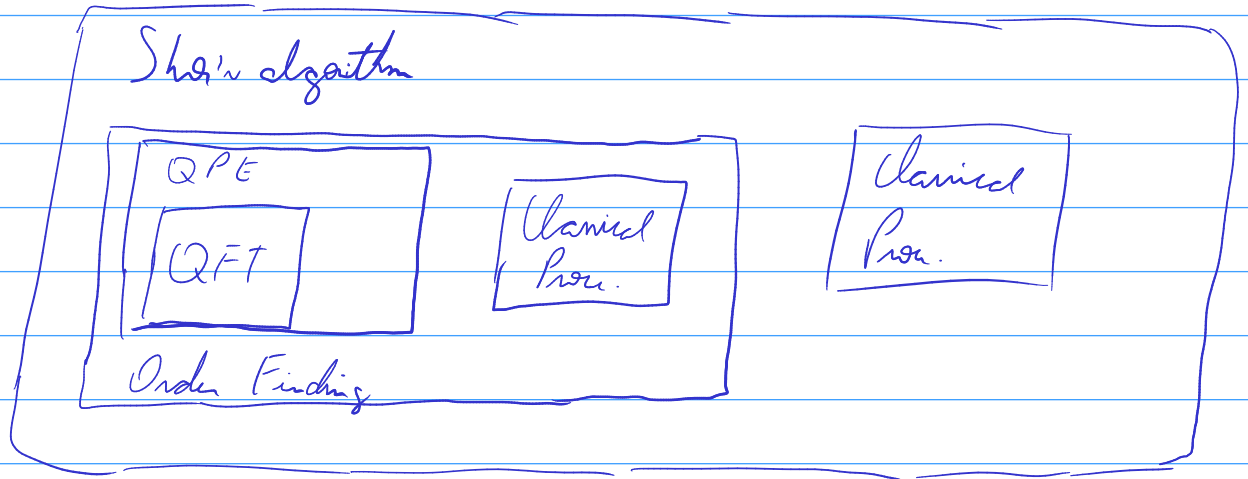


Shor's algorithm

→ Take an integer N and returns its prime factors



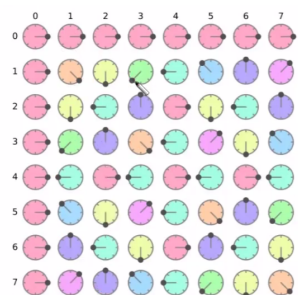
Discrete Fourier Transformation (DFT)

$$\frac{1}{\sqrt{N}} \begin{pmatrix} 1 & 1 & 1 & 1 & \dots & 1 \\ 1 & \omega & \omega^2 & \omega^3 & \dots & \omega^{N-1} \\ 1 & \omega^2 & \omega^4 & \omega^6 & \dots & \omega^{2(N-1)} \\ 1 & \omega^3 & \omega^6 & \omega^9 & \dots & \omega^{3(N-1)} \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & \omega^{N-1} & \omega^{2(N-1)} & \omega^{3(N-1)} & \dots & \omega^{(N-1)(N-1)} \end{pmatrix}$$

$$\omega = e^{i\frac{2\pi}{N}} \leftarrow \text{some rotation}$$

$$\text{DFT}_{jk} = \frac{\omega^{jk}}{\sqrt{N}}$$

← it takes a input vector $\vec{x} = (x_0, x_1, \dots, x_{N-1})$
 $x_i \in \mathbb{C}$ and returns $\text{DFT}(\vec{x})$



← a rotation matrix with $N=8$

→ Using $\vec{y} = \text{DFT } \vec{v}$ you find some periods
in \vec{v}

$$\vec{v} = (1, 0, 0, 1, 0, 0, 1, 0, 0)$$

period $r = 3$
 $N = 9$

↑ if we use DFT here, will get some peaks
at multiples of $\frac{N}{r}$

if we have $\vec{v} \rightarrow$ period r_1
 $\vec{w} \rightarrow$ period r_2

DFT $(\vec{v} + \vec{w})$ will picture peaks at multiple
of $\frac{N}{r_1}$ and $\frac{N}{r_2}$

QFT

→ DFT on quantum states
 $|\psi\rangle \xrightarrow{\text{QFT}} \text{DFT } |\psi\rangle$

$$|\psi\rangle = \sum_{j=0}^{N-1} x_j |j\rangle$$

$$\text{DFT } |\psi\rangle \rightarrow \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} \sum_{j=0}^{N-1} x_j \omega^{kj} |k\rangle$$

