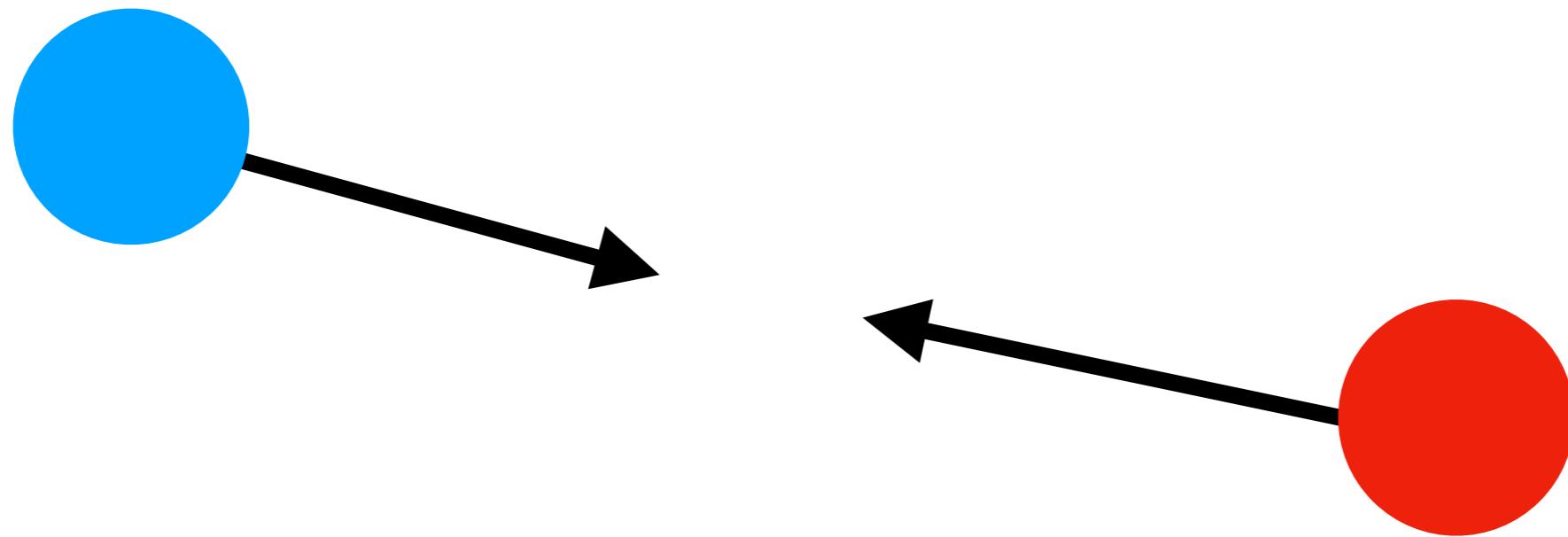




# Lecture 18:

# Monte Carlo methods

# What do we need for Simulation?

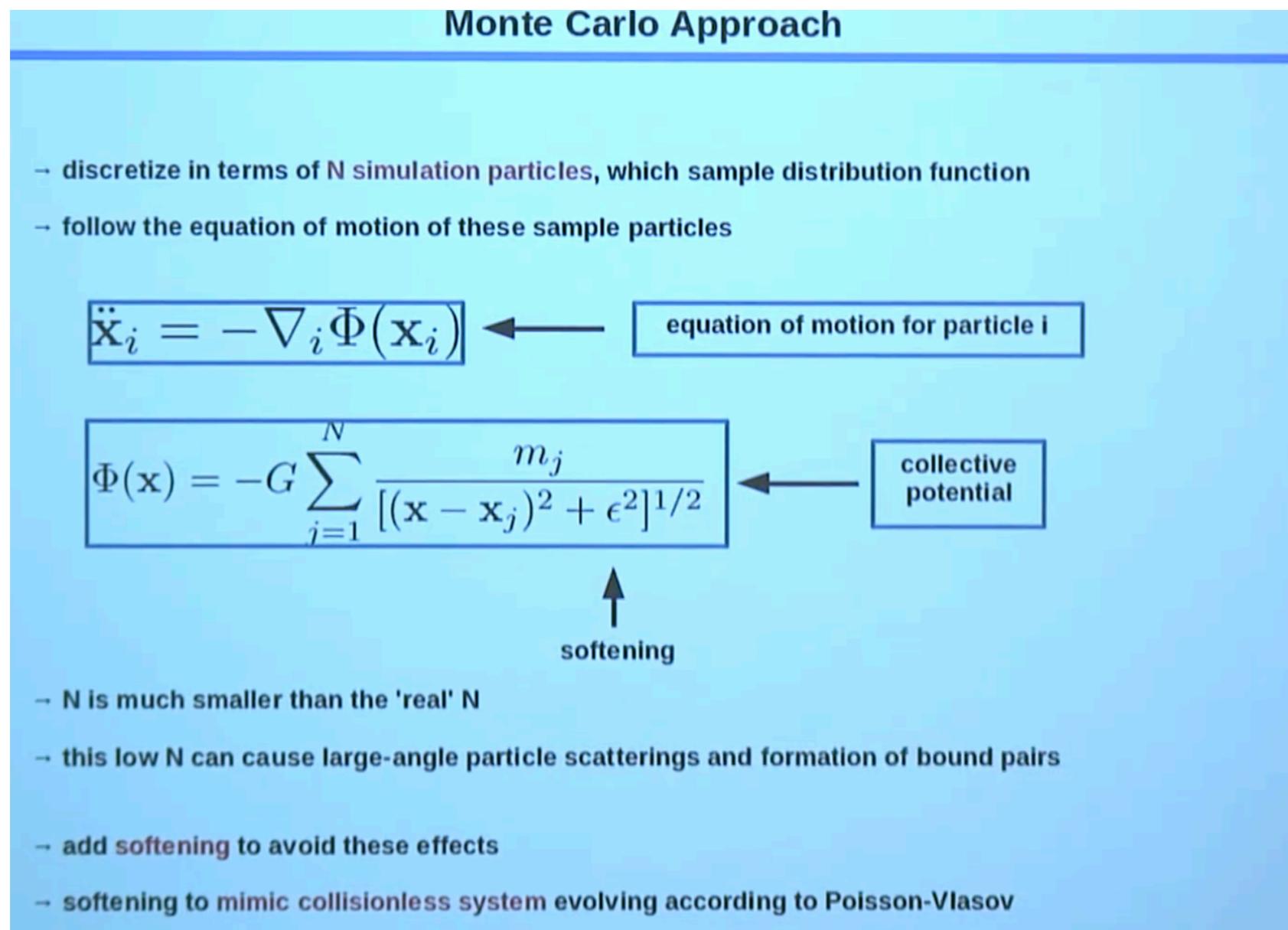


$$F = \frac{Gm_1m_2}{|\vec{r}_1 - \vec{r}_2 + \epsilon|^3} (\vec{r}_1 - \vec{r}_2)$$

- To simulate this on a computer we will add a “softening” term

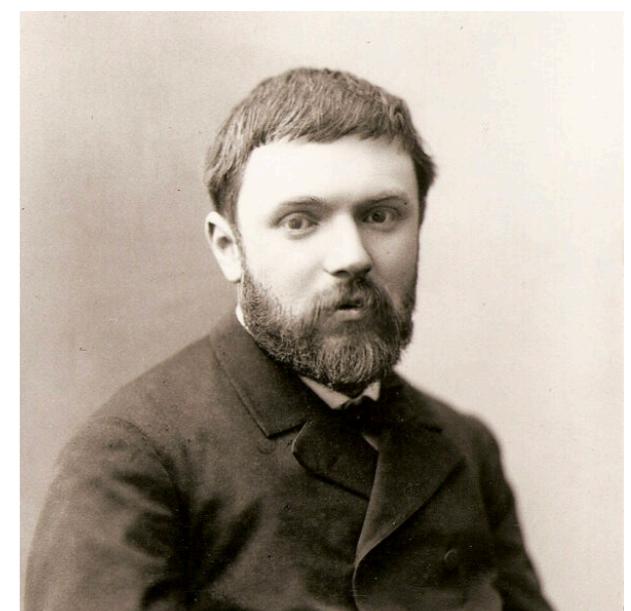
# This is how DM/galaxies<sup>3</sup> are modeled

- Instead of treating matter as a fluid
  - Discretize matter into chunks and solve n-body problem



# 3 body Problem: History

- The original two body problem was solved in 18th century
  - Work done by Netwon, Bernoulli Bros, Euler, Laplace,...
  - All started on the 3 body problem and built on the two body
- King Oscar II decided to make a competition:
  - For his 60th birthday he bestowed a prize on who could solve the n-body problem

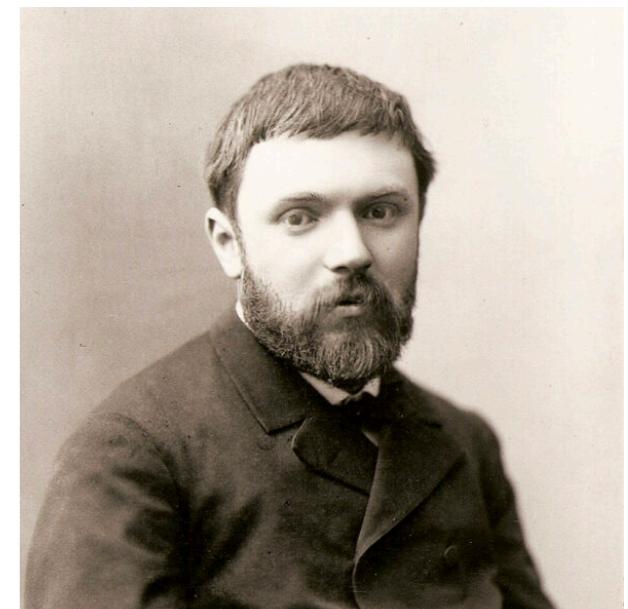


# 3 body Problem: History

- The original two body problem was solved in 18th century
  - Work done by Netwon, Bernoulli Bros, Euler, Laplace,...
  - All started on the 3 body problem and built on the two body
- King Oscar II (king of Sweden) decided to make a competition:
  - For his 60th birthday he bestowed a prize on who could solve the n-body problem

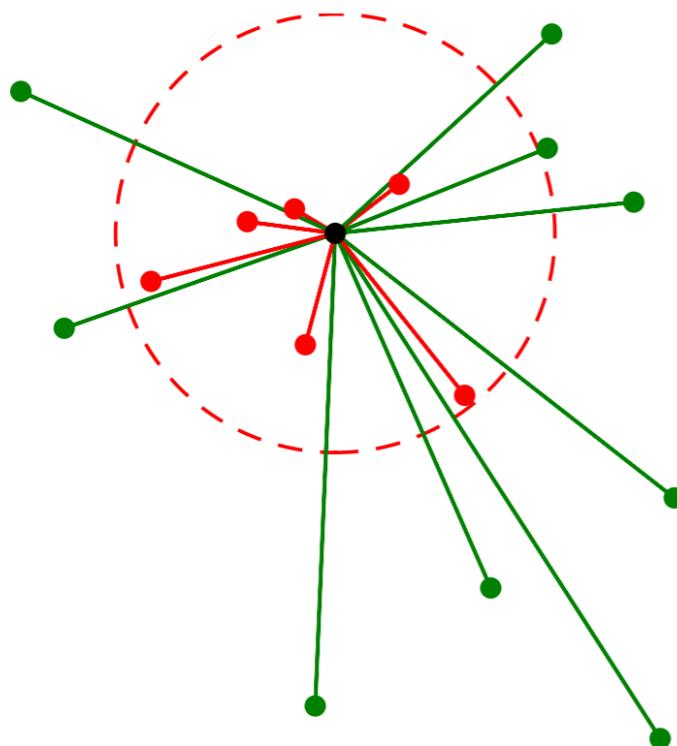
Henri Poincare →

He proved no solution existed!



# Going to N-body

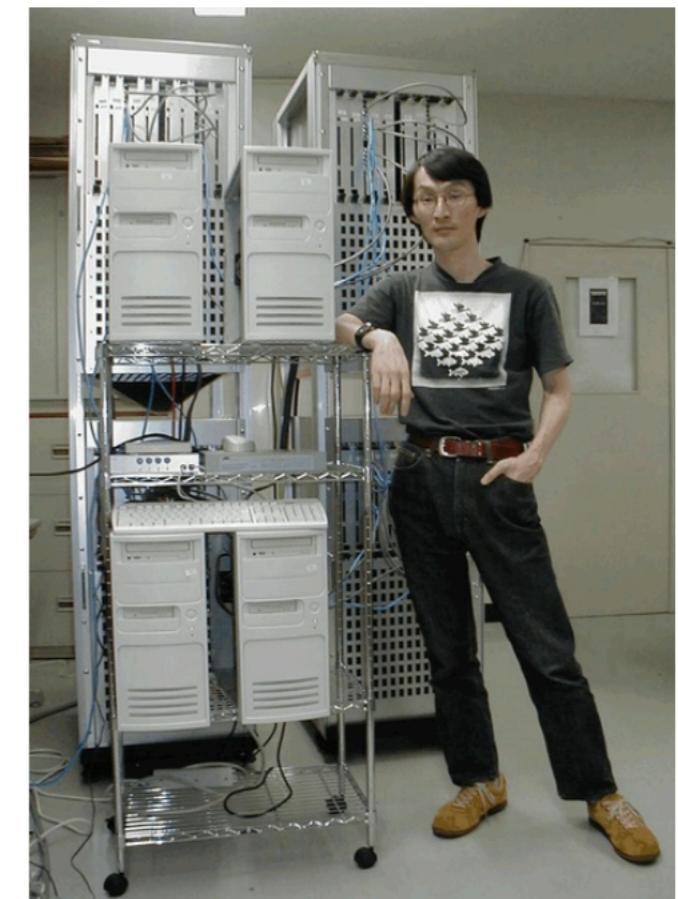
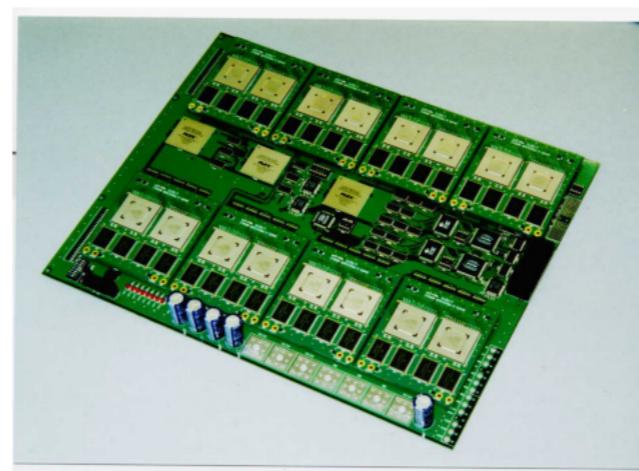
- The challenge of solving this numerically for n-body
  - This scales with the number of bodies  $N^2$
  - Requires the computations of all pairwise distances



- For  $N=1000$  (1 Million/step computations)
- For  $N=10^6$  Something ridiculous

