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Natural Widths of Atomic K and L Levels, $K\alpha$ X-Ray Lines and Several KLL Auger Lines

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Semi-empirical values of the natural widths of K, L_1 , L_2 and L_3 levels, $K\alpha_1$ and $K\alpha_2$ x-ray lines, and KL_1L_1 , KL_1L_2 and KL_2L_3 . Auger lines for the elements $10 \le Z \le 110$ are presented in tables and graphs. Level width Γ_i (i = K, L_1 , L_2 , L_3) is obtained from the relation $\Gamma_i = \Gamma_{R,i}/\omega_i$, using the theoretical radiative rate $\Gamma_{R,i}$ from Scofield's relativistic, relaxed Hartree-Fock calculation and the fluorescence yield ω_i from Krause's evaluation. X-ray and Auger line widths are calculated as the sums of pertinent level widths. This tabulation of natural level and line widths is internally consistent, and is compatible with all relevant experimental and theoretical information. Present semi-empirical widths, especially those of $K\alpha_1$ and $K\alpha_2$ x-rays, are compared with measured widths. Uncertainties of semi-empirical values are estimated.

Key words: Auger effect; Auger line width; K shell: L shell: level width; line width; natural width; x-ray line width.

1. Introduction

In the preceding paper [1]¹ we presented an evaluation of fluorescence, Coster-Kronig and Auger yields for K and L shells. That evaluation utilized experimental and theoretical data up to October 1977 and led to sets of values that are internally consistent and are compatible with all relevant information. It is possible to use these fluorescence yield values, in conjunction with theoretical radiative rates, to derive reliable semi-empirical values of natural widths of K and L levels and, from these, natural widths of $K\alpha$ x-ray lines and KLL Auger lines. Such values are of importance in many areas of basic and applied science, most notably in the various branches of electron and x-ray spectrometry.

Level and line widths presented in this paper are natural atomic widths. In principle, natural widths are smaller than the actual widths observable in an experiment, because several effects, especially multiplet splitting and multiple ionization, may broaden a level or an emission line. However, these broadening effects are in general small, except for light elements, and will be discussed in some detail in section 5.

The most recent tabulations of x-ray line widths were made by Salem and Lee [2] and by Pessa [3]. To arrive at their recommendations of $K\alpha_1$ and $K\alpha_2$ x-ray line widths, Salem and Lee applied least-squares fits to experimental x-ray data and Pessa used x-ray and some level width data. However, line widths for $Z \leq 30$ were not reduced to natural atomic widths in either tabulation. Instead, these line widths reported in references 2 and 3 contain contributions from various broadening effects and pertain to solid systems normally encountered in practice.

Tables of K level widths, and in several cases L level widths, have been presented by Parratt [4], Blokhin [5], Meisel and Nefedov [6], Sevier [7], and Salgueiro et al. [8]. Level widths in reference 8 are based on theoretical radiative rates and semi-empirical fluorescence yields, those in references 4 to 7 are based on the relatively scarce measurements of x-ray absorption edges or absorption lines and on the more plentiful measurements of x-ray emission lines suitably adjusted for the width of the final level reached in an x-ray transition. A compilation of level width data up to about 1971 can be found in Sevier [7]; and recent theoretical results are shown in graphic form in reference 1, where also the original calculations of the partial level widths are cited.

No tabulation of *KLL* Auger line widths has been made to date.

It should be emphasized that, except for the theory-based $A(Z + B)^4$ relation [4,5,8], no other pertinent information has been included in past empirical efforts to systematize measured level and line widths as a function of atomic number Z and to achieve internal consistency.

2. Equations

In this work we derive the natural level width Γ_i from the radiative width (or rate) $\Gamma_{R,i}$ and the fluorescence yield ω_i of level i using the relation

$$\Gamma_i = \Gamma_{\rm R,i}/\omega_i \tag{1}$$

where $i = K, L_1, L_2, L_3$.

Natural x-ray and Auger line widths are calculated as the sums of the widths of levels involved in the transitions:

$$\Gamma(K\alpha_k) = \Gamma(K) + \Gamma(L_\ell), \tag{2}$$

¹ Figures in brackets indicate literature references at the end of this paper.