→ QFT is exponentially faster than DFT

Binary expansion of integers

→ Transform an int. to a sum of powers of 2

$$9 = 1 \cdot 2^3 + 0 \cdot 2^2 + 0 \cdot 2^1 + 1 \cdot 2^0 = 1001$$
$$0,75 = \frac{1}{2} + \frac{1}{4} = 1 \cdot 2^{-1} + 1 \cdot 2^{-2} = 0,11$$

$$9 + 0,75 = 1001,11$$

$$2 \cdot 1001,11 \rightarrow 1 \cdot 2^4 + 0 \cdot 2^3 + 0 \cdot 2^2 + 1 \cdot 2^1 + 1 \cdot 2^0 + 1 \cdot 2^{-1}$$
$$\rightarrow 10011,1 \leftarrow \text{shift left}$$

Integer binary:

$$i = 2^n \sum_{l=1}^n 2^{-l} i_l$$

↳ the biggest factor is being mult. by the index l

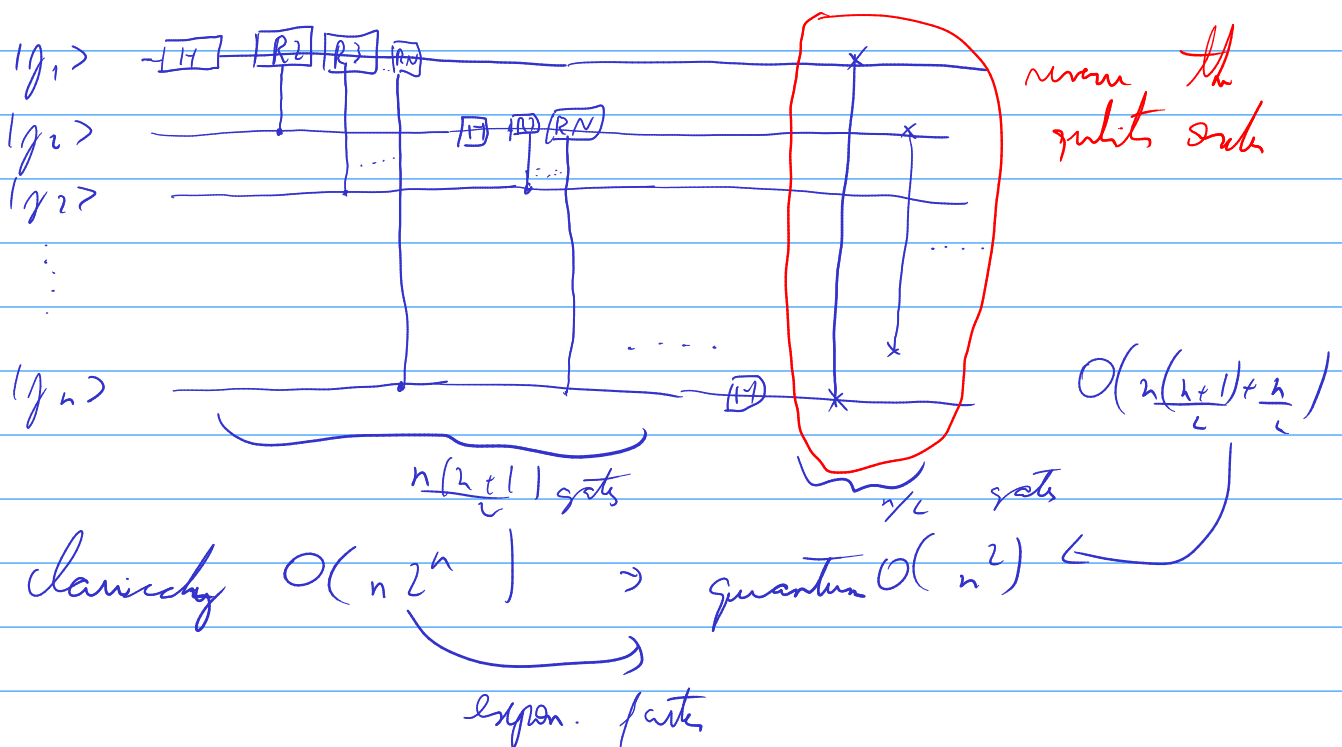
QFT on basis states

$$\text{QFT} |j_1 j_2 \dots j_n\rangle = \frac{1}{2^{n/2}} \left((|0\rangle + e^{2\pi i j_1 / 2} |1\rangle) \otimes (|0\rangle + e^{2\pi i j_2 / 2} |1\rangle) \otimes \dots \right)$$

QFT circuit

$$\text{gate } R_k = \begin{pmatrix} 1 & 0 \\ 0 & e^{2\pi i / 2^k} \end{pmatrix}$$

$$CR_k = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & e^{2\pi i / 2^k} \end{pmatrix}$$



Inverse QFT

- QFT is unitary so $\text{QFT} \text{QFT}^\dagger = I$
- run QFT backwards
- get the period, invert and invert the input vector

