

## 第二章

# 均匀物质的热力学性质

## § 2.1 内能、焓、自由能和吉布斯函数的全微分

主要目的：

利用数学方法  热力学函数间的微分关系

### 一、数学定义

函数  $f(x, y)$  的全微分

$$df = \left( \frac{\partial f}{\partial x} \right)_y dx + \left( \frac{\partial f}{\partial y} \right)_x dy$$

自变量



状态参量( $p, S, V, T$ )

函数



热力学函数（态函数）( $U, H, F, G$ )

$$U = U(S, V) \quad H = H(S, p)$$

$$F = F(T, V) \quad G = G(T, p)$$

热力学的基本微分方程  $dU = TdS - pdV$

## 二、热力学量表示为偏导数

1、函数关系：  $U = U(S, V)$

全微分： 
$$dU = \left( \frac{\partial U}{\partial S} \right)_V dS + \left( \frac{\partial U}{\partial V} \right)_S dV$$

热力学基本方程 
$$dU = TdS - pdV$$

对比得： 
$$T = \left( \frac{\partial U}{\partial S} \right)_V \quad -p = \left( \frac{\partial U}{\partial V} \right)_S$$

## 2、函数关系：

$$H = H(S, p)$$

全微分：
$$dH = \left( \frac{\partial H}{\partial S} \right)_p dS + \left( \frac{\partial H}{\partial p} \right)_S dp$$

根据焓的定义 
$$H = U + pV$$

全微分：
$$\begin{aligned} dH &= dU + pdV + Vdp \\ &= TdS - pdV + pdV + Vdp = TdS + Vdp \end{aligned}$$

对比得：
$$T = \left( \frac{\partial H}{\partial S} \right)_p \quad V = \left( \frac{\partial H}{\partial p} \right)_S$$

### 3、函数关系：

$$F = F(T, V)$$

全微分：
$$dF = \left( \frac{\partial F}{\partial T} \right)_V dT + \left( \frac{\partial F}{\partial V} \right)_T dV$$

根据自由能的定义

$$F = U - TS$$

全微分：
$$\begin{aligned} dF &= dU - TdS - SdT \\ &= TdS - pdV - TdS - SdT = -SdT - pdV \end{aligned}$$

对比得：
$$-S = \left( \frac{\partial F}{\partial T} \right)_V \quad -p = \left( \frac{\partial F}{\partial V} \right)_T$$

4、函数关系:

$$G = G(T, p)$$

全微分: 
$$dG = \left( \frac{\partial G}{\partial T} \right)_p dT + \left( \frac{\partial G}{\partial p} \right)_T dp$$

根据吉布斯函数的定义 
$$G = U - TS + pV$$

全微分: 
$$\begin{aligned} dG &= dU - TdS - SdT + pdV + Vdp \\ &= TdS - pdV - TdS - SdT + pdV + Vdp \\ &= -SdT + Vdp \end{aligned}$$

对比得: 
$$-S = \left( \frac{\partial G}{\partial T} \right)_p \quad V = \left( \frac{\partial G}{\partial p} \right)_T$$

### 三、麦氏关系

求偏导数的次序可以交换  $\frac{\partial^2 f}{\partial x \partial y} = \frac{\partial^2 f}{\partial y \partial x}$

在函数关系  $U = U(S, V)$  中得到:

$$T = \left( \frac{\partial U}{\partial S} \right)_V \quad -p = \left( \frac{\partial U}{\partial V} \right)_S$$

$$\Rightarrow \left( \frac{\partial}{\partial V} T \right)_S = \frac{\partial}{\partial V} \left[ \left( \frac{\partial U}{\partial S} \right)_V \right]_S = \frac{\partial^2 U}{\partial V \partial S}$$
$$\left( \frac{\partial}{\partial S} (-p) \right)_V = \frac{\partial}{\partial S} \left[ \left( \frac{\partial U}{\partial V} \right)_S \right]_V = \frac{\partial^2 U}{\partial S \partial V}$$

$$\Rightarrow \left( \frac{\partial T}{\partial V} \right)_S = - \left( \frac{\partial p}{\partial S} \right)_V$$



