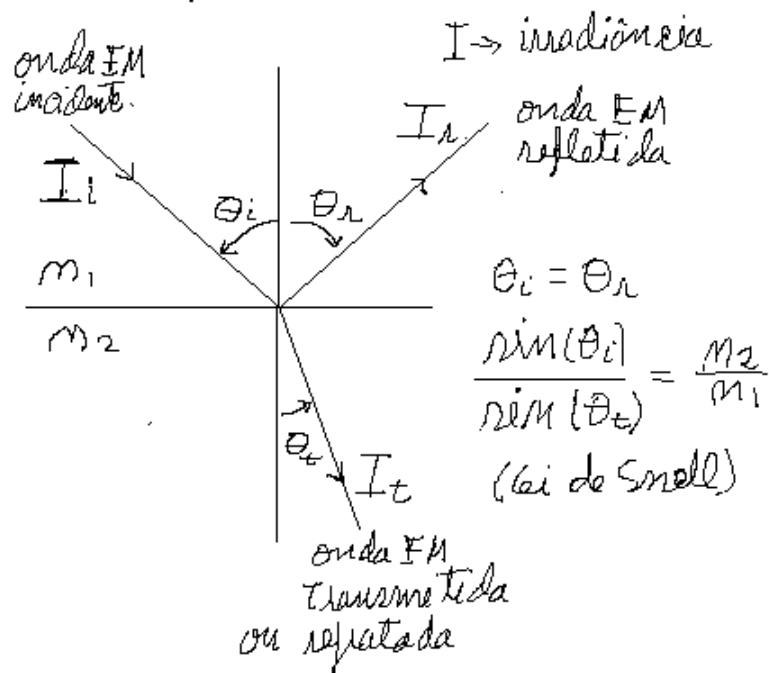


Equações de Fresnel, verificação experimental.

LET 2021

Reflexão e refração de ondas eletromagnéticas



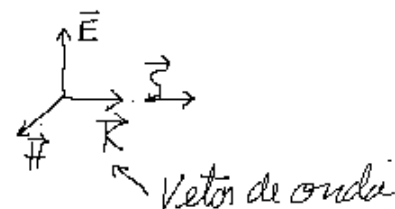
$$I_r (I_i) = ?$$

$$I_t (I_i) = ?$$

$$I = \langle \vec{S} \rangle, \quad \vec{S} = \vec{E} \times \vec{H} \quad [I] = [S] = \text{W m}^{-2}$$

Vetor de Poynting

onda EM:



Equação de ondas:

$$\nabla \times \nabla \times \vec{E} = -\frac{\partial}{\partial t} (\nabla \times \mu \vec{H}) = -\mu \frac{\partial}{\partial t} (\vec{J} + \frac{\partial \vec{D}}{\partial t}) = -\mu \frac{\partial \vec{J}}{\partial t} - \mu \epsilon \frac{\partial^2 \vec{E}}{\partial t^2}$$

Com $\vec{J} = 0 \rightarrow$ (meio não condutor)

$$\nabla \times \nabla \times \vec{E} = -\mu \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} \rightarrow \vec{\nabla} (\vec{\nabla} \cdot \vec{E}) - \nabla^2 \vec{E} = -\mu \epsilon \frac{\partial^2 \vec{E}}{\partial t^2}, \quad \vec{\nabla} \cdot \vec{D} = \rho \Rightarrow \vec{\nabla} \cdot \epsilon \vec{E} = \rho$$

Com $\rho = 0$ (meio não eletrizado)

$$\nabla^2 \vec{E} - \mu \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} = 0 \rightarrow \text{Equação de ondas p/ meio isolante não condutor e não eletrizado.}$$

Eq. Ondas admite soluções do tipo: $\vec{E} = \vec{E}_0 e^{i(\omega t - \vec{k} \cdot \vec{r})}$, verifiquemos em que condições.

Equações de Maxwell:

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}, \quad \vec{\nabla} \cdot \vec{B} = 0$$

$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}, \quad \nabla \cdot \vec{D} = \rho$$

Equações constitutivas:

$$\vec{B} = \mu \vec{H}, \quad \vec{D} = \epsilon \vec{E}, \quad \vec{J} = \sigma_e \vec{E}$$