# Historical Solutions

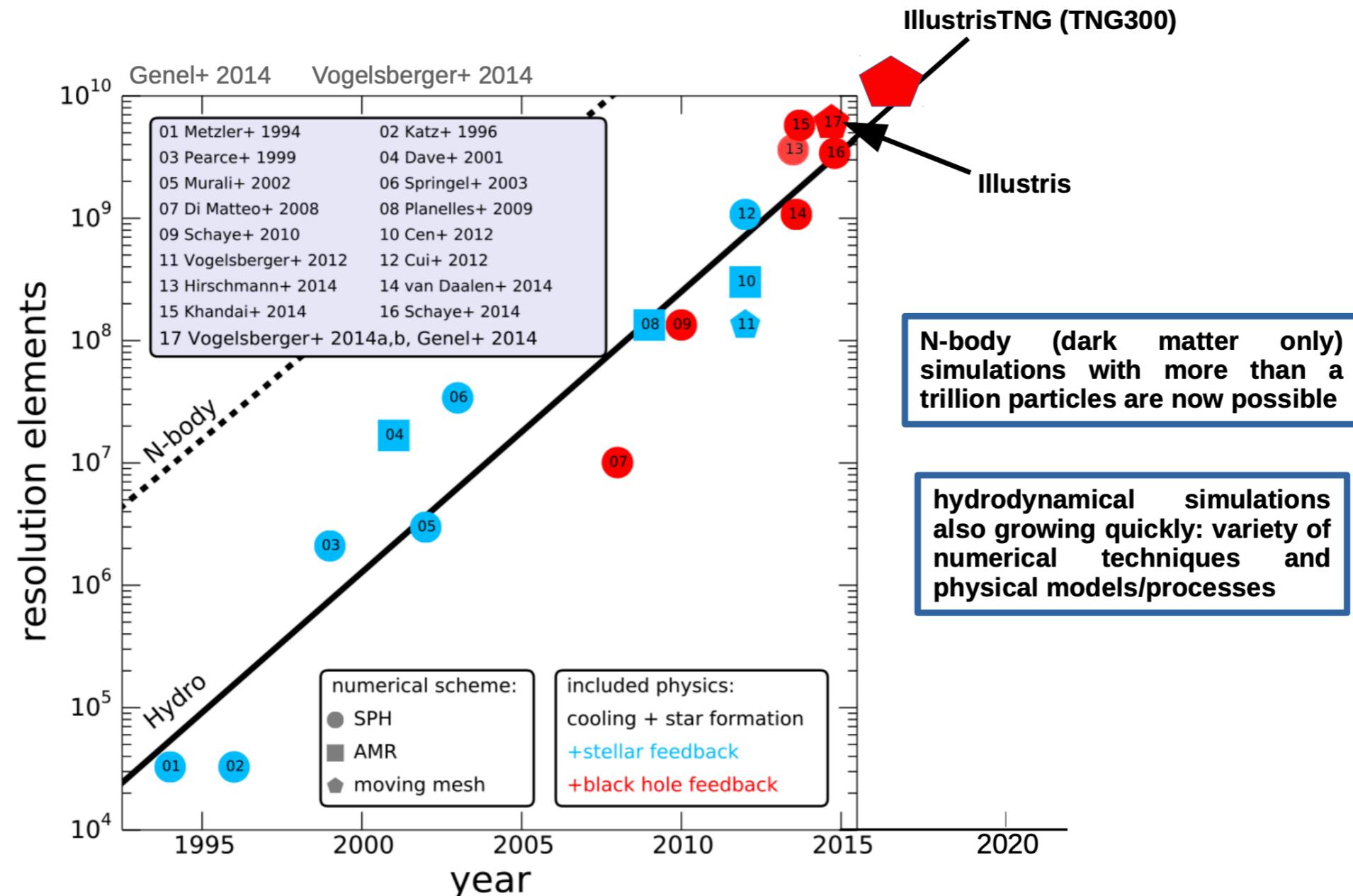
- One approach has been to build dedicated computing hardware
  - Dedicated hardware that can do large-scale parallel computation
  - Focused specifically on n-body simulation
  - GRAPE boards (GRAvity PipelinE)
- Now done with GPUs



Jun Makino with GRAPE-6

# Scaling of n-body

## The Evolution of Large-Scale Simulations

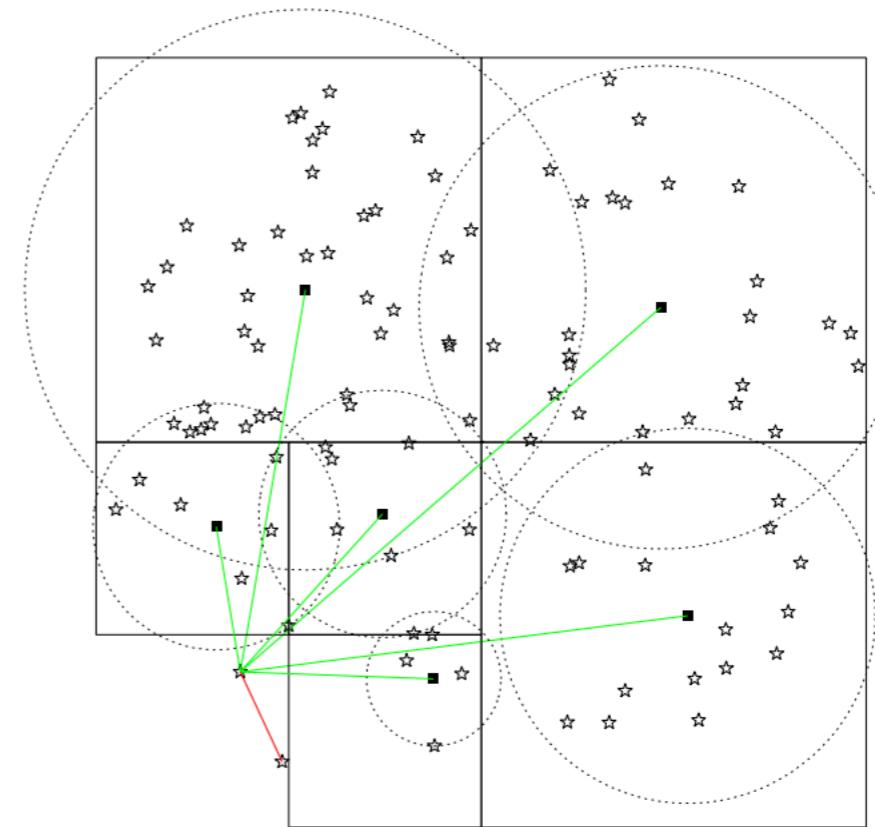
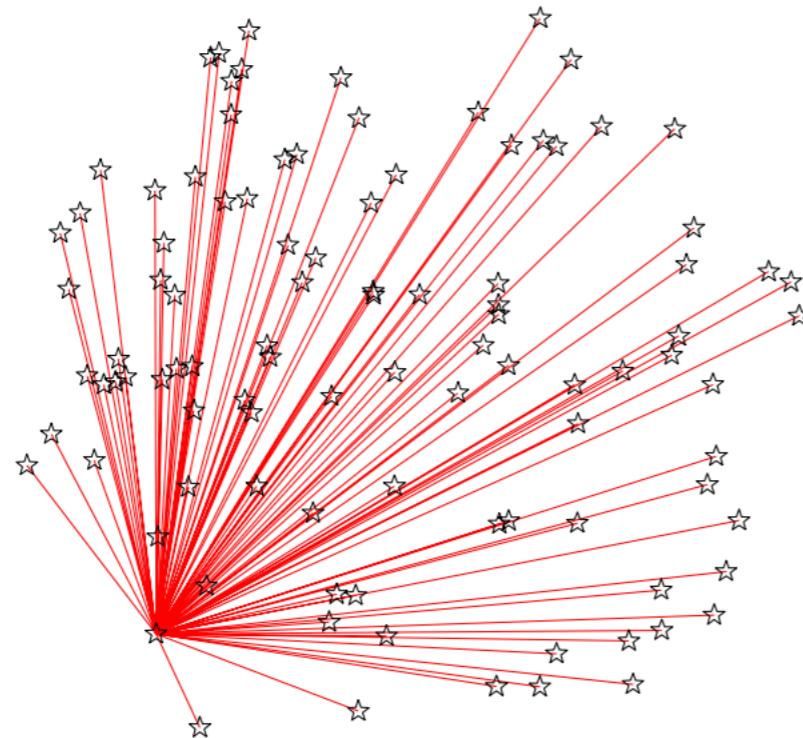


- M. Vogelsberger([https://indico.cern.ch/event/736594/contributions/3184103/attachments/1738225/2812076/talk\\_vogelsberger.pdf](https://indico.cern.ch/event/736594/contributions/3184103/attachments/1738225/2812076/talk_vogelsberger.pdf))

# How<sup>9</sup>

# do you deal with N-body?

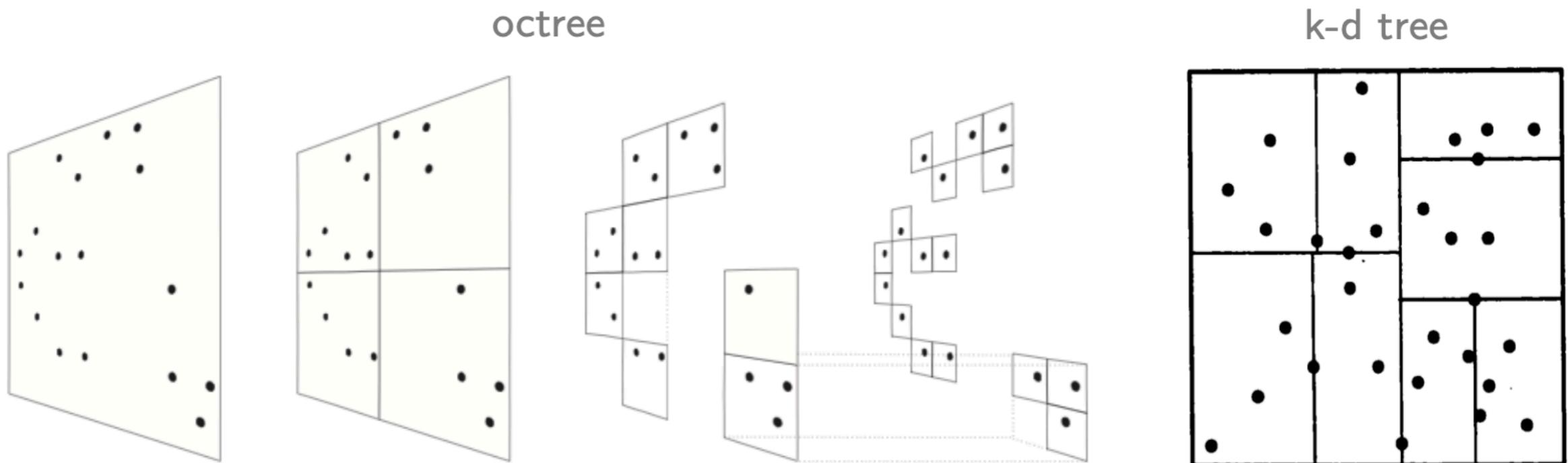
- Barnes-Hut Algorithm re-rank stars into a tree structure
  - Structure is a grid over the whole space



[from Dehnen & Read 2011]

# Tree Construction

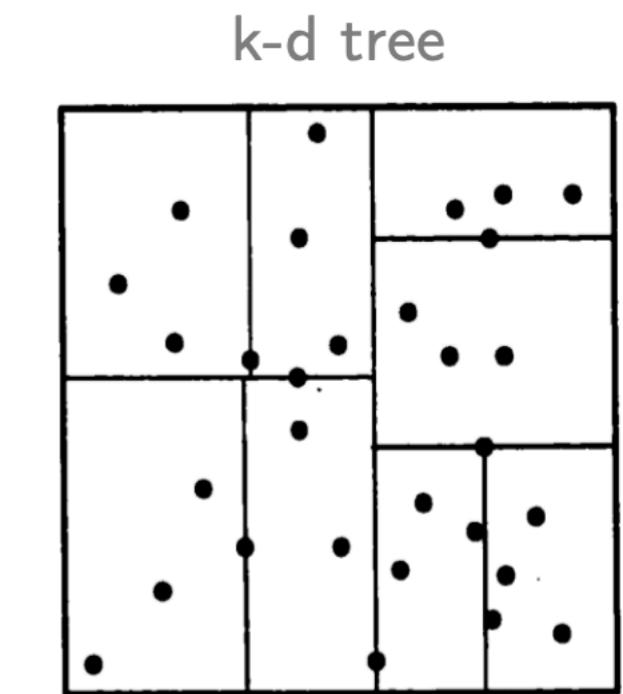
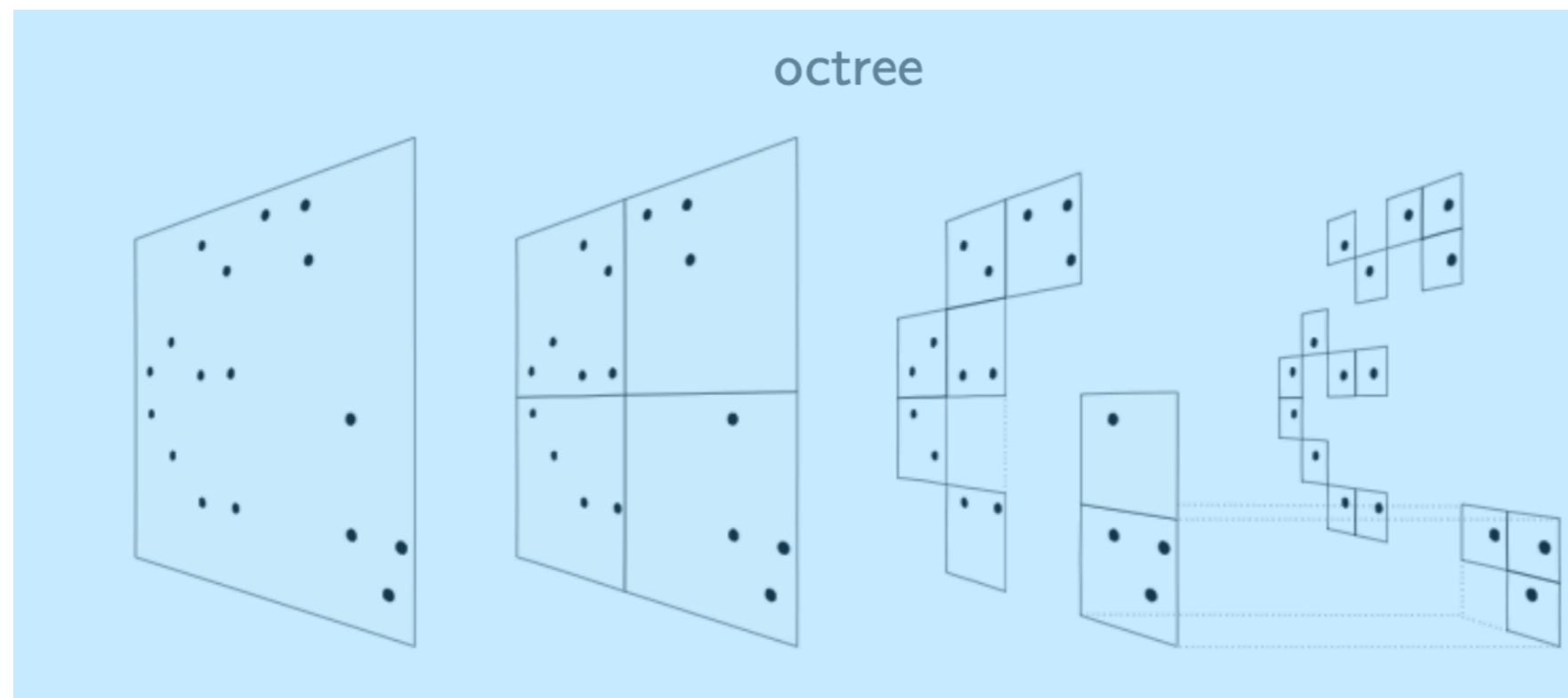
- QuadTree/OctTree
  - Split each square(cube) into 4(8) sub regions
- KD Tree
  - Use the data to draw equal numbered regions in space



