

Constraints on Beyond Standard Model vector boson interactions in an effective field theory using differential cross section of the production of two jets associate with four charged leptons at $\sqrt{s} = 13$ TeV with the ATLAS detector

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ABSTRACT: In this report, the production of a pair on-shell Z bosons associate with a pair of jets in an electroweak enhanced phase space region was measured at centre-of-mass energy $\sqrt{s} = 13$ TeV with 139 fb^{-1} of pp collision data collected by the ATLAS detector. The result is compatible with the Standard Model prediction. The differential cross sections as functions of two CP-sensitive observables were obtained by removing the detector effect from the measurement. The differential measurements were then reinterpreted in a dimension-6 Standard Model Effect Field Theory and constraints on relevant CP-odd and CP-even new physics operators were obtained.

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1 Introduction

The Standard Model of particle physics (SM)[1] has been a very successful theory in predicting the nature of fundamental particles and their interactions. However, the SM is not yet complete, for example, the enormous matter-antimatter asymmetry[2] observed in the universe is not fully described by the SM. Being sensitive to Beyond the SM (BSM) physics is one of the main goals of the Large Hadron Collider (LHC)[3]. At current energy scale, the BSM effects may manifest themselves as small corrections to the SM, searching for deviations from the SM prediction; therefore a general way to search the new physics phenomena.

Since the discovery of the last particle predicted by the SM, the Higgs boson, in 2012[4], the nature of its interactions with other particles has been carefully investigated in order to seek any deviations from the SM. The Vector Boson Scattering (VBS) process[5], an important pure electroweak process in the LHC, enables us to probe the nature of Higgs boson interactions with vector bosons and to constrain BSM hypothesis altering the nature of the Higgs boson and vector bosons which may

provide alternative Electroweak Symmetry Breaking (EWSB) mechanism[6] and additional sources of CP violation[7].

In the LHC, VBS events are produced by a pair of vector bosons radiated from the colliding partons scattered to another pair of vector bosons subsequently. At the detector-level, the signature of VBS events are two jets from the initial partons associate with the decay products of two vector boson, and we denote such events VVjj. The evidence of electroweak production of $W^\pm W^\pm jj$ (EW $W^\pm W^\pm jj$) was first observed during the first run of LHC with the ATLAS detector[8]; and the observation of this very process has been confirmed by the CMS[9] collaboration during the second run of LHC in 2018[10], while the observation in the EW ZZjj channel was reported in 2019[11]. In spite the endeavour has been put on the VBS process, the model-independent measurement of the differential cross section has not been made so far.

With a total of 139 fb^{-1} proton proton collision data at $\sqrt{s} = 13\text{ TeV}$ accumulated in the ATLAS detector[12] during the second run of the LHC, a great opportunity to scrutinise the SM at much higher accuracy has been provided. In this report, the fiducial differential cross section of ZZjj process measured at $\sqrt{s} = 13\text{ TeV}$ using the whole run II data is reported and the results has been used to constrain a BSM effective field theory[13] which provides alternative Higgs boson interactions, vector boson interactions and additional source of CP violation in the electroweak sector.

2 Recap

In the last semester[14], an iterative unfolding procedure to remove the detector effect[15] was performed on the data corresponding to an integrated luminosity of 36 fb^{-1} using the ATLAS detector in a VBS enhanced fiducial phase space. Three distributions which are possibly sensitive to the BSM physics effect were unfolded, namely $\Delta\phi_{jj}$, M_{jj} and M_{4l} , detailed discussion on these observables can be found in Ref. [14]. The unfolding procedure also gives the covariance matrices which take the bin to bin migration into account. With these infrastructures in hand, the differential cross section in various distributions can be easily measured; and the constraints on BSM physics theories can be set using the covariance matrices and the corresponding differential cross section, the detail of which will be discussed in section 4.3.

3 ZZjj differential cross section

3.1 ZZjj processes and backgrounds

The ATLAS detector in the LHC is a layered designed general-purpose particle detector, the technical detail of which can be found in Ref. [12]. The detector signatures of VVjj events in the ATLAS detector have a variety of patterns, which depend on

the kind of final state vector bosons, namely W or Z boson; and the decay channel of them[16]. ZZjj events with both Z bosons decay leptonically are denoted ZZjj4l, which is the selected channel in this analysis. Although the branching ratio for both Z bosons decay leptonically is around 10^{-3} [17], the channel has the minimal hadronic background and no missing momenta.

At the leading order, the ZZjj4l events can be separated to two categories, the electroweak production with six electroweak vertices (EW ZZjj4l), and the QCD production including two strong-interaction vertices and four electroweak vertices (QCD ZZjj4l), at the current energy scale, the QCD ZZjj4l production has a cross section approximately 100 times larger than the cross section of EW production, which is around 1 fb. The QCD ZZjj4l can be further separated by the initial species of initial partons. If both initial partons are gluon, the process is classified as ggZZ4l, otherwise classified as qqZZ4l. Some example leading order Feynman diagrams of these processes are shown in Fig. 1.

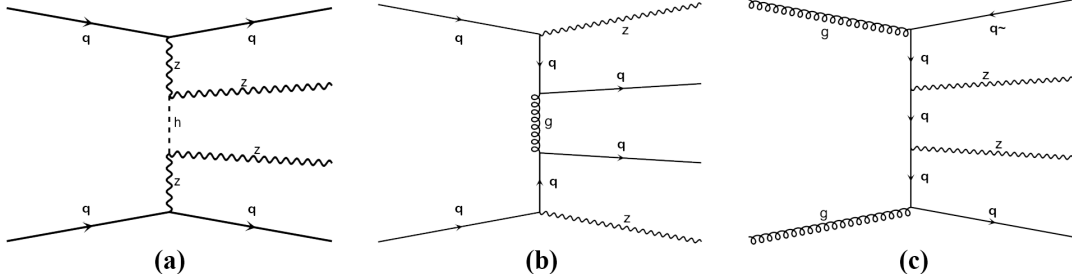


Figure 1. The leading order SM Feynman diagrams contributing to ZZjj processes. (a) EW ZZjj process, (b) QCD qqZZjj process and (c) is QCD ggZZjj process.

