

## Estudo do Condensador

- Carga e descarga do condensador num circuito RC
- Variação com a frequência da permitividade elétrica de um dielétrico com perdas

# Sumário

- Conceitos básicos
- Tipos de condensadores
- Carga e descarga do condensador num circuito RC
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- Análise da permitividade elétrica, equação de Clausius-Mossotti
- Modelo elementar para a polarização eletrónica
- Permitividade elétrica complexa e condensador com perdas

## Estudo do Condensador, conceitos básicos

Dispositivo composto por 2  
condutores em presença:



$$C = \frac{Q}{V_1 - V_2}$$

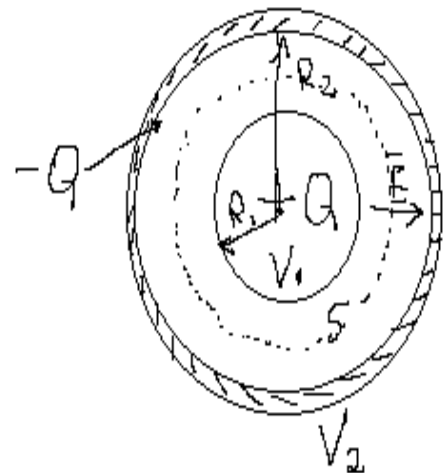
O condensador armazena carga e  
energia elétrica:

$$W_e = \frac{1}{2} \sum_{i=1}^2 Q_i V_i = \frac{1}{2} (Q_1 V_1 + Q_2 V_2) = \frac{1}{2} Q (V_1 - V_2) = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} C (V_1 - V_2)^2$$

# Diferentes tipos de Condensadores

## Condensador esférico

$$C = \frac{Q}{V_1 - V_2}$$



$$\oint_S \vec{E} \cdot \vec{N} dS = \frac{Q}{\epsilon_0} \quad (\text{I Gauss})$$

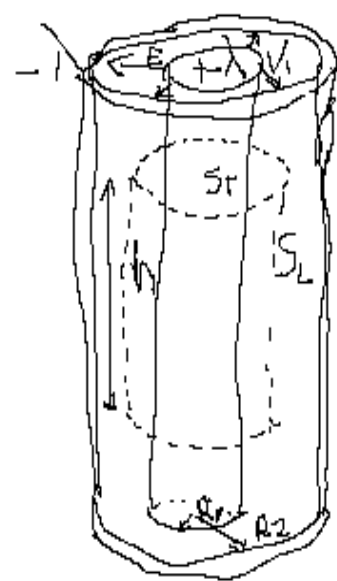
$$E \cdot S = \frac{Q}{\epsilon_0}$$

$$E = \frac{Q}{4\pi\epsilon_0} \frac{1}{r^2}$$

$$V_1 - V_2 = \int_{R_1}^{R_2} \vec{E} \cdot d\vec{p} = \int_{R_1}^{R_2} E \cdot dr = \frac{Q}{4\pi\epsilon_0} \left[ -\frac{1}{R_2} + \frac{1}{R_1} \right]$$

$$C = \frac{Q}{\frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{R_1} - \frac{1}{R_2} \right]} = \frac{4\pi\epsilon_0}{\frac{1}{R_1} - \frac{1}{R_2}}$$

## Condensador Cilindrico



$$C/L = \frac{Q/L}{V_1 - V_2}$$

$$\int_{S_L + S_T} \vec{E} \cdot \vec{n} ds = \frac{\lambda h}{\epsilon_0}$$

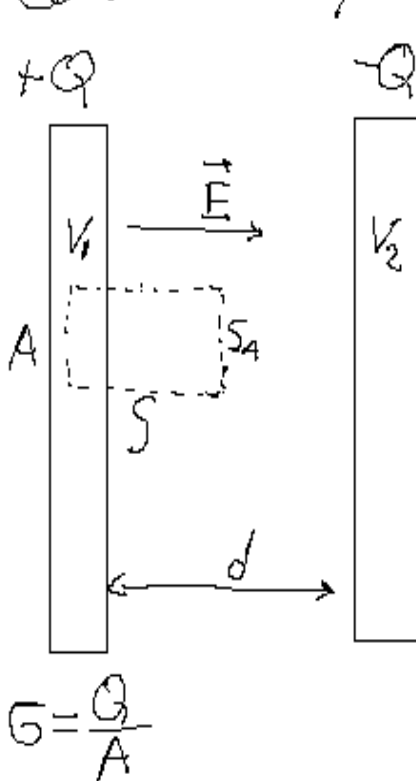
$$E S_L = \frac{\lambda h}{\epsilon_0}$$

