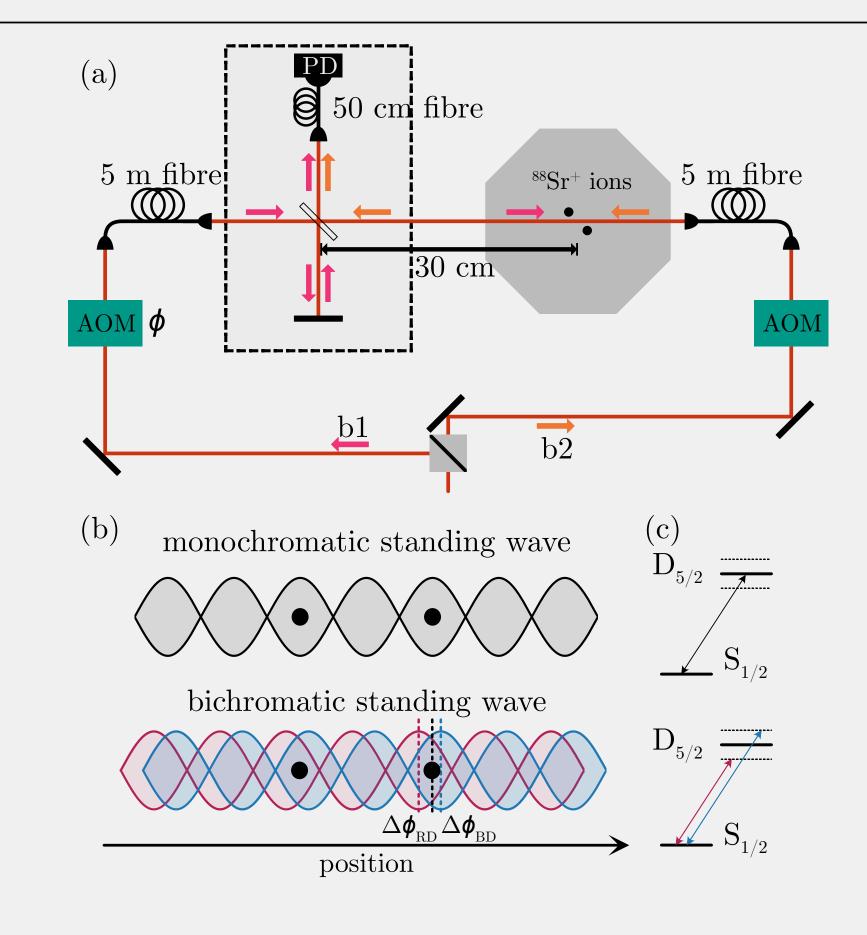
Standing Wave Single and MS Gates

- Using standing wave gives complete control over phase visible to ions.
- This extra freedom allows fast gates by preventing the saturation effect seen in the travelling MS.
- Bichromatic standing wave Hamiltonian where ions are seperated by $n\lambda/2$:

 $\hat{H}_{MS-SW} = \hbar 2\Omega \hat{S}_{\phi} \cos{(\delta t)} \sin{(\Delta \phi/2)} + \hbar 2\Omega \eta \hat{S}_{\phi} \cos{(\delta t)} (\hat{a}e^{-i\omega_z t} + \hat{a}^{\dagger}e^{i\omega_z t}) \cos{(\Delta \phi/2)}$