In addition to the large QCD ZZjj background, other none-ZZjj processes may produce the same or similar final states as the VBS processes; these background processes, therefore, mimic the detector signature of VBS process. none-ZZjj background processes considered in this analysis are the production of one vector boson plus a Higgs boson, the Higgs boson then decays to four charged leptons (VH4l); the production of three vector bosons (VVV); and the production of two charged leptons and a pair of $t\bar{t}$ quarks (ttll), the decay of the top quark via W bosons give birth to another pair of charged leptons. All three backgrounds are order six pure EW processes at leading order, which is the same order as VBS process. The example Feynman diagrams of these background processes at leading order are shown in Fig. 2. Contribution of other processes in the studied phase space were neglected.

Despite the extreme rarity and enormous background of EW ZZjj4l process, there exist region in the phase space where the EW ZZjj4l is enhanced while other background processes are suppressed. With $36fb^{-1}$ of data, the significance of the EW ZZjj4l process was measured to be 1.5σ . The detailed discussion can be found in the first-semester report[14].

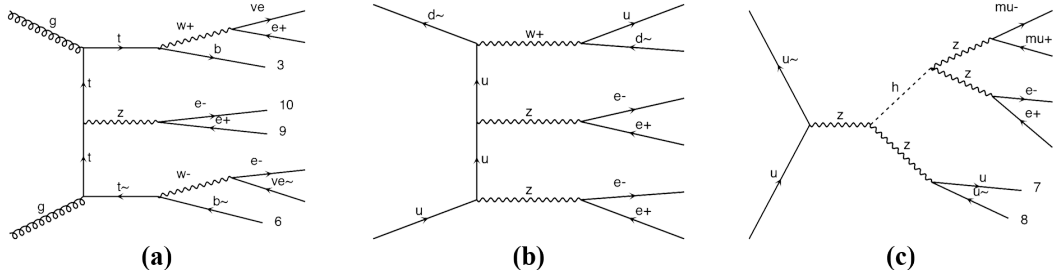


Figure 2. The leading order SM feynman diagrams contributing to none-ZZjj background processes, subsequent decay of the intermediate particles are also shown. (a) ttll process, (b) ZZW^+ process and (c) ZH process.

3.2 Data and Simulations

This analysis was performed using the proton-proton collision data recorded in the ATLAS detector during the run II of LHC with centre-of-mass energy $\sqrt{s} = 13 \text{ TeV}$. These events were required to pass the online trigger system[18], which has an overall efficiency of 96% to 99%[19]. The total accumulated data sample corresponding to an integrated luminosity of 139 fb^{-1} was used in this analysis.

Monte Carlo (MC) Simulated events[20] are used in this analysis to correct the detector-level events for detector effects, and to estimate the expected contribution of signal and background processes. The ATLAS collaboration has switched to a newer version of software to generate MC simulations, which gives better Next-to-Leading-Order (NLO) estimations and updated jets reconstruction algorithm compared to MC simulations used last semester[14].

The EW ZZjj4l process with no on-shell Higgs boson intermediate state was modelled by the SHERPA 2.2.2[21] event generator with NNPDF3.0NNLO Parton Distribution Functions (PDF)[22] at leading order (LO), parton showering was handled by the SHERPA CSSHOWER[23]. The Higgs EW ZZjj4l, in other word Vector Boson Fusion (VBF), events were simulated at NLO precision in QCD using POWHEG-BOX[24] interfaced to PYTHIA 8[25], detailed NLO correction of the VBF process can be found in Ref. [26]. For $gg \rightarrow H \rightarrow ZZ \rightarrow 4l$ events, POWHEG-BOX event generator and similar NLO precision QCD estimation[27] were chosen. Other $ggZZ4l$ processes were modelled using SHERPA 2.2.2 with NNPDF3.0NNLO PDF at LO; the results were merged with CSSHOWER parton showering. The QCD qqZZ4l events were also modelled separately, $t\bar{t} \rightarrow H \rightarrow ZZ \rightarrow 4l$ events were generated at NLO[28] using POWHEG-BOX, while other qqZZ4l events were modelled by SHERPA 2.2.2 generator with NNPDF3.0NNLO PDF at LO, and CSSHOWER parton showering was used to model hadronization of final-state quarks. The VH background events were generated at NLO precision with MiNLO[29] using POWHEG-BOX event generator with NNPDF3.0 PDF and PYTHIA 8 showering. The VVV and ttll processes were modelled by SHERPA 2.2.2 with NNPDF3.0NNLO PDF at LO.

Each particle-level MC simulation event goes through detailed ATLAS detector simulation[30] based on the GEANT 4[31] framework to produce detector-level simulation events. The detector simulation of Every process gives three “generations” of simulation files, each generation corresponds to a specific detector status, neighbouring bunch crossings (pile-up), integrated luminosity in one year of operation of the LHC.

3.3 Events selection

Object	fiducial region
Electrons	“Loose” ID criterion
	“FixedCutPflowLoose” criterion
	$p_T > 7 \text{ GeV}, \eta < 2.47$
	$ d_0/\sigma_{d_0} < 5, z_0 \times \sin\theta < 0.5 \text{ mm}$
Muons	“Loose” ID criterion
	“FixedCutPflowLoose” criterion
	$p_T > 7 \text{ GeV}, \eta < 2.7$
	$ d_0/\sigma_{d_0} < 3, z_0 \times \sin\theta < 0.5\text{mm}$
Jets	$P_T > 30 \text{ GeV}$ and $J_{vt} > 0.6$ for $ \eta < 2.4$
	$P_T > 40 \text{ GeV}$ for $2.4 < \eta < 4.5$
OCSF lepton pairs	two pairs of OCSF lepton with M_Z closest to Z mass
	$p_T > 20, 20, 10 \text{ GeV}$ for leading 3 leptons
	$\Delta R_{l+l-} > 0.2$
	$70 \text{ GeV} < M_Z < 110 \text{ GeV}$
Dijets system	leading two jets with $y_{j_1} \times y_{j_2} < 0$ and $ y_{j_1} - y_{j_2} > 2$
	dijet invariant mass $M_{jj} > 200 \text{ GeV}$
ZZjj system	$P_T \text{ balance} < 0.5$

Table 1. Summary of selections applied to both simulations and observaed events.

The selection of ZZjj4l events relies on the final state particles of the ZZjj4l processes, namely electrons, muons and jets. All the selections discussed here were implimented using the CERN ROOT framework[32].

Electrons were identified by the energy deposit in the electromagnetic calorimeter with corresponding tracks in the inner tracking detector (ID). Muons were reconstructed by the tracks in the muon spectrometer (MS) matched to tracks in the ID, muons can also be identified by the MS alone in the region where the ID does not cover. Jets are clustered by the energy deposit in the hadronic calorimeter using the particle flow (PFlow)[33] algorithm.