Eq. de ondas EM:

$$\nabla^2 \vec{E} - \epsilon \mu \frac{\partial^2 \vec{E}}{\partial t^2} = 0 = \frac{\partial^2 \vec{E}}{\partial x^2} + \frac{\partial^2 \vec{E}}{\partial y^2} + \frac{\partial^2 \vec{E}}{\partial z^2} - \epsilon \mu \frac{\partial^2 \vec{E}}{\partial t^2}$$



$$(-ik_x)^2 \vec{E} + (-ik_y)^2 \vec{E} + (-ik_z)^2 \vec{E} - \epsilon \mu (i\omega)^2 \vec{E} = 0$$

$$(-k_x^2 - k_y^2 - k_z^2 + \epsilon \mu \omega^2) \vec{E} = 0$$

$$-k^2 + \epsilon \mu \omega^2 = 0$$

$$k = \sqrt{\epsilon \mu} \omega$$

$$k = \sqrt{\epsilon \mu} k v$$

$$v = \frac{1}{\sqrt{\epsilon \mu}} \Rightarrow k = \frac{\omega}{v}, \quad m = \frac{e}{v},$$

$$k = \frac{2\pi}{\lambda} = m \frac{2\pi}{\lambda_0} \quad k = m \frac{\omega}{c} = m k_0 \quad k_0 = \frac{\omega}{c} = \frac{2\pi}{\lambda_0} = \frac{2\pi}{\lambda_0}$$

$$\lambda = \frac{\lambda_0}{m}$$

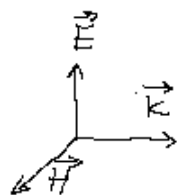
$$i(\omega t - \vec{k} \cdot \vec{r})$$

Relação entre \vec{E} e \vec{H} : $\vec{H} = \vec{H}_0 e$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} = -\mu \frac{\partial \vec{H}}{\partial t} = -\mu (i\omega) \vec{H}$$

$$\vec{\nabla} \times \vec{E} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & E_z \end{vmatrix} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ -ik_x & -ik_y & -ik_z \\ E_x & E_y & E_z \end{vmatrix} = -i \vec{k} \times \vec{E}$$

$$-i \vec{k} \times \vec{E} = -i \mu \omega \vec{H} \Rightarrow \vec{k} \times \vec{E} = \mu \omega \vec{H}$$



$$\Rightarrow H = \frac{k}{\mu \omega} E$$

Teste de solução tipo $\vec{E}(\vec{r}, t) = \vec{E}_0 e^{i(\omega t - \vec{k} \cdot \vec{r})}$
onda plana:

$$\frac{\partial \vec{E}}{\partial x} = \vec{E}_0 e^{i(\omega t - k_x x - k_y y - k_z z)} (-ik_x)$$

$$\frac{\partial^2 \vec{E}}{\partial x^2} = (-ik_x)^2 \vec{E}_0 e^{i(\omega t - \vec{k} \cdot \vec{r})} = (-ik_x)^2 \vec{E}$$

$$\frac{\partial^2 \vec{E}}{\partial t^2} = (i\omega)^2 \vec{E}_0 e^{i(\omega t - \vec{k} \cdot \vec{r})} = (i\omega)^2 \vec{E}$$

$$\phi = \omega t - \vec{k} \cdot \vec{r} \quad \frac{d\phi}{dt} = 0 = \omega - \vec{k} \cdot \frac{d\vec{r}}{dt}$$

$$\omega = \vec{k} \cdot \vec{v} \rightarrow \omega = k v$$

$$\omega = k v$$

Relação entre \vec{k} e \vec{E} :

$$\vec{\nabla} \cdot \vec{E} = \vec{\nabla} \cdot \left(\frac{1}{\epsilon} \vec{D} \right) = \frac{1}{\epsilon} \vec{\nabla} \cdot \vec{D} = 0 \quad \left(\text{meio não eletrizado} \right)$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = E_{0x} e^{i(\omega t - \vec{k} \cdot \vec{r})} (-ik_x) +$$

$$+ E_{0y} e^{i(\omega t - \vec{k} \cdot \vec{r})} (-ik_y) + E_{0z} e^{i(\omega t - \vec{k} \cdot \vec{r})} (-ik_z) =$$

$$= -i (E_x k_x + E_y k_y + E_z k_z) = -i (\vec{k} \cdot \vec{E}) = 0 \Rightarrow \vec{k} \perp \vec{E}$$

Verificação que $\vec{E} = \vec{E}_0 e^{i(\omega t - \vec{k} \cdot \vec{r})}$

é uma onda plana: encontrar lugar geométrico dos pontos tal que $\vec{E}(0, \vec{r}) = \vec{E}_0$:

$$\vec{E}(0, \vec{r}) = \vec{E}_0 e^{i \vec{k} \cdot \vec{r}} = \vec{E}_0 \Rightarrow \vec{k} \cdot \vec{r} = 2m\pi \Rightarrow k_x x + k_y y + k_z z = 2m\pi \rightarrow \text{equação de família de planos}$$