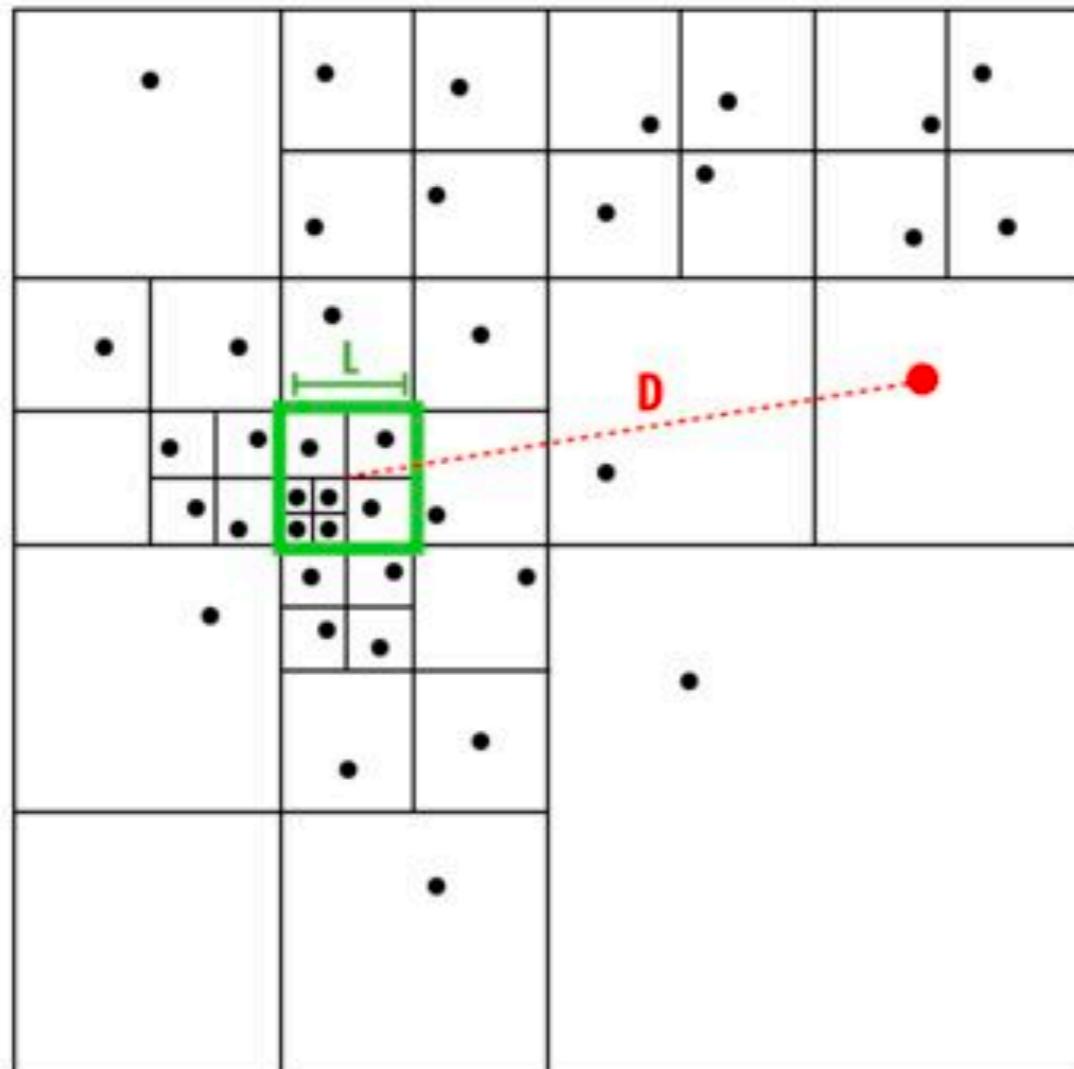
# Tree Construction

- QuadTree/OctTree
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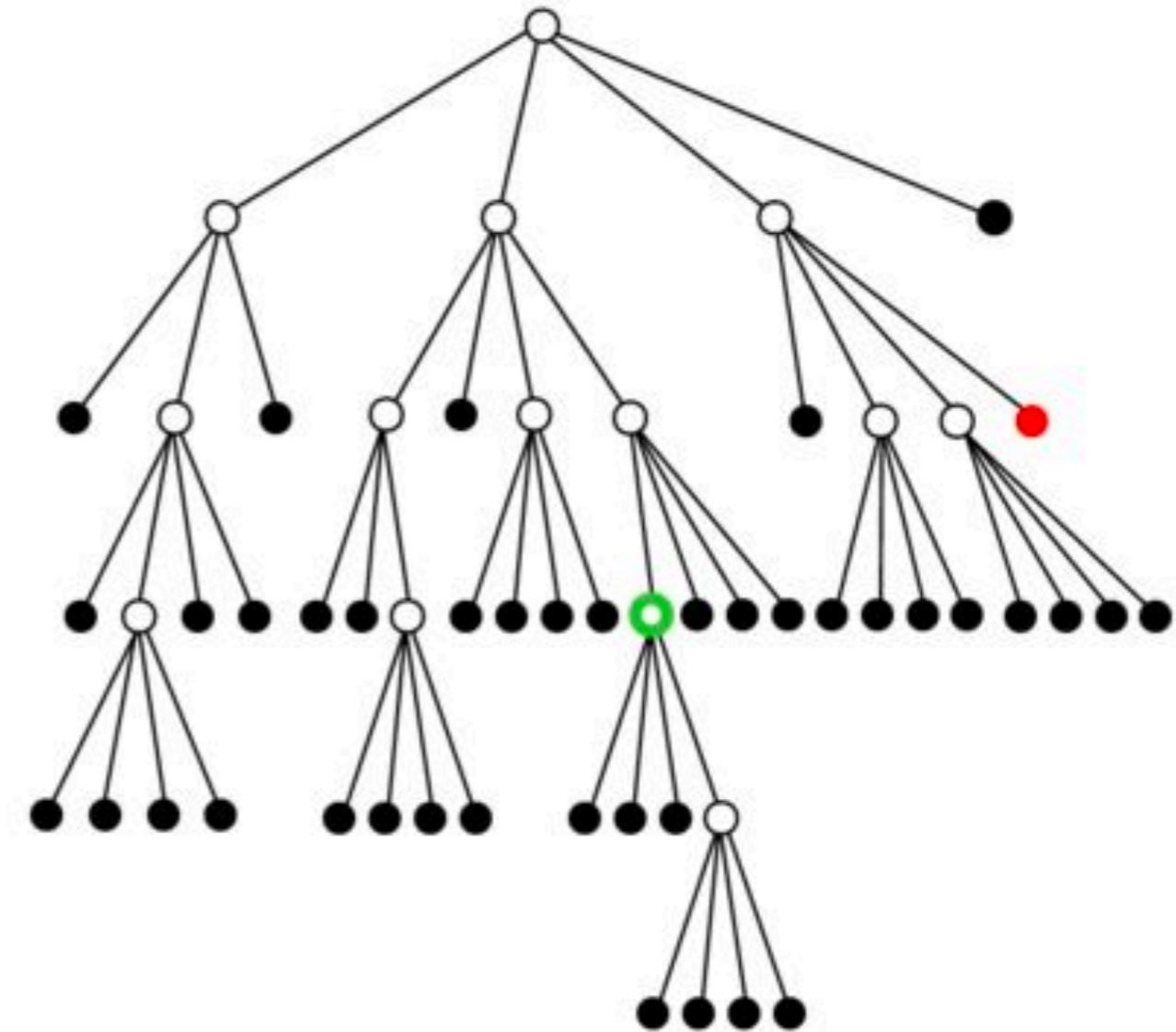
**Just requires that we know the bounds of the space**



# Visualizing Tree



Spatial Domain



Quad-Tree Representation

Note that for the big grids, this equates to Gauss' law style approach  
Treat each square as a star w/total mass at mass weighted center

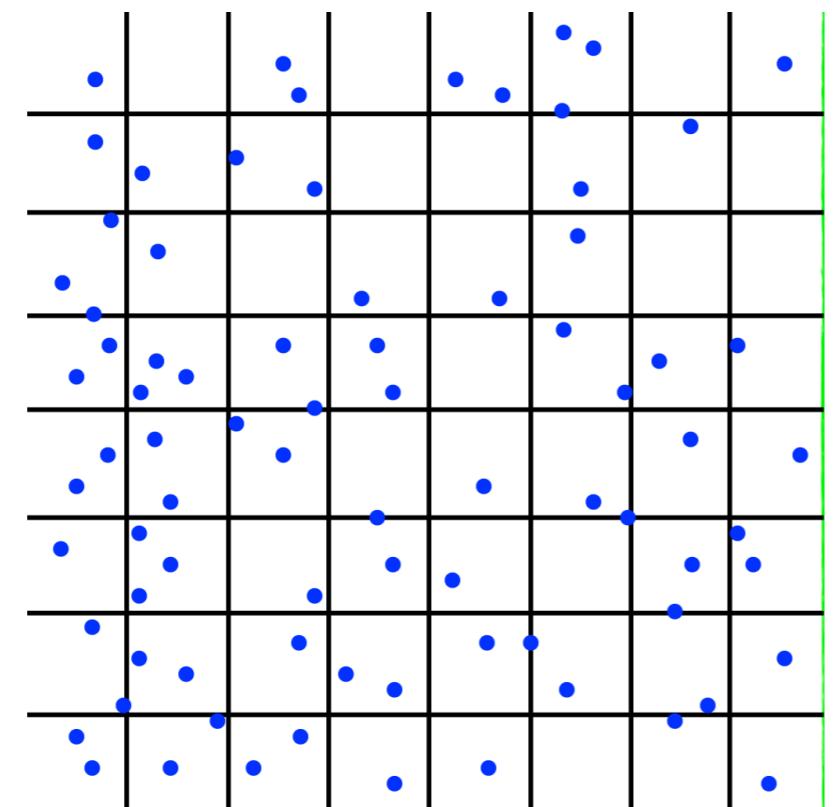
# Barnes-Hut Algorithm

- We can follow a step by step construction of this :
  - 1. Construct tree structure with bounds
  - 2. Loop over stars and fill tree structure
  - 3. Loop over stars and compute distance
    - Full n-body computation for nearby trees only
  - 4. Step forward everything
- The above process is  $N \log(N)$  in computational time

# Larger Scale Concept

- For the really large scale n-body dynamics
  - To get global motion
  - Grid up the whole space and Fourier transform it
  - Evolve the system in Fourier space

FFT(



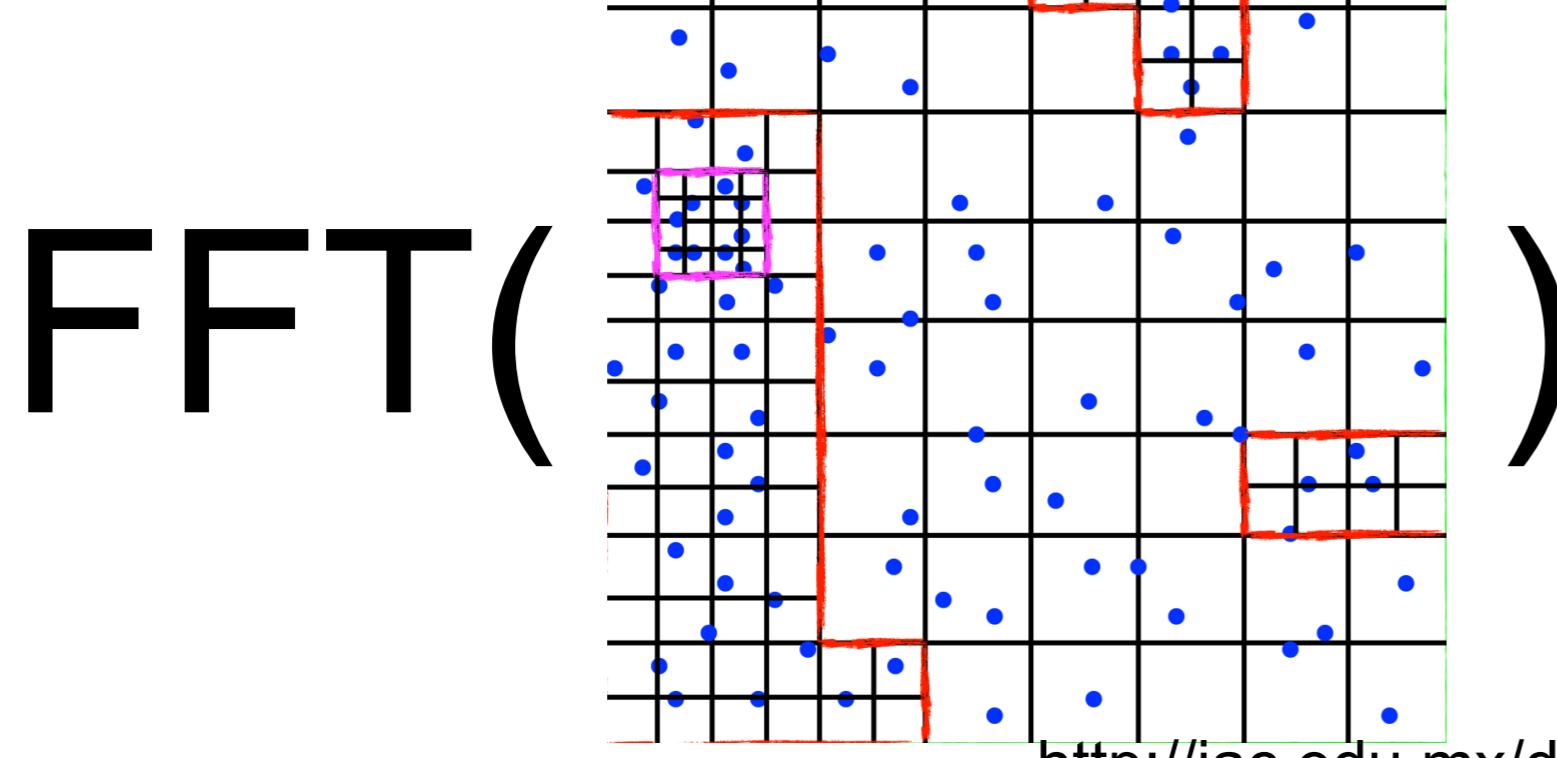
)

Solve for the Gravational  
Potential in Fourier space

Step Potential

# Larger Scale Concept

- For the really large scale n-body dynamics
  - To get global motion
  - Grid up the whole space and Fourier transform it
  - Evolve the system in Fourier space



Solve for the Gravational Potential in Fourier space

Step Potential



# Lecture 18:

# Monte Carlo methods

# Monaco



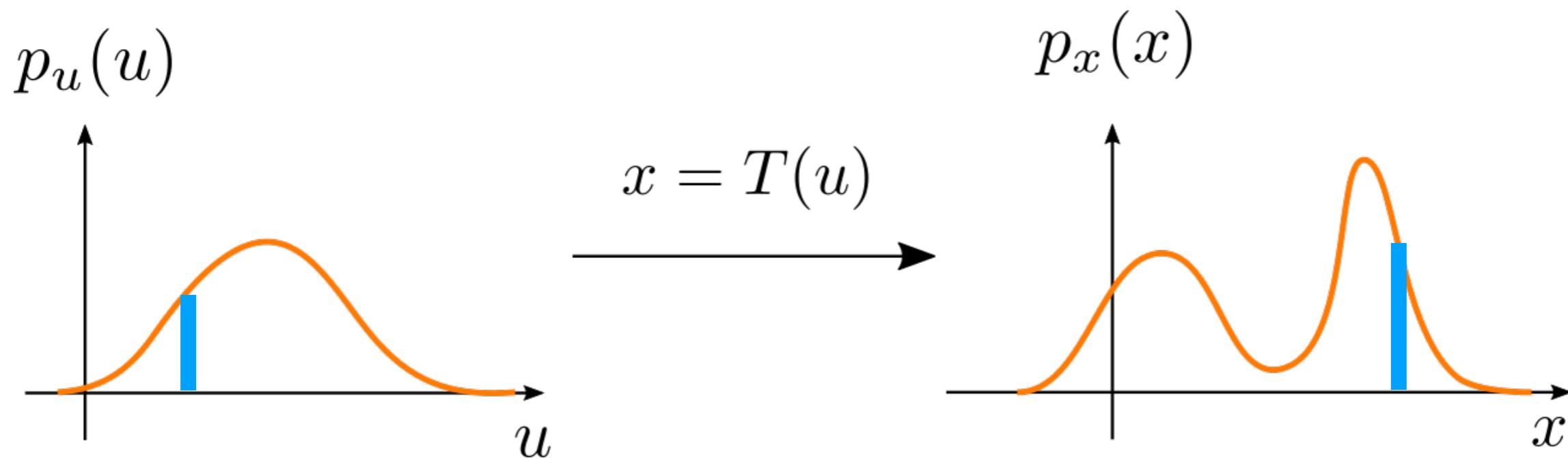
# Monte Carlo(MC)

- Have been seeing Monte Carlo methods throughout class
  - Any time we randomly sample that's an MC method
  - Effectively we are just rolling the die



# Monte Carlo vs Integration

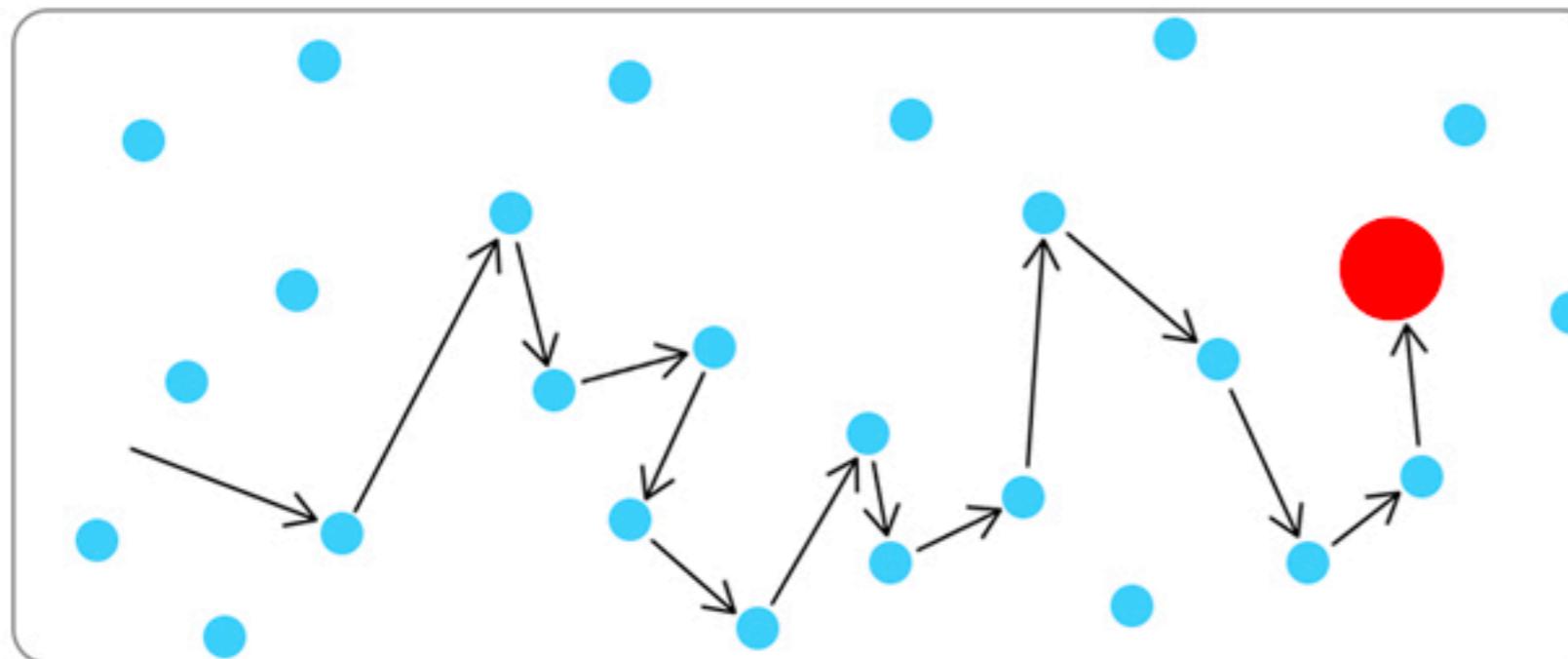
- Monte Carlo is a form of integrator
  - However non-deterministic and varies over distribution