The selecting criteria are mostly unchanged from last semester, which can be found in Ref. [14], the only differences are that last semester one of the Z boson can be off-shell, this time the mass of the second pair of Opposite Charge, Same Flavour (OCSF) lepton pair was required to have an invariant mass above 70 GeV, which means both Z bosons are on-shell Z boson, and a jet vertex tagger (Jvt)[34] was applied to jets with $|\eta| < 2.4$ this semester to select jets with higher possibility come from hard interactions. The selection requirements applied are summarised in Table 1, definitions of abbreviations and symbols can be found in Ref. [14].

Both the predicted particle-level and detector-level SM events were required to satisfy all criteria listed above, apart from selections that not applicable to particle-level events, for example, the lepton ID criterion and Jvt selection. The real-world data events were required to pass the same selections as well. The selection of the BSM EFT particle-level events was implemented using the RIVET[35] framework and the jets reconstruction algorithm chosen was *anti-k_t*[36] with radius parameter $R = 0.4$; other criteria were the same as the SM particle-level selections.

The observed and SM predicted event yields in the fiducial region are listed in Table 2, where the “Other” row denotes the overall contribution of VVV , VH and $t\bar{t}l\bar{l}$ processes. The observed event yield shows no significant deviation from the SM prediction. The significance of EW ZZjj process is measured to be 2.18σ , which corresponds to a p-value = 8.7×10^{-4} . Some selections, for example, M_{jj} were deliberately raised or abolished compared to the last semester[14] to let more QCD ZZjj events into the fiducial region, which is essential for BSM physics models limit setting, details will be discussed in section 4.3.

Process	Events yield
EW ZZjj4l	22.9 ± 0.1
QCD qqZZ4l	101.3 ± 0.5
QCD ggZZ4l	26.9 ± 0.2
Other	7.9 ± 0.1
Total	158.9 ± 0.6
Data	164

Table 2. The SM predicted and observed fiducial event yields for different processes. Uncertainties quoted are statistical only.

3.4 Chosen Observables

To be sensitive to the CP-odd BSM physics model, the distributions of observables selected should be sensitive to CP-violating effects; in this analysis, several CP-sensitive observables have been investigated.

The most potent observable in probing CP-violation is the signed azimuthal angle between the leading and subleading jets[37],

$$\Delta\phi_{jj} = \phi_1 - \phi_2 \quad (3.1)$$

where the ordering of jets is determined by the rapidity of them so that j_1 is the jet with higher rapidity (more forward in space). By artificially selecting one direction, the $\Delta\phi_{jj}$ is by definition sensitive to the Parity violation[38]. The other chosen observable, θ_1 [39], is defined as the angle between the leading Z boson and the third lepton measured in the ZZ rest frame, i.e.

$$\theta_1 = \cos^{-1}\left(\frac{\mathbf{Z}_1 \cdot \mathbf{l}_3}{|\mathbf{Z}_1||\mathbf{l}_3|}\right) \quad (3.2)$$

where the bold character denotes the three-vector of particles in ZZ rest frame. Other observables mentioned in Ref. [39] were investigated as well, but θ_1 was chosen for its best sensitivity to the CP-even BSM EFT operators, more on that in section 4.3.

These distributions were initially finely binned and later re-binned to variable bin sizes, the bin width were determined by requiring at least 14 expected detector-level ZZjj events each bin. The expected and observed detector-level event yields as functions of the two selected observables are shown in Fig. 3. The observation is compatible with the SM predictions.

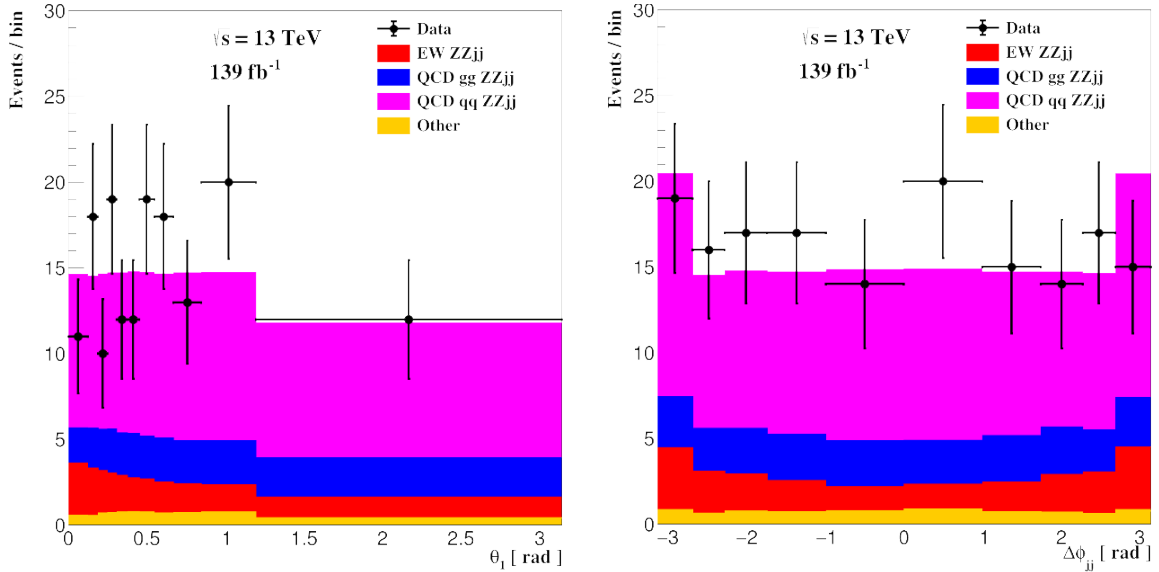


Figure 3. Observed and expected event yield as function of θ_1 and $\Delta\phi_{jj}$ with variable bin width.

3.5 Unfolding

The measured $\Delta\phi_{jj}$ and θ_1 spectra were unfolded to remove detector effects, including the efficiency of the trigger system, and limited resolution and efficiency of the

tracking detectors and calorimeters[40]. The unfolding procedure also gives the covariance matrices account for the correlation between bins in a given distribution. By performing such correction, the results become model-independent so that they can be directly compared with BSM particle-level predictions.