Verificação: $\vec{k} \cdot \vec{r}_1 = 2m\pi$
 $\vec{k} \cdot \vec{r}_2 = 2m\pi$

$$\vec{k} \cdot (\vec{r}_1 - \vec{r}_2) = 0 \Rightarrow \vec{k} \perp (\vec{r}_1 - \vec{r}_2) \rightarrow \text{vetor do plano}$$

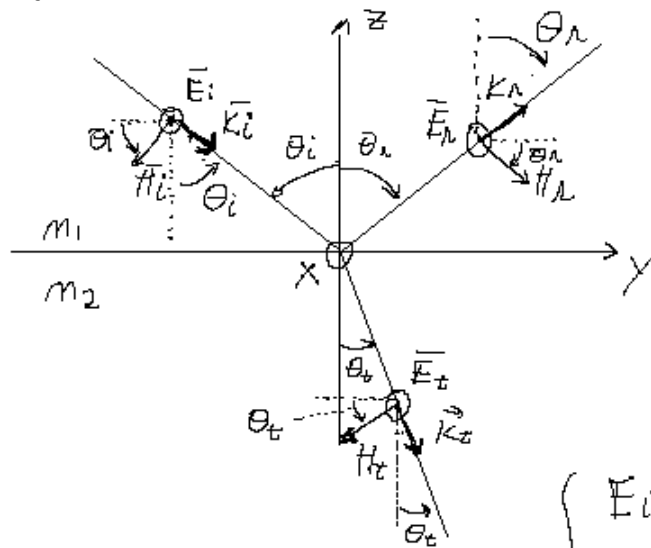
Equações de Fresnel:

Condições fronteira dos campos \vec{E} e \vec{H}

$$E_{\text{tangencial } 1} = E_{\text{tangencial } 2}$$

$$H_{\text{tangencial } 1} = H_{\text{tangencial } 2}$$

Polarização S ($\vec{E} \perp$ plano de incidência)



$E_m z=0$ podemos escrever:

$$\begin{cases} E_i + E_r = E_t \\ -H_i \cos \theta_i + H_r \cos \theta_r = -H_t \cos \theta_t \end{cases}$$

$$H = \frac{k}{\mu \omega} E$$

$$\begin{cases} E_{i0} e^{i(\omega t - k_y y)} + E_{r0} e^{i(\omega t - k_y y)} = E_{t0} e^{i(\omega t - k_y y)} \\ -\cos \theta_i H_{i0} e^{i(\omega t - k_y y)} + \cos \theta_r H_{r0} e^{i(\omega t - k_y y)} = -\cos \theta_t H_{t0} e^{i(\omega t - k_y y)} \end{cases}$$

$$\begin{cases} E_{i0} + E_{r0} = E_{t0} \\ -\cos \theta_i \frac{k_i}{\mu_1 \omega} E_{i0} + \cos \theta_r \frac{k_r}{\mu_1 \omega} E_{r0} = -\cos \theta_t \frac{k_t}{\mu_2 \omega} E_{t0} \end{cases}$$

$$\Rightarrow \begin{cases} \omega_i = \omega_r = \omega_t \\ k_{iy} = k_{ry} = k_{ty} \end{cases}$$