$$E = \frac{\lambda h}{2\pi \lambda \epsilon_0}$$

$$V_1 - V_2 = \int_{R_1}^{R_2} \vec{E} \cdot d\vec{R} = \frac{\lambda}{2\pi \epsilon_0} \ln\left(\frac{R_2}{R_1}\right)$$

$$C/L = \frac{\lambda}{\frac{\lambda}{2\pi \epsilon_0} \ln\left(\frac{R_2}{R_1}\right)} = \frac{2\pi \epsilon_0}{\ln\left(\frac{R_2}{R_1}\right)}$$

## Condensador plano



$$C = \frac{Q}{V_1 - V_2}$$

$$\int_S \vec{E} \cdot \vec{n} ds = \frac{Q S_A}{\epsilon_0}$$

$$E S_A = \frac{Q S_A}{\epsilon_0}$$

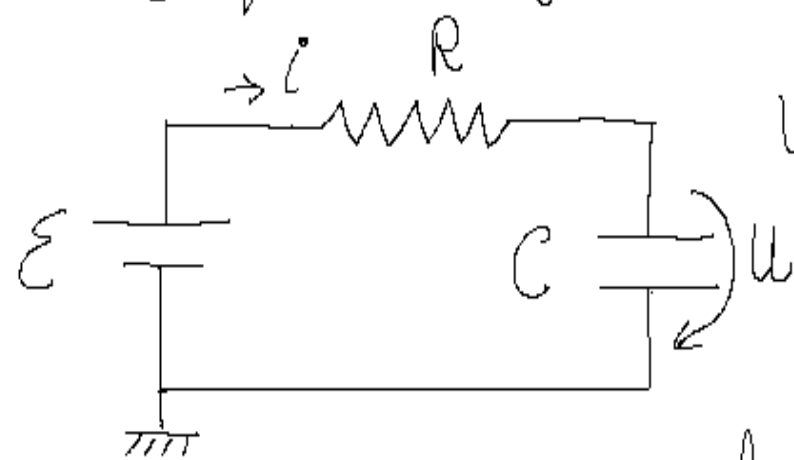
$$E = \frac{Q}{\epsilon_0}$$

$$V_1 - V_2 = \int_1^2 \vec{E} \cdot d\vec{R} = E d$$

$$V_1 - V_2 = \frac{Q d}{\epsilon_0}$$

$$C = \frac{Q A}{\frac{Q d}{\epsilon_0}} = \frac{\epsilon_0 A}{d}$$

# Carga e descarga do condensador num circuito RC



Lei de Kirchhoff:

$$C = \frac{Q}{U} \quad i = \frac{dQ}{dt} = C \frac{dU}{dt}$$

$$\varepsilon = Ri + U \rightarrow \varepsilon = RC \frac{dU}{dt} + U$$

Eq. Homogénea  $0 = RC \frac{dU}{dt} + U \rightarrow U = U_0 e^{-\frac{t}{RC}}$

Eq. Característica  $0 = RC \lambda + 1$   
 $\rightarrow \lambda = -\frac{1}{RC}$

$$U = U_0 e^{-t/RC} + \varepsilon$$

$$U(0) = 0 \Rightarrow U_0 = -\varepsilon$$

$$U = \varepsilon (1 - e^{-t/RC})$$

$$i = C \frac{dU}{dt} = \frac{C\varepsilon}{RC} e^{-t/RC}$$

## Balanco energético

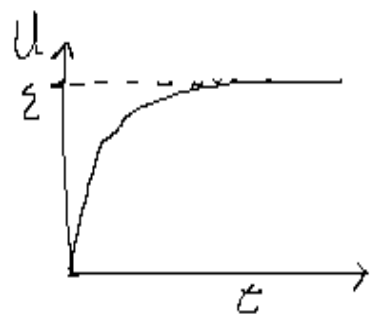
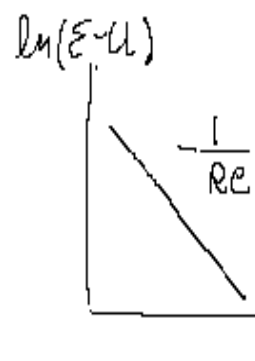
$$W_\varepsilon = \int_0^\infty \varepsilon i dt = \varepsilon \int_0^\infty i dt =$$