- Monte Carlo typically used when
  - we can't model things analytically any more
  - Replace a whole distribution with just an event (small region)

# Brownian Motion

## Brownian Motion



Fluid molecule

Suspended particle

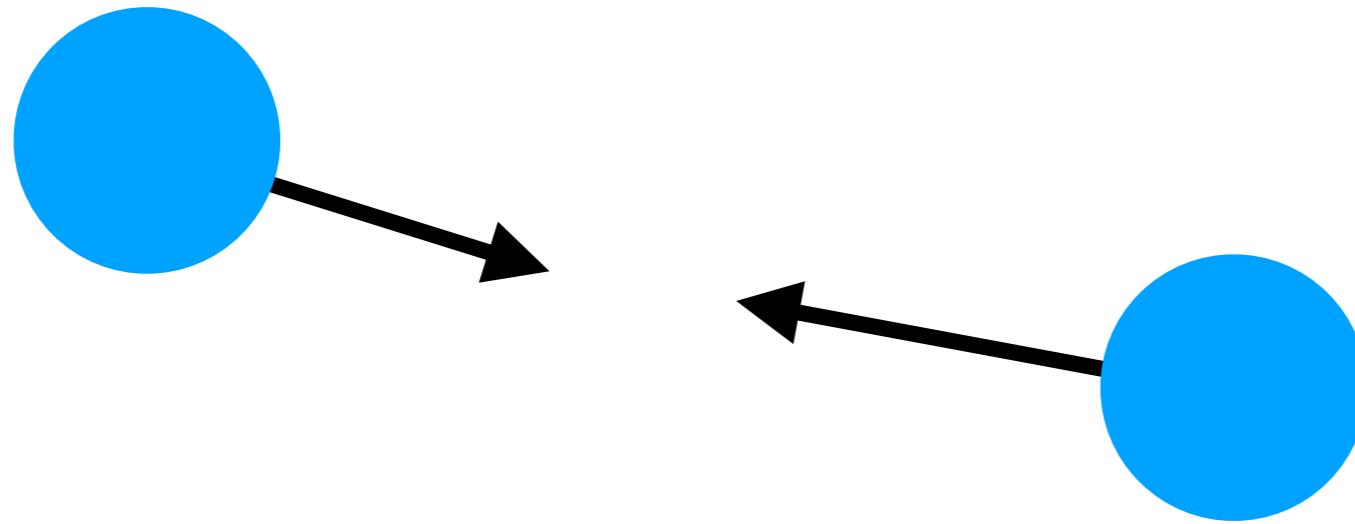
 ScienceFacts.net

$$\begin{aligned}
 f(v_x, v_y, v_z) &= \left[ \frac{m}{2\pi kT} \right]^{3/2} e^{-m(v_x^2 + v_y^2 + v_z^2)/2kT} \\
 &= \left[ \frac{m}{2\pi kT} \right]^{3/2} e^{-mv^2/2kT}
 \end{aligned}$$

using  $v^2 = v_x^2 + v_y^2 + v_z^2$

- At each step
  - We just randomly sample the velocity from a Gaussian
  - We can do this many times to look at overall motion

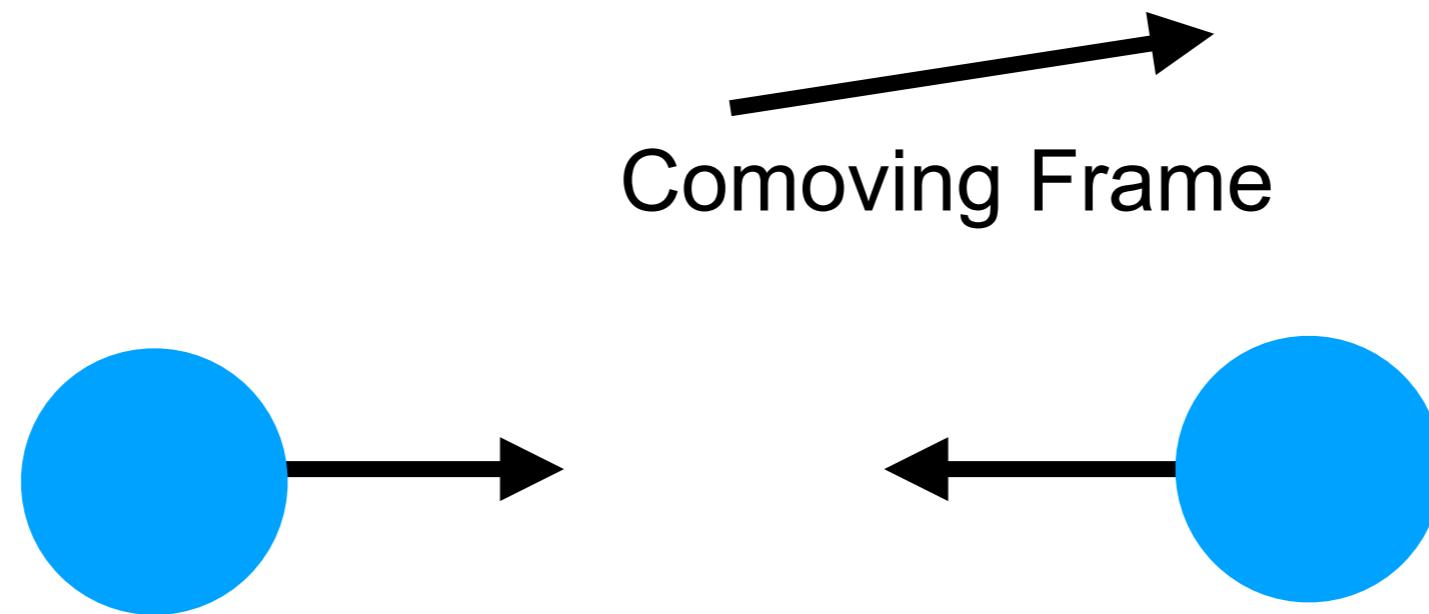
# The motion at each step



## Elastic Collision

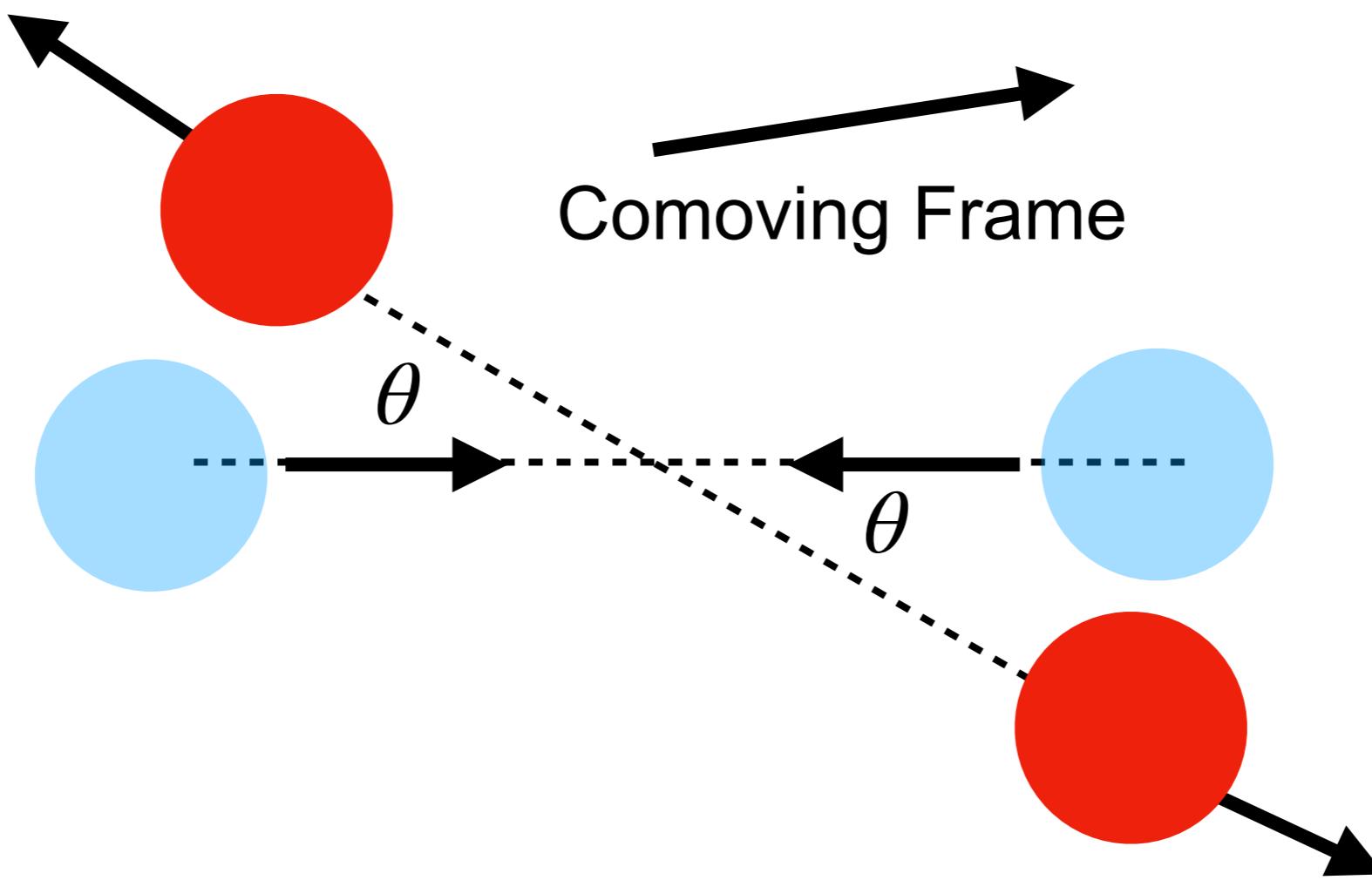
- Just sample particle collisions at each step

# The motion at each step



Elastic Collision  
In COM Frame

# The motion at each step



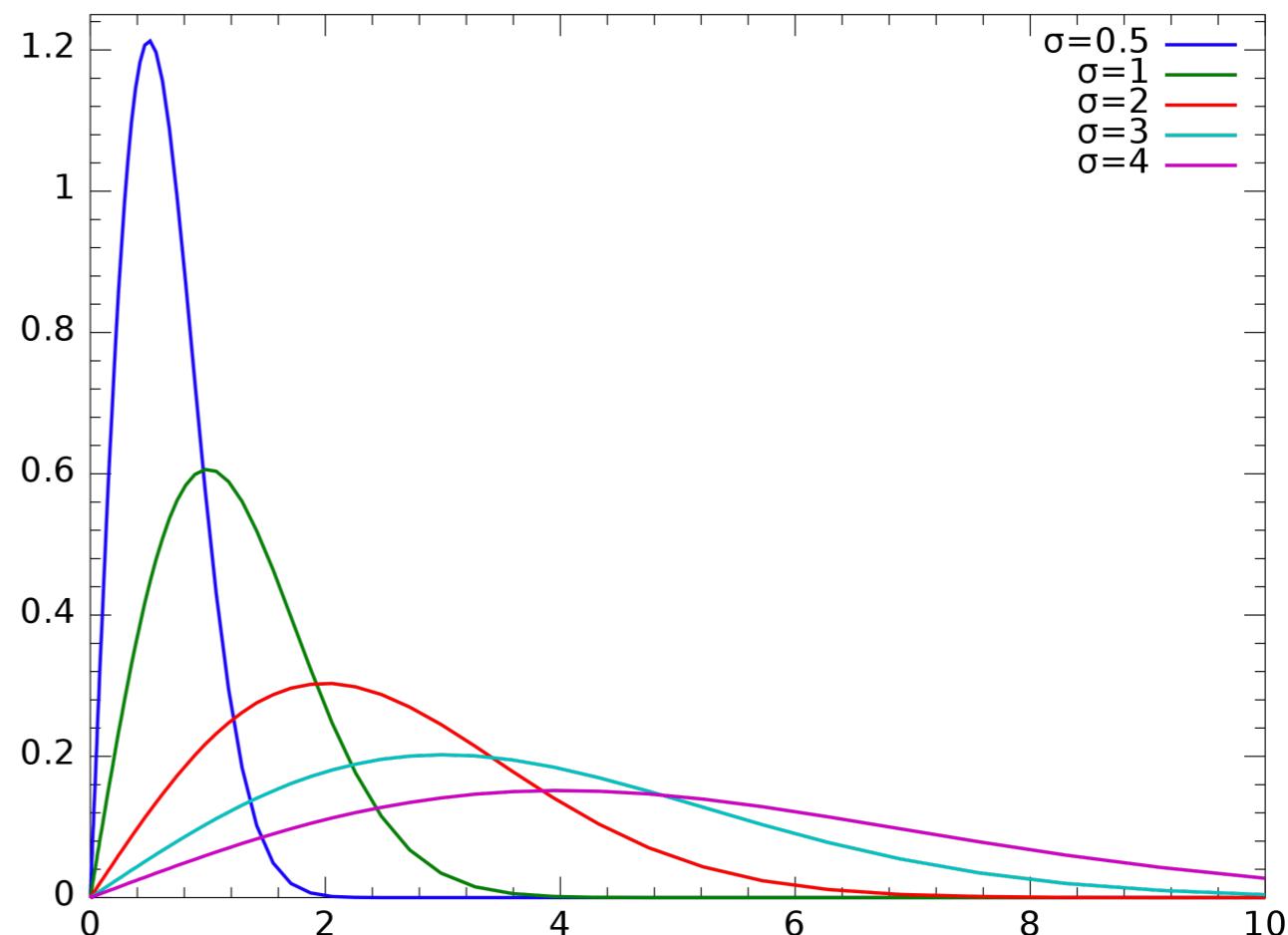
Elastic Collision  
In COM Frame

# Rayleigh Distribution

Rayleigh is a distribution of the radius in a 2D Gaussian

$$f_U(x; \sigma) = f_V(x; \sigma) = \frac{e^{-x^2/(2\sigma^2)}}{\sqrt{2\pi\sigma^2}}. \quad f(x; \sigma) = \frac{x}{\sigma^2} e^{-x^2/(2\sigma^2)}, \quad x \geq 0,$$

$$F_X(x; \sigma) = \frac{1}{2\pi\sigma^2} \int_0^{2\pi} \int_0^x r e^{-r^2/(2\sigma^2)} dr d\theta = \frac{1}{\sigma^2} \int_0^x r e^{-r^2/(2\sigma^2)} dr.$$