$$k_i \sin \theta_i = k_r \sin \theta_r = k_t \sin \theta_t$$

$$m_1 k_0 \sin \theta_i = m_1 k_0 \sin \theta_r$$

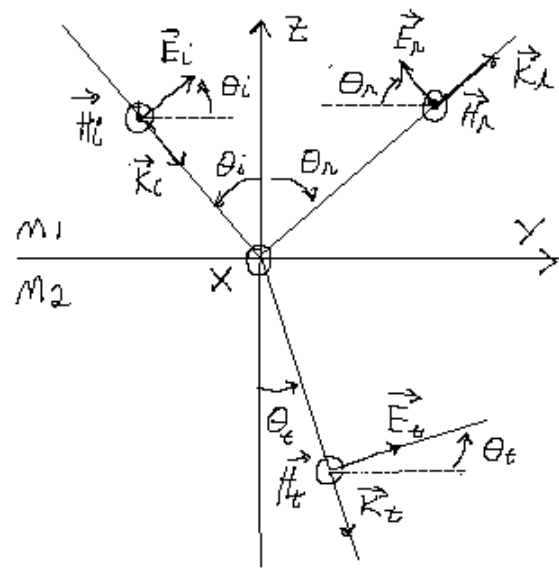
$$\theta_i = \theta_r$$

$$\text{Lei de Snell} \rightarrow m_1 \sin \theta_i = m_2 \sin \theta_t$$

$$\begin{cases} E_{i0} + E_{r0} = E_{t0} \\ -\cos\theta_i \frac{\mu_1}{\mu_2} E_{i0} + \cos\theta_r \frac{\mu_1}{\mu_2} E_{r0} = -\cos\theta_t \frac{\mu_1}{\mu_2} E_{t0} \end{cases} \quad \begin{cases} E_{i0} + E_{r0} = E_{t0} \\ (E_{i0} - E_{r0}) \cos\theta_i \frac{\mu_1}{\mu_2} = E_{t0} \cos\theta_t \frac{\mu_1}{\mu_2} \end{cases}$$

$$\begin{cases} \frac{E_{r0}}{E_{i0}} = \frac{\frac{\mu_1}{\mu_2} \cos\theta_i - \frac{\mu_2}{\mu_1} \cos\theta_t}{\frac{\mu_1}{\mu_2} \cos\theta_i + \frac{\mu_2}{\mu_1} \cos\theta_t} \\ \frac{E_{t0}}{E_{i0}} = \frac{2 \frac{\mu_1}{\mu_2} \cos\theta_i}{\frac{\mu_1}{\mu_2} \cos\theta_i + \frac{\mu_2}{\mu_1} \cos\theta_t} \end{cases} \quad \text{equações de Fresnel para a polarização S}$$

Equações de Fresnel (polarização P)



Condições fronteira dos campos \vec{E} e \vec{H} $E_{\text{tangencial 1}} = E_{\text{tangencial 2}}$

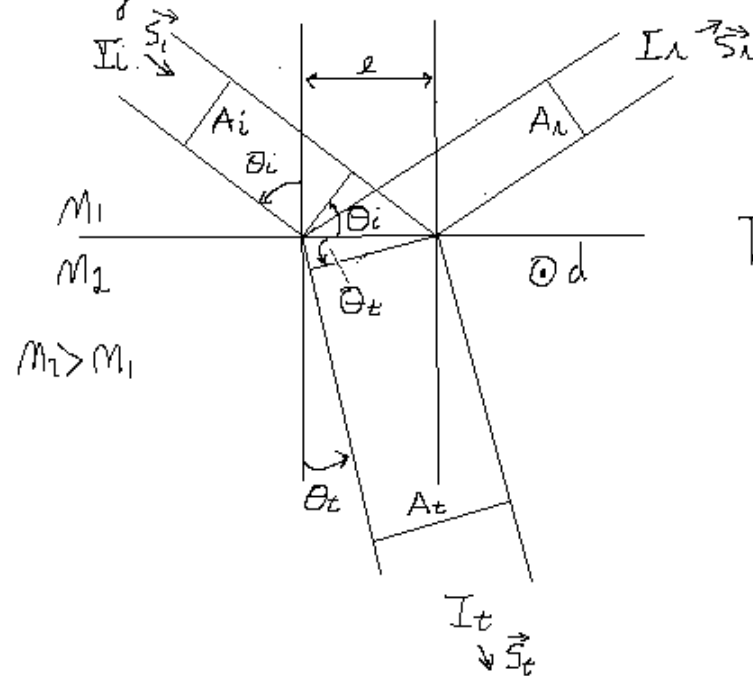
$H_{\text{tangencial 1}} = H_{\text{tangencial 2}}$

Em $z=0$ podemos escrever:

$$\begin{cases} H_i + H_r = H_t \\ E_i \cos\theta_i - E_r \cos\theta_r = E_t \cos\theta_t \end{cases} \quad \begin{cases} \frac{\mu_1}{\mu_2} E_i + \frac{\mu_1}{\mu_2} E_r = \frac{\mu_1}{\mu_2} E_t \\ E_i \cos\theta_i - E_r \cos\theta_r = E_t \cos\theta_t \end{cases}$$

$$\begin{cases} \frac{\mu_1}{\mu_2} (E_{i0} + E_{r0}) = \frac{\mu_1}{\mu_2} E_{t0} \\ E_{i0} \cos\theta_i - E_{r0} \cos\theta_r = E_{t0} \cos\theta_t \end{cases} \quad \begin{cases} \frac{E_{r0}}{E_{i0}} = \frac{\frac{\mu_2}{\mu_1} \cos\theta_i - \frac{\mu_1}{\mu_2} \cos\theta_t}{\frac{\mu_2}{\mu_1} \cos\theta_i + \frac{\mu_1}{\mu_2} \cos\theta_t} \\ \frac{E_{t0}}{E_{i0}} = \frac{2 \frac{\mu_1}{\mu_2} \cos\theta_i}{\frac{\mu_2}{\mu_1} \cos\theta_i + \frac{\mu_1}{\mu_2} \cos\theta_t} \end{cases} \quad \text{equações de Fresnel polarização P}$$