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² The subscripts L_1, L_2 , and L_3 are also written simply as 1,2,3: e.g., $\omega_{L_2} = \omega_2$

where k = 1, 2; $\ell = 3, 2$; and

$$\Gamma(KL_kL_\ell) = \Gamma(K) + \Gamma(L_k) + \Gamma(L_\ell), \tag{3}$$

where $L_k L_\ell = L_1 L_1$, $L_1 L_2$, $L_2 L_3$.

While eq (2) is exact, eq (3) is an approximation to the exact expression

$$\Gamma(KL_kL_\ell) = \Gamma(K) + \Gamma(L_kL_\ell) \tag{4}$$

and implies invariance of the decay width regardless of the number of holes present.

3. Source Material

Semi-empirical values presented in table 1 and figures I to 4 were obtained from eqs (1) to (3) with theoretical radiative rates $\Gamma_{R,i}$ from Scofield's relativistic, relaxed Hartree-Fock calculation [9-11] and fluorescence yields ω_i from Krause's evaluation [1]. As a consequence, the present semi-empirical level and line widths utilize the same source material as the fluorescence yields. Specifically, these widths, like the fluorescence yields, are compatible with available information on (a) fluorescence, Auger, and Coster-Kronig yields, (b) theoretical radiative and radiationless transition rates, (c) level widths, (d) x-ray and Auger line widths, (e) x-ray and Auger spectra, and (f) Coster-Kronig energies. In the evaluation of the fluorescence yields, only those level and line width data were used for which non-lifetime broadening could be presumed negligible. For further details on the evaluation procedure, the reader is referred to reference 1.

4. Estimate of Uncertainties

Uncertainties in the level width values are the quadratically added uncertainties in fluorescence yields ω_i and radiative rates $\Gamma_{R,i}$. We neglect correlation between the radiative rate and the fluorescence yield introduced through the evaluation process [1]. Generally, uncertainties are dominated by those assigned to ω_i [1]. Uncertainty of the theoretical K-shell rate, $\Gamma_{R,K}$, is assumed to be 3-5%, and of the L-subshell rates, $\Gamma_{R,L_{\ell}}$, to be 4-7%. These estimates rely on a comparison of theoretical rates with the following experimental data: first, the intensity ratios of x-ray transitions to levels with different principal quantum numbers [12], and second, the $K\alpha$ x-ray widths at higher Z where the radiative width predominates [2,3]. Earlier, Scofield surmised an uncertainty of about 2 % for $\Gamma_{R,K}$ on the basis of K x-ray branching ratios [11].

Uncertainties in *line* widths are the combined errors in the components.

Results are summarized in table 2. As expected, the uncertainty in the widths of the $K\alpha$ and KL_2L_3 lines are dominated by the K level width. Uncertainties are relatively large at low Z, because of large uncertainties in the fluorescence yields ω_i [1], which were obtained

from the small number of experimental data and realistic calculations [13] available in this Z range. The estimates given in table 2 for light elements are conservative, at least for closed shell atoms, as can be seen from table 3 in which our semi-empirical values are compared with experimental data [14–26].

5. Effects Influencing the Widths

The natural level width, which is governed solely by the lifetime of a hole in that level, may not be observable in an experiment because of various effects influencing the width of a level. Similarly, natural line widths of transitions interconnecting levels may be difficult to delineate. General treatises of the various effects have been given in several places [4,5,7]; detailed discussions have been presented in many specific studies [27-45]. Here, only a brief outline of the most important effects and their consequences will be given in regard to the K and L shells.

Effects can be placed into two general categories: lifetime effects and non-lifetime effects. Either category can be advantageously divided into two Z ranges: (a) a range in which the levels under consideration are shallow core levels, usually in the penultimate shell, and (b) a range in which the levels are deep core levels, separated from the peripheral shell by at least one completely filled principal shell.

The lifetime of a hole in the level i is altered if additional holes are present in other levels. This lifetime effect can be substantial for shallow levels, that is in light atoms [39-41], but is negligible for deep core levels, that is in heavier atoms, provided the additional holes are in the outer shells [42]. A similar situation prevails for atoms in chemical or solid state matrices since chemical bonding entails removal or addition of peripheral charge.

If the level width is measured by photoelectron spectrometry, the widths corresponding to various hole states can generally be determined separately, for the various photolines are separated in energy [33,44]. The same holds true for x-ray lines and, especially, Auger lines of the light elements [33,44,45], unless too many hole states are present as in the case of most ion-atom collisions. A drastic change in the lifetime of the $L_{2.3}$ level of argon in the presence of an $M_{2.3}$ hole has been noted in an Auger spectrum [16,23,41], in which the single-hole, diagram line is well separated in energy from the double-hole, satellite line.

