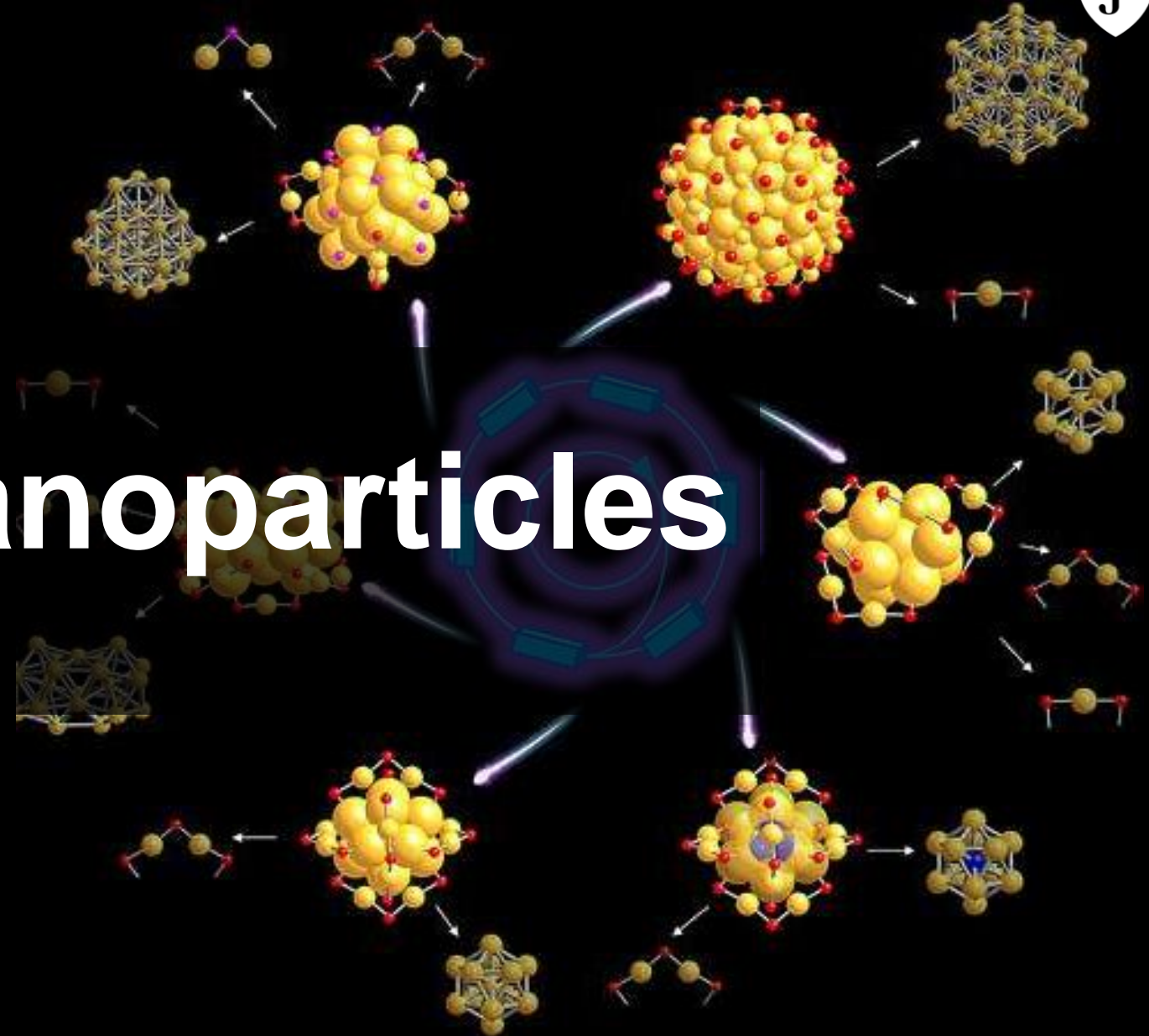


The background features a central, glowing, hourglass-shaped nanostructure with a bright yellow-orange light source at its narrowest point. This central structure is surrounded by several smaller, spherical, textured particles in shades of blue and teal, all set against a dark, gradient background.