Refletância e Transmittância na superfície de refração entre dois meios LHI não condutores



Refletância: $R \equiv \frac{P_r}{P_i} = \frac{A_r I_r}{A_i I_i} = \frac{I_r}{I_i} = \frac{\langle |\vec{S}_r| \rangle}{\langle |\vec{S}_i| \rangle} = \frac{\langle |\vec{E}_r \times \vec{H}_r| \rangle}{\langle |\vec{E}_i \times \vec{H}_i| \rangle} = \frac{\langle E_r H_r \rangle}{\langle E_i H_i \rangle} = \left(\frac{E_{r0}}{E_{i0}} \right)^2$

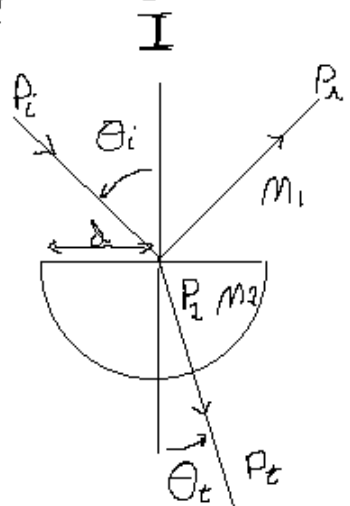
Transmittância: $T \equiv \frac{P_t}{P_i} = \frac{A_t I_t}{A_i I_i} = \frac{d \cos \theta_t I_t}{d \cos \theta_i I_i} = \frac{\cos \theta_t}{\cos \theta_i} \frac{I_t}{I_i} = \frac{\cos \theta_t \langle E_t H_t \rangle}{\cos \theta_i \langle E_i H_i \rangle} \Rightarrow$

$$T = \frac{\cos \theta_t}{\cos \theta_i} \frac{n_2}{n_1} \left(\frac{E_{t0}}{E_{i0}} \right)^2$$

$T + R = 1 \rightarrow$ Conservação da energia

Lei de Lambert-Beer: $I(d) = I_0 e^{-\alpha d}$
(meio com perdas)

Experiência laboratorial:

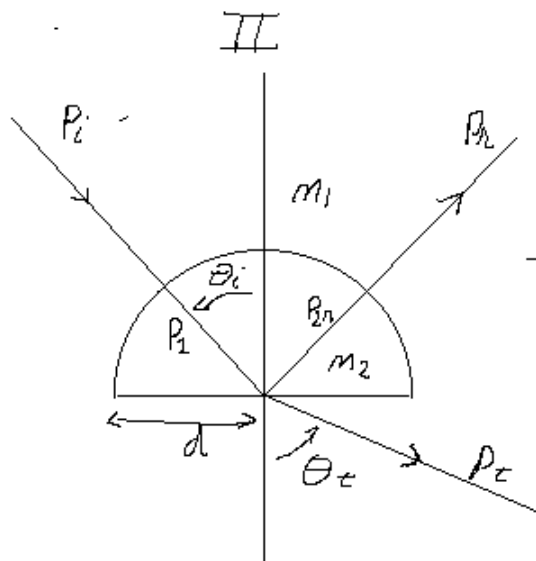


$$R_T = \frac{P_r}{P_i} = R_{12}(\theta_i)$$

$$T_T = \frac{P_t}{P_i} = \frac{P_t}{P_2} \times \frac{P_2}{P_i} = T_{21}(0) \times T_{12}(\theta_i)$$

$$R_T + T_T < 1 \quad T_{12}(0) = T_{21}(0) = \frac{4n_1 n_2}{(n_1 + n_2)^2} \Big|_{n_1 = n_2}$$