Diagram lines coincide with satellite lines in medium to heavy atoms in the case of x rays [32,36,37] and in heavy atoms in the case of Auger effect [33] assuming the additional holes are in outer levels. In these cases the line widths are affected in other ways, even though the lifetimes for the different lines may be, and usually are [42], identical.

If the peripheral electron configuration of an atom is changed, whether by multiple ionization or by chemical bonding, level energies change and Coster-Kronig decay channels may either open or close. When this occurs, the effect on the natural level or line width is usually dramatic because of the importance of Coster-Kronig transitions. A notable example can be found in the 3d transition metal series, where the L_2 - L_3 X Coster-Kronig transition, which is forbidden in free atoms, becomes allowed in bound atoms with $26 \le Z \le 30$ [38] and possibly down to Z = 22 [1]. Hence, the L_2 level, $K\alpha_2$ x-ray and KL_2L_3 Auger lines are broader in the metal than in free atoms due to a lifetime effect. Use the ω_2 values in table 1 of the preceding paper [1] in eq (1) to obtain these natural width values pertaining to the metals of these elements.

Non-lifetime effects are caused by a myriad of effects [7,27-37], of which exchange coupling (multiplet splitting) is the most conspicuous one in the light elements [27,28]. These effects always lead to a level broadening. For example, if the energy spacings of multiplet compo nents are less than the lifetime governed width, level broadening will result. In 3d elements, specifically $22 \le Z \le 28$, the multiplet splitting of the L_2 and L_3 levels is of such a degree that these levels as well as the pertinent x-ray and Auger lines are considerably broadened [27-29]. Although observed only in bound atoms so far, this effect is also present in free atoms. Since exchange coupling becomes weaker with increasing separation of the partially filled shells [43], K and L level widths and interconnecting transitions will experience negligible broadening from this source in the heavier elements.

Another non-lifetime effect is associated with the occurrence of multiple vacancies [32,33,36,37]. If, in addition to the hole in a deep core level, outer shell vacancies are created, as for example by photon and electron excitation, the resulting satellites will appear within the breadth of the single-hole, diagram emission line. Because the coincidence is not exact and because the satellite lines are subject to multiplet splitting, emission lines will be broadened. This has been demonstrated recently for $K\alpha$ x-ray lines in iron [37] and copper [36]. However, as Z increases this effect loses in significance compared with the natural width of the diagram line.

For detailed descriptions of other effects contributing to the widths of levels and lines, the reader is referred to references 7, 30, 34 and 35. The possibility that decay widths of multiplet components vary within a multiplet has been considered theoretically for light atoms [39,40].

The widths presented in table 1 and figures 1 to 4 are natural widths, in atoms with a single initial hole, on the basis of the data and procedures used (secs 2 and 3), except that small, and probably negligible, admixtures from other sources including those discussed in this section could not be completely eliminated in the evaluation process of the fluorescence yields [1].

The discussion of the most important effects influencing the widths shows that level and line widths are virtually identical with the natural widths in atoms in which K and L levels are deep core levels. On the other hand, level and line widths can deviate from the natural

widths in atoms in which K and L levels are shallow levels. The dividing lines between the "deep" and "shallow" regions can be placed near $Z\approx 13$ for the K level and near Z=30 for the L levels and for transitions between K and L levels.

6. Comparison With Experimental Data

In this section, we compare our semi-empirical values with various experimental data on level and line widths. Such comparisons provide an overview of the present state of the art and may pinpoint areas where improvements are needed.

In table 3, free-atom data [14-26], mostly of recent origin, are compared with semi-empirical values. Generally, satisfactory accord is noted except for the KL_1L_1 and KL_1L_2 widths of argon. These discrepancies in the argon Auger line widths may be due, at least in part, to the use of the approximate expression eq (3).

Semi-empirical natural line widths of $K\alpha_1$ and $K\alpha_2$ x rays are compared in figures 2, 3, 5 and 6 with both fitted and selected original experimental data. As seen in figures 2 and 3, semi-empirical widths show the expected smooth increase with atomic number, while fitted data [2,3] exhibit some undulations. Part of the increased width observed in experiments with atoms in solids [2,3] between Z=21 and 28 is due to exchange interactions of a 2p hole with the incomplete 3d shell [7,27,28]. Starting with Z=29 [30], x-ray widths are essentially lifetime governed.