在函数关系  $H = H(S, p)$  中得到:

$$T = \left( \frac{\partial H}{\partial S} \right)_p \quad V = \left( \frac{\partial H}{\partial p} \right)_S$$

$$\Rightarrow \left( \frac{\partial}{\partial p} T \right)_S = \frac{\partial}{\partial p} \left[ \left( \frac{\partial H}{\partial S} \right)_p \right]_S = \frac{\partial^2 H}{\partial p \partial S}$$

$$\left( \frac{\partial}{\partial S} V \right)_p = \frac{\partial}{\partial S} \left[ \left( \frac{\partial H}{\partial p} \right)_S \right]_p = \frac{\partial^2 H}{\partial S \partial p}$$

$$\Rightarrow \left( \frac{\partial T}{\partial p} \right)_S = \left( \frac{\partial V}{\partial S} \right)_p$$

在函数关系  $F = F(T, V)$  中得到:

$$-S = \left( \frac{\partial F}{\partial T} \right)_V \quad -p = \left( \frac{\partial F}{\partial V} \right)_T$$

$$\Rightarrow \left( \frac{\partial}{\partial V} (-S) \right)_T = \frac{\partial}{\partial V} \left[ \left( \frac{\partial F}{\partial T} \right)_V \right]_T = \frac{\partial^2 F}{\partial V \partial T}$$

$$\left( \frac{\partial}{\partial T} (-p) \right)_V = \frac{\partial}{\partial T} \left[ \left( \frac{\partial F}{\partial V} \right)_T \right]_V = \frac{\partial^2 F}{\partial T \partial V}$$

$$\Rightarrow \left( \frac{\partial S}{\partial V} \right)_T = \left( \frac{\partial p}{\partial T} \right)_V$$

在函数关系  $G = G(T, p)$  中得到:

$$-S = \left( \frac{\partial G}{\partial T} \right)_p \quad V = \left( \frac{\partial G}{\partial p} \right)_T$$

$$\Rightarrow \left( \frac{\partial}{\partial p} (-S) \right)_T = \frac{\partial}{\partial p} \left[ \left( \frac{\partial G}{\partial T} \right)_p \right]_T = \frac{\partial^2 G}{\partial p \partial T}$$

$$\left( \frac{\partial}{\partial T} V \right)_p = \frac{\partial}{\partial T} \left[ \left( \frac{\partial G}{\partial p} \right)_T \right]_p = \frac{\partial^2 G}{\partial T \partial p}$$

$$\Rightarrow \left( \frac{\partial S}{\partial p} \right)_T = - \left( \frac{\partial V}{\partial T} \right)_p$$

# 热力学微分关系

热力学函数	热力学基本方程	热力学偏导数	麦克斯韦关系
$U$	$dU = TdS - pdV$	$T = \left( \frac{\partial U}{\partial S} \right)_V$ $p = - \left( \frac{\partial U}{\partial V} \right)_S$	$\left( \frac{\partial T}{\partial V} \right)_S = - \left( \frac{\partial p}{\partial S} \right)_V$
$H = U + pV$	$dH = TdS + Vdp$	$T = \left( \frac{\partial H}{\partial S} \right)_p$ $V = \left( \frac{\partial H}{\partial p} \right)_S$	$\left( \frac{\partial T}{\partial p} \right)_S = \left( \frac{\partial V}{\partial S} \right)_p$
$F = U - TS$	$dF = -SdT - pdV$	$S = - \left( \frac{\partial F}{\partial T} \right)_V$ $p = - \left( \frac{\partial F}{\partial V} \right)_T$	$\left( \frac{\partial p}{\partial T} \right)_V = \left( \frac{\partial S}{\partial V} \right)_T$
$G = H - TS$	$dG = -SdT + Vdp$	$S = - \left( \frac{\partial G}{\partial T} \right)_p$ $V = \left( \frac{\partial G}{\partial p} \right)_T$	$- \left( \frac{\partial V}{\partial T} \right)_p = \left( \frac{\partial S}{\partial p} \right)_T$