The unfolding procedure relies on the thorough understanding of the ATLAS detector, in other words, how would particle-level final-state particles interact with the detector and what detector-level signals are expected. The imperfection of the detector results in events migration, fake signals and missing particles. Since the direct access to particle-level events is impossible, those three effects are estimated by the particle-level and detector-level simulations and represented by a respond matrix R_{ij} ; more details can be found in Ref. [14]. An iterative unfolding algorithm[15] using the particle-level predicted distributions as an initial hypothesis, implemented by RooUnfold framework[41] was used to achieve the unfolding. One iteration was found to be a good balance between bias and statistical uncertainties for both distributions.

Before the unfolding procedure, the none-ZZjj background contributions were removed from the observed data, leaving only ZZjj contributions; the faked events were then removed, and the iterative method was performed; finally, a bin-by-bin efficiency correction was applied.

3.6 Results

After applying the unfolding technique, the event yields can be easily be reinterpreted to the differential cross section in the fiducial region. The measured and expected fiducial differential cross section as functions of selected observables are shown in Fig. 4, showing no significant deviation from the SM predictions

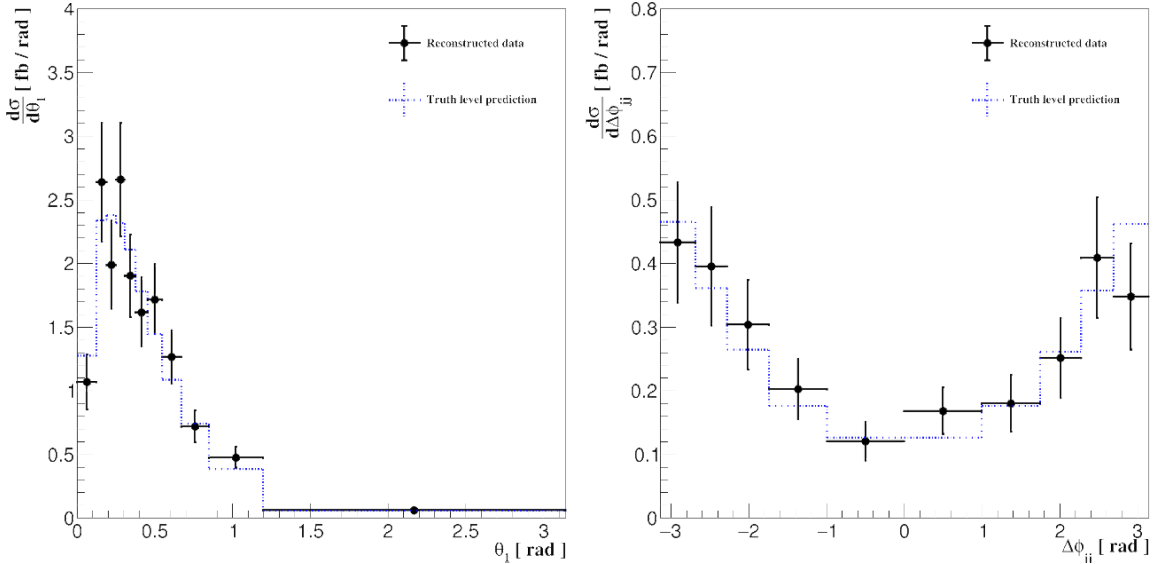


Figure 4. observed and expected event yield as functions of θ_1 and $\Delta\phi_{jj}$ with variable bin width.

4 Reinterpretation in an effective field theory

4.1 EFT theory

In the SM, no CP-violation is predicted in the electroweak interactions; however, as stated, the SM is incomplete. In the SMEFT, the SM is treated as a low-energy (infrared) leading order approximation of a yet to be determined grand theory when new physics phenomena reside in a much higher energy scale, in this case, $\Lambda \gg \bar{v}_T$. One can expand the SM by constructing a series of expansion in the order of the ratio to new physics scale. With the SM Lagrangian defined in Ref. [42], the EFT Lagrangian constructed using $SU(3) \times SU(2) \times U(1)$ symmetry with dimension-6 BSM operators can be written as

$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \mathcal{L}^{(6)}, \quad \mathcal{L}^{(6)} = \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i \quad (4.1)$$

where c_i is the Wilson coefficient, and \mathcal{O}_i is the corresponding dimension-6 BSM operator, higher dimensional Lagrangian was neglected. The structure of the Lagrangian allows us to split the quantum amplitude, square of matrix element, to the SM and BSM part, i.e.

$$|\mathcal{M}|^2 = |\mathcal{M}_{SM}|^2 + 2\text{Re}(\mathcal{M}_{SM}^\dagger \mathcal{M}_{EFT}) + |\mathcal{M}_{EFT}|^2 + \mathcal{O}(\Lambda^{-6}) \quad (4.2)$$

where $2\text{Re}(\mathcal{M}_{SM}^\dagger \mathcal{M}_{EFT})$ is proportional to Λ^{-2} , which corresponds to the interference between the SM and EFT, and $|\mathcal{M}_{EFT}|^2$ is proportional to Λ^{-4} , which is the direct EFT contribution. In this report, only one BSM operator is discussed at a time, so we hereby denote the cross section corresponding to each component of the amplitude. The total cross section can be written as

$$\sigma_{EFT}(C) = \sigma_{SM} + \frac{C}{\Lambda^2} \sigma_{Linear} + \frac{C^2}{\Lambda^4} \sigma_{Quad} \quad (4.3)$$

where the *Quad* and *Linear* indicate the corresponding cross section are proportional to c^2 and c respectively.

The VBS process typically contains VVV, HVV and Vqq interaction vertices, EFT operators provide alternative interactions to these vertices, for example, by exchanging unknown massive fermion whose mass is much higher than the Higgs boson. Studied EFT operators are summarised in Table 3, in which the ψ is the Higgs field, W and B are the $SU(3) \times SU(2) \times U(1)$ electroweak field tensor and τ^I is the generator of the $SU(2)$ symmetry group. The operators with $\tilde{W}_{\mu\nu}^I = \epsilon^{IJK} W_{\mu\nu}^{JK}/2$ provide CP-violating effect through the interference with the SM, while the operators without tilde are CP-even. For CP-odd operators, the integrated σ_{Linear} vanish[43] so that can only be probed by CP-sensitive observables, whereas the σ_{Quad} can be probed by the total cross section deviation from the SM prediction.