The deviation of the fitted values by Salem and Lee [2] from the present semi-empirical values (fig 5) occurs in several, but nonidentical, regions for $K\alpha_1$ and $K\alpha_2$ lines and might be attributed to their use of inaccurate input data. Values tabulated by Pessa [3] agree satisfactorily with our semi-empirical values except in the vicinity of Z=90. Measurements by Nelson and Saunders [46], though exhibiting relatively larger scatter, agree satisfactorily with present values if the average percent-deviation ($\lesssim 2\%$) is considered.

In figure 6, level width differences $\Delta\Gamma(L_{2,3}) = \Gamma(L_2)$ $\Gamma(L_3)$ taken from various sources are plotted for Z > 30. We note small fluctuations in the present values, probably due to uncorrelated errors in $\Gamma(L_2)$ and $\Gamma(L_3)$, and find good agreement with $M(I_{2,0}) = \Gamma(I_{2}M_{2}R_{1}) - \Gamma(I_{2}M_{2}R_{1})$ of Salem and Lee's [2] fitted values. To be rigorous, the $L_3M_4\alpha_2$ line width should be used for calculating the difference $\Delta\Gamma(L_{2,3})$: but we preferred to base comparison on the $L_3M_5lpha_1$ line width since the $Llpha_1$ line is the much stronger line and the width of the M4 level is nearly the same³ as the width of the M_5 level except for $Z \ge 90$. While the width differences $\Delta\Gamma(L_{2,3})$ derived from L xray emission lines agree well with the differences obtained in this work, those derived from the $K\alpha$ doublet, $\Delta\Gamma(L_{2,3}) = \Gamma(KL_2\alpha_2) - \Gamma(KL_3\alpha_1)$ are in very poor agreement. In fact, the large scatter in the values based on Kα width determinations indicates a need for improved measurements of $K\alpha$ line widths.

^aLike L Coster-Kronig rates (see ref (1)), the Coster-Kronig rate $M_4 \sim M_3 X$ is probably greatly overestimated by theory [47].

Table 1. Natural widths of K and L levels, $K\alpha$ x-ray lines, and several KLL Auger lines in eV. (FWHM)^a