$$= \frac{\varepsilon^2}{R} \int_0^\infty e^{-t/RC} dt = \frac{\varepsilon^2}{R} \left[ -\frac{1}{1/RC} e^{-t/RC} \right]_0^\infty = \frac{\varepsilon^2 C}{2}$$

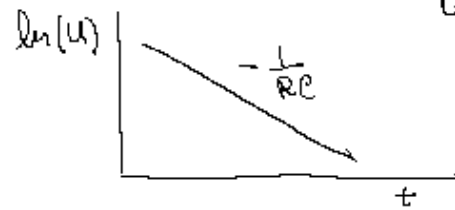
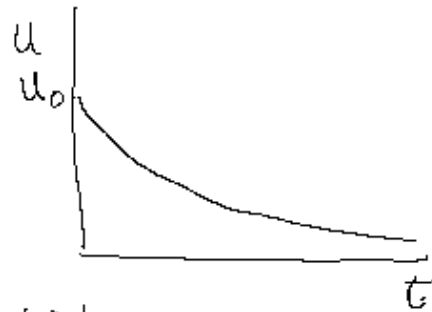
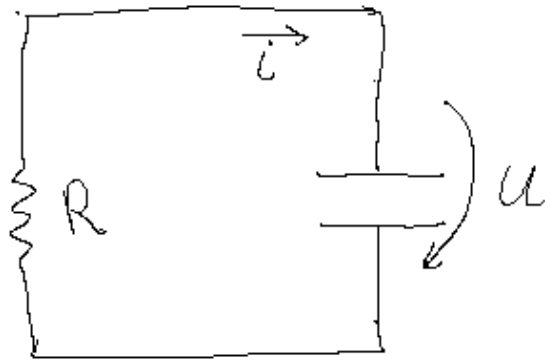
$$W_C = \frac{1}{2} C U^2(t=\infty) = \frac{1}{2} C \varepsilon^2$$

$$W_R = \int_0^\infty R i^2 dt = R \frac{\varepsilon^2}{R^2} \int_0^\infty e^{-\frac{2t}{RC}} dt = \frac{1}{2} C \varepsilon^2$$