## 说明:

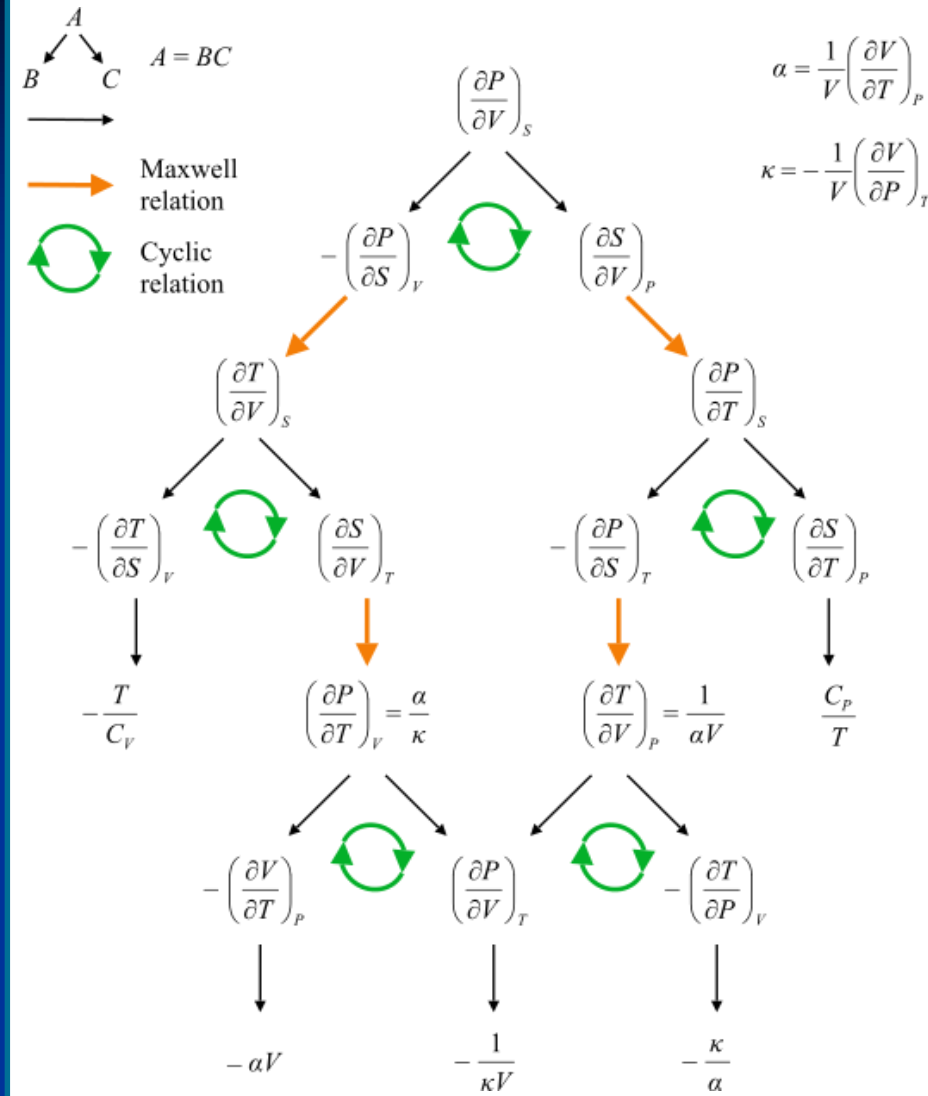
- 1、表中这套热力学关系是从热力学基本方程  $dU = TdS - pdV$  导出的，从变量变换的角度看，只可能导出其它三个基本方程。
- 2、利用表中关系，加上  $C_p$ 、 $C_V$  和附录一中的几个偏微分学公式，就可以研究均匀闭系的各种热力学性质。
- 3、表中关系是解决热力学问题的基础，应熟记它们。

简单记忆麦克斯韦关系的一种方法，如下：

$$\begin{array}{cc} T \longrightarrow V \\ p \longrightarrow S \end{array} \longrightarrow -\left(\frac{\partial T}{\partial V}\right)_S = \left(\frac{\partial p}{\partial S}\right)_V \quad -\left(\frac{\partial V}{\partial T}\right)_p = \left(\frac{\partial S}{\partial p}\right)_T$$

$$\begin{array}{cc} T \downarrow \\ p \end{array} \quad \begin{array}{cc} V \downarrow \\ S \end{array} \longrightarrow \left(\frac{\partial T}{\partial p}\right)_S = \left(\frac{\partial V}{\partial S}\right)_p \quad \left(\frac{\partial p}{\partial T}\right)_V = \left(\frac{\partial S}{\partial V}\right)_T$$

# 麦氏关系



[https://en.wikipedia.org/wiki/Maxwell\\_relations](https://en.wikipedia.org/wiki/Maxwell_relations)