•		Le	vel		X-ray line		Auger line		
Element	K	L_1	L ₂	L ₃	Kα ₁	Kα ₂	KL_1L_1	KL ₁ L ₂	KL ₂ L ₃
10 Ne	0.24	(<0.1)	(0.000)	(0.000)	0.24	0.24	(<0.4)	(<0.3)	0.24
11 Na	0.30	(0.2)	(0.000)	(0.000)	0.30	0.30	(0.7)	(0.5)	0.30
12 Mg	0.36	0.41	0.001	0.001	0.36	0.36	1.19	0.77	0.36
13 A1	0.42	0.73	0.004	0.004	0.43	0.43	1.89	1.16	0.43
14 Si	0.48	1.03	0.015	0.014	0.49	0.49	2.54	1.52	0.51
15 P	0.53	1.26	0.032	0.033	0.57	0.56	3.04	1.82	0.60
16 S	0.59	1.49	0.054	0.054	0.65	0.64	3.56	2.13	0.70
17 CI	0.64	1.58	0.083	0.087	0.72	0.72	3.80	2.30	0.81
18 Ar	0.68	1.63	0.126	0.128	0.81	0.80	3.93	2.43	0.93
19 K	0.74	1.92	0.152	0.156	0.89	0.89	4.57	2.81	1.05
20 Ca	0.81	2.07	0.17	0.17	0.98	0.98	4.94	3.04	1.15
21 Sc	0.86	2.21	0.19	0.19	1.05	1.06	5.27	3.26	1.25
22 Ti	0.94	2.34	0.24	0.22	1.16	1.18	5.62	3.52	1.40
23 V	1.01	2.41	0.26	0.24	1.26	1.28	5.84	3.69	1.52
24 Cr 25 Mn	1.08 1.16	2.54	0.29	0.27	1.35	1.37	6.15	3.90	1.64
26 Fe	1.16	2.62 2.76	0.34 0.37	0.32	1.48	1.50	6.40	4.12	1.82
27 Co	1.33	2.79	0.37	0.36	1.61	1.62	6.77	4.38	1.99
28 Ni	1.44	2.89	0.52	0.43	1.76 1.94	1.76	6.91	4.56	2.19
29 Cu	1.55	3.06	0.62	0.56	2.11	1.96 2.17	7.22 7.67	4.85 5.23	2.44 2.73
30 Zn	1.67	3.28	0.72	0.65	2.32	2.39	8.22	5.66	3.04
31 Ga	1.82	3.38	0.83	0.76	2.59	2.66	8.59	6.04	3.42
32 Ge	1.96	3.53	0.95	0.82	2,78	2.92	9.02	6.45	3.73
33 As	2.14	3.79	1.03	0.94	3.08	3.17	9.71	6.96	4.11
34 Se	2.33	3.94	1.13	1.00	3.33	3.46	10.2	7.40	4.46
35 Br	2,52	4.11	1.21	1.08	3.60	3.73	10.8	7.84	4.81
36 Kr	2.75	4.28	1.31	1.17	3.92	4.06	11.3	8.34	5.23
37 Rb	2.99	4.44	1.43	1.27	4.26	4.92	11.9	8.85	5.69
38 Sr	3.25	4.67	1.54	1.39	4.63	4.79	12.6	9.45	6.17
39 Y	3.52	4.71	1.65	1.50	5.02	5.18	13.0	9.89	6.68
40 Zr 41 Nb	3.84	4.78	1.78	1.57	5.40	5.62	13.4	10.4	7.18
41 No 42 Mo	4.14 4.52	3.94	1.87	1.66	5.80	6.01	12.0	9.9	7.67
42 MO 43 Tc	4.32	4.25 4.36	1.97	1.78	6.31	6.49	13.0	10.7	8.28
44 Ru	5.33	4.58	2.08 2.23	1.91	6.82	6.99	13.6	11.4	8.90
45 Rh	5.77	4.73	2.23	2.00	7.33	7.56	14.5	12.1	9.55
46 Pd	6.24	4.93	2.43	2.13	7.90	8.12	15.2	12.9	10.2
47 Ag	6.75	4.88	2.43	2.25 2.40	8.49	8.67	16.1	13.6	10.9
48 Cd	7.28	4.87	2.62	2.50	9.16 9.79	9.32	16.5	14.2	11.7
49 In	7.91	5.00	2.72	2.65	10.56	9.91 10.63	17.0 17.9	14.8 15.6	12.4 13.3
50 Sn	8.49	2.97	2.84	2.75	11.2	11.3	14.4	14.3	
51 Sb	9.16	3.13	3.00	2.87	12.0	12.2	15.4	15.3	14.1 15.0
52 Te	9.89	3.32	3.12	2.95	12.8	13.0	16.5	15.3	15.0
53 I	10.6	3.46	3.25	3.08	13.7	13.8	17.5	17.3	16.9
54 Xe	11.4	3.64	3.40	3.13	14.6	14.8	18.7	18.5	18.0
55 Cs	12.3	3.78	3.51	3.25	15.6	15.8	19.9	19.6	19.1
56 Ba	13.2	3.92	3.57	3.32	16.5	16.8	21.1	20.7	20.1
57 La	14.1	4.06	3.68	3.41	17.6	17.8	22.3	21.9	21.2
58 Ce	15.1	4.21	3.80	3.48	18.6	18.9	23.6	23.2	22.5
59 Pr	16.2	4.34	3.89	3.60	19.8	20.1	24.9	24.4	23.7

Table 1. Natural widths of K and L levels, $K\alpha$ x-ray lines, and several KLL Auger lines in eV. $(FWHM)^a$ — continued.