$$W_\varepsilon = W_C + W_R$$



Descarga do condensador num circuito RC



$$U = -Ri = -R \frac{dQ}{dt}$$

$$U = -RC \frac{dU}{dt}$$

$$RC \frac{dU}{dt} + U = 0 \rightarrow U = U_0 e^{-\frac{t}{RC}}$$

$$RC D + 1 = 0$$

$$D = -\frac{1}{RC}$$

$$\ln(U) = -\frac{t}{RC} + \ln(U_0)$$

$$i = C \frac{dU}{dt} = -\frac{U_0}{R} e^{-\frac{t}{RC}}$$

Balanco energético

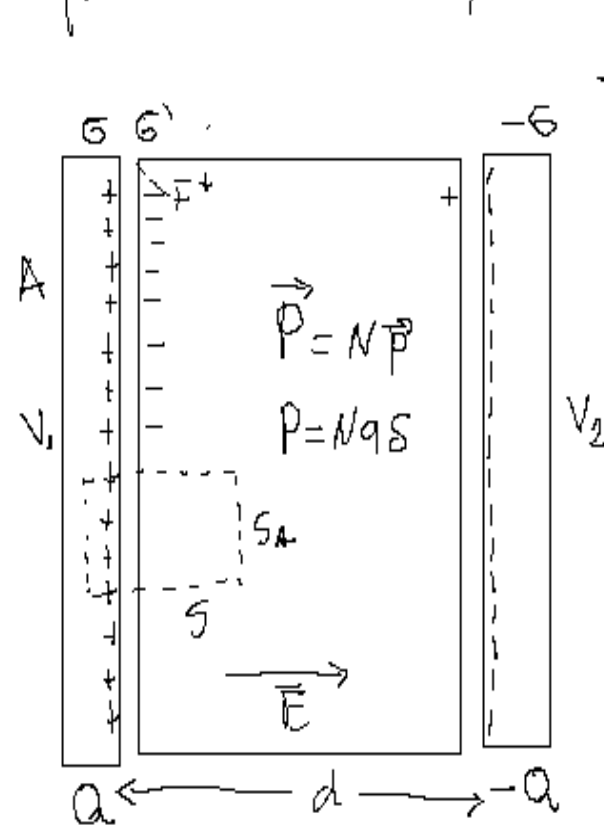
$$W_C = \frac{1}{2} C U_0^2$$

$$W_R = \int_0^{\infty} R i^2 dt = R \frac{U_0^2}{R^2} \int_0^{\infty} e^{-2t/RC} dt$$

$$W_A = \frac{1}{2} C U_0^2$$

$$\boxed{W_C = W_R}$$

# Campo elétrico na presença de matéria, condensador plano com dielétrico



$$G' = \frac{Q'}{A} = \frac{-qN}{A} = \frac{-qN\delta A}{A} = -qN\delta$$

$$G' = -P \quad G' = \vec{P} \cdot \vec{N}$$

$$ES_A = \frac{G + G'}{\epsilon_0} S_A \rightarrow E = \frac{G - P}{\epsilon_0}$$

$$\epsilon_0 E + P = G \rightarrow \epsilon_0 \vec{E} + \vec{P} = \vec{D}$$

$$\vec{P} = \epsilon_0 \chi_e \vec{E} \rightarrow \vec{D} = (\epsilon_0 + \epsilon_0 \chi_e) \vec{E} = \epsilon \vec{E} = \epsilon_0 (1 + \chi_e) \vec{E}$$

$\chi_e \rightarrow$  Suscetibilidade elétrica (adimensional)

$\epsilon = \epsilon_0 (1 + \chi_e) \rightarrow$  Permissividade elétrica ( $F m^{-1}$ )

$$\vec{E} = \frac{G - \epsilon_0 \chi_e \vec{E}}{\epsilon_0} \rightarrow \epsilon_0 \vec{E} + \epsilon_0 \chi_e \vec{E} = G \rightarrow E = \frac{G}{\epsilon_0 (1 + \chi_e)} = \frac{G}{\epsilon} \rightarrow C = \frac{Q}{V_1 - V_2} = \frac{GA}{Ed} = \frac{\epsilon A}{d}$$

$$\epsilon E = G = D$$



# Análise da permissividade elétrica

$$\epsilon = \epsilon_0(1 + \chi_e) \quad \vec{P} = \epsilon_0 \chi_e \vec{E}, \quad \vec{P} = N \vec{p}, \quad \vec{p} = \alpha \vec{E}_i$$

↑ número de dipolos  
↑ por unidade de volume
 ↑ polarizabilidade  
↑ molecular
 ↑ campo local

## Relação entre $E_i$ e $E$

Sistemas pouco densos (gases)  
 $E_i = E$

Sistemas densos (líquidos e sólidos)

$$E_i = E + \frac{P}{3\epsilon_0}$$

$$P = N\alpha E_i = \begin{cases} N\alpha E \Rightarrow \epsilon_0 \chi_e = N\alpha \Rightarrow \epsilon = \epsilon_0 + N\alpha & (\text{gases}) \\ N\alpha \left( E + \frac{P}{3\epsilon_0} \right) \Rightarrow P = \frac{N\alpha}{1 - \frac{N\alpha}{3\epsilon_0}} E \Rightarrow \epsilon_0 \chi_e = \frac{N\alpha}{1 - \frac{N\alpha}{3\epsilon_0}} \Rightarrow \epsilon = \epsilon_0 + \frac{N\alpha}{1 - \frac{N\alpha}{3\epsilon_0}} \end{cases}$$