## § 2.2 麦氏关系的简单应用

上节导出了麦氏关系：

$$(1) \left( \frac{\partial T}{\partial V} \right)_S = - \left( \frac{\partial p}{\partial S} \right)_V$$

$$dU = TdS - pdV$$

$$(2) \left( \frac{\partial T}{\partial p} \right)_S = \left( \frac{\partial V}{\partial S} \right)_p$$

$$dH = TdS + Vdp$$

$$(3) \left( \frac{\partial p}{\partial T} \right)_V = \left( \frac{\partial S}{\partial V} \right)_T$$

$$dF = -SdT - pdV$$

$$(4) \left( \frac{\partial V}{\partial T} \right)_p = - \left( \frac{\partial S}{\partial p} \right)_T$$

$$dG = -SdT + Vdp$$

麦氏关系给出了热力学量的偏导数之间的关系。利用麦氏关系，可以把一些不能直接从实验测量的物理量以物态方程（或体胀系数和等温压缩系数）和热容量等可以直接从实验测量的物理量表达出来表示出来。

一、选 $T$ 、 $V$ 为状态参量，熵为： $S = S(T, V)$

内能为： $U = U(S, V) = U(S(T, V), V) = U(T, V)$

$$\text{全微分: } dS = \left( \frac{\partial S}{\partial T} \right)_V dT + \left( \frac{\partial S}{\partial V} \right)_T dV$$

$$dU = \left( \frac{\partial U}{\partial T} \right)_V dT + \left( \frac{\partial U}{\partial V} \right)_T dV$$

$$\begin{aligned} dU = TdS - pdV &= T \left[ \left( \frac{\partial S}{\partial T} \right)_V dT + \left( \frac{\partial S}{\partial V} \right)_T dV \right] - pdV \\ &= T \left( \frac{\partial S}{\partial T} \right)_V dT + \left[ T \left( \frac{\partial S}{\partial V} \right)_T - p \right] dV \end{aligned}$$



利用麦氏关系： $\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial p}{\partial T}\right)_V$

可得： $dU = T\left(\frac{\partial S}{\partial T}\right)_V dT + \left[T\left(\frac{\partial p}{\partial T}\right)_V - p\right]dV$

对比得： $C_V = \left(\frac{\partial U}{\partial T}\right)_V = T\left(\frac{\partial S}{\partial T}\right)_V$

$$\left(\frac{\partial U}{\partial V}\right)_T = T\left(\frac{\partial p}{\partial T}\right)_V - p$$

公式  $\left(\frac{\partial U}{\partial V}\right)_T = T\left(\frac{\partial p}{\partial T}\right)_V - p$  的意义:

对于理想气体:  $pV_m = RT$

→ 
$$\left(\frac{\partial U_m}{\partial V_m}\right)_T = T \frac{R}{V_m} - p = \frac{TR - pV_m}{V_m} = 0$$

对于理想气体, 内能只是温度的函数。

焦耳定律

对于范氏气体：

$$\left(p + \frac{a}{V_m^2}\right)(V_m - b) = RT$$

由

$$\left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial p}{\partial T}\right)_V - p$$



$$\left(\frac{\partial U_m}{\partial V_m}\right)_T = \frac{RT}{V_m - b} - p = \frac{a}{V_m^2}$$

温度不变时，范氏气体的内能随体积的变化率。  
与体积有关。

二、选 $T$ 、 $p$ 为状态参量，熵为： $S = S(T, p)$

焓为： $H = H(S, p) = H(S(T, p), p) = H(T, p)$

全微分： $dS = \left(\frac{\partial S}{\partial T}\right)_p dT + \left(\frac{\partial S}{\partial p}\right)_T dp$      $dH = \left(\frac{\partial H}{\partial T}\right)_p dT + \left(\frac{\partial H}{\partial p}\right)_T dp$

热力学基本方程： $dH = TdS + Vdp$

$$dH = T \left(\frac{\partial S}{\partial T}\right)_p dT + \left[ T \left(\frac{\partial S}{\partial p}\right)_T + V \right] dp$$

对比得： $C_p = \left(\frac{\partial H}{\partial T}\right)_p = T \left(\frac{\partial S}{\partial T}\right)_p$

利用麦氏关系：

$$\left(\frac{\partial S}{\partial p}\right)_T = - \left(\frac{\partial V}{\partial T}\right)_p$$

$$\left(\frac{\partial H}{\partial p}\right)_T = T \left(\frac{\partial S}{\partial p}\right)_T + V = V - T \left(\frac{\partial V}{\partial T}\right)_p$$

