Today's outline - February 16, 2023



- Mermin's interpretation of parallelism
- Simon's algorithm
- Distributed computation
- The Fourier transform

Reading Assignment: Reiffel: 7.8, 8.1–8.2 Wong: 7.6–7.7

Homework Assignment #04: due Friday, February 17, 2023

Homework Assignment #05: due Thursday, March 02, 2023

Mermin's interpretation



David Mermin proposed a simpler interpretation for how quantum algorithms and the solution to the Bernstein-Vazirani problem, in particular

Consider a C_{not} acting on the Hadamard basis

$$\begin{split} C_{not}|++\rangle &= C_{not} \frac{1}{2} (|00\rangle + |01\rangle + |10\rangle + |11\rangle) = \frac{1}{2} (|00\rangle + |01\rangle + |11\rangle + |10\rangle) = |++\rangle \\ C_{not}|+-\rangle &= C_{not} \frac{1}{2} (|00\rangle - |01\rangle + |10\rangle - |11\rangle) = \frac{1}{2} (|00\rangle - |01\rangle + |11\rangle - |10\rangle) = |--\rangle \\ C_{not}|-+\rangle &= C_{not} \frac{1}{2} (|00\rangle + |01\rangle - |10\rangle - |11\rangle) = \frac{1}{2} (|00\rangle + |01\rangle - |11\rangle - |10\rangle) = |-+\rangle \\ C_{not}|--\rangle &= C_{not} \frac{1}{2} (|00\rangle - |01\rangle - |10\rangle + |11\rangle) = \frac{1}{2} (|00\rangle - |01\rangle - |11\rangle + |10\rangle) = |+-\rangle \end{split}$$

If we then apply the Hadamard transform to each bit the resulting truth table becomes

This is simply a C_{not} gate applied to the high order qubit controlled by the low order qubit

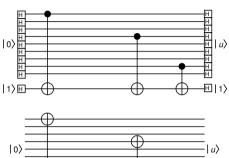
Mermin's interpretation



This insight leads to a simple way to look at the black box for U_{f_u}

- 1. Prepare an *n*-qubit register $|0\rangle_n$
- 2. Prepare an ancilla qubit $|a\rangle = |1\rangle$
- 3. Apply the Hadamard gate to all qubits
- **4**. Place a $C_{not}|u_i\rangle|a\rangle$ for each $u_i=1$
- 5. Apply the Hadamard gate to all qubits

The net effect is to have the ancilla bit "turn on" each qubit in the unknown, $C_{not}|a\rangle|u_i\rangle$ where $u_i=1$



From this perspective there is no quantum parallelism but simply a discrete circuit which produces the desired outcome

Of course, this presupposes that one knows what $|u\rangle$ is so we are peering into the black box

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Simon's problem – description



| Suppose we have a 2-to-1 function $f(x)$ such that $f(x) = f(x \oplus a)$ where a is secret and both x and a are n bit strings | X | f(x) |
|--|-----|------|
| For example, when $n=3$ we might have the table | 000 | 111 |
| There are 4 values for $f(x)$, each appearing twice, once in the top | 001 | 000 |
| half of the table and once in the bottom | | 110 |
| The goal of the algorithm is to find the the secret string a | 011 | 010 |
| Classically, this can be done by querying the function until we obtain two identical values for $f(x)$ and then calculate $a = x_0 \oplus x_1$ This can take up to $2^{n-1} + 1$ queries so the computation is $O(2^n)$ | | 000 |
| | | 111 |
| | | 010 |
| | | 110 |

In this case, we can see that $a=010\oplus 111=101$ and this holds for all matched pairs in the table

In contrast, Simon's quantum algorithm is a calculation which is O(n)