$T_{Tp} = T_T e^{-\alpha d}$ com perdas
 $R_{Tp} = R_T$



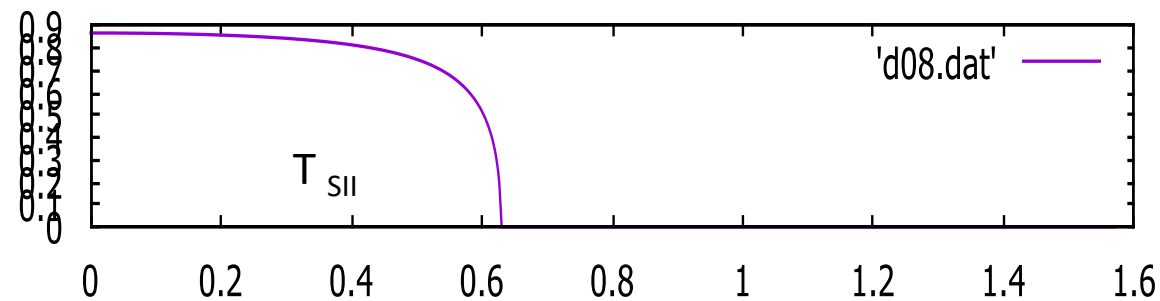
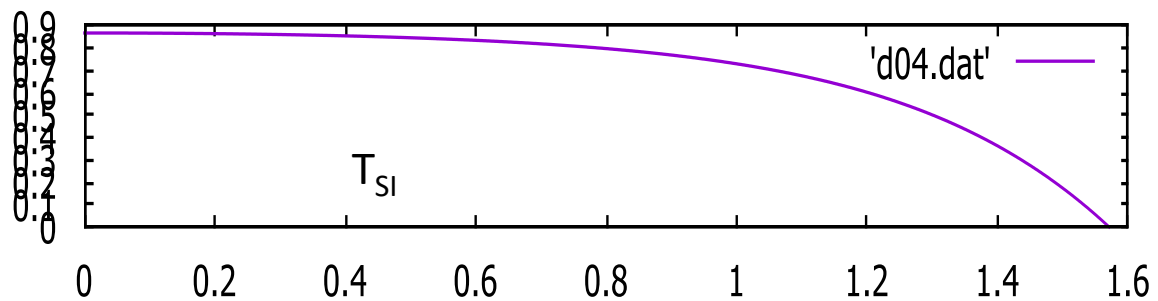
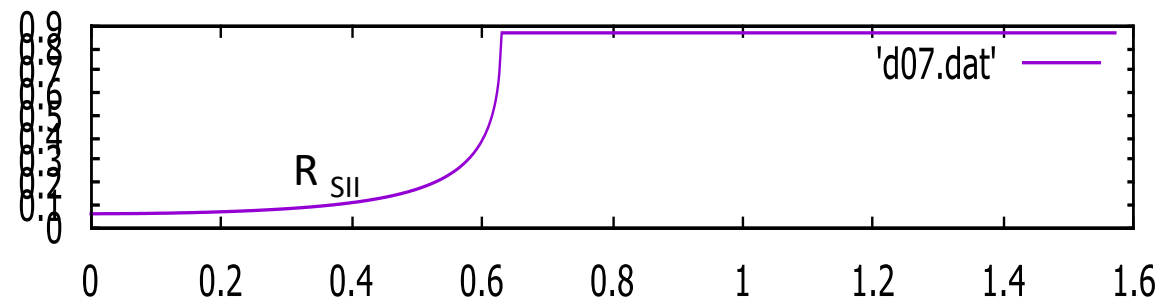
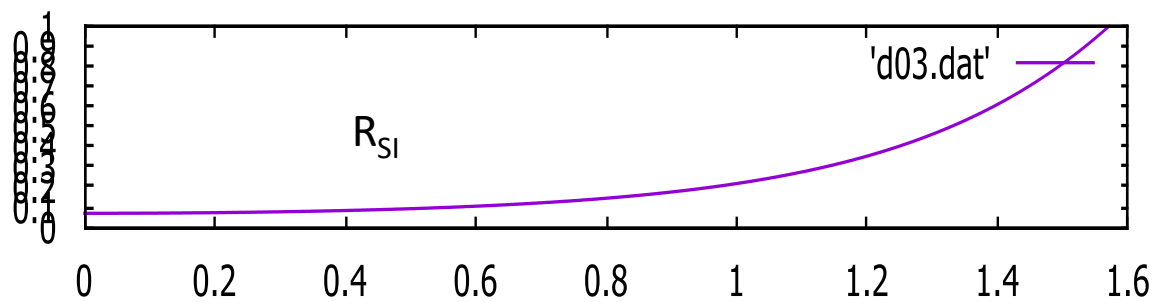