↑  
 Eq. de Clausius-Mossotti  
 ou Lorentz-Lorentz

Fontes de polarização: {  
 eletrônica  
 orientação de dipolos  
 iônica

# Análise da Permittividade elétrica

## Modelo elementar para a polarização eletrónica:

Eq. movimento:  
p/ electrão no  
potencial do  
átomo

$$m \frac{d^2 x}{dt^2} + \gamma' \frac{dx}{dt} + kx = q E_i$$

$$\frac{d^2 x}{dt^2} + \frac{\gamma'}{m} \frac{dx}{dt} + \frac{k}{m} x = \frac{q}{m} E_i$$

$$-\omega^2 \bar{x} + i\omega \gamma \bar{x} + \omega_0^2 \bar{x} = \frac{q}{m} \bar{E}_i \rightarrow \bar{x} = \frac{q/m \bar{E}_i}{\omega_0^2 - \omega^2 + i\omega \gamma}, \quad \bar{p} = q \bar{x} = \underbrace{\frac{q^2/m}{\omega_0^2 - \omega^2 + i\omega \gamma}}_{\alpha} \bar{E}_i$$

$$\gamma = \frac{\gamma'}{m}, \quad \omega_0^2 = \frac{k}{m}$$

$$\alpha = \alpha(\omega) = \frac{q^2/m}{\omega_0^2 - \omega^2 + i\omega \gamma} \quad (\text{complexo}) \Rightarrow \epsilon = \epsilon_0 + \frac{N \alpha(\omega)}{1 - \frac{N \alpha(\omega)}{3 \epsilon_0}} = \epsilon(\omega) \rightarrow \text{complexo}$$

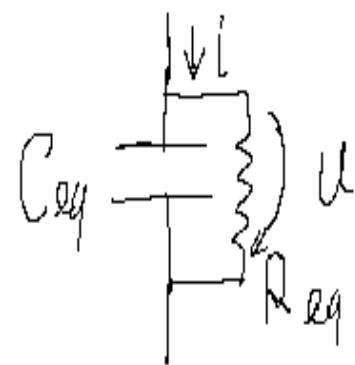
$$E_i = \text{Re} \left\{ E_{i0} e^{i(\omega t + \phi_E)} \right\} = \text{Re} \left\{ \bar{E}_i e^{i\omega t} \right\} \quad \bar{E}_i = E_{i0} e^{i\phi_E}$$

$$x = \text{Re} \left\{ x_0 e^{i(\omega t + \phi_x)} \right\} = \text{Re} \left\{ \bar{x} e^{i\omega t} \right\} \quad \bar{x} = x_0 e^{i\phi_x}$$

Condensador plano com dielétrico de permissividade elétrica complexa

$$C = \frac{\epsilon A}{d} = \frac{A}{d} (\epsilon_r + j\epsilon_i) = \frac{A}{d} \epsilon_r + j \frac{A}{d} \epsilon_i = C_r + jC_i \quad j \equiv \sqrt{-1}$$

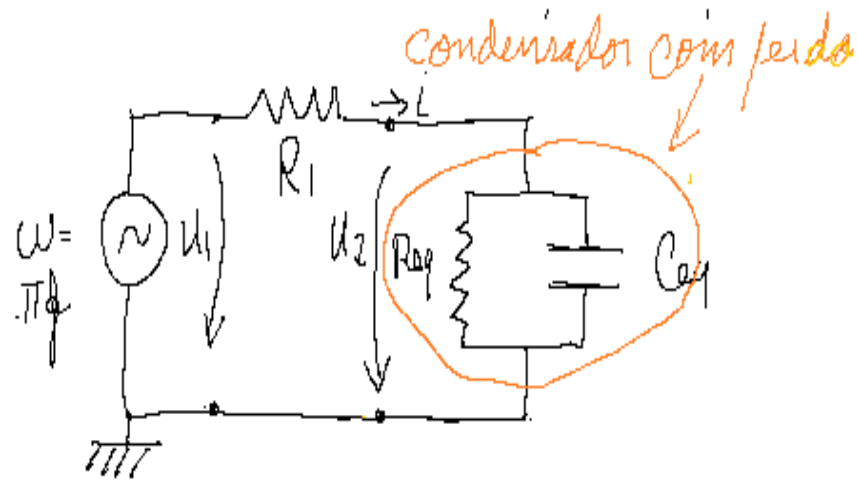
$$u \begin{cases} \downarrow +Q \\ -Q \end{cases} \quad i = \frac{dQ}{dt} = C \frac{du}{dt} \rightarrow \bar{I} = j\omega \bar{Q} = j\omega C \bar{U} = j\omega \bar{U} (C_r + jC_i) = j\omega C_r \bar{U} - \omega C_i \bar{U}$$