Simon's algorithm – quantum circuit



The problem requires two registers of n bits each which we designate with $|0\rangle_n$ and $|0\rangle_n$ as input and output registers, respectively

$$|\phi_{0}\rangle = |0\rangle_{n}|0\rangle_{n}$$

$$|\phi_{1}\rangle = |W|\otimes I(|0\rangle_{n}|0\rangle_{n}) = \frac{1}{\sqrt{2^{n}}} \sum_{x=0}^{2^{n}-1} |x\rangle|0\rangle_{n}$$

$$|\phi_{2}\rangle = \frac{1}{\sqrt{2^{n}}} \sum_{x=0}^{2^{n}-1} |x\rangle|f(x)\rangle$$

$$|\phi_{3}\rangle = \frac{1}{\sqrt{2}} (|x_{0}\rangle + |x_{0} \oplus a\rangle) |f(x_{0})\rangle$$

$$\begin{aligned} |\phi_{4}\rangle &= \frac{\sqrt{2}}{W} \otimes I \left[\frac{1}{\sqrt{2}} \left(|\mathbf{x}_{0}\rangle + |\mathbf{x}_{0} \oplus \mathbf{a}\rangle \right) |f(\mathbf{x}_{0})\rangle \right] = \frac{1}{\sqrt{2^{n}}} \frac{1}{\sqrt{2}} \sum_{y=0}^{2^{n}-1} \left[(-1)^{\mathbf{x}_{0} \cdot y} + (-1)^{(\mathbf{x}_{0} \oplus \mathbf{a}) \cdot y} \right] |\mathbf{y}\rangle |f(\mathbf{x}_{0})\rangle \\ &= \frac{1}{\sqrt{2^{n+1}}} \sum_{y=0}^{2^{n}-1} (-1)^{\mathbf{x}_{0} \cdot y} \left[1 + (-1)^{\mathbf{a} \cdot y} \right] |\mathbf{y}\rangle |f(\mathbf{x}_{0})\rangle \end{aligned}$$

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Simon's algorithm – quantum circuit



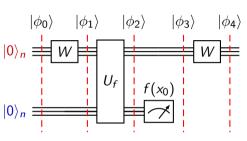
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Dropping the $|f(x_0)\rangle$ as it has already been measured, we have

$$|\phi_4\rangle = \frac{1}{\sqrt{2^{n+1}}} \sum_{y=0}^{2^n-1} (-1)^{x_0 \cdot y} \left[1 + (-1)^{a \cdot y}\right] \frac{|y\rangle}{|y\rangle}$$

There are two cases to consider for the modulo 2 scalar product $a \cdot y$

$$y \cdot a \neq 0 \quad \longrightarrow \quad |\phi_4\rangle \equiv 0$$



The second case is for $a \cdot y = 0$, in which case

$$|\phi_4\rangle = \frac{1}{\sqrt{2^{n+1}}} \sum_{y=0}^{2^n-1} (-1)^{x_0 \cdot y} [1+1] |y\rangle = \frac{1}{\sqrt{2^{n-1}}} \sum_{y=0}^{2^n-1} (-1)^{x_0 \cdot y} |y\rangle$$

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This is a superposition of 2^n possible states, one of which will be observed when $|\phi_4\rangle$ is measured

If n-1 linearly independent $|y\rangle$ are measured, it is possible to solve $y \cdot a = 0$

| Simon's algorithm – example | X | f(x) |
|---|------|------|
| Suppose a system with $n=4$ and $a=1001$, $f(x)$ has the truth table | 0000 | 1111 |
| | 0001 | 0001 |
| $ \phi_0\rangle = 0\rangle 0\rangle = 0000\rangle 0000\rangle$ $ \phi_0\rangle = \phi_1\rangle = \phi_2\rangle = \phi_3\rangle = \phi_4\rangle$ | 0010 | 1110 |
| $1 \sum_{i=1}^{15} \frac{1}{i} \frac{1}{$ | 0011 | 1101 |
| $ \phi_1\rangle = \frac{1}{4} \sum_{15} \mathbf{x}\rangle 0000\rangle$ $ 0\rangle_4 = W$ | 0100 | 0000 |
| x=0 | 0101 | 0101 |
| $ \phi_2\rangle = \frac{1}{4} \sum_{k=0}^{15} x\rangle f(x)\rangle$ | 0110 | 1010 |
| $ \phi_2\rangle = \frac{1}{4} \sum_{x=0}^{ x\rangle r(x)\rangle} 0\rangle_4$ | 0111 | 1001 |
| X=0 | 1000 | 0001 |
| $ \phi_3\rangle=rac{1}{\sqrt{2}}\left[\mathbf{x}_0\rangle+ \mathbf{x}_0\oplus\mathbf{a}\rangle\right] f(\mathbf{x}_0)\rangle$ For example, suppose $f(\mathbf{x}_0)=1010$ | 1001 | 1111 |
| | 1010 | 1101 |
| $ \phi_3\rangle=rac{[0110\rangle+ 1111\rangle]}{\sqrt{2}} f(x_0)\rangle$ now apply the Walsh transformation | 1011 | 1110 |
| v = | 1100 | 0101 |
| $ \phi_4 angle = rac{[0000 angle - 0010 angle - 0100 angle + 0110 angle + 1001 angle - 1011 angle - 1101 angle + 1111 angle]}{\sqrt{8}}$ | 1101 | 0000 |
| · | 1110 | 1001 |
| Note that any value of $ f(x_0)\rangle$ measured will result in these 8 $ x_0\rangle$ | 1111 | 1010 |