T		L	evel		X-ra	y line	Auger line		
Element	K	L ₁	L ₂	L ₃	Kα ₁	Kα ₂	$\overline{KL_1L_1}$	KL ₁ L ₂	KL ₂ L ₃
60 Nd	17.3	4.52	3.97	3.65	20.9	21.3	26.3	25.8	24.9
61 Pm	18.5	4.67	4.06	3.75	22.2	22.5	27.8	27.2	26.3
62 Sm	19.7	4.80	4.15	3.86	23.6	23.8	29.3	28.6	27.7
63 Eu	21.0	4.91	4.23	3.91	24.9	25.2	30.8	30.1	29.1
64 Gd	22.3	5.05	4.32	4.01	26.4	26.7	32.4	31.7	30.7
65 Tb	23.8	5.19	4.43	4.12	27.9	28.2	34.1	33.4	32.3
66 Dy	25.2	5.25	4.55	4.17	29.4	29.8	35.7	35.0	34.0
67 Ho	26.8	5.33	4.66	4.26	31.1	31.5	37.5	36.8	35.7
68 Er	28.4	5.43	4.73	4.35	32.7	33.1	39.2	38.5	37.5
69 Tm	30.1	5.47	4.79	4.48	34.6	34.9	41.0	40.3	39.3
70 Yb	31.9	5.53	4.82	4.60	36.5	36.7	42.9	42.2	41.3
71 Lu	33.7	5.54	4.92	4.68	38.4	38.7	44.8	44.2	43.4
72 Hf	35.7	5.63	5.02	4.80	40.5	40.7	47.0	46.3	45.5
73 Ta	37.7	5.58	5.15	4.88	42.6	42.9	48.9	48.5	47.8
74 W	39.9	5.61	5.33	4.98	44.9	45.2	51.1	50.8	50.2
75 Re	42.1	6.18	5.48	5.04	47.2	47.6	54.5	53.8	52.6
76 Os	44.4	7.25	5.59	5.16	49.6	50.0	58.9	57.2	55.1
77 Ir	46.8	8.30	5.69	5.25	52.1	52.5	63.4	60.8	57.7
78 Pt	49.3	9.39	5.86	5.31	54.6	55.2	68.1	64.6	60.5
79 A u	52.0	10.5	6.00	5.41	57.4	58.0	72.9	68.4	63.4
80 Hg	54.6	11.3	6.17	5.50	60.1	60.8	77.3	72.1	66.3
81 T1	57.4	12.0	6.32	5.65	63.1	63.8	81.4	75.7	69.4
82 Pb	60.4	12.2	6.48	5.81	66.2	66.8	84.8	79.1	72.7
83 Bi	63.4.	12.4	6.67	5.98	69.4	70.1	88.2	82.5	76.1
84 Po	66.6	12.6	6.83	6.13	72.7	73.4	91.9	86.1	79.6
85 At	69.8	12.8	7.01	6.29	76.1	76.8	95.5	89.7	83.1
86 Rn	73.3	13.1	7.20	6.41	79.7	80.5	99.4	93.5	86.9
87 Fr	76.8	13.3	7.47	6.65	83.4	84.2	103	97.5	90.9
88 Ra 89 Ac	80.4 84.1	13.4 13.6	7.68 7.95	6.82 6.98	87.2 91.1	88.1 92.0	107 111	101	94.9
90 Th	88.0	13.7	8.18	7.13				106	99.0
91 Pa	91.9	14.3	8.75	7.13	95.2	96.2	115	110	103
92 U	96.1	14.0	9.32	7.33 7.43	99.3	100.7	121	115	108
93 Np	100	14.0	9.91	7.43 7.59	103.5	105.4	124	119	113
94 Pu	105	13.5	10.5	7.82	108	110	128	124	118
95 Am	109	13.3	10.9	8.04	113	115	132	129	123
96 Cm	114	13.6	11.4		117	120	136	133	128
97 Bk	119	13.8	11.8	8.26 8.55	122	125	141	139	133
98.Cf	124	14.0	12.2	8.75	127 132	130	146	144	138
99 Es	129	14.3	12.7	9.04	132	136 141	151 157	150 155	144 150
100 Fm	135	14.4	13.1	9.33	144	148	164	162	
101 Md	140	14.8	13.6	9.61	150	153	170	162	157
102 No	145	15.1	14.0	9.90	155	159	176		163
103 Lr	150	15.9	14.4	10.1	161	165	182	175	169
104	156	16.2	14.9	10.5	167	171	182	181	175
105	162	16.5	15.5	10.8	173	177	195	188	182
106	168	16.8	15.9	11.2	179	184		194	188
107	174	17.0	16.4	11.6	185	190	201	200	195
108	181	17.2	17.0	12.0	193	198	208	207	202
109	187	17.6	17.5	12.3	200	205	215 222	215 222	210 217
110	193	18.1	18.1	12.7	206	211	229	229	224

^a Semi-empirical values (eqs. (1) to (3)); values in parentheses are estimated. Uncertainties are listed in Table 2.