$$i = \frac{U}{R_{eq}} + \frac{dQ}{dt} = \frac{U}{R_{eq}} + C_{eq} \frac{dU}{dt} \rightarrow \bar{I} = \frac{\bar{U}}{R_{eq}} + j\omega C_{eq} \bar{U}$$

$$\begin{cases} C_{eq} = C_r \\ \frac{1}{R_{eq}} = -\omega C_i \end{cases} \rightarrow \begin{cases} \epsilon_r \frac{A}{d} = C_{eq} \\ \epsilon_i \frac{A}{d} = -\frac{1}{\omega R_{eq}} \end{cases}$$

Determinação experimental de  $\epsilon(\omega) = \epsilon_r + j\epsilon_i$  com condensador plano



$$\begin{cases} U_1 = R_1 \dot{I}_1 + U_2 \\ \dot{I}_1 = \frac{U_2}{R_{eq}} + C_{eq} \frac{dU_2}{dt} \end{cases} \rightarrow \begin{cases} \frac{\bar{U}_1 - \bar{U}_2}{R_1} = \bar{I}_1 \\ \bar{I}_1 = \frac{\bar{U}_2}{R_{eq}} + j\omega C_{eq} \bar{U}_2 \end{cases} \begin{cases} \frac{\bar{U}_1 - \bar{U}_2}{R_1} = \frac{\bar{U}_2}{R_{eq}} + j\omega C_{eq} \bar{U}_2 \end{cases}$$

$$\frac{\bar{U}_1}{R_1} = \bar{U}_2 \left( \frac{1}{R_1} + \frac{1}{R_{eq}} + j\omega C_{eq} \right) \rightarrow \frac{\bar{U}_1}{\bar{U}_2} = 1 + \frac{R_1}{R_{eq}} + j\omega R_1 C_{eq}$$

$$\frac{U_{1ef}}{U_{2ef}} = \left| \frac{\bar{U}_1}{\bar{U}_2} \right| = \sqrt{\left(1 + \frac{R_1}{R_{eq}}\right)^2 + \omega^2 R_1^2 C_{eq}^2} \rightarrow C_{eq} = \frac{1}{R_1 \omega} \sqrt{\left(\frac{U_{1ef}}{U_{2ef}}\right)^2 - \left(1 + \frac{R_1}{R_{eq}}\right)^2} \quad (2)$$

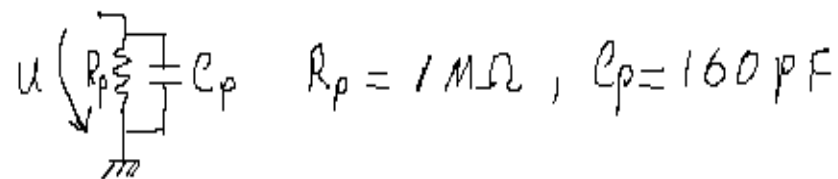
Potência dissipada  
no conversor com  
perdas: P:

$$P = \langle u_2 i \rangle = \left\langle \frac{u_2^2}{R_{eq}} \right\rangle = \frac{u_{2ef}^2}{R_{eq}} = \left\langle u_2 \frac{u_1 - u_2}{R_1} \right\rangle = \frac{\langle u_2 u_1 \rangle - \langle u_2^2 \rangle}{R_1} \rightarrow$$