Nanostructured Materials and Nanotechnology 2022-2023

ermelinda.macoas@tecnico.ulisboa.pt

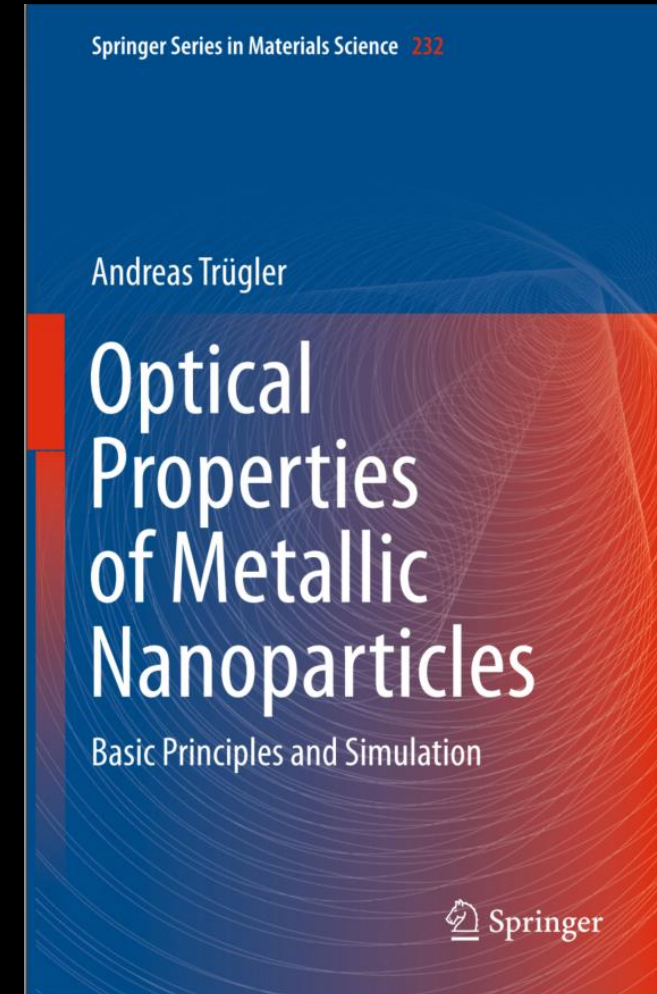
Metal Nanoparticles



Bibliography

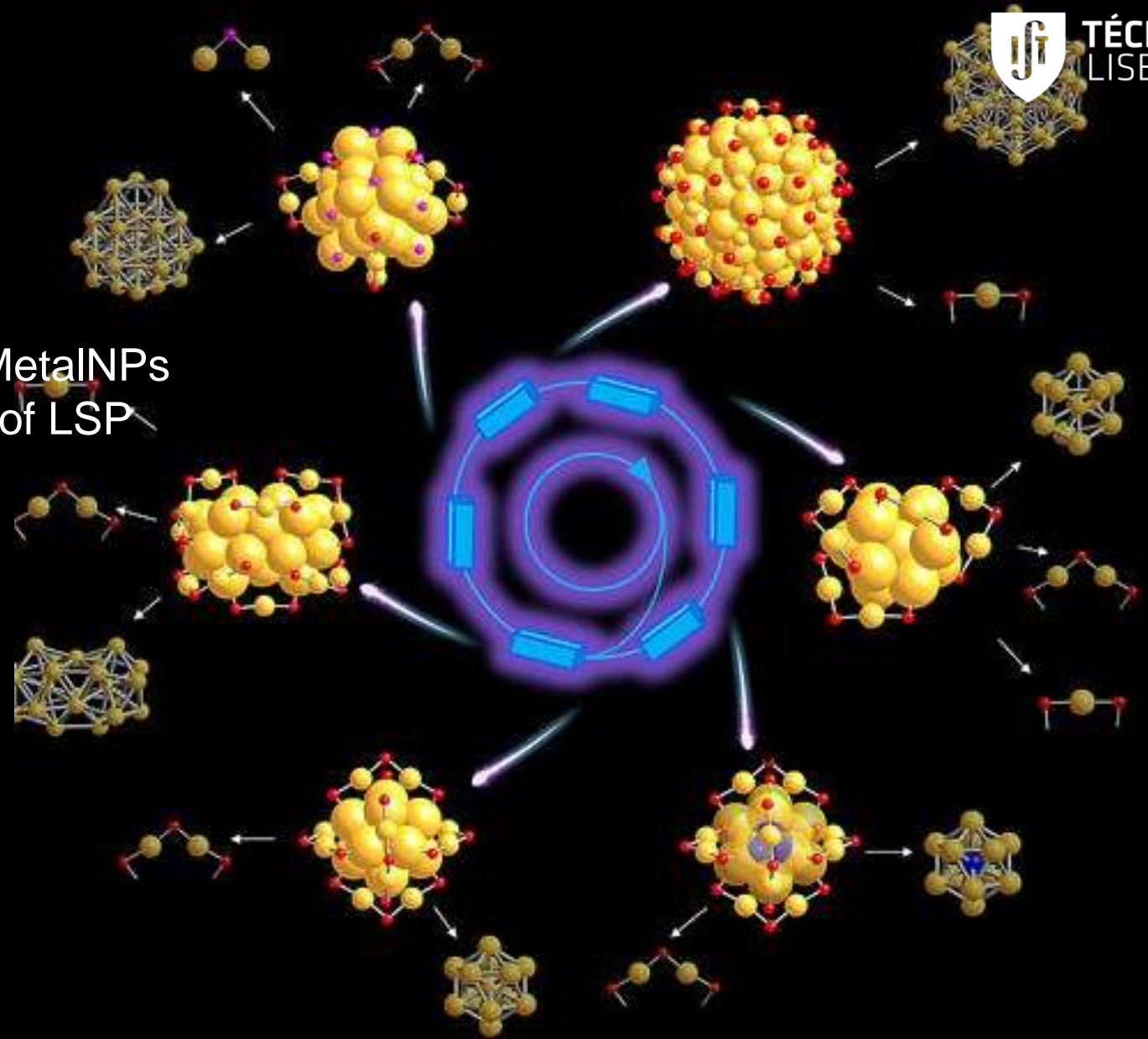
1. “Surface plasmon resonance in gold nanoparticles: a review”
Vincenzo Amendola et al 2017 J. Phys.: Condens. Matter 29
203002, <https://doi.org/10.1088/1361-648X/aa60f3>
2. Gold nanorods: Synthesis, characterization and applications
Pérez-Juste et al, Coordination Chemistry Reviews 249 (2005)
1870–1901 doi:[10.1016/j.ccr.2005.01.030](https://doi.org/10.1016/j.ccr.2005.01.030)

Extra reading



Content

- Localized surface plasmon
- Extinction cross-section for spherical MetalNPs
- Resonance conditions for observation of LSP
- Composition Effect on LSPR
- Size Effect on LSPR
- Medium Effect on LSPR
- Shape Effect
- Field enhancement
- Preparation of MetalNPs
- Applications of MetalNPs



Lateral Flow Immunoassay

COVID-19
(SARS-CoV-2)

C
G
M

A

B

Lycurgus Cup, Roman glass, 4th c.AD.
 Nanoparticles (70 nm) of gold and silver are
 responsible for the dichroic effect (red/green)