三、选  $p$ 、 $V$  为状态参量，熵为：  $S = S(p, V)$

$$U = U(S, V) = U(S(p, V), V) = U(p, V)$$

$$dS = \left( \frac{\partial S}{\partial p} \right)_V dp + \left( \frac{\partial S}{\partial V} \right)_p dV \quad dU = \left( \frac{\partial U}{\partial p} \right)_V dp + \left( \frac{\partial U}{\partial V} \right)_p dV$$

$$dU = TdS - pdV = T \left[ \left( \frac{\partial S}{\partial p} \right)_V dp + \left( \frac{\partial S}{\partial V} \right)_p dV \right] - pdV$$

$$= T \left( \frac{\partial S}{\partial p} \right)_V dp + \left[ T \left( \frac{\partial S}{\partial V} \right)_p - p \right] dV$$

利用麦氏关系：  $\left( \frac{\partial p}{\partial T} \right)_V = \left( \frac{\partial S}{\partial V} \right)_T$   $\left( \frac{\partial V}{\partial T} \right)_p = - \left( \frac{\partial S}{\partial p} \right)_T$

$$dU = -T \left( \frac{\partial V}{\partial T} \right)_s dp + \left[ T \left( \frac{\partial p}{\partial T} \right)_s - p \right] dV$$

对比得：

$$\left( \frac{\partial U}{\partial p} \right)_V = -T \left( \frac{\partial V}{\partial T} \right)_s \quad \left( \frac{\partial U}{\partial V} \right)_p = T \left( \frac{\partial p}{\partial T} \right)_s - p$$


## 四、计算任意简单系统的定压热容量与定容热容量之差

$$\text{由 } C_p - C_V = T \left( \frac{\partial S}{\partial T} \right)_p - T \left( \frac{\partial S}{\partial T} \right)_V$$

$$S(T, p) = S(T, V(T, p))$$

$$f(x, z) = f(x, y(x, z))$$

$$\left( \frac{\partial f}{\partial x} \right)_z = \left( \frac{\partial f}{\partial x} \right)_y + \left( \frac{\partial f}{\partial y} \right)_x \left( \frac{\partial y}{\partial x} \right)_z$$

$$\left( \frac{\partial S}{\partial T} \right)_p = \left( \frac{\partial S}{\partial T} \right)_V + \left( \frac{\partial S}{\partial V} \right)_T \left( \frac{\partial V}{\partial T} \right)_p$$


$$C_p - C_V = T \left( \frac{\partial S}{\partial V} \right)_T \left( \frac{\partial V}{\partial T} \right)_p = T \left( \frac{\partial p}{\partial T} \right)_V \left( \frac{\partial V}{\partial T} \right)_p$$

对于理想气体

$$= nR$$

利用麦氏关系:  $\left( \frac{\partial S}{\partial V} \right)_T = \left( \frac{\partial p}{\partial T} \right)_V$

对于任意简单系统

$$= TpV \frac{1}{p} \left( \frac{\partial p}{\partial T} \right)_V \left( \frac{1}{V} \left( \frac{\partial V}{\partial T} \right)_p \right) = TpV \beta \alpha = \frac{VT \alpha^2}{\kappa_T}$$

$\alpha = \kappa_T \beta p$

固体与液体的  $C_V$  很难测量，通过  $C_p$  计算之。

## 附 雅可比行列式

$x, y$  是状态参量,  $u$  和  $v$  是热力学函数:

$$u(x, y), \quad v(x, y).$$

### 雅可比行列式定义

$$\frac{\partial(u, v)}{\partial(x, y)} = \begin{vmatrix} \frac{\partial u}{\partial x} & \frac{\partial u}{\partial y} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y} \end{vmatrix} = \frac{\partial u}{\partial x} \frac{\partial v}{\partial y} - \frac{\partial u}{\partial y} \frac{\partial v}{\partial x}$$

性质:

$$1) \quad \left( \frac{\partial u}{\partial x} \right)_y = \frac{\partial(u, y)}{\partial(x, y)} = \begin{vmatrix} \frac{\partial u}{\partial x} & \frac{\partial u}{\partial y} \\ \frac{\partial y}{\partial x} & \frac{\partial y}{\partial y} \end{vmatrix} = \frac{\partial u}{\partial x} \cdot 1 - \frac{\partial u}{\partial y} \cdot 0$$

$$2) \quad \frac{\partial(u, v)}{\partial(x, y)} = - \frac{\partial(v, u)}{\partial(x, y)} \quad \left| \begin{array}{cc} \frac{\partial u}{\partial x} & \frac{\partial u}{\partial y} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y} \end{array} \right| = - \left| \begin{array}{cc} \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y} \\ \frac{\partial u}{\partial x} & \frac{\partial u}{\partial y} \end{array} \right|$$