# Proton Therapy



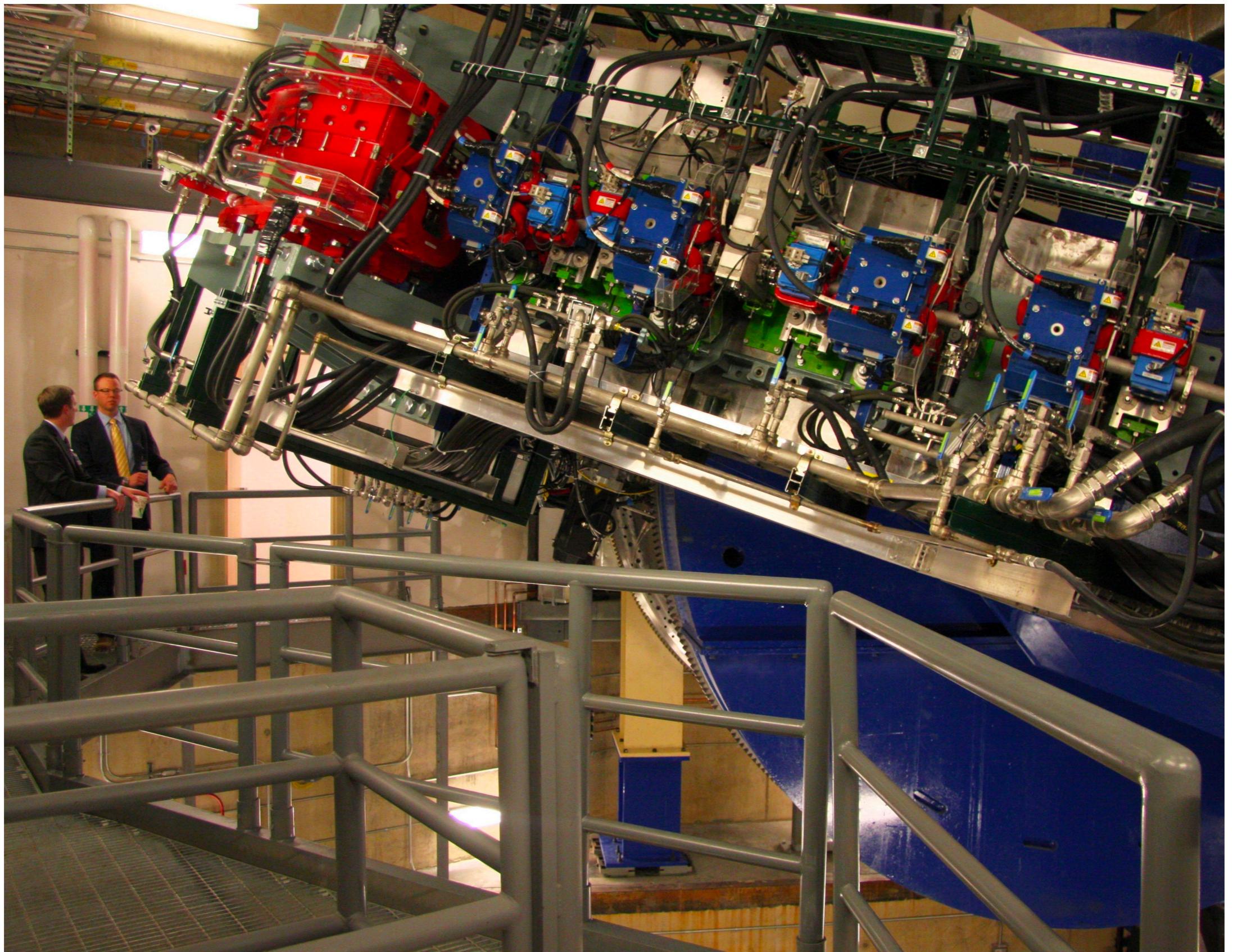
Proton Therapy Center at MGH

# Typical Device

## Particle Therapy Centre

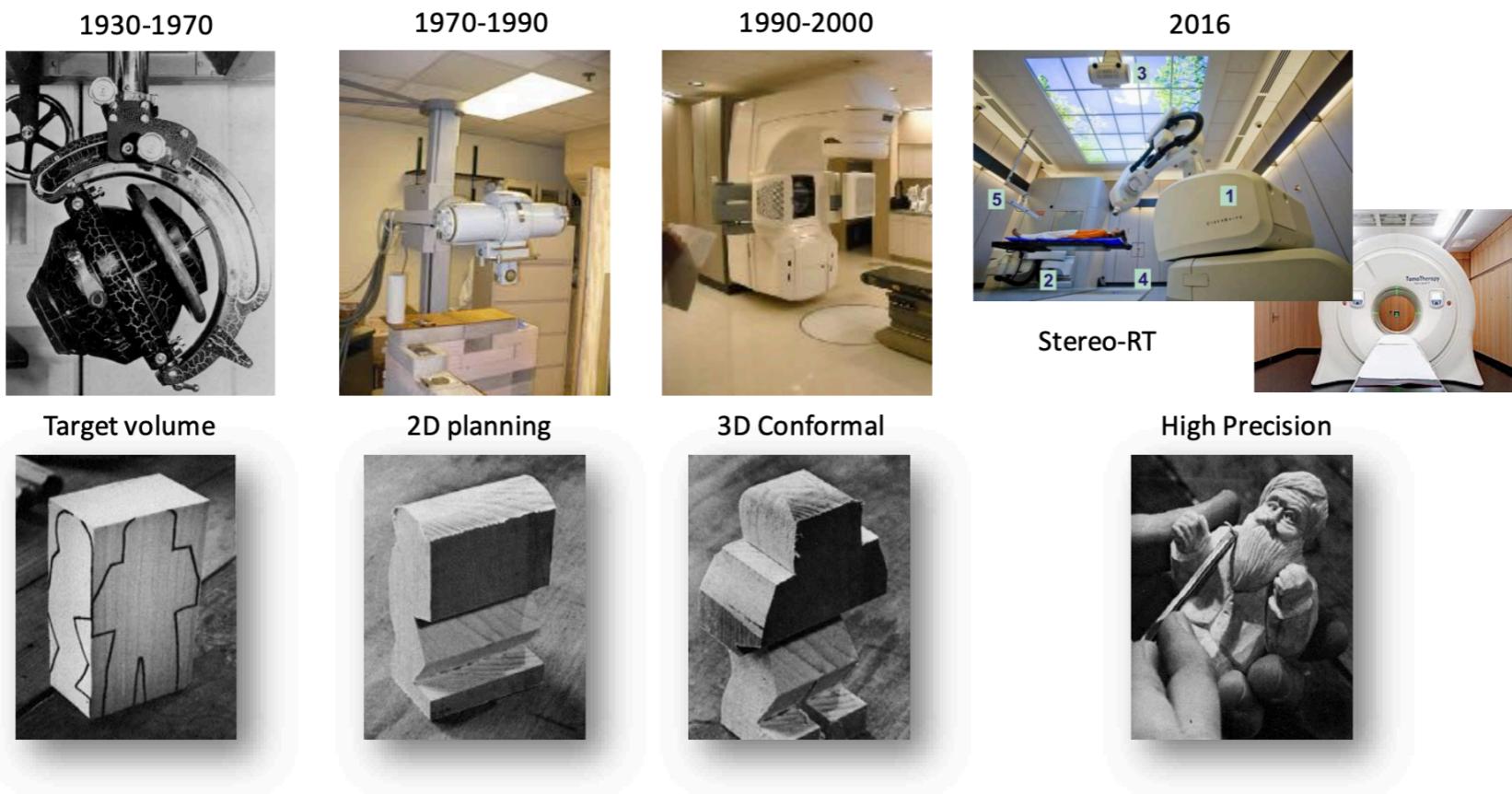


# Mayo Clinic



# Radiation Therapy

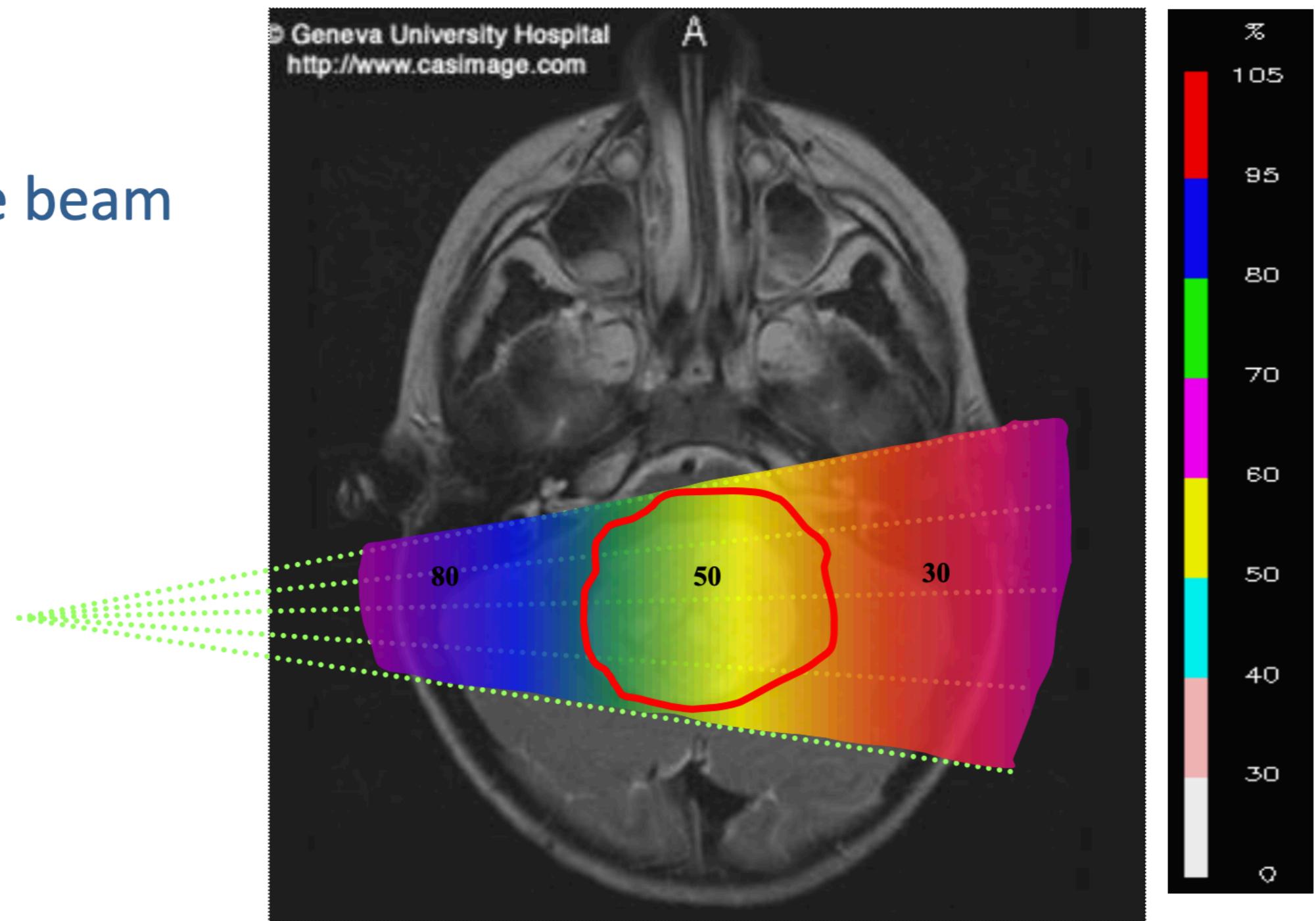
## Fractionation and Enhanced precision



- To fight Cancer
- Radiation therapy has had a long history of usage
- Radiation is sent to a tumor to kill it
- Critical when you can't cut the tumor out

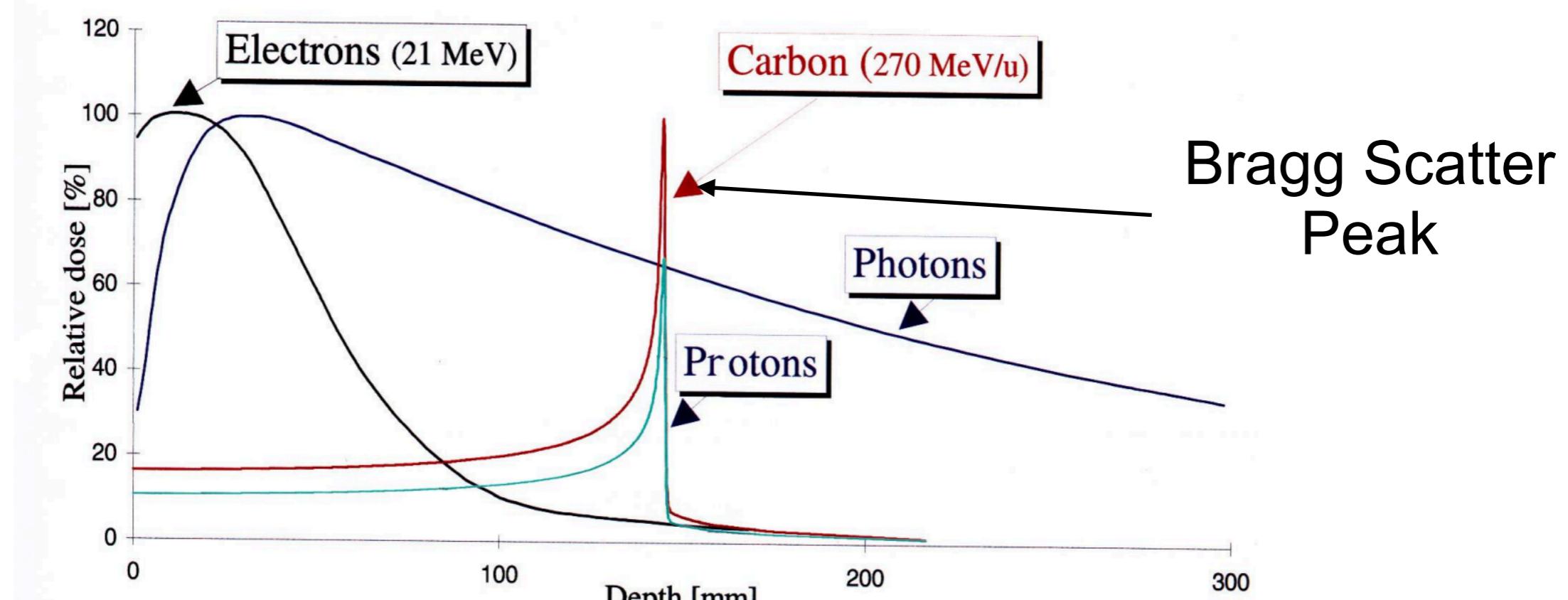
# Classical Radiotherapy with X-rays

single beam

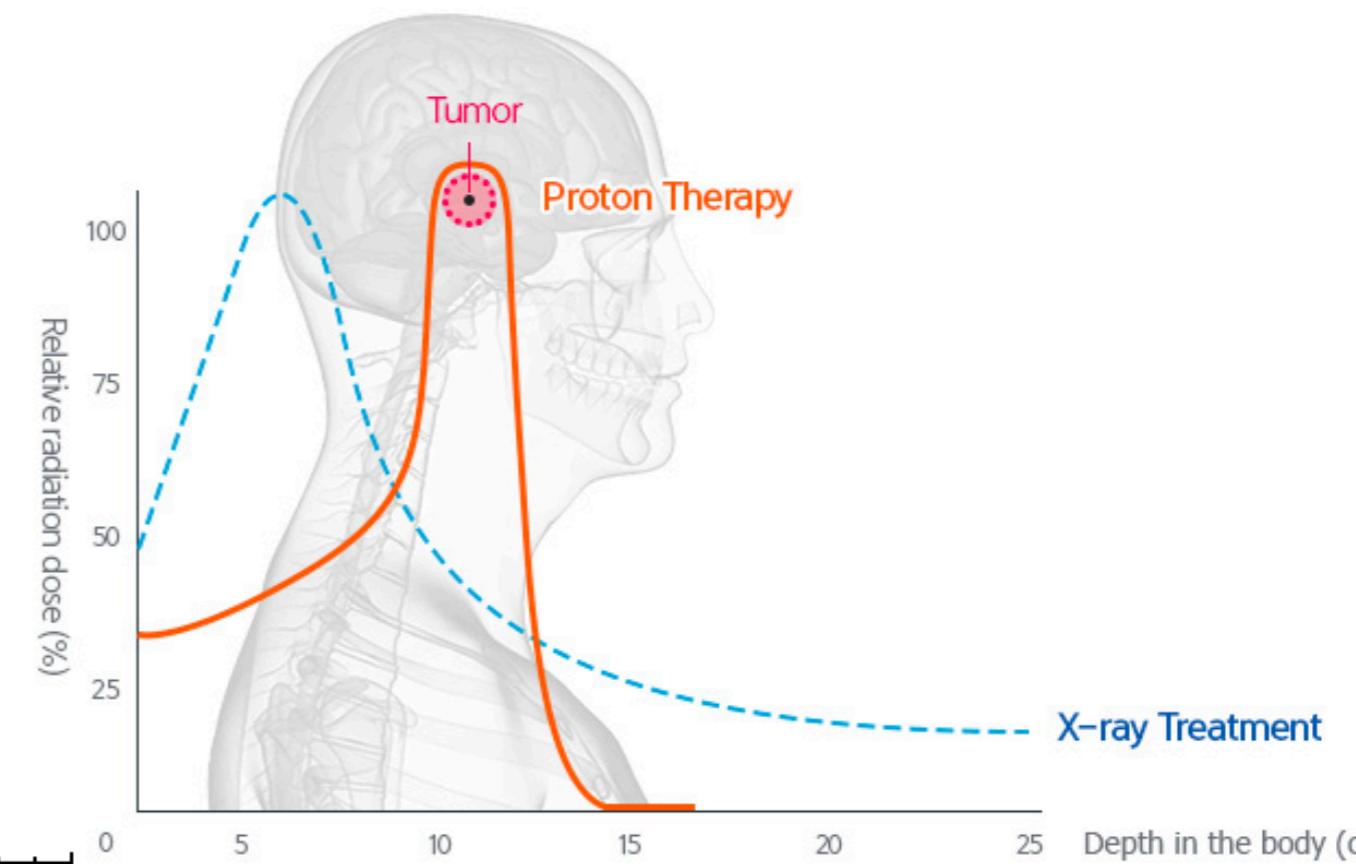
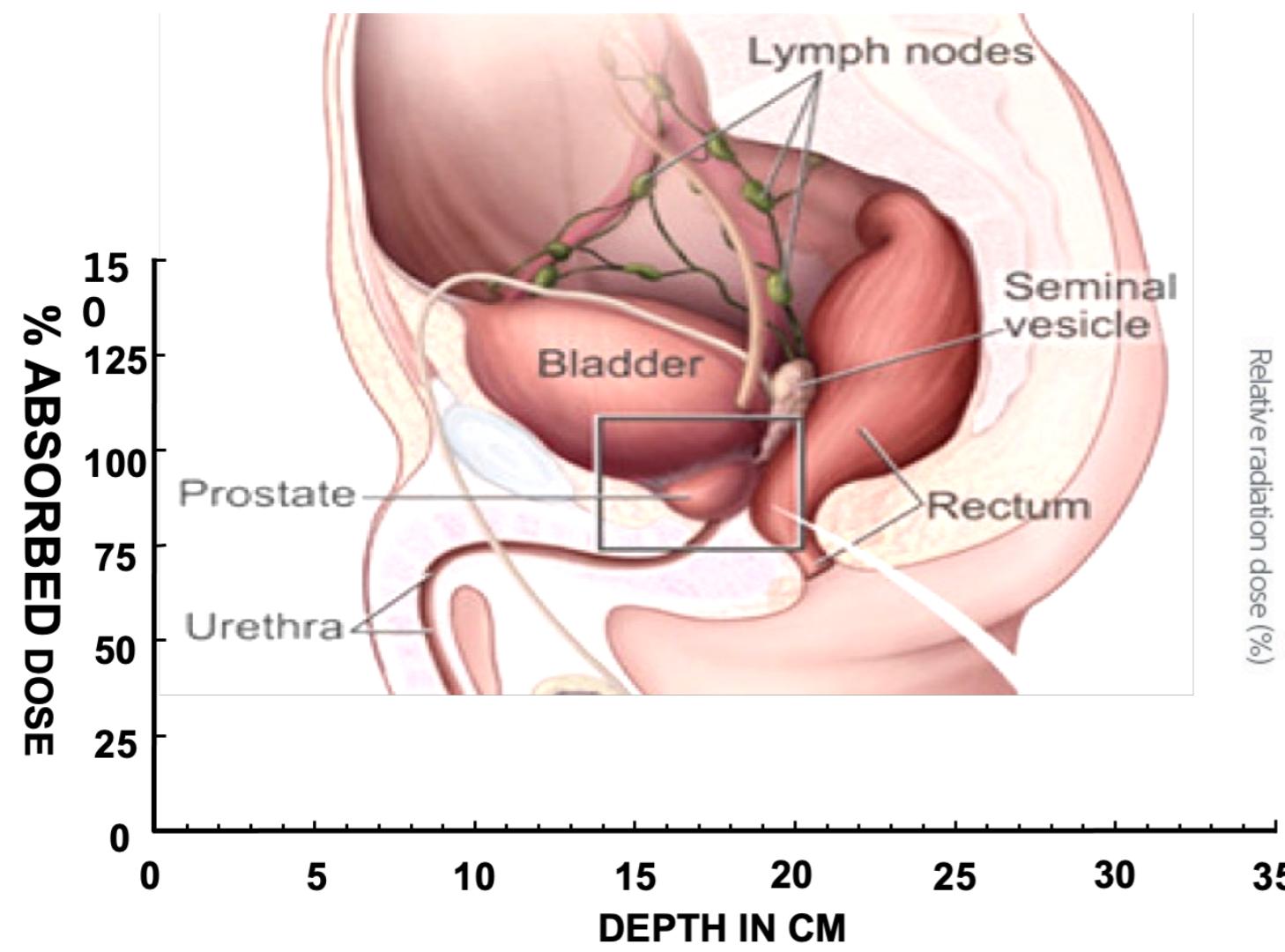