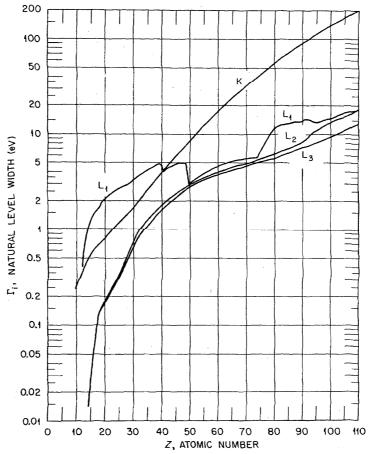


FIGURE 1. Natural widths of K and L levels. Graphic display of table 1.

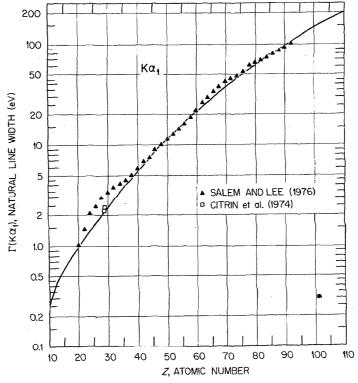


Figure 2. Natural widths of $K\alpha_1$ x-ray lines. Graphic display of table 1. Values of reference 2 are also shown.

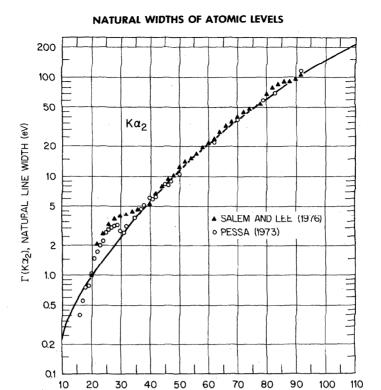


FIGURE 3. Natural widths of $K\alpha_2$ x-ray lines. Graphic display of table 1. Values of references 2 and 3 are also shown.

Z, ATOMIC NUMBER

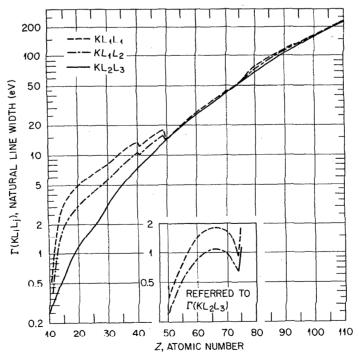


FIGURE 4. Natural widths of three KLL Auger lines. Graphic display of table 1. Note magnification of width differences shown in inset for $50 \le Z \le 75$.

Table 2. Estimated uncertainties (in percent) of semi-empirical values of natural level and line widths.

Level or Line										
	10-20	20-30	30-40	40-50	50-60	60-70	70-80	80-90	90-100	100-110
K	10ª	5	5	4	4	3	3	3	4	5
L_1	30 ^a	30a	30 ^b	25 ^b	20	15	15	15	20	20
L_2	25ª	25ª	25 ^b	20	10	10	10	10	10 ^b	15
L ₃	25a	25ª	20	15	10	10	. 8	8	8	10
Kα _{1.2} .	10 ^a	.7ª	6	5	4	4	3 -	3	4	5
$KL_1L_1^c$	25a	25ª	20 ^b	10 ^b	8	5	5	5	5	6
$KL_1L_2^c$	25°	20 ^a	15 ^b	10 ^b	5	5	5	4	4	6
$KL_2L_3^c$	10 ^a	10a	10	8	5	5	4	3	4	6

Table 3. Experimental free-atom data compared with semi-empirical natural level and line widths (in eV)