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Simon's algorithm – example



$$|\phi_4\rangle = \frac{1}{\sqrt{8}} [|0000\rangle - |0010\rangle - |0100\rangle + |0110\rangle + |1001\rangle - |1011\rangle - |1101\rangle + |1111\rangle]$$

The result of the final measurement, $|y\rangle$ will be one of these eight values and each of them should satisfy the linear equation $a \cdot y = 0$

Since we know that $a=|1001\rangle$ for this example, we can check this identity and the other 6 have the same properties

$$|1001\rangle \cdot |0000\rangle = 1 \cdot 0 + 0 \cdot 0 + 0 \cdot 0 + 1 \cdot 0 = 0$$

$$|1001\rangle \cdot |1001\rangle = 1 \cdot 1 + 0 \cdot 0 + 0 \cdot 0 + 1 \cdot 1 = 2 = 0$$

It is now necessary to collect n-1=3 independent values of $|y\rangle$ to solve for a

| Trial | $ y\rangle$ | Indep.? |
|-------|----------------|---------|
| 1 | 0000⟩ | No |
| 1 | $ 0010\rangle$ | Yes |
| 1 | $ 0100\rangle$ | Yes |
| 1 | $ 0110\rangle$ | No |
| 1 | $ 1001\rangle$ | Yes |

Create a matrix from the $y \cdot a = 0$ equation and the three independent values obtained

$$\begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} a_3 \\ a_2 \\ a_1 \\ a_0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

Simon's algorithm – example



Solve this matrix equation by Gaussian elimination

$$\begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} a_3 \\ a_2 \\ a_1 \\ a_0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

Convert the matrix to an upper triangular form by swapping rows 1 and 3

$$\begin{bmatrix} 1 & 0 & 0 & 1 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} a_3 \\ a_2 \\ a_1 \\ a_0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

Since the bottom row of the matrix is all zeros, a_0 can be either 0 or 1

$$a_0 = 0$$

 $a_1 = 0, \quad a_2 = 0, \quad a_3 + a_0 = 0$
 $a_3 = 0 \longrightarrow \quad a = |0000\rangle$

trivial, incorrect solution

$$a_0 = 1$$
 $a_1 = 0, \quad a_2 = 0, \quad a_3 + a_0 = 0$ $a_3 = -1 = 1 \quad \longrightarrow \quad a = |1001\rangle$

correct solution

Distributed computation



Alice and Bob are each provided with an $N=2^n$ bit number, u and v respectively

Alice must compute an n-bit number a and Bob must compute an n-bit number b such that

$$d_H(u, v) = 0 \longrightarrow a = b$$
 $d_H(u, v) = N/2 \longrightarrow a \neq b$
else \longrightarrow no condition on a and b

This is a challenging problem because u and v are exponentially larger than a and b

A classical solution requires a communication of at least N/2 bits but with enough entangled pairs, no additional communication is needed in a quantum solution

Start with *n* entangled pairs of particles, (a_i, b_i) in states $\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$ with

$$a_0, a_1, \ldots, a_{n-1}, b_0, b_1, \ldots, b_{n-1} \longrightarrow |\psi\rangle = \frac{1}{\sqrt{N}} \sum_{i=0}^{N-1} |i, i\rangle$$