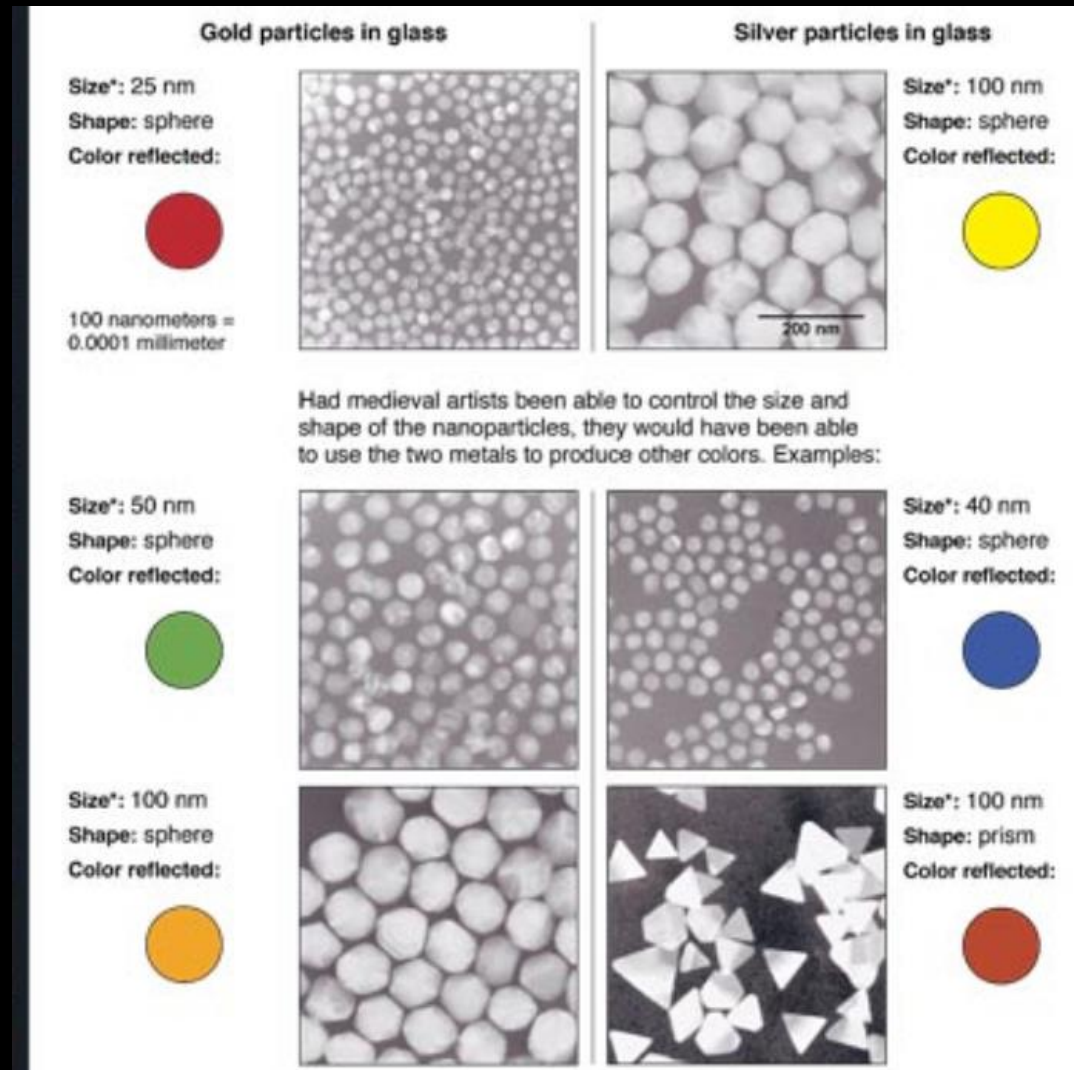


Img credit|Science direct, picture provided by The British Museum



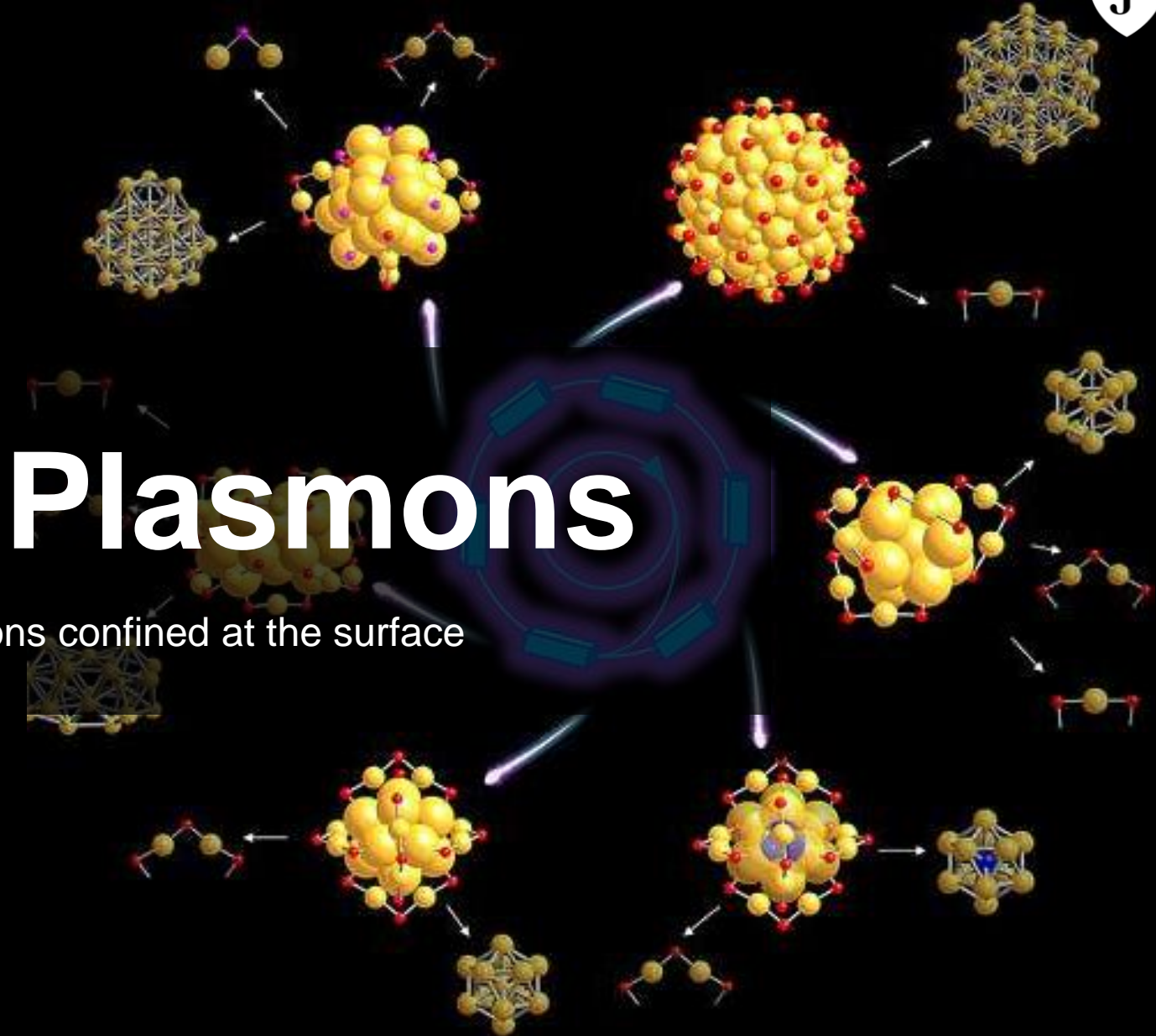
The First Nanotechnologists

Ancient stained-glass makers knew that by putting varying, tiny amounts of gold and silver in the glass, they could produce the red and yellow found in stained-glass windows. Similarly, today's scientists and engineers have found that it takes only a small amounts of nanoparticles, precisely placed, to change a material's physical properties