QPE

→ any U has an eigenvector

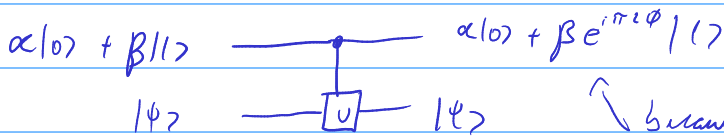
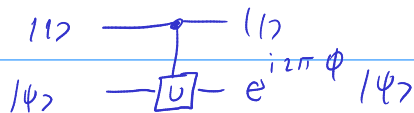
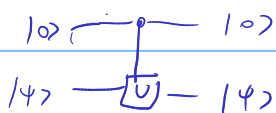
$$U|\lambda\rangle = \lambda|\lambda\rangle$$

$$\lambda_k = e^{i\theta} \Leftrightarrow e^{i2\pi\phi_k}$$

$$0 \leq \phi_k < 1$$

→ take an U and an eigenvector of $U|\psi\rangle \rightarrow U|\psi\rangle = e^{i2\pi\phi}|\psi\rangle$ it returns ϕ

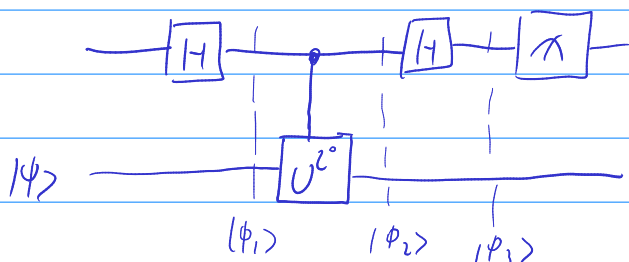
idea behind QPE



because

$$(\alpha|0\rangle + \beta|1\rangle)|\psi\rangle \xrightarrow{U} \alpha|0\rangle|\psi\rangle + \beta e^{i2\pi\phi}|1\rangle|\psi\rangle$$

$|\psi\rangle$ is not affected



$$|\phi_1\rangle = |+\rangle|\psi\rangle \quad |\phi_2\rangle = \frac{1}{\sqrt{2}}(|0\rangle + e^{i2\pi\phi}|1\rangle)|\psi\rangle$$

$$|\phi_3\rangle = \frac{1}{\sqrt{2}}(|+\rangle + e^{i\pi\phi}|-\rangle)|\psi\rangle$$

$$\frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) + e^{i2\pi\phi} \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle) \right) |\psi\rangle$$

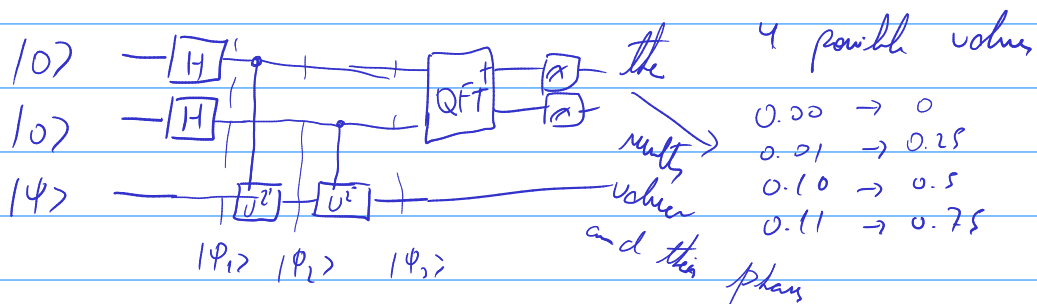
$$\frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{2}} (|0\rangle + |1\rangle + e^{i2\pi\phi}(|0\rangle - |1\rangle)) \right) |\psi\rangle$$

$$\frac{1}{2} (|0\rangle + |1\rangle + e^{i2\pi\phi}|0\rangle - e^{i2\pi\phi}|1\rangle) |\psi\rangle$$

$$\frac{1}{2} ((1 + e^{i2\pi\phi})|0\rangle + (1 - e^{i2\pi\phi})|1\rangle) |\psi\rangle$$

if $\phi = 0 \rightarrow |0\rangle|\psi\rangle$
 $\phi = \frac{1}{2} \rightarrow |1\rangle|\psi\rangle$

the only 2 possible values for 1 bit
 $0.0 \rightarrow 0$
 $0.1 \rightarrow \frac{1}{2} (2^{-1})$



$$|\phi_1\rangle = |+\rangle|+\rangle|\psi\rangle$$

$$|\phi_2\rangle = \left(\frac{1}{\sqrt{2}} (|0\rangle + e^{i4\pi\phi}|1\rangle) \right) |+\rangle|\psi\rangle$$

$$|\phi_3\rangle = (|0\rangle + e^{i4\pi\phi}|1\rangle) (|0\rangle + e^{i2\pi\phi}|1\rangle) |\psi\rangle$$