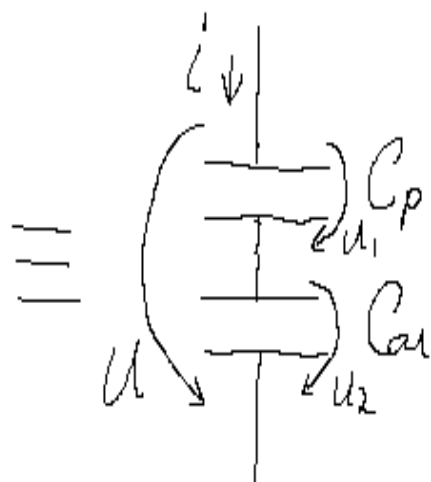
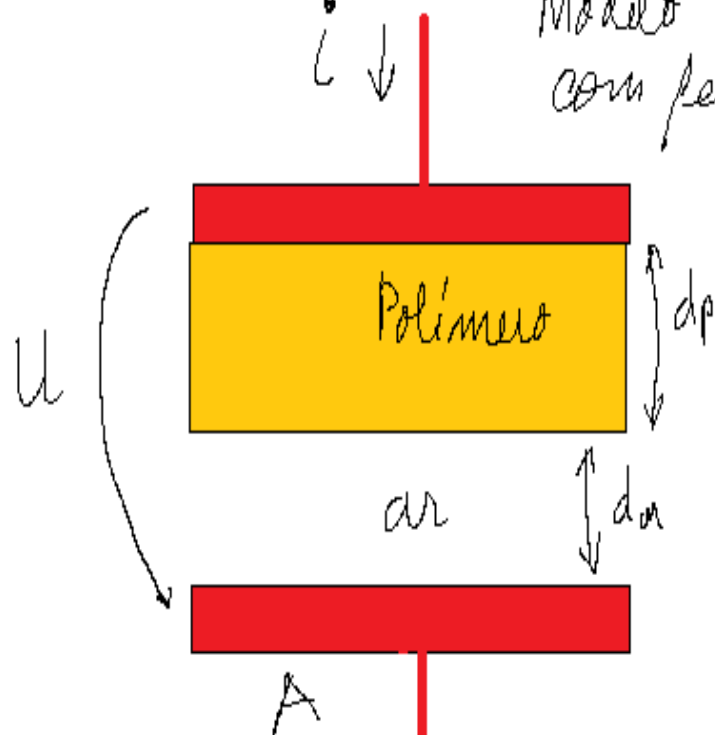
$$\rightarrow R_{eq} = R_1 \frac{u_{2ef}^2}{\langle u_1 u_2 \rangle - u_{2ef}^2} \quad (1)$$

$$\langle x \rangle \equiv \frac{1}{T} \int_0^T x dt, \quad x_{ef}^2 \equiv \frac{1}{T} \int_0^T x^2 dt$$

Nota: Circuito equivalente  
da ponta de prova do osciloscópio:



Modelo mais completo p/ descrever o condensador com perdas em estudo no lab.



$$i = C_T \frac{dU}{dt} = C_T \left( \frac{dU_1}{dt} + \frac{dU_2}{dt} \right), \quad i = C_p \frac{dU_1}{dt}$$

$$i = C_{ar} \frac{dU_2}{dt}$$

$$i = C_T \left( \frac{i}{C_p} + \frac{i}{C_{ar}} \right)$$

$$\frac{1}{C_T} = \frac{1}{C_p} + \frac{1}{C_{ar}} \rightarrow \frac{1}{C_{T\lambda} + jC_{Ti}} - \frac{1}{C_{ar}} = \frac{1}{C_{p\lambda} + jC_{pi}}$$

$$C_{p\lambda} + jC_{pi} = \frac{C_{ar} (C_{T\lambda} + jC_{Ti}) (C_{ar} - C_{T\lambda} + jC_{Ti})}{(C_{ar} - C_{T\lambda})^2 + C_{Ti}^2}$$

$$C_{ar} = \epsilon_0 \frac{A}{d_{ar}}$$

$$C_{T\lambda} = C_{eq}$$

$$C_{Ti} = \frac{1}{\omega R_{eq}}$$

$$\begin{cases} C_{p\lambda} = \epsilon_{\lambda p} \frac{A}{d_p} \\ C_{pi} = \epsilon_{ip} \frac{A}{d_p} \end{cases}$$

FIM