Surface Plasmons

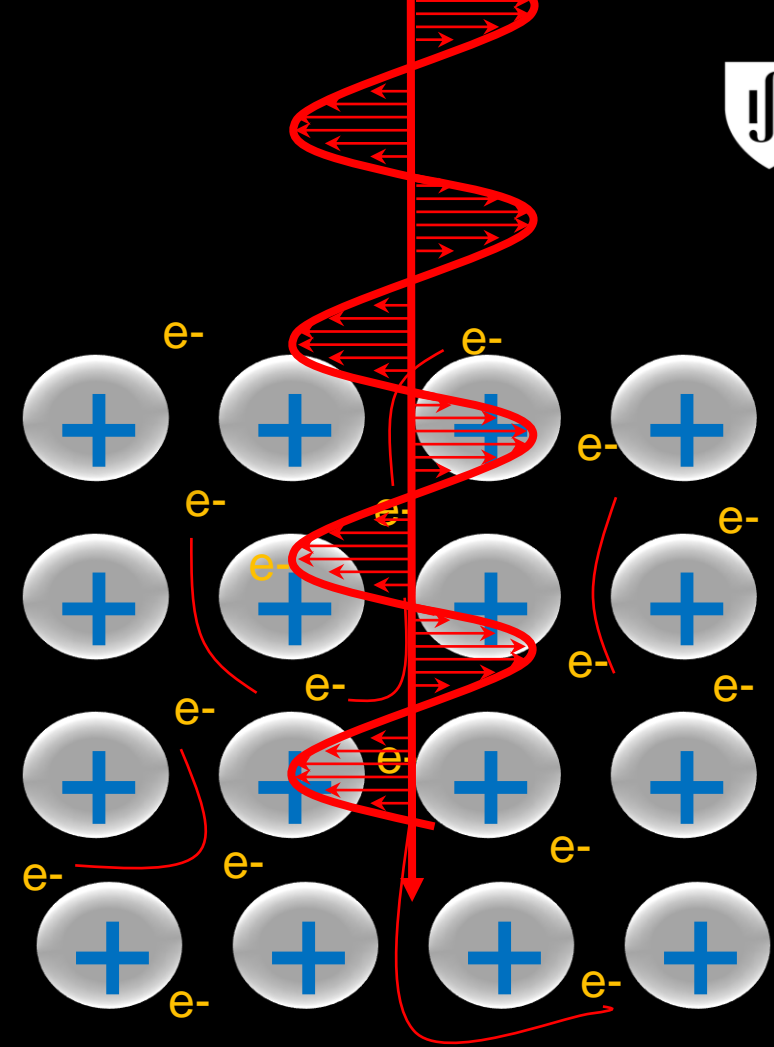
collective electron oscillations confined at the surface



Electron gas

A conductor (metal) can be modeled by a regular array of atoms where each atom gives up electrons that are shared amongst the atoms in the material forming a free e^- gas (e^- sea).

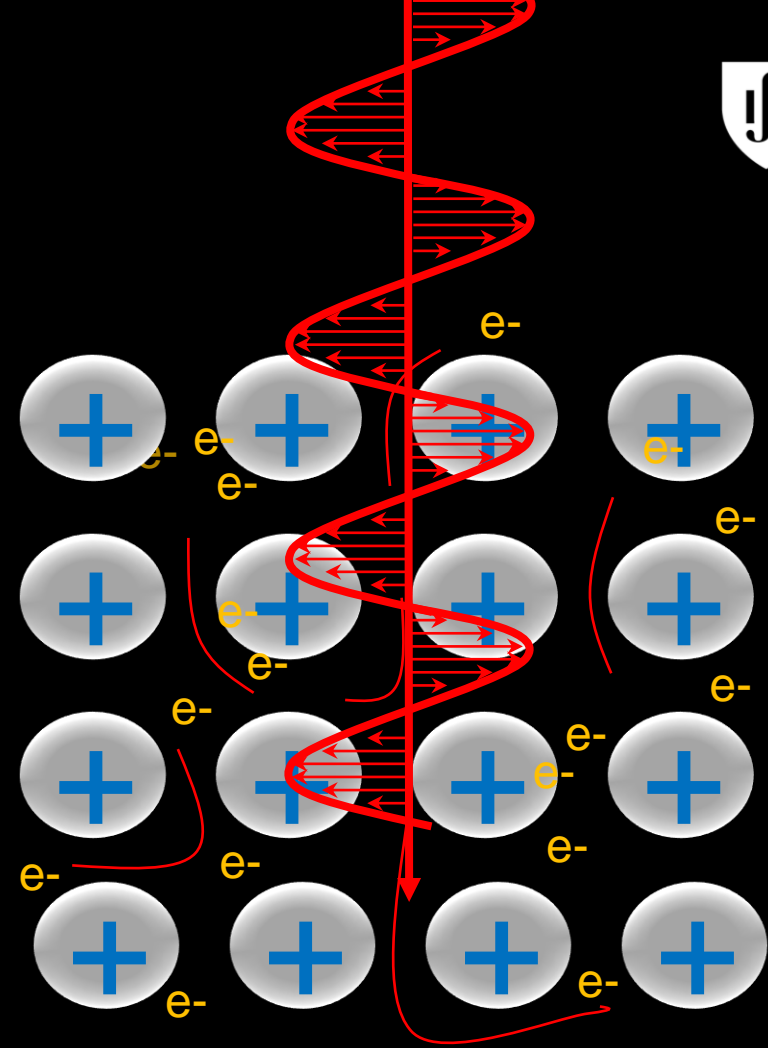
Quantitized system with conduction e^- moving in the periodic potential of the nuclei



Electron gas

A conductor (metal) can be modeled by a regular array of atoms where each atom gives up electrons that are shared amongst the atoms in the material forming a free e^- gas (e^- sea).

Quantitized system with conduction e^- moving in the periodic potential of the nuclei



Electric field exerts forces on charged particles

$$F=q.E$$

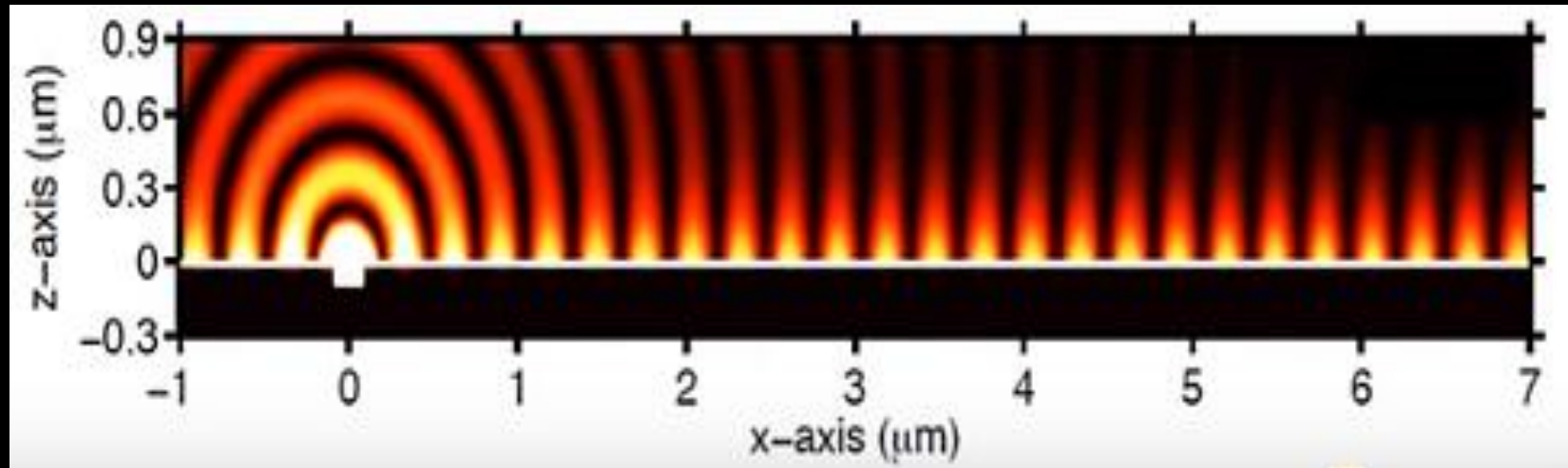
Dielectric

Metal

out of phase

z

SPP Intensity is maximal at the interface and it decays exponentially into both the metal and the dielectric



Dielectric

Metal

$\delta_{\text{diel}} \approx 250\text{-}1000 \text{ nm}$ Evanescent field
length in the dielectric

$\delta_{\text{SPP}} \approx 2\text{-}20 \text{ }\mu\text{m}$ at $\lambda \approx 500 \text{ nm}$ Surface plasmon
up to 1 mm at $\lambda \approx 1550 \text{ nm}$ propagation length

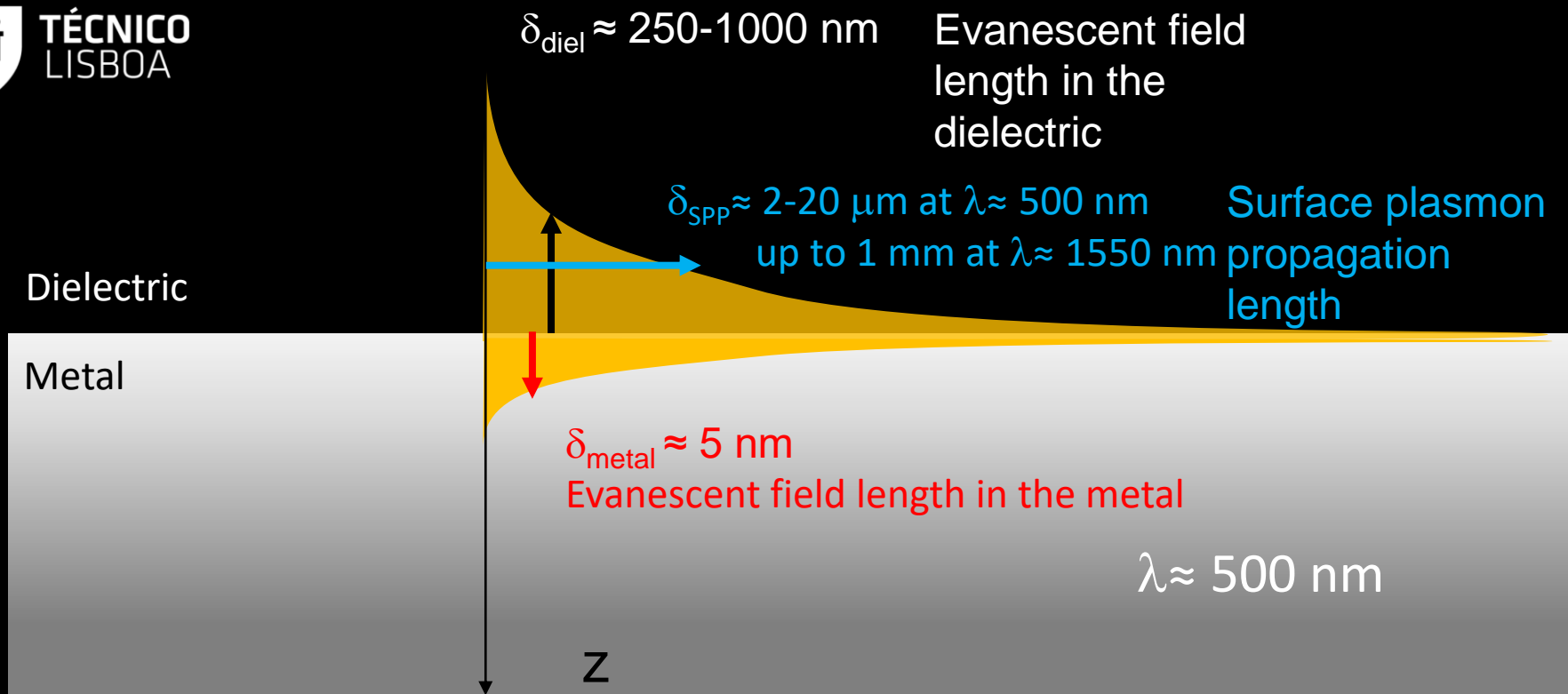
$\delta_{\text{metal}} \approx 5 \text{ nm}$
Evanescent field length in the metal

$\lambda \approx 500 \text{ nm}$

z

SPP Intensity is
maximal at the interface
and it decays
exponentially into both
the metal and the
dielectric

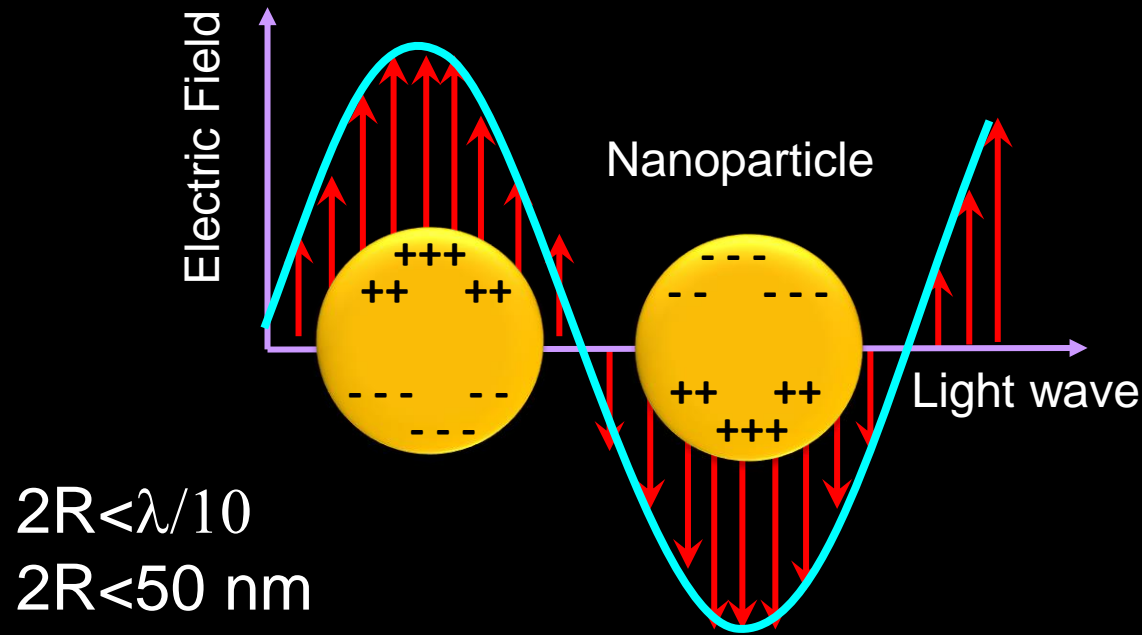




SPP Intensity is maximal at the interface and it decays exponentially into both the metal and the dielectric

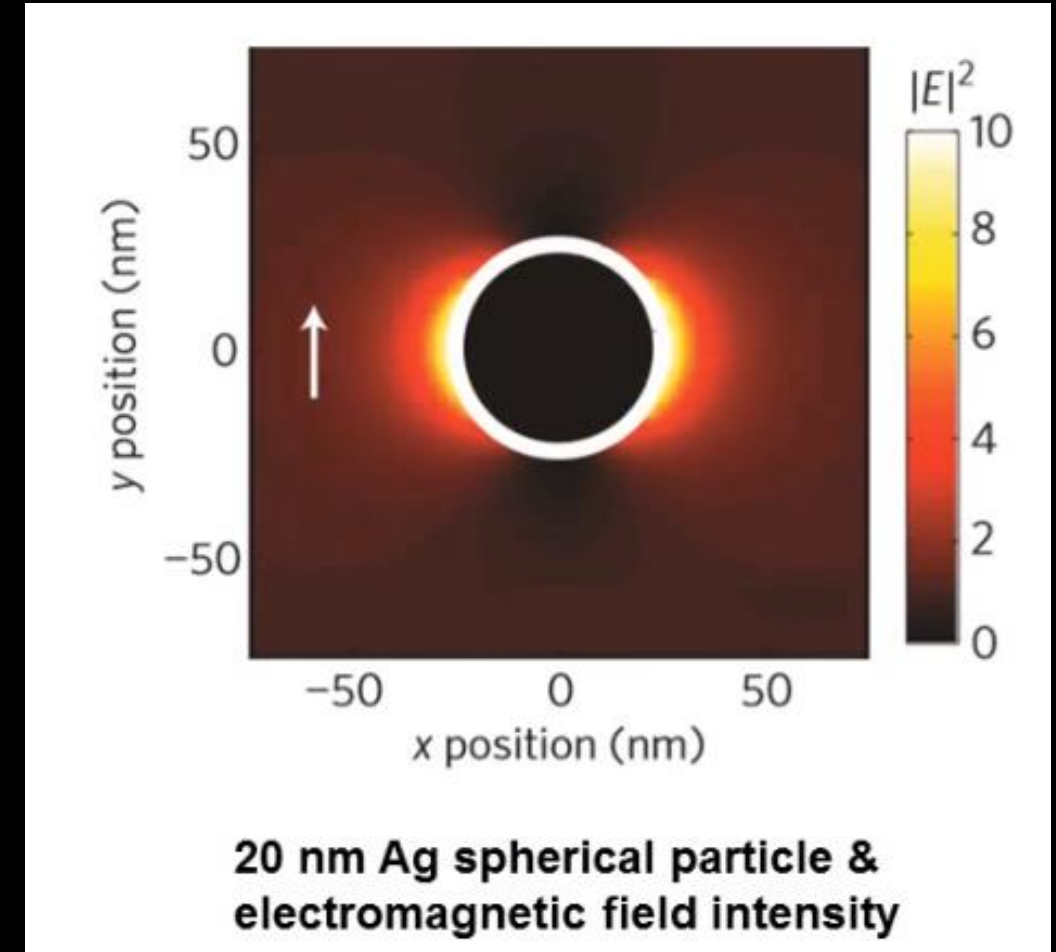
If the lateral dimensions of the interface become much smaller than the surface plasmon propagation length the surface plasmon is said to be localized
LSPP= localized surface plasmon polariton

In NPs the plasmon field can couple with the incident electric field resulting in field enhancement



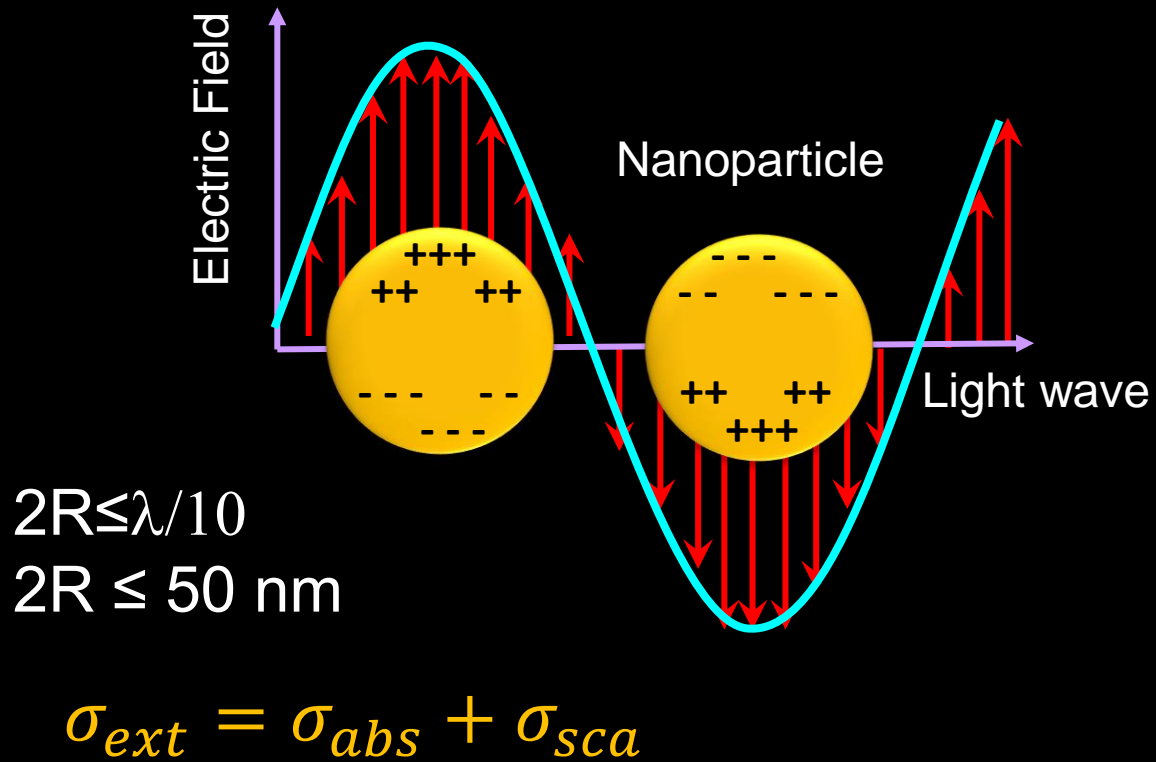
Field is uniform inside the NPs
NPs are macroscopic dipole

The charge displacement inside the NP generates a field that opposes the incident field.

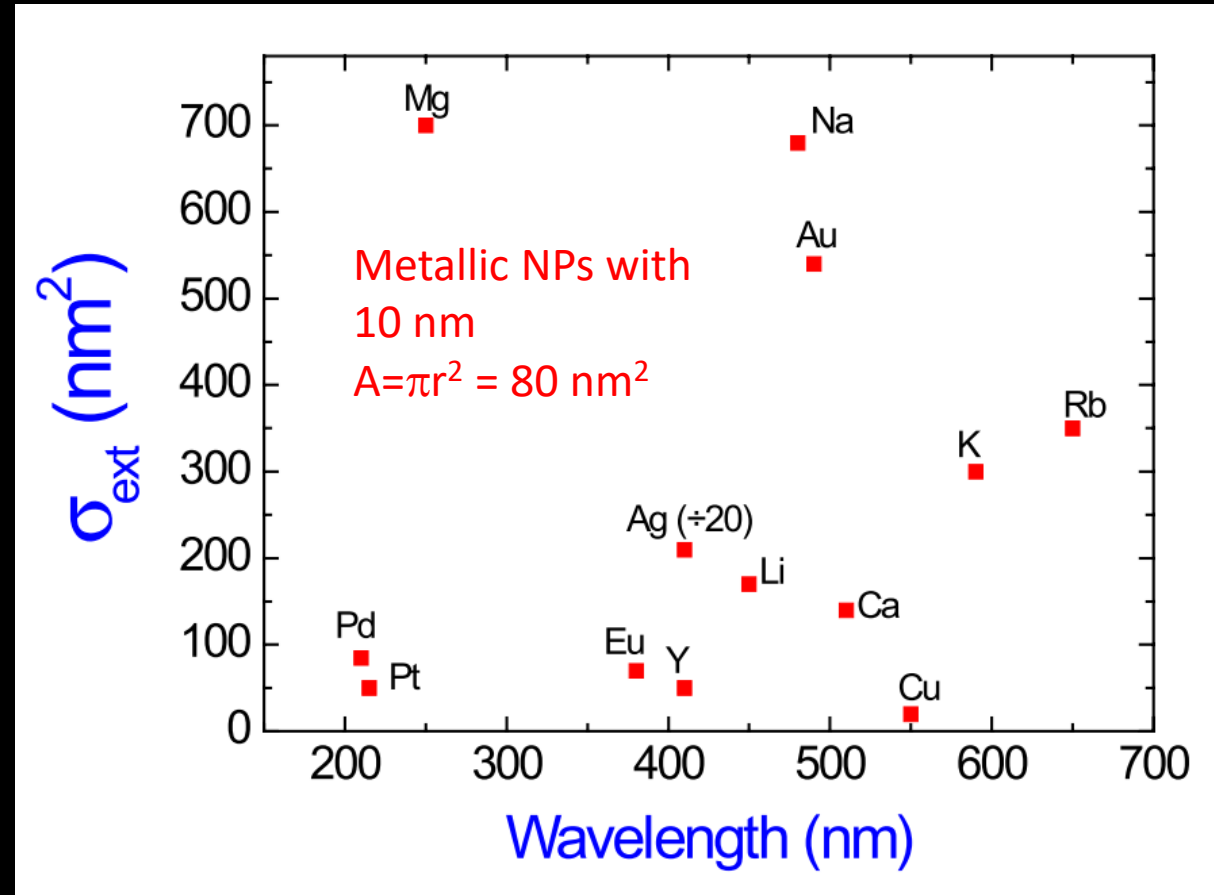


δ_{diel} 20-40 nm for LSPP

LSPP vs SPP in sensing



According to Mie's scattering theory (dipolar approximation), particles with radius much smaller than the wavelength of radiation have extinction cross section (σ_{ext}) that can be bigger than the particle

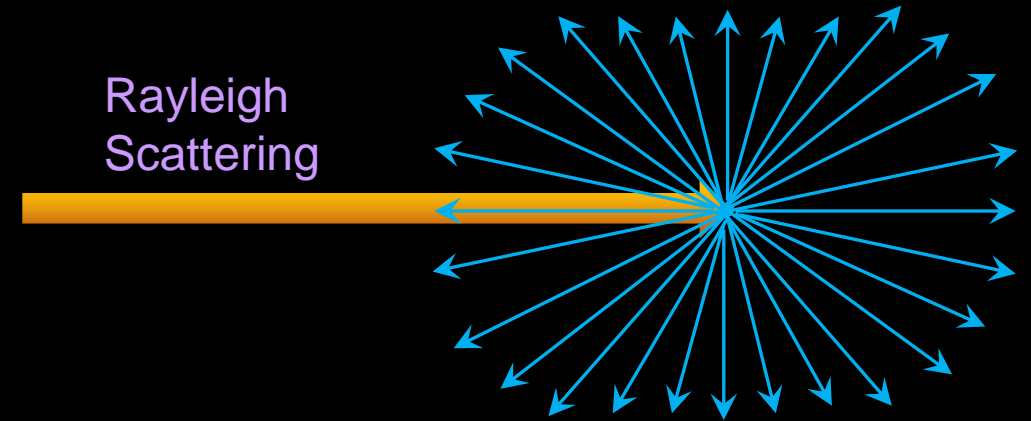


[Garcia, J. Phys. D: Appl. Phys. 44 \(2011\) 283001](#)

For a given particle, the extinction coefficient is a complex function of particle size, wavelength of light, and complex refractive index of the particle.

Different approximations can be employed for certain particle size ranges.

$2R \ll \lambda/10$ $R < 10$ nm Rayleigh scattering (λ^{-4})



The sky is blue due to
Rayleigh scattering
(λ^{-4}) by small
particle



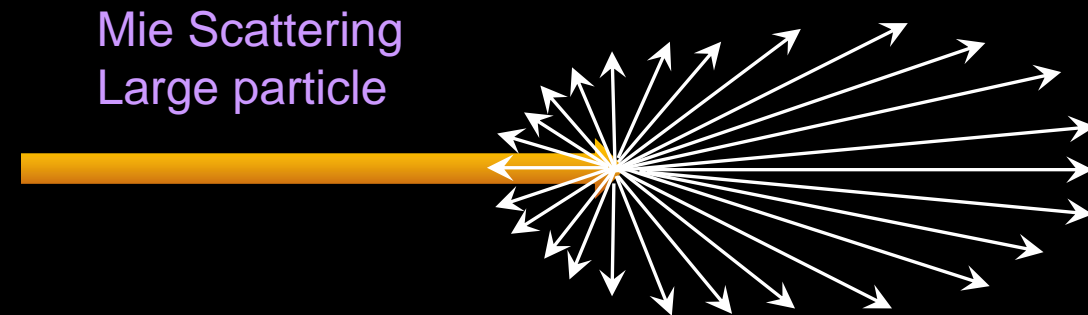