Distributed computation



Alice uses the phase change subroutine with $f(i) = u_i$

Bob uses the phase change subroutine with $f(i) = v_i$

$$\sum_{i=0}^{N-1} |i
angle \longrightarrow \sum_{i=0}^{N-1} (-1)^{u_i} |i
angle$$

$$\sum_{i=0}^{N-1} |i
angle \longrightarrow \sum_{i=0}^{N-1} (-1)^{\mathsf{v}_i} |i
angle$$

They each apply the Walsh transformation to get a common global state

$$|\psi
angle = rac{1}{\sqrt{N}}\sum_{i=0}^{N-1}(-1)^{u_i\oplus v_i}\left(rac{\mathcal{W}}{|i
angle}
ight)\otimes \mathcal{W}|i
angle
ight) = rac{1}{N\sqrt{N}}\sum_{i=0}^{N-1}\sum_{i=0}^{N-1}\sum_{k=0}^{N-1}(-1)^{u_i\oplus v_i}(-1)^{i\cdot j}(-1)^{i\cdot k}|jk
angle$$

The probability that the measurement results in a=x=b is the modulus squared of $\langle x,x|\psi\rangle$

$$\langle x, x | \psi \rangle = \frac{1}{N\sqrt{N}} \sum_{i=0}^{N-1} (-1)^{u_i \oplus v_i} (-1)^{i \cdot x} (-1)^{i \cdot x} = \frac{1}{N\sqrt{N}} \sum_{i=0}^{N-1} (-1)^{u_i \oplus v_i}$$

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Distributed computation



The probability that Alice and Bob measure the same n bit value, x, is given by

If u = v, then $(-1)^{u_i \oplus v_i} = 1$ so when summed over the N possible values of x, $P_{xx} = 1$ and Alice and Bob will measure a = b with probability 1

For $d_H(u, v) = N/2$ there will be exactly the same number of 1 and -1 values in the sum so $P_{xx} = 0$ and Alice and Bob will measure a = b with probability 0

$$P_{\mathsf{x}\mathsf{x}} = |\langle \mathsf{x}, \mathsf{x} | \psi \rangle|^2 = \left| \frac{1}{N\sqrt{N}} \sum_{i=0}^{N-1} (-1)^{u_i \oplus v_i} \right|^2$$

$$\langle \mathbf{x}, \mathbf{x} | \psi \rangle = \frac{1}{\sqrt{N}}$$

$$\langle \mathbf{x}, \mathbf{x} | \psi \rangle = \mathbf{0}$$

Quantum parallelism



The action of the quantum operator U_f on a maximally superposed state appears to do more computation than a classical computation of f(x)

$$U_f|x,0\rangle = \frac{1}{\sqrt{N}}\sum_{x=0}^{N-1}|x,f(x)\rangle$$

This is not the case, as a single measurement of an m qubit system can only result in a single m-qubit value

Exponential speedups through the use of quantum computation are not possible in general and in many cases there is no speedup at all

Even an efficient quantum algorithm cannot probe the vast space of an $n \otimes m$ qubit system

The value of quantum computing lies in two general techniques

Amplification of outputs of interest: By transforming the state in a way that the output values of interest have a higher probability of being measured

Measuring properties of the set of all f(x): An example is using the quantum Fourier Transform to determine the periodicity of f(x)

Discrete Fourier transform



The quantum Fourier transform is an important building block for many quantum algorithms

In order to develop the efficient implementation of the quantum Fourier transform, it is useful to start with the classical discrete and fast Fourier transforms

The discrete Fourier transform (DFT) is a linear transformation which takes a discrete column vector a(k) to a column vector of Fourier coefficients, A(x), where $0 \le k, x \le N-1$

$$A(x) = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} a(k) e^{2\pi i k x/N}$$

The DFT operator is an $N \times N$ matrix with elements

$$F_{xk} = \frac{1}{\sqrt{N}} e^{2\pi i k x/N}$$

Assume $a(k) = e^{-2\pi i u k/N}$ is a function of frequency u < N which evenly divides N

Computing the Fourier coefficients,

$$A(x) = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} a(k) e^{2\pi i k x/N} = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{-2\pi i u k/N} e^{2\pi i k x/N} = \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{2\pi i k (x-u)/N}$$

All are zero except for when $x - u = 0 \mod N$ so the only term which survives is A(u)