$$3) \quad \frac{\partial(u, v)}{\partial(x, y)} = \frac{\partial(u, v)}{\partial(x, s)} \frac{\partial(x, s)}{\partial(x, y)} \quad \left| \begin{array}{cc} \frac{\partial u}{\partial x} & \frac{\partial u}{\partial s} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial s} \end{array} \right| \left| \begin{array}{cc} \frac{\partial x}{\partial x} & \frac{\partial x}{\partial y} \\ \frac{\partial s}{\partial x} & \frac{\partial s}{\partial y} \end{array} \right| = \left| \begin{array}{ccc} \frac{\partial u}{\partial x} & \frac{\partial u}{\partial s} & \frac{\partial s}{\partial y} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial s} & \frac{\partial s}{\partial y} \end{array} \right| = \left| \begin{array}{cc} \frac{\partial u}{\partial x} & \frac{\partial u}{\partial y} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y} \end{array} \right|$$

$$4) \quad \frac{\partial(u, v)}{\partial(x, y)} = 1 / \frac{\partial(x, y)}{\partial(u, v)}$$



例1、求证绝热压缩系数与等温压缩系数之比等于定容热容量与定压热容量之比。

$$\kappa_S = -\frac{1}{V} \left( \frac{\partial V}{\partial p} \right)_S, \quad \kappa_T = -\frac{1}{V} \left( \frac{\partial V}{\partial p} \right)_T$$

证明：


$$\begin{aligned} \frac{\kappa_S}{\kappa_T} &= \frac{-\frac{1}{V} \left( \frac{\partial V}{\partial p} \right)_S}{-\frac{1}{V} \left( \frac{\partial V}{\partial p} \right)_T} = \frac{\frac{\partial(V, S)}{\partial(p, S)}}{\frac{\partial(V, T)}{\partial(p, T)}} = \frac{\frac{\partial(V, S)}{\partial(V, T)}}{\frac{\partial(V, S)}{\partial(p, T)}} \\ &= \frac{\left( \frac{\partial S}{\partial T} \right)_V}{\left( \frac{\partial S}{\partial T} \right)_p} = \frac{T \left( \frac{\partial S}{\partial T} \right)_V}{T \left( \frac{\partial S}{\partial T} \right)_p} = \frac{C_V}{C_p} \end{aligned}$$

$\left( \frac{\partial u}{\partial x} \right)_y = \frac{\partial(u, y)}{\partial(x, y)}$

## 例2、求证

$$C_p - C_V = -T \frac{\left(\frac{\partial p}{\partial T}\right)_V^2}{\left(\frac{\partial p}{\partial V}\right)_T}$$

证明:  $C_p = T \left(\frac{\partial S}{\partial T}\right)_p = T \frac{\partial(S, p)}{\partial(T, p)}$       利用麦氏关系:  $\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial p}{\partial T}\right)_V$



$$= T \frac{\frac{\partial(S, p)}{\partial(T, V)}}{\frac{\partial(T, p)}{\partial(T, V)}} = T \frac{\left(\frac{\partial S}{\partial T}\right)_V \left(\frac{\partial p}{\partial V}\right)_T - \left(\frac{\partial S}{\partial V}\right)_T \left(\frac{\partial p}{\partial T}\right)_V}{\left(\frac{\partial p}{\partial V}\right)_T} = C_V - T \frac{\left(\frac{\partial p}{\partial T}\right)_V^2}{\left(\frac{\partial p}{\partial V}\right)_T}$$

例3、考虑一理想气体，其熵为

$$S = \frac{n}{2} \left\{ \sigma + 5R \ln \frac{U}{n} + 2R \ln \frac{V}{n} \right\}$$

其中， $n$ 为摩尔数， $R$ 为气体常数， $U$ 为能量， $V$ 为体积， $\sigma$ 为常数，求出定压和定容热容量。

解：  $dU = TdS - pdV$

温度 $T$ 为  $T = \left( \frac{\partial U}{\partial S} \right)_V$   $\left( \frac{\partial S}{\partial U} \right)_V = \frac{1}{T} = \frac{5}{2} nR \frac{1}{U}$

$$U = \frac{5}{2} nRT \quad C_V = \frac{5}{2} nR$$

$$C_p = C_V + nR = \frac{7}{2} nR$$