Coefficient	Operator	Relevant vertices
c_W	$\epsilon^{IJK} W_\mu^{I,\nu} W_\nu^{J,\rho} W_\rho^{K,\mu}$	VVV
$c_{\tilde{W}}$	$\epsilon^{IJK} \tilde{W}_\mu^{I,\nu} W_\nu^{J,\rho} W_\rho^{K,\mu}$	VVV
c_{HW}	$\psi^\dagger \psi W_{\mu\nu}^I W^{I,\mu\nu}$	HVV
$c_{H\tilde{W}}$	$\psi^\dagger \psi \tilde{W}_{\mu\nu}^I W^{I,\mu\nu}$	HVV
c_{HWB}	$\psi^\dagger \tau^I \psi W_{\mu\nu}^I B^{\mu\nu}$	HVV and Vqq
$c_{H\tilde{W}B}$	$\psi^\dagger \tau^I \psi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	HVV and Vqq

Table 3. Studied Wilson coefficients and the corresponding operators and the interaction vertices they alter.

4.2 particle-level EFT event generation

The parton-level EW ZZjj EFT events were generated at centre-of-mass energy of 13 TeV using MADGRAPH 5[44] event generator, with the SMEFTsim[13] EFT implementation. The chosen SMEFTsim package is a simplified $U(3)^5$ EFT model which dramatically reduces the number of Wilson coefficients (69[13] versus 2499 for full SMEFT[45]) thus reduces the need for computation power. The $\{M_W, M_Z, G_F\}$ input scheme was chosen for certain theoretical advantages and better precision over the $\{\alpha_{EW}, M_Z, G_F\}$ input scheme[13]. The parton-level results were then passed to PYTHIA 8[25] to model the decay of Z boson and the hadronization of final-state quarks. The particle-level events then go through the same selections described in section 3.3. The EFT events the selections were reimplemented using the RIVET framework[35].

The linear and quadratic EFT events of EW jj4l process were generated separately with only one Wilson coefficient set to $c = 5$, while others set to zero, and the scale factor was set to $\Lambda = 1$ TeV. One can scale the resulting spectra using the proportional rule for σ_{Linear} and σ_{Quad} , see equation 4.3, to get distributions corresponding to other values of the Wilson coefficient.

Due to the limitation of computing resources, I am not able to generate adequate EW jj4l events to minimise the statistical fluctuation. As a result, only processes with two on-shell Z bosons were generated to keep generating time reasonable. The predicted SM differential cross-sections as a function of $\Delta\phi_{jj}$ generated with and without the off-shell Z boson in the fiducial region are shown in Fig. 5, indicating little difference between the two samples in the fiducial region.

In fig. 6 and 7, the ratio of EFT differential cross sections with $C = 1$ as functions of $\Delta\phi_{jj}$ and θ_1 to the SM prediction are shown. the $\Delta\phi_{jj}$ spectrum shows excellent sensitivity to the CP-odd EFT operators, and θ_1 spectrum shows great sensitivity to CP-even EFT operators. The most considerable EFT effect is the contributions corresponding to Wilson coefficients C_W and $C_{\tilde{W}}$, while little sensitivities were shown for C_{HWB} and $C_{H\tilde{W}B}$. To quantify the inclusive CP-violation, a geometrical asymmetry

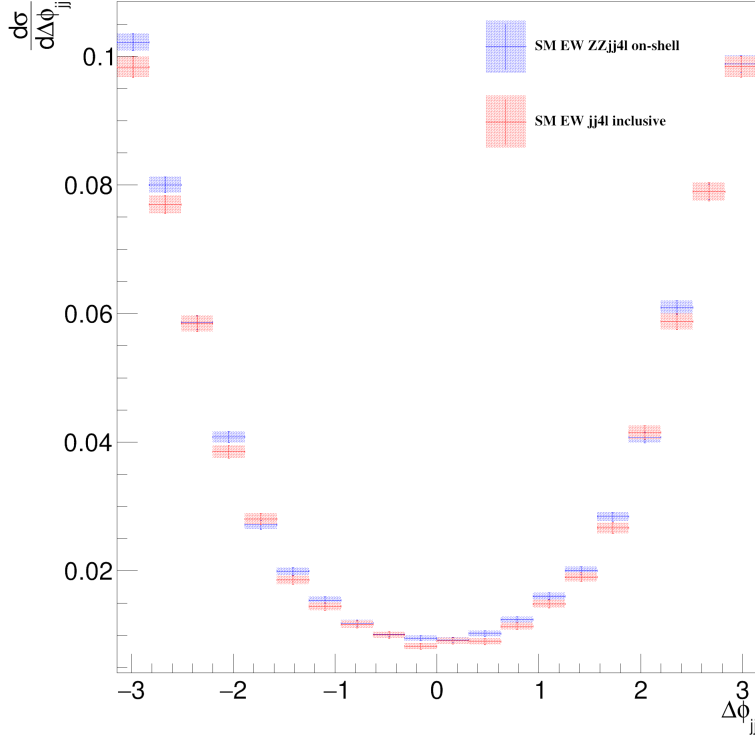


Figure 5. The predicted SM differential cross-section as function of $\Delta\phi_{jj}$ generated with and without the off-shell Z boson in the fiducial region, indicating no obvious mismodelling by not generating processes with off-shell Z bosons.

A is defined as

$$A = \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-} \quad (4.4)$$

where σ_{\pm} is the cross section in the positive and negative $\Delta\phi_{jj}$ part in the fiducial region. The change of fiducial cross sections and induced Asymmetries by each EFT operator with $C = 1$ are summarised in Table 4, indicating the CP-odd operators do not affect the total cross section at linear order but do have contributions to the Asymmetry. Note that at $C = 1$ the quadratic contributions are similar to the linear contributions, at larger Wilson coefficient the quadratic correction can dominate the EFT effect.