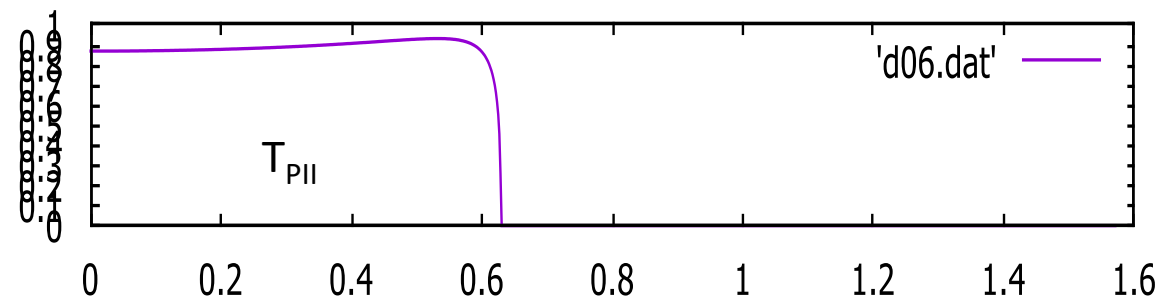
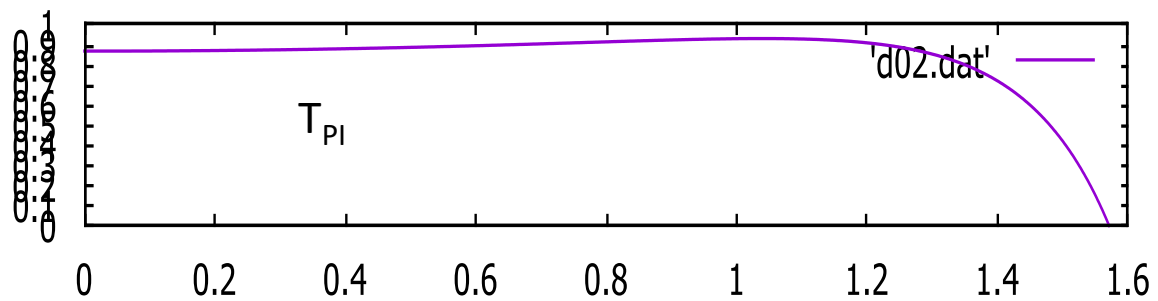
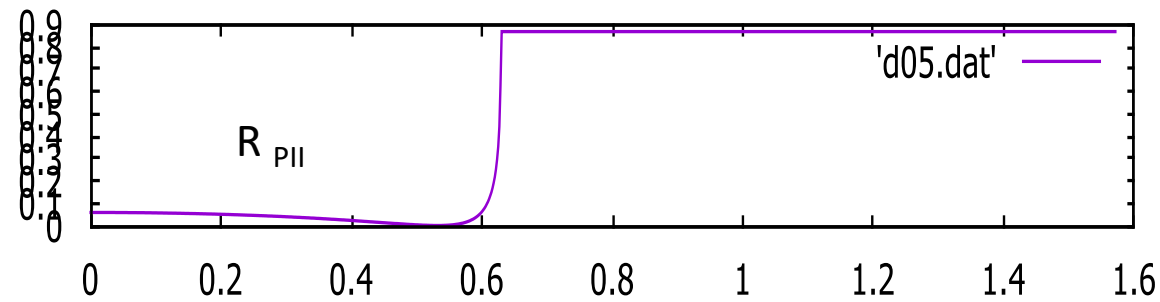
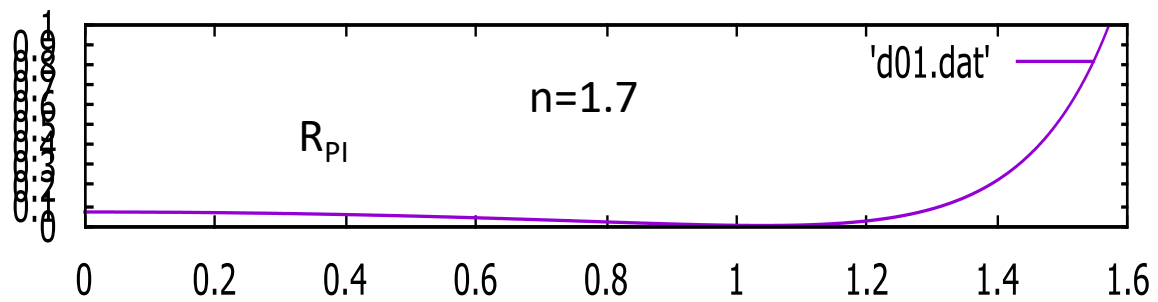
$$R_T = \frac{P_r}{P_i} = \frac{P_r}{P_{2L}} \times \frac{P_{2L}}{P_2} \times \frac{P_2}{P_i}$$