# Hadron Therapy

- Therapy
  - Hadrons allow you to control deposit
  - Can vary the depth of the hadrons through Bragg scatter

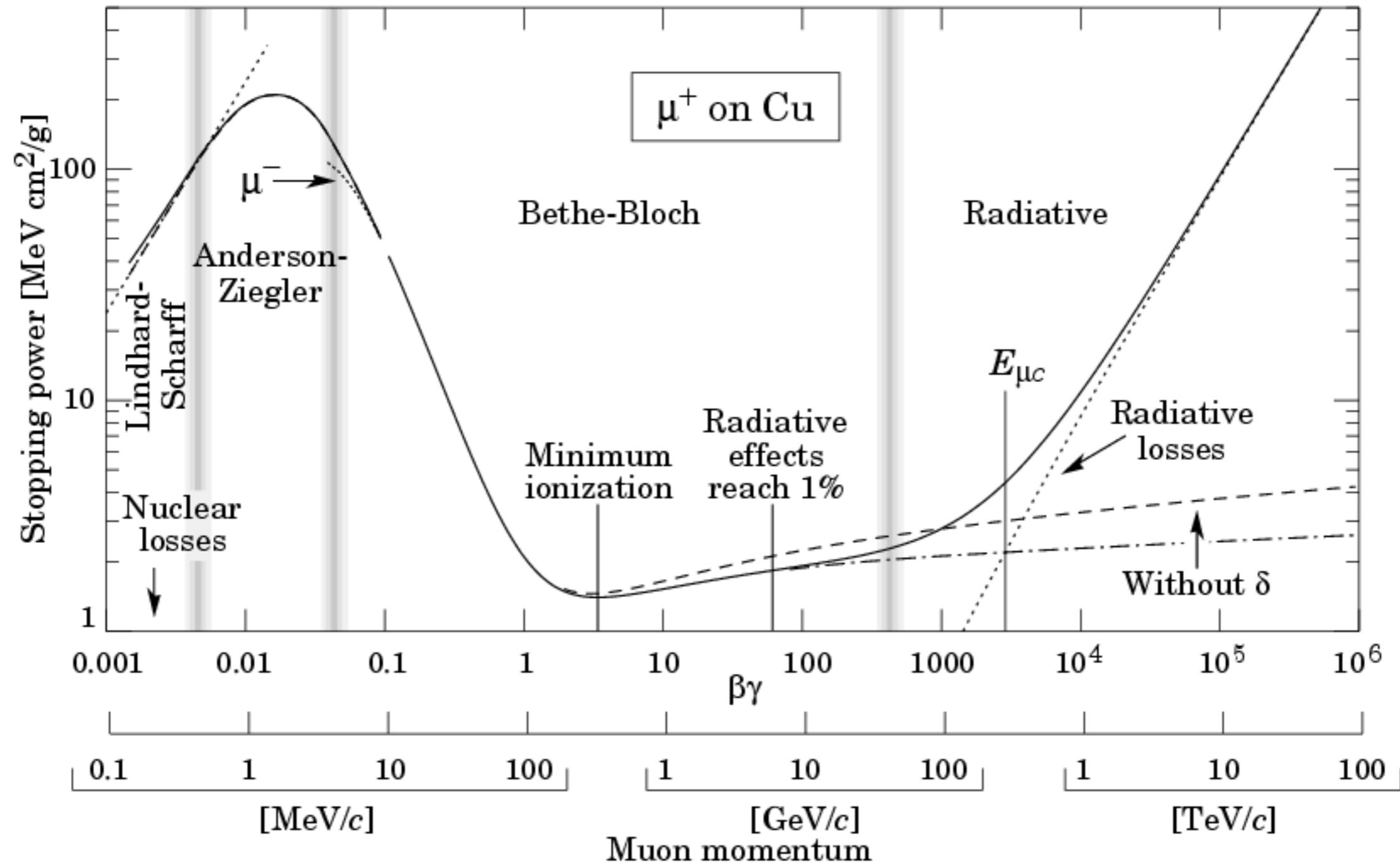


# Proton Therapy



# Bethe-Bloch Equation

- Charged Particles in matter are governed by this equation



# Protons Governed

$$-\left\langle \frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[ \frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 T_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right] [\cdot \rho]$$

density

$$K = 4\pi N_A r_e^2 m_e c^2 = 0.307 \text{ MeV g}^{-1} \text{ cm}^2$$

$$T_{\max} = 2m_e c^2 \beta^2 \gamma^2 / (1 + 2\gamma m_e/M + (m_e/M)^2)$$

[Max. energy transfer in single collision]

$z$  : Charge of incident particle

$M$  : Mass of incident particle

$Z$  : Charge number of medium

$A$  : Atomic mass of medium

$I$  : Mean excitation energy of medium

$\delta$  : Density correction [transv. extension of electric field]

$$N_A = 6.022 \cdot 10^{23}$$

[Avogardo's number]

$$r_e = e^2 / 4\pi \epsilon_0 m_e c^2 = 2.8 \text{ fm}$$

[Classical electron radius]

$$m_e = 511 \text{ keV}$$

[Electron mass]

$$\beta = v/c$$

[Velocity]

$$\gamma = (1 - \beta^2)^{-1/2}$$

[Lorentz factor]

Validity:

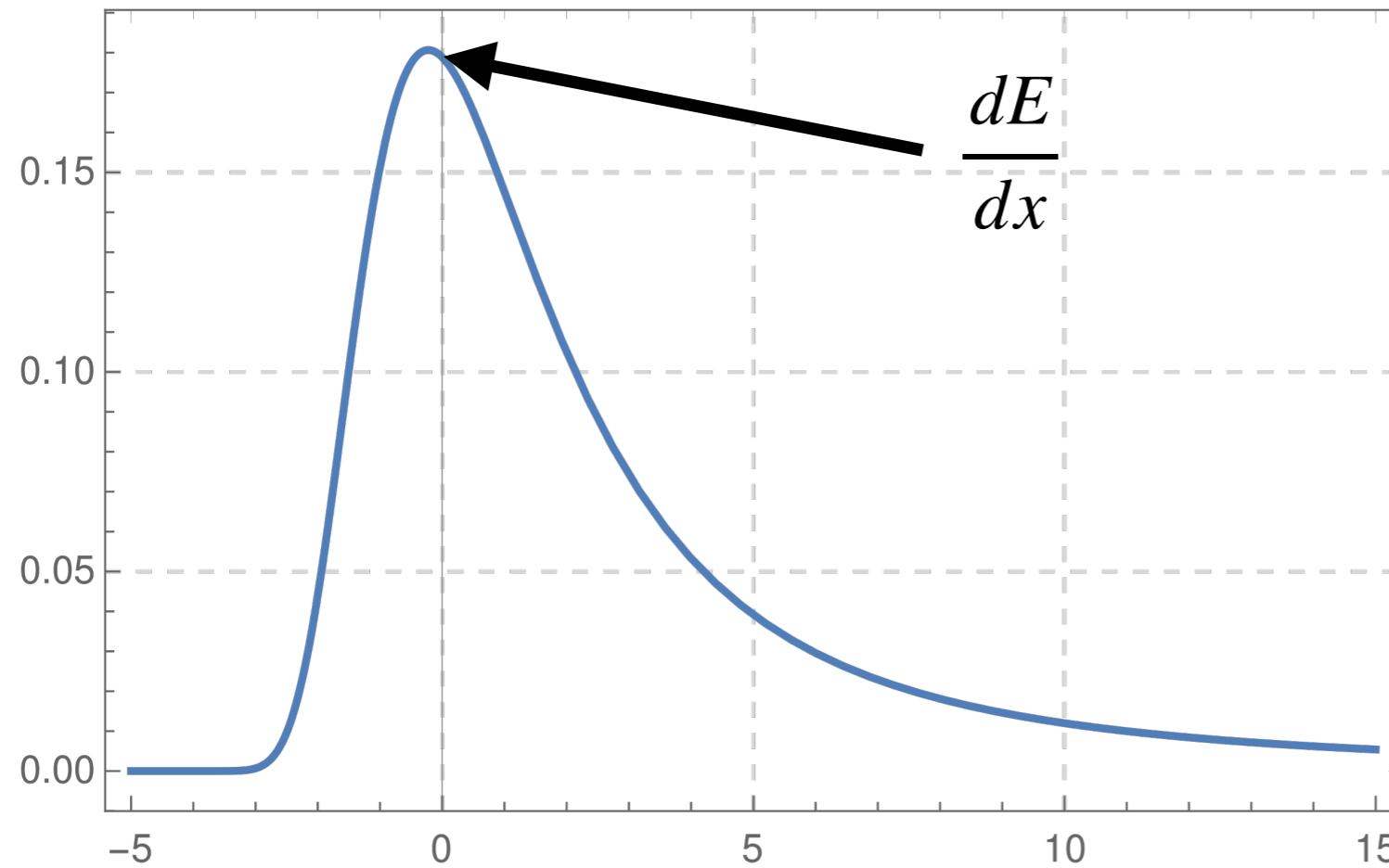
$$0.05 < \beta\gamma < 500$$

$$M > m_\mu$$

# Actual Energy Loss

- As we step along we lose energy by the Landau distribution

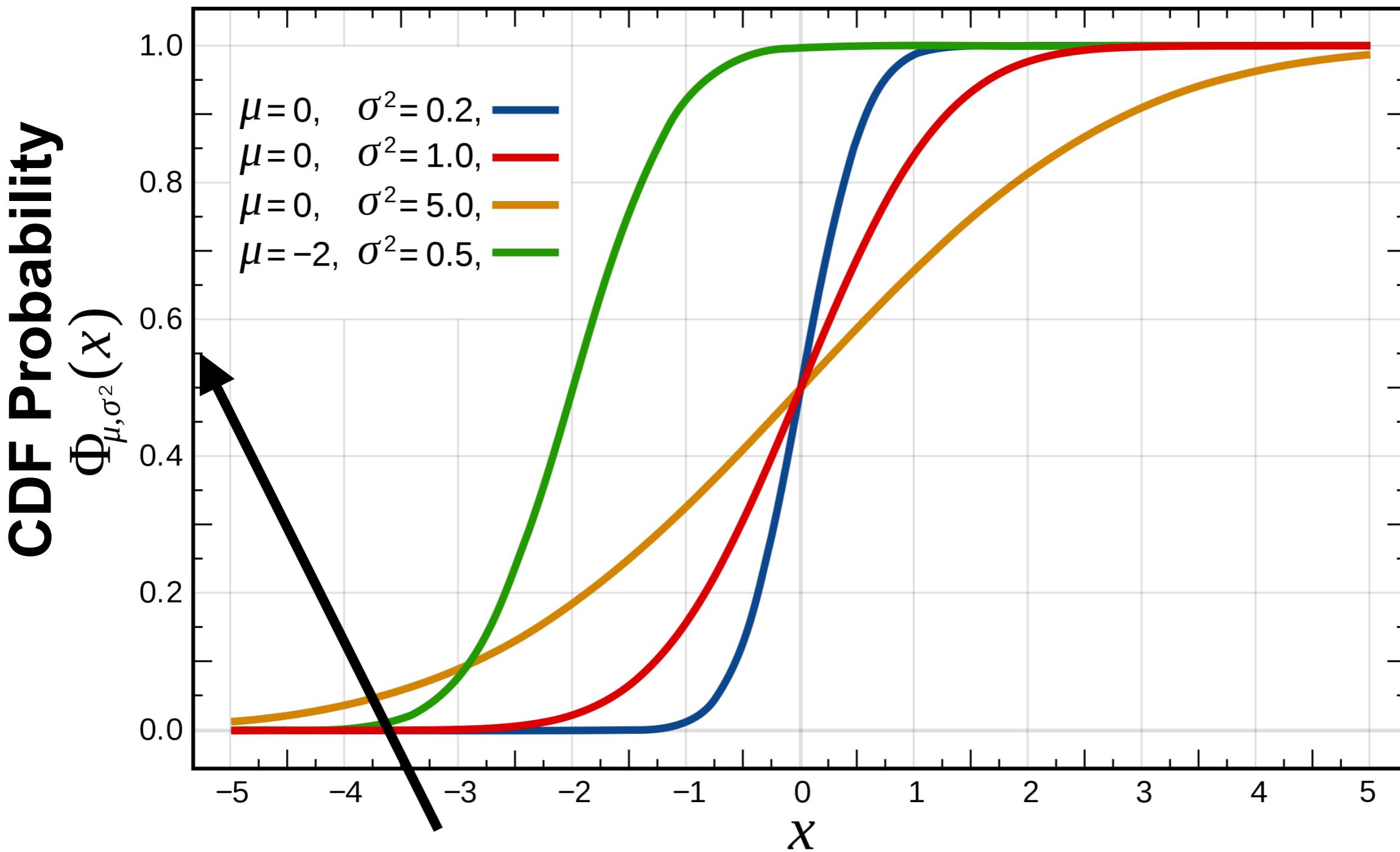
$$p(x) = \frac{1}{2\pi i} \int_{a-i\infty}^{a+i\infty} e^{s \log(s) + xs} ds,$$