4.3 Limit setting procedure

With relatively high observed statistics in each bin of the two studied distributions, the uncertainties can be modelled by Gaussian. The χ^2 function constructed to set limits on the Wilson coefficients is defined as

$$\chi^2 = (\vec{\sigma}_{data} - \vec{\sigma}_{EFT})^T \mathbf{Cov}^{-1} (\vec{\sigma}_{data} - \vec{\sigma}_{EFT}) \quad (4.5)$$

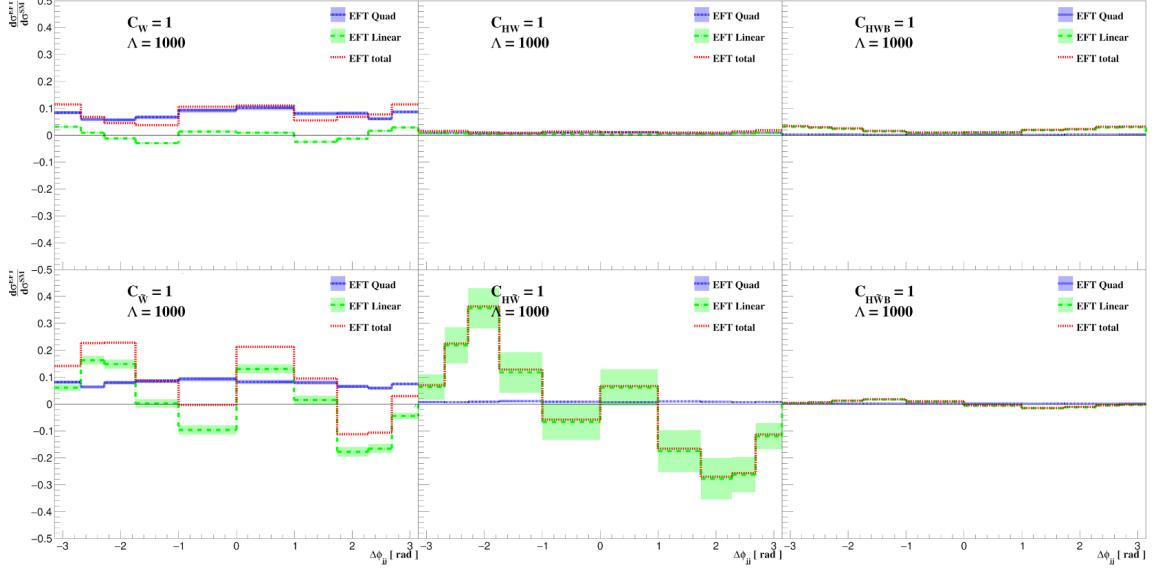


Figure 6. The ratio of EFT differential cross section with $C = 1$ as function of $\Delta\phi_{jj}$ to the SM predicted cross sections, showing that the $\Delta\phi_{jj}$ spectrum have great sensitivity to CP-odd EFT operators.

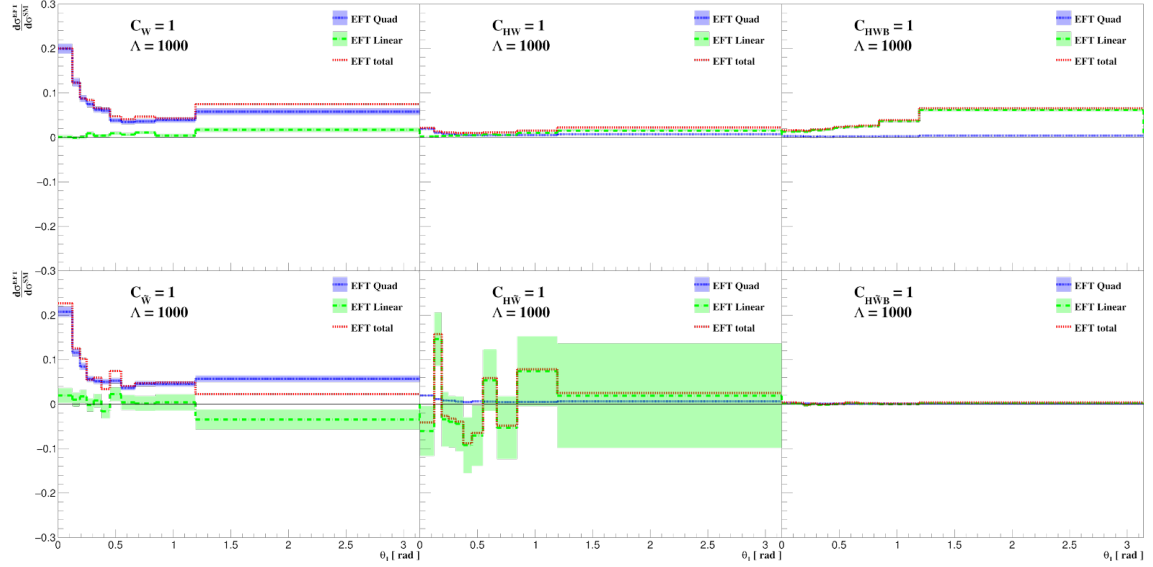


Figure 7. The ratio of EFT differential cross section with $C = 1$ as function of θ_1 to the SM predicted cross sections, showing that the θ_1 spectrum have great sensitivity to CP-even EFT operators.

where $\vec{\sigma}_{data}$ and $\vec{\sigma}_{EFT}$ are the vectors of observed and predicted EFT cross sections in every bin of a particular distribution, \mathbf{Cov} is the statistical covariance matrix given by the the unfolding procedure. Each time two of the Wilson coefficients C_1 , C_2 was scanned and a pair of $\{C_1^{min}, C_2^{min}\}$ which minimise the χ^2 was determined.

Type	σ_{Linear} [fb]	σ_{Quad} [fb]	σ_{EFT}/σ_{SM}	Asymmetry
SM	-	-	-	0.002 ± 0.003
c_W	0.008 ± 0.002	0.119 ± 0.004	0.082 ± 0.003	0.003 ± 0.003
$c_{\tilde{W}}$	0.006 ± 0.008	0.117 ± 0.004	0.080 ± 0.006	-0.059 ± 0.006
c_{HW}	0.0083 ± 0.0005	0.0117 ± 0.0003	0.0130 ± 0.0004	-0.002 ± 0.003
$c_{H\tilde{W}}$	-0.02 ± 0.03	0.0117 ± 0.0003	-0.00 ± 0.02	-0.15 ± 0.02
c_{HWB}	0.0342 ± 0.0008	0.00313 ± 0.00007	0.0242 ± 0.0005	-0.002 ± 0.003
$c_{H\tilde{W}B}$	0.0002 ± 0.0007	0.00175 ± 0.00004	0.0012 ± 0.0005	-0.010 ± 0.003

Table 4. The predicted particle level EFT fiducial cross section, their ratio to the SM prediction and the Asymmetries at $C = 1$ and $\Lambda = 1$ TeV. The uncertainties quoted are statistical only.