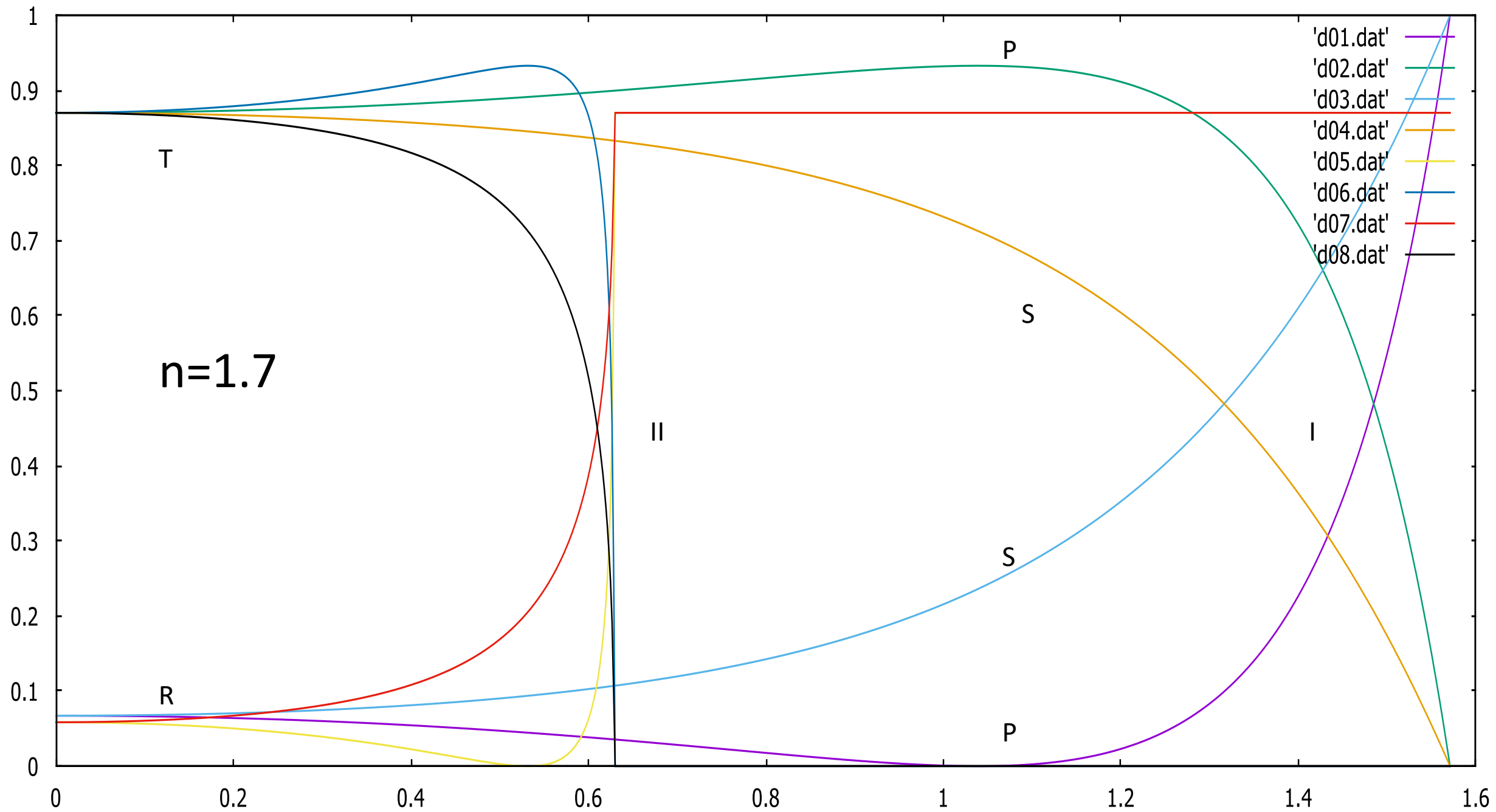
$$R_T = T_{21}(0) R_{21}(\theta_i) T_{12}(0)$$

$$T_T = \frac{P_t}{P_i} = \frac{P_t}{P_2} \times \frac{P_2}{P_i} = T_{21}(\theta_i) \times T_{12}(0)$$

$$R_T + T_T < 1$$

$R_{Tp} = R_T e^{-2\alpha d}$
 $T_{Tp} = T_T e^{-\alpha d}$ } com perdas





FIM