F1	Level or	Natural wi	dth				
Element	Line ^a	Semi-empirical ^b	Expt.	Notes and references			
10 Ne	K	0.24	0.27(1)	Photoelectron; Svensson et al., Ref. 14			
	$K\alpha_{1,2}$	0.24°	0.27(5)	Electron exc.; Agren et al., Ref. 15			
	KL_2L_3	0.24	0.21(2)	Proton exc.; Stolterfoht, Ref. 16			
11 Na	$L_1 L_{2.3} M_1$	0.2	0.24(2)	Electron exc.; Breuckmann et al., Ref. 17			
18 Ar	K	0.68	0.68(3)	X-ray absorption; Watanabe, Ref. 18			
	KL_1L_1	3.93	7.5(9)	Electron exc.; Krause, Ref. 19			
	KL_1L_2	2.43	3.6(3)	Electron exc.; Krause, Ref. 19			
	KL_2L_3	0.93	1.0(1)	Electron exc.; Krause, Ref. 19			
	KL_2L_3	0.93	0.92	Electron exc.; Asplund et al., Ref. 20			
	$L_1 M_{2,3}$	1.63	1.3	Electron exc.; Nordgren et al., Ref. 21			
	L ₁	1.63	1.8(2)	Electron exc.; Mehlhorn, Ref. 22			
	$L_{2,3}M_1$	0.127	0.12	Electron exc.; Nordgren et al., Ref. 21			
	$L_{2,3}M_{2,3}M_{2,3}$	0.127	0.10(2)	Proton exc.; Ridder et al., Ref. 23			
	2,5 2,5 2,5		0.10(2)	Electron exc.; Kondow et al., Ref. 24			
			0.16(2)	Electron exc.; Mehlhorn, Ref. 22			
30 Zn	$L_{2,3}M_{4,5}M_{4,5}$	0.70	0.50(15)	Electron exc.; Aksela et al., Ref. 25			
36 Kr	$L_{2}^{-2,3}$ 4,3 4,3	1.31	1.3	X-ray absorption; Wuilleumier, Ref. 26			
	L ₃	1.17	1.2	X-ray absorption; Wuilleumier, Ref. 26			

[&]quot;Widths of M levels are negligibly small.

 $[^]a$ Uncertainties from neglect of multi-body effects not included. b Near Coster-Kronig discontinuities, uncertainties may exceed those listed; cutoffs or onsets may be shifted for KL_1L_1 Auger transitions.

^cErrors due to the use of approximate formula eq. 3 are not included.

 $^{^{\}mathrm{b}}$ This work.

^cMultiply by ≈ 1.15 to include splitting of doublet.

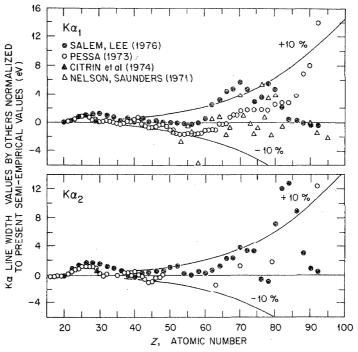


Figure 5. $K\alpha$ line widths recommended by Salem and Lee [2] and by Pessa [3], and reported by Citrin et al., [30], and by Nelson and Saunders [46] are normalized to present semi-empirical values. For clarity, only the $K\alpha_1$ values of reference 46 are shown. Note the arbitrarily chosen \pm 10% uncertainty limits.

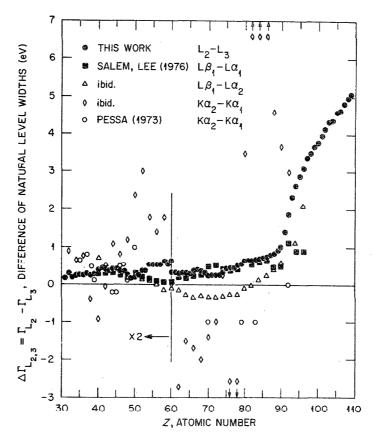


Figure 6. Differences of natural level widths, $\Delta\Gamma(L_{2,3})=\Gamma(L_2)-\Gamma(L_3),$ from table 1 and other sources [2,3].

Acknowledgements

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Note Added in Proof

Since submitting this manuscript, the following data on level and line widths have come to my attention: $ArL_{2.3}=0.127(10)$ eV [G. C. King et al., J. Phys. B12, 2479 (1978)]; $EuK\alpha_1=24.8(4)$ eV, $TmK\alpha_1=34.4(7)$ eV, $TaK\alpha_1=42.9(1.1)$ eV [K. C. Wang et al., Phys. Rev. A17, 1735 (1978)]; $NeK\alpha_{1,2}=0.24$ eV [T. P. Tooman and R. J. Liefeld, $Int.\ Conf.\ Phys.\ X-Ray\ Spectra$, Extended abstracts, R. D. Deslattes, editor, pp. 352–354, NBS (1976)]; and $N(K\to\pi 2p)=0.128(6)$ eV in N_2 [G. C. King et al., Chem. Phys. Lett. 52, 50 (1977)]. On the basis of the line width of 0.04 eV measured for $Ar(L_{2.3}M_{2.3}\to M_1M_{2.3})$ [21], the widths of NeK reported in reference 16 and $ArL_{2.3}$ reported in reference 23 should be changed to 0.23(3) eV and 0.12(2) eV, respectively, [D. Ridder, private communication (1978)].