The $\Delta\chi^2$ was constructed by

$$\Delta\chi^2(C_1, C_2) = \chi^2(C_1, C_2) - \chi_{min}^2 \quad (4.6)$$

According to the Wilk's theorem[46], the $\Delta\chi^2$ also follows the χ^2 distribution, whose corresponding Number of Degrees of Freedom (NDF) equals the remaining free parameters, in this case, 2. The Confidence Level (CL) of each point on the plane was calculated by

$$1 - CL = \int_{\Delta\chi^2}^{\infty} f_{k=2}(x) dx \quad (4.7)$$

where the χ_{min}^2 is the minimum χ^2 on the plane and $f_k(x)$ is the χ^2 distribution function with NDF $k = 2$.

4.4 Results

The 68% and 95% expected and observed limits on two of the studied CP-odd Wilson coefficients with other coefficients set to 0 calculated using the $\Delta\phi_{jj}$ distribution are shown in the left panel right panels in Figure 8, the asymmetries on the plane are also overlayed. The observed $\Delta\chi^2$ are shown in the left panels.

The 68% and 95% expected and observed limits on two of the studied CP-even Wilson coefficients with other coefficients set to 0 calculated using the θ_1 distribution are shown in the left panel right panels in Figures 9. The observed $\Delta\chi^2$ are shown in the left panels.

The one-dimensional observed and expected constraints of individual Wilson coefficient are also calculated, in this case the NDF of the $\Delta\chi^2$ distribution was set to 1 so that the resulting limits are quite tighter than the two-dimensional limits. The results are summarised in Table 5. As expected, the $\Delta\phi_{jj}$ distribution is more sensitive to CP-odd operators, whereas the θ_1 distribution gives tighter limits on CP-even Wilson coefficients.