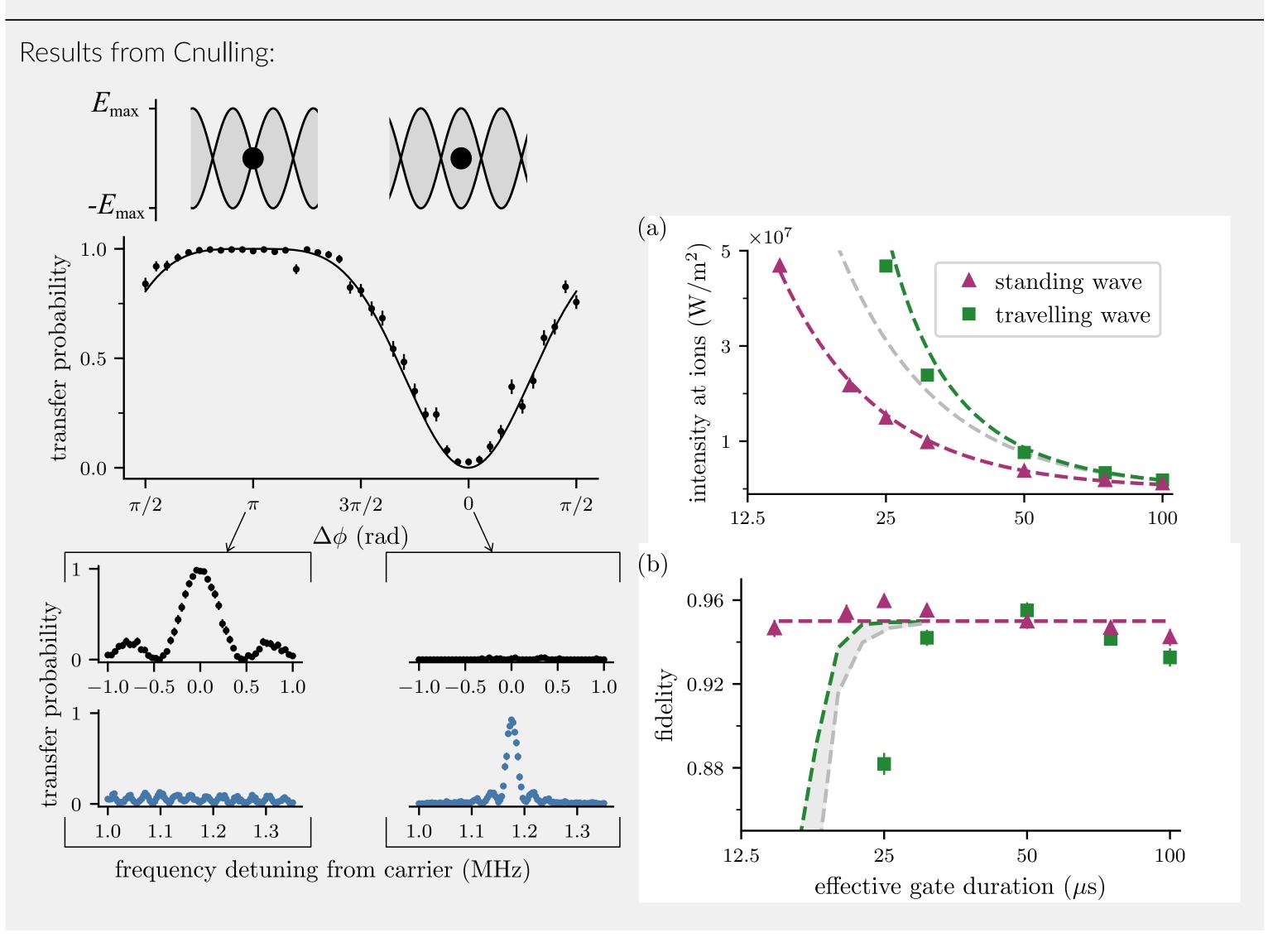
- Setting $\Delta \phi = 0$ (ions sitting at antinodes) we suppress the carrier term and maximise sideband coupling.
- However standing waves require phase stabilisation to perform coherent interactions.

Phase Stabilization



- Passively using enclosure and actively by feedback process to AOM.
- 1) Fast drifts removed utilising interference of light from the two branches on a photodiode (PD). X XX diagram rather than words. 2) As PD lockpoint is 30 cm away from ions, a second feedback loop using the ion as a sensor is required. We do this by performing a $\pi/2$ -pulse using arm 1 ($\pi/2$, b1) followed immediately by a $\pi/2$ -pulse using arm 2 ($\pi/2$, b2). This is equivalent to a zero-delay Ramsey sequence, which gives a signal sensitive on the difference in phase between the two pulses, hence the relative phase be- tween the branches.
- No detriment using lattice: Standard Randomized Benchmarking gave quality of our single qubit rotations per gate to be $\epsilon=0.173(3)\%$ for a travelling wave and $\epsilon=0.144(3)\%$ for the standing wave.

Experimental Results



New Platform

- New ion trap experiment in development for exploring fast gate regime.
- Two high NA lenses allow optical access on both faces of the trap to create an array of singly addressing standing waves.
- 3D segmented trap design from NPL facilitates low heating rates whilst enabling ion shuttling and crystal rotations.
- Calcium 40
- quadropole transition used.
- MuMetal shield and permanent magnets.
- fig: solidworks of new experiment