Average of this distribution  
gives Bethe-Bloch

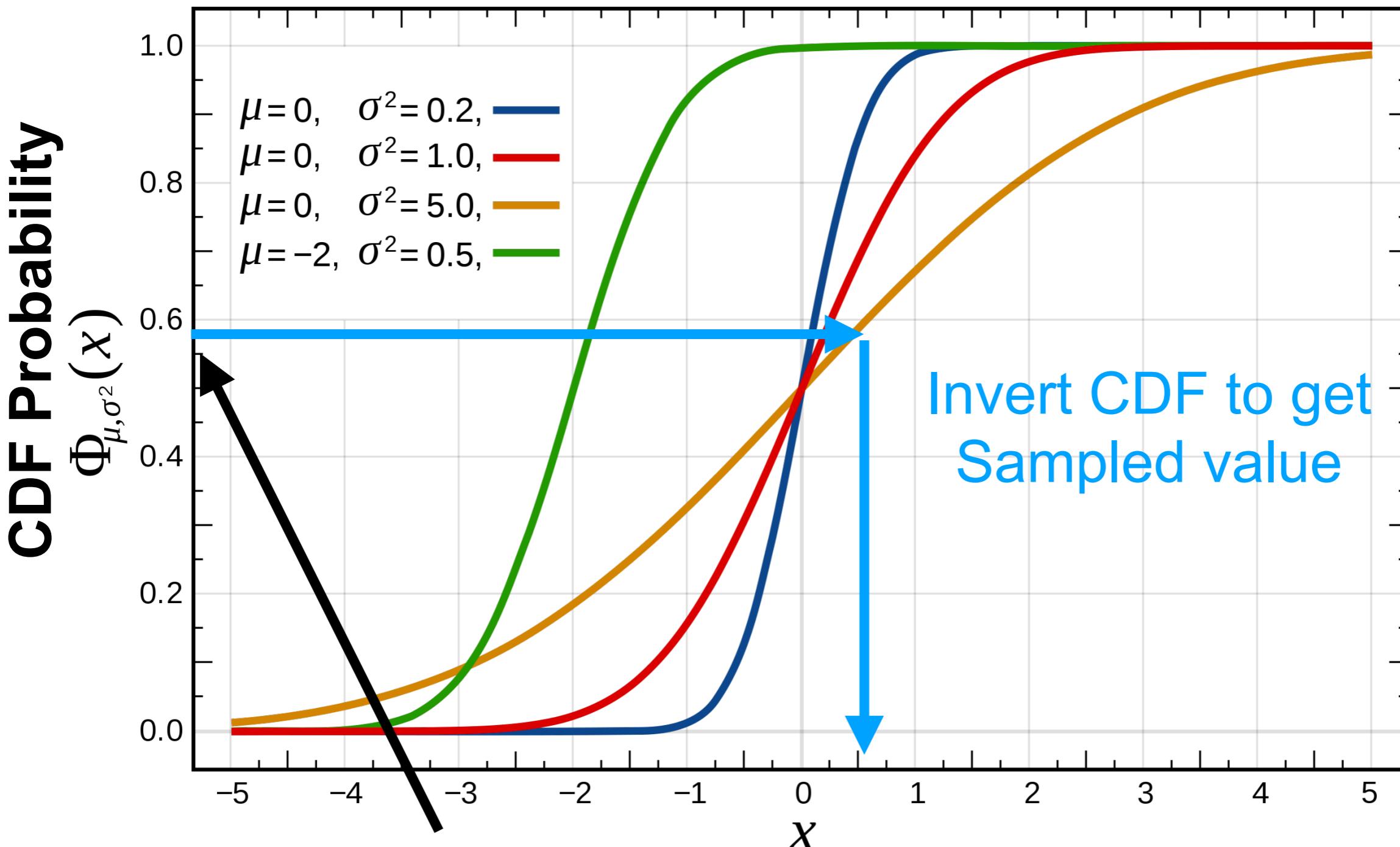
We can sample this  
At each step

# Sampling a Distribution

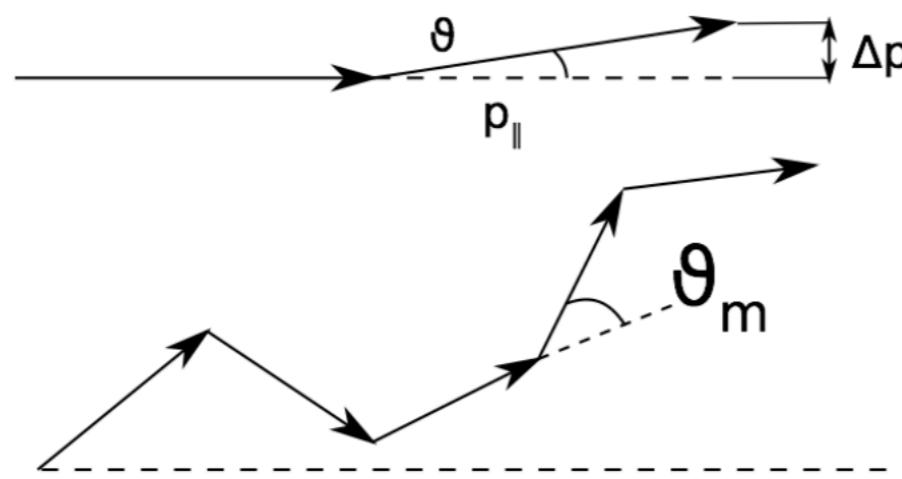


Sample from a p-value from 0 to 1 (flat 0 to 1)

# Sampling a Distribution



# Multiple Scatter Particles



after  $k$  collisions

$$\begin{aligned}\theta &\simeq \frac{\Delta p_{\perp}}{p_{\parallel}} \simeq \frac{\Delta p_{\perp}}{p} \\ &= \frac{2Zze^2}{b} \frac{1}{pv}\end{aligned}$$

$$\langle \theta_k^2 \rangle = \sum_{m=1}^k \theta_m^2 = k \langle \theta^2 \rangle$$

- Single collision (thin absorber): Rutherford scattering  $d\sigma/d\Omega \propto \sin^{-4} \theta / 2$
- Few collisions: difficult problem
- Many ( $>20$ ) collisions: statistical treatment “Molière theory”

# Multiple Scatter Particles

$$\theta \simeq \frac{\Delta p_\perp}{p} \simeq \frac{\Delta p_\perp}{p}$$

Obtain the **mean deflection angle in a plane** by averaging over many collisions and integrating over  $b$ :

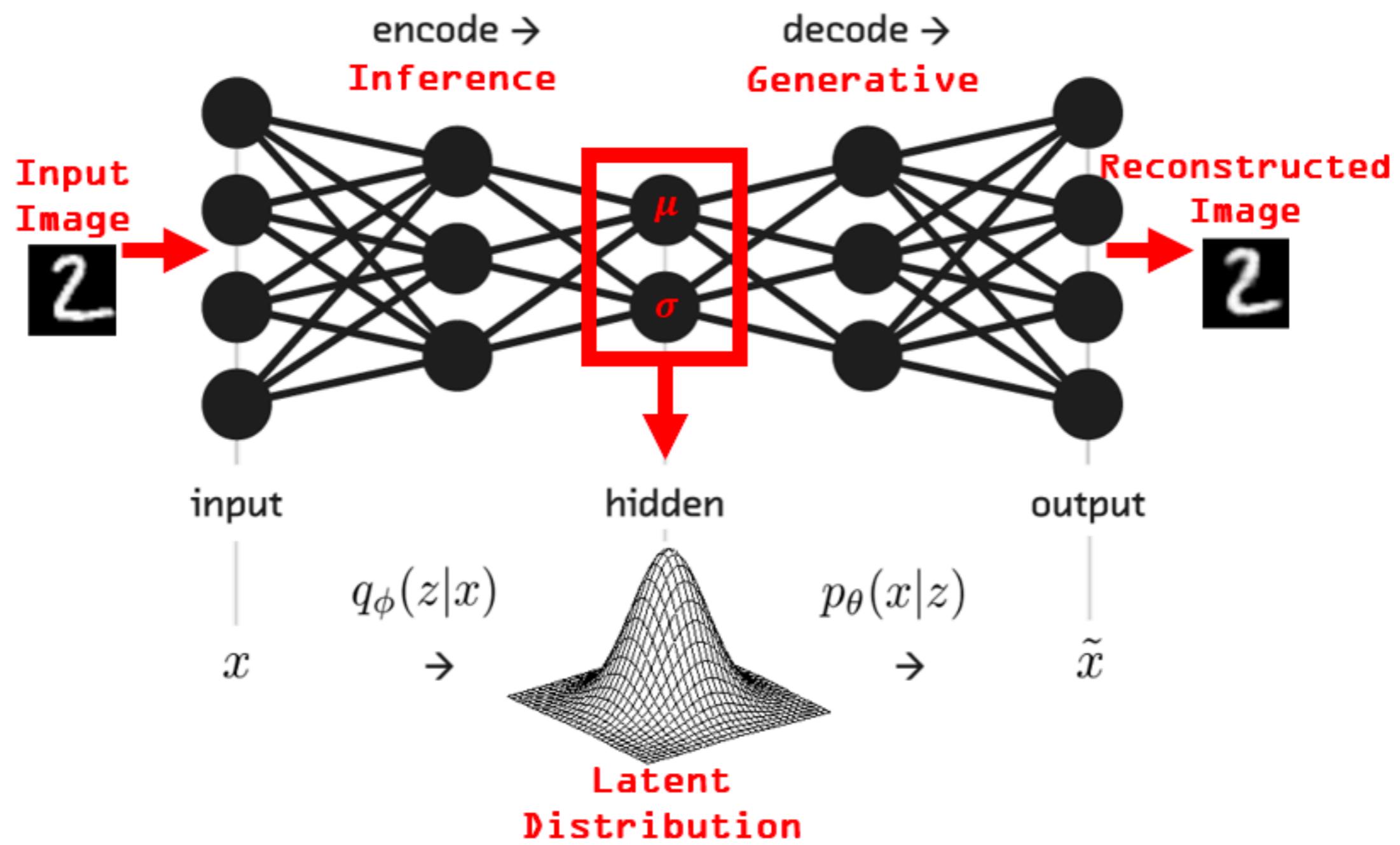
$$\sqrt{\langle \theta^2(x) \rangle} = \theta_{\text{rms}}^{\text{plane}} = \frac{13.6 \text{ MeV}}{\beta pc} z \sqrt{\frac{x}{X_0}} \left( 1 + 0.038 \ln \frac{x}{X_0} \right)$$

- Material constant  $X_0$ : radiation length
- $\propto \sqrt{x} \rightarrow$  use thin detectors
- $\propto 1/\sqrt{X_0} \rightarrow$  use light detectors
- $\propto 1/\beta p \rightarrow$  serious problem at low momenta

In 3 dimensions:  $\theta_{\text{rms}}^{\text{space}} = \sqrt{2} \theta_{\text{rms}}^{\text{plane}}$        $13.6 \rightarrow 19.2$

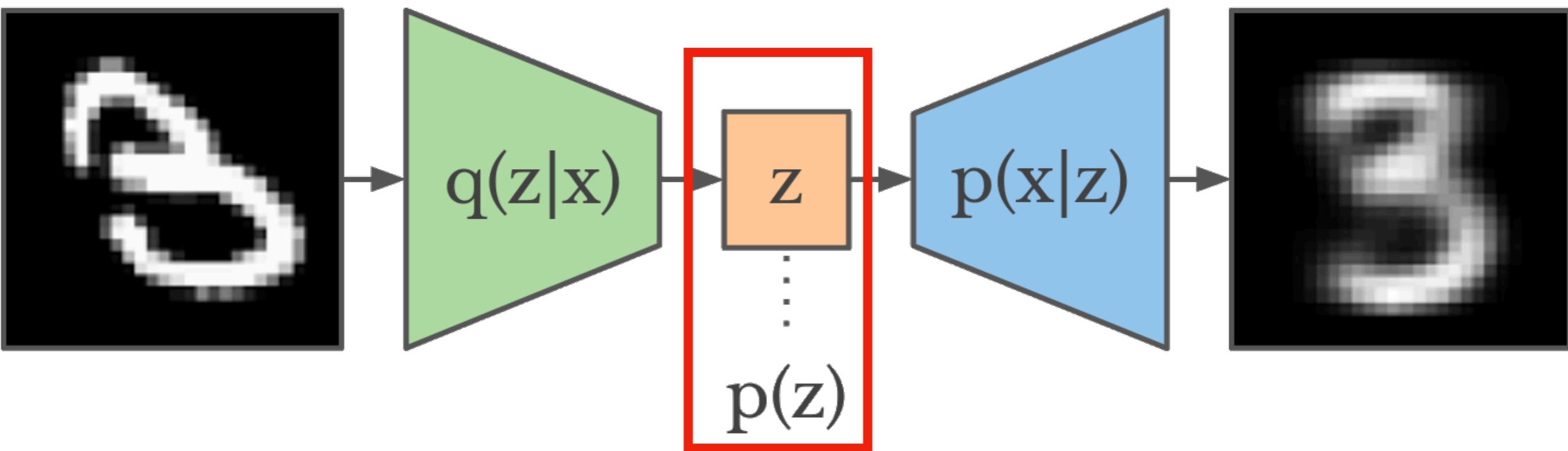
# VAE

- Variational Autoencoder is a great way to model objects

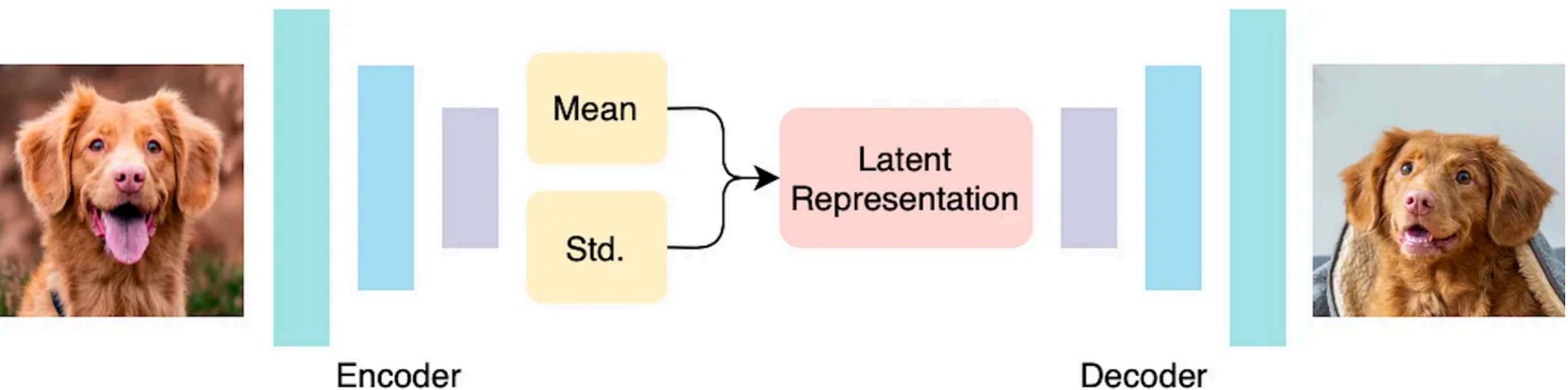


# VAE

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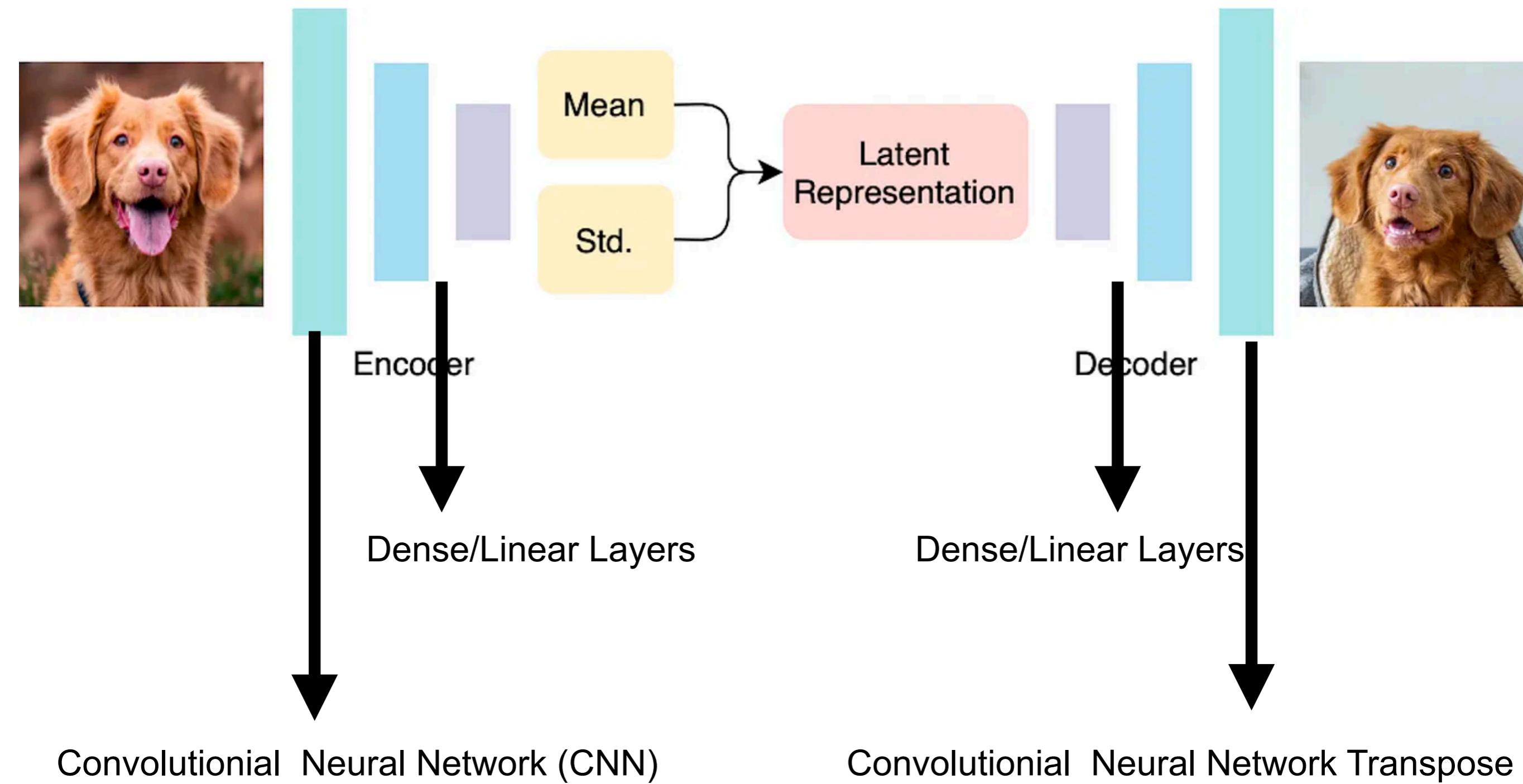


Randomly sample a normal distribution in this space

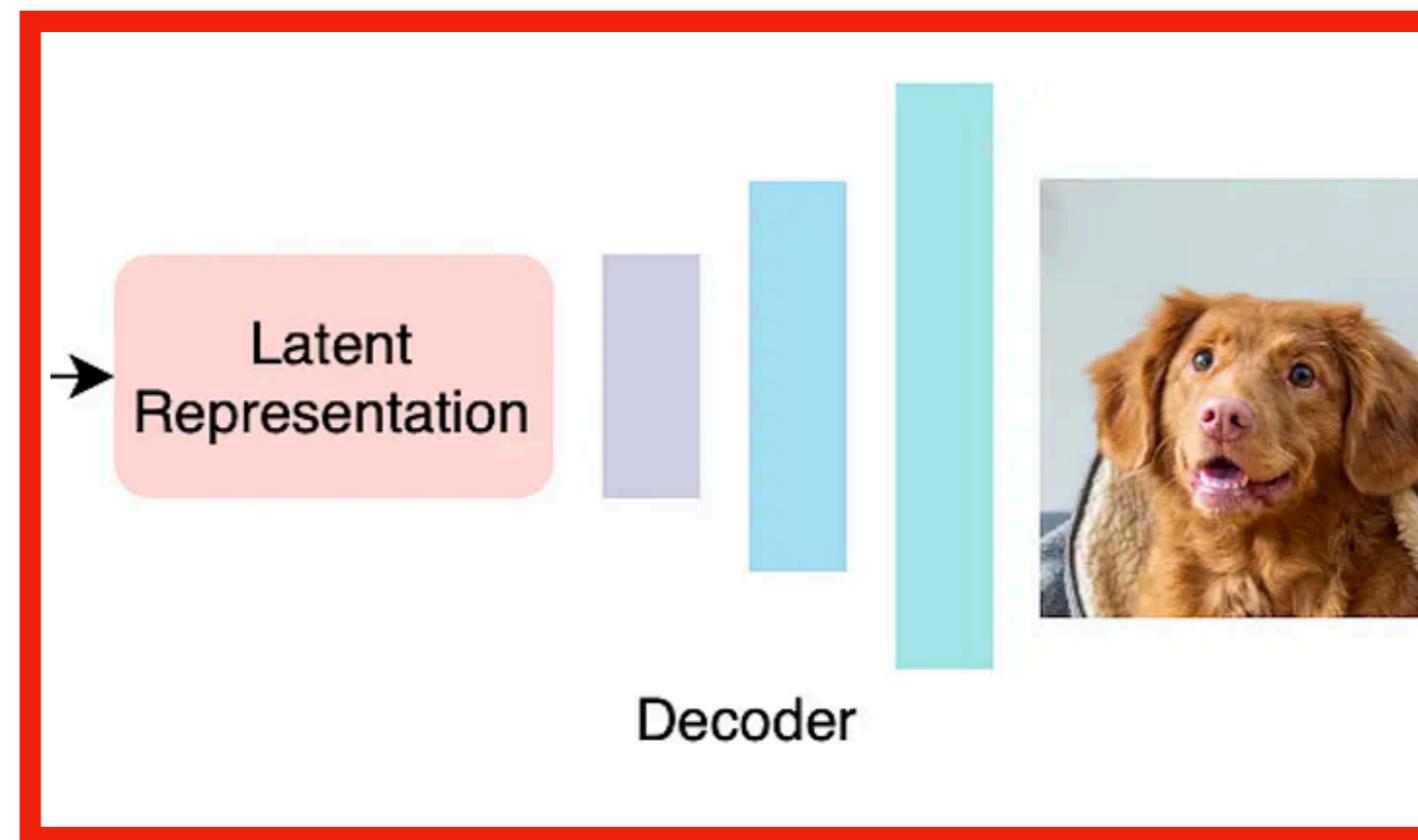
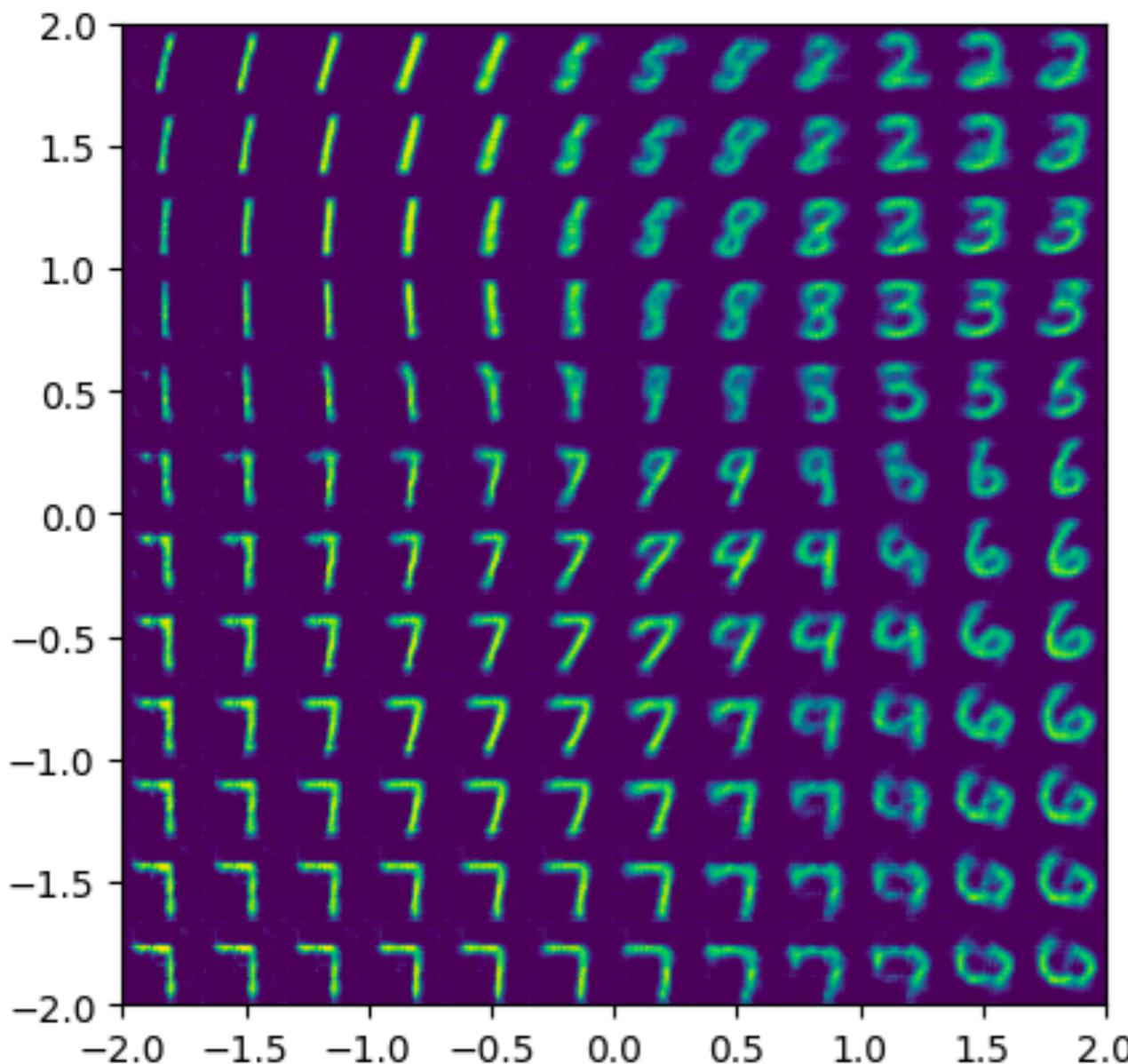


# MNIST VAE encoder

- We will use a CNN to encode the data and process it

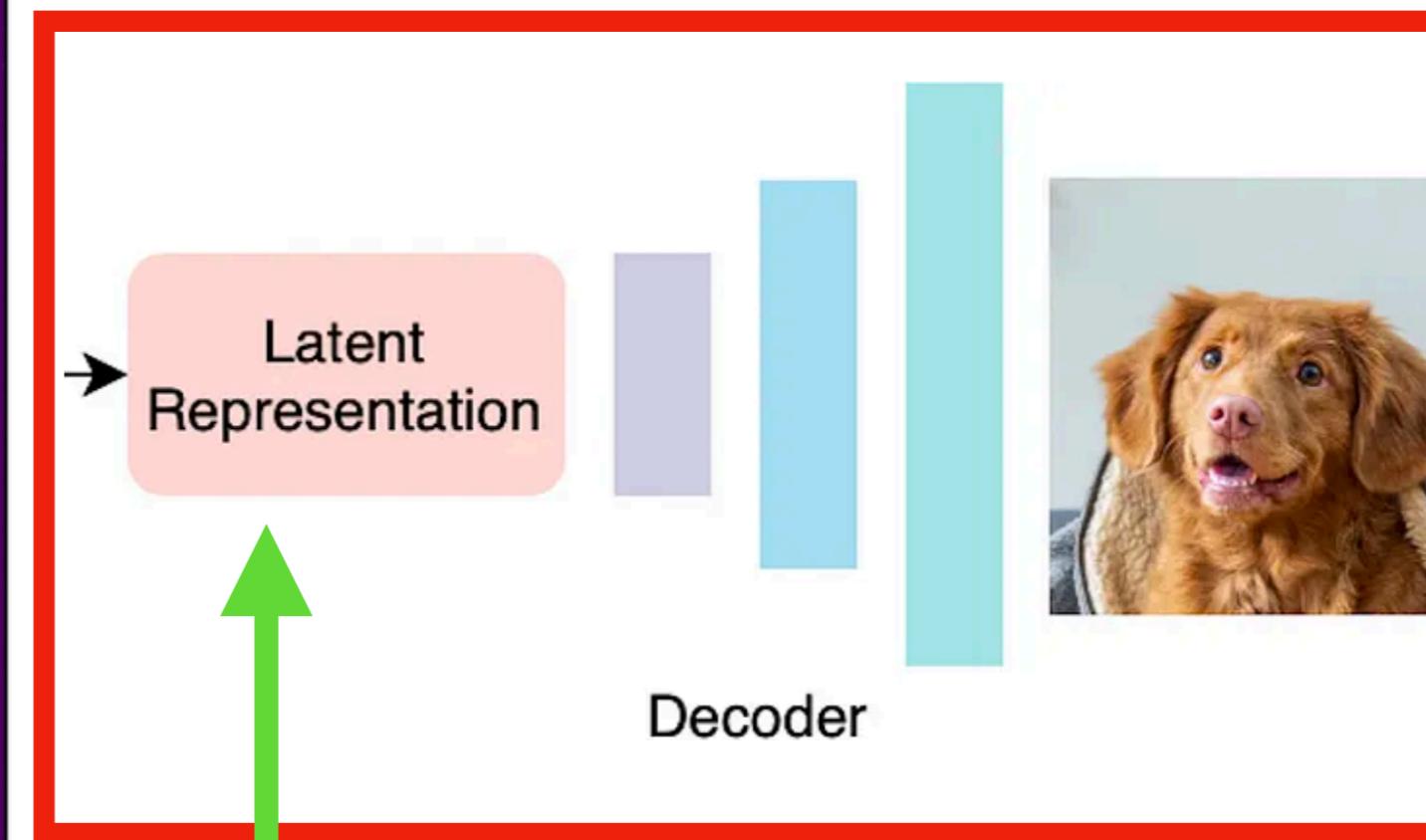
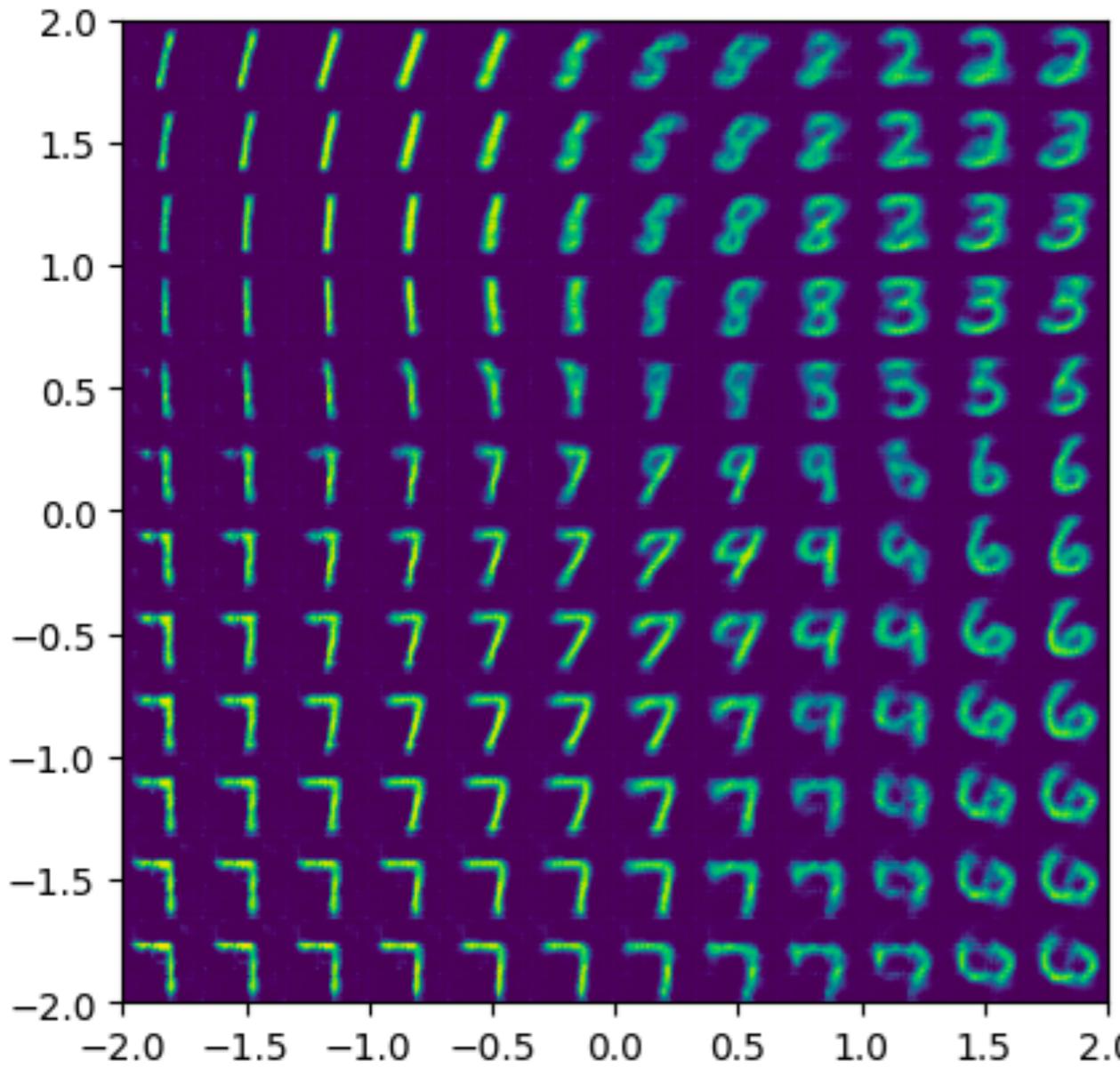


# Exploring the latent<sup>43</sup> space?



- We can sample the latent space as a generator

# Conditional VAE



**Generator**

Variables that we would like to condition on

- Force known inputs into the VAE
  - That way our latent space has explicit knowledge of what is going on