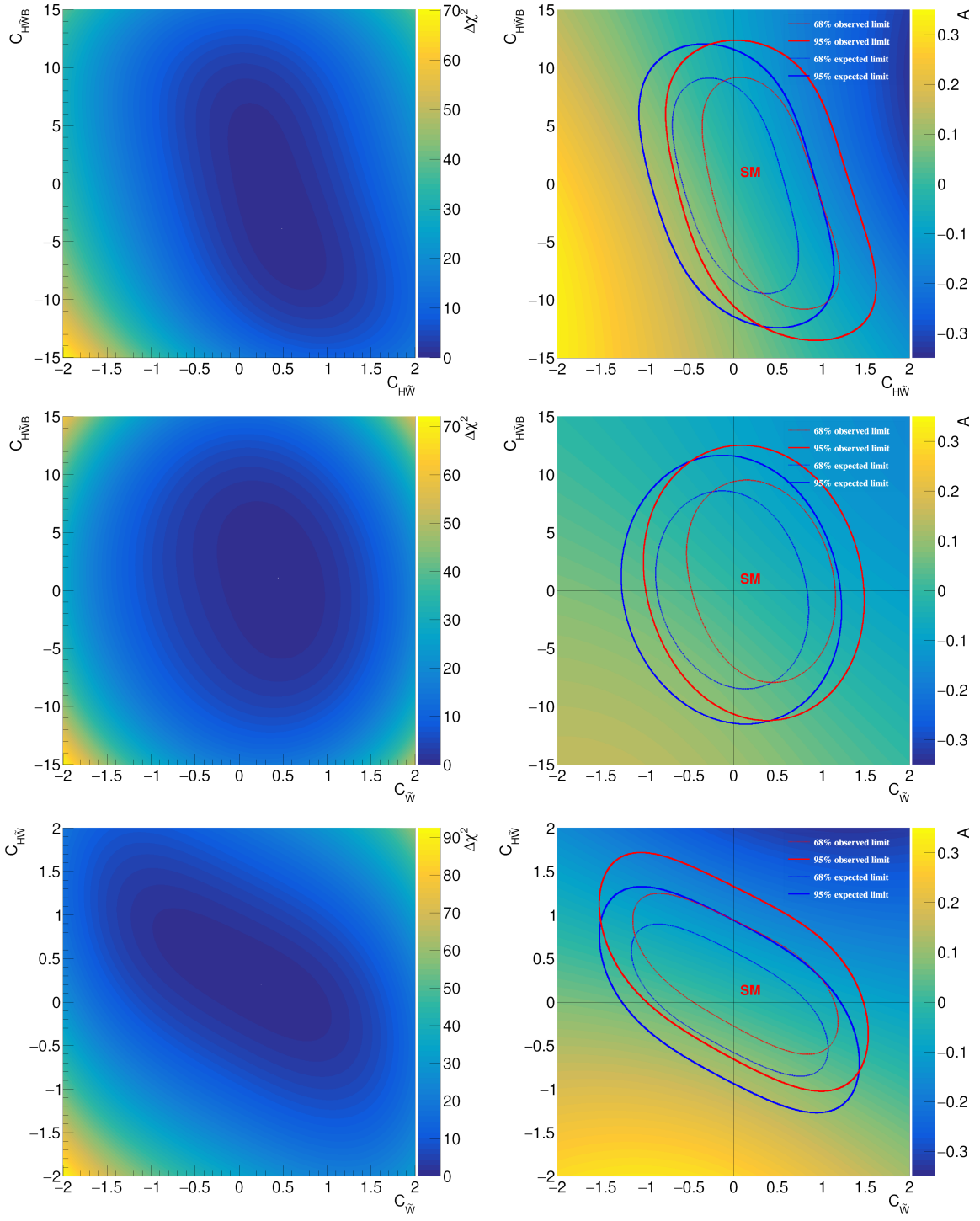


Figure 8. *TODO*

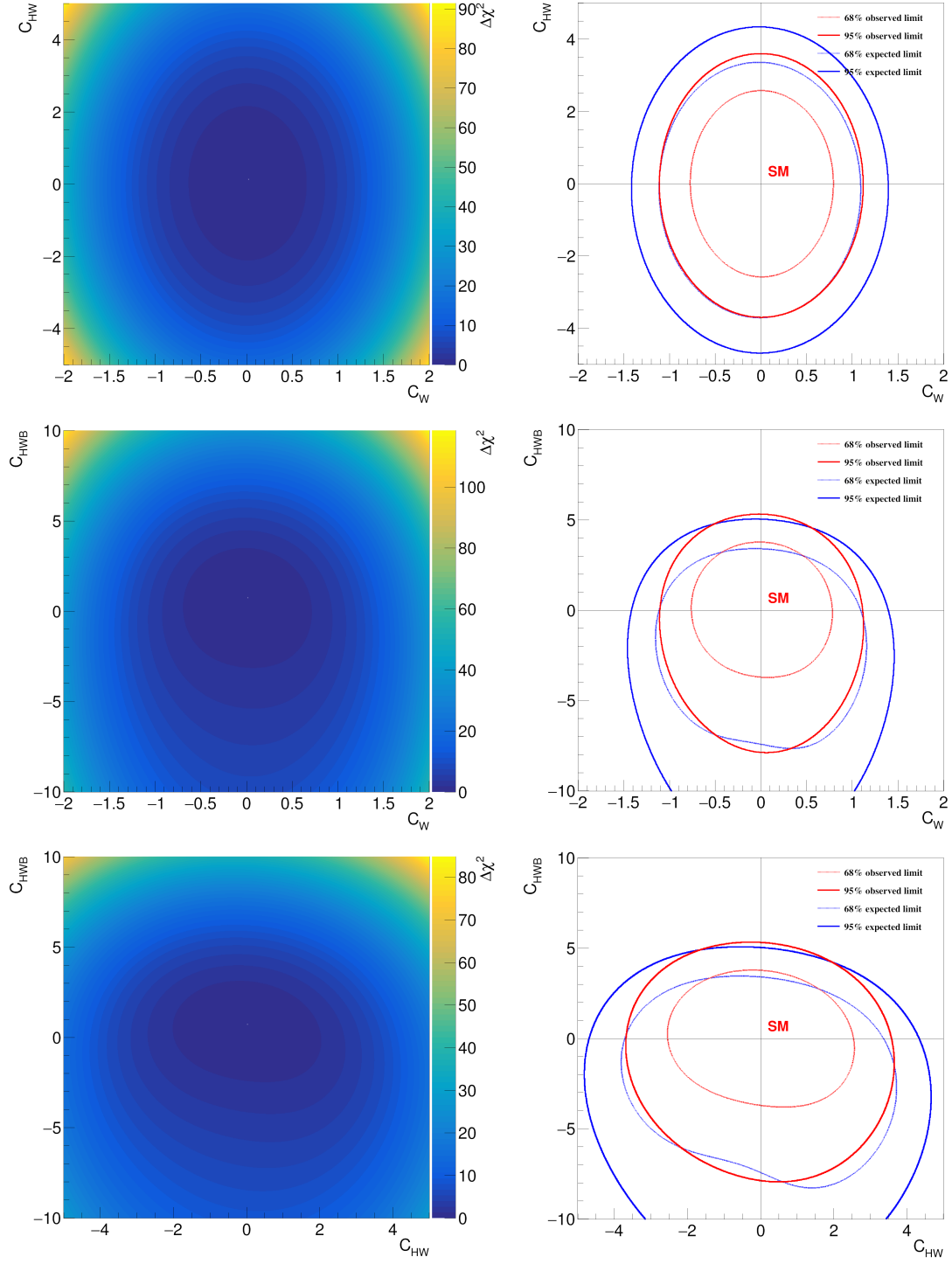


Figure 9. *TODO*

5 Discussion

The process of EW ZZjj with on-shell Z bosons shows good sensitivity to selected EFT operators in the fiducial region. Compared to the $H \rightarrow Z^*Z \rightarrow 4l$ process[37],

Coeff.	Distro.	68% obs.	95% obs.	68% exp.	95% exp.
C_W	θ_1	[-0.37, 0.41]	[-0.66, 0.68]	[-0.64, 0.62]	[-0.89, 0.88]
	$\Delta\phi_{jj}$	[-0.80, 0.39]	[-1.06, 0.82]	[-0.72, 0.65]	[-1.01, 0.94]
$C_{\tilde{W}}$	θ_1	[-0.41, 0.36]	[-0.69, 0.63]	[-0.64, 0.58]	[-0.90, 0.84]
	$\Delta\phi_{jj}$	[-0.06, 0.56]	[-0.38, 0.81]	[-0.31, 0.31]	[-0.59, 0.58]
C_{HW}	θ_1	[-1.24, 1.33]	[-2.22, 2.17]	[-2.16, 1.83]	[-3.01, 2.66]
	$\Delta\phi_{jj}$	[-2.78, 1.55]	[-3.65, 2.67]	[-2.62, 1.94]	[-3.53, 2.85]
$C_{H\tilde{W}}$	θ_1	[0.30, 0.99]	[-0.03, 1.33]	[-0.32, 0.33]	[-0.64, 0.65]
	$\Delta\phi_{jj}$	[-0.03, 0.37]	[-0.23, 0.57]	[-0.19, 0.19]	[-0.37, 0.38]
C_{HWB}	θ_1	[-0.96, 1.62]	[-2.58, 2.65]	[-1.54, 1.28]	[-3.55, 2.35]
	$\Delta\phi_{jj}$	[-11.4, 1.06]	[-13.1, 2.35]	[-2.04, 1.51]	[-11.9, 2.72]
$C_{H\tilde{W}B}$	θ_1	[-2.17, 4.60]	[-5.24, 6.68]	[-5.52, 5.18]	[-7.79, 7.44]
	$\Delta\phi_{jj}$	[-2.38, 4.96]	[-5.51, 7.35]	[-3.52, 3.58]	[-6.14, 6.24]

Table 5. *todo*

constraints on operators corresponding to C_HW , $C_{H\tilde{W}}$ and C_{HWB} , $C_{H\tilde{W}B}$ are tighter. The model-independent technique shows good robustness and the CP-sensitive observables are proved useful.

However, there are a few limitations. Firstly, although large sets of EFT events were generated, due to the low branching ratio only around 2000 events were eventually selected in the fiducial region, which can lead to rather large statistical fluctuations. These uncertainties were not taken into account in the limit setting procedure.

Secondly, only statistical uncertainties were considered, systematical and theoretical uncertainties in both EFT and SM distributions were not taken into account. Lower sensitivities are expected if these uncertainties are being considered.

Thirdly, The NDF of the χ_2 distributions should be determined experimentally using large sets of pseudo-experiments, however, in this analysis, they are set to 1 or 2 depending on the number of involved Wilson coefficients.

Finally, if more than one Wilson coefficients were non-zero values, the Λ^{-4} correction to the cross section should also contain the interference between these Wilson coefficients, which can only be ignored when C_1C_2/Λ^{-4} is small compared to other Λ^{-4} contributions, in other words one of the involving Wilson coefficients is small compared to the other one. These corrections was ignored regardlessly in this report.

Future development of this analysis can be carried out based on these limitations, constraints on other new physics models can also be easily estimated using the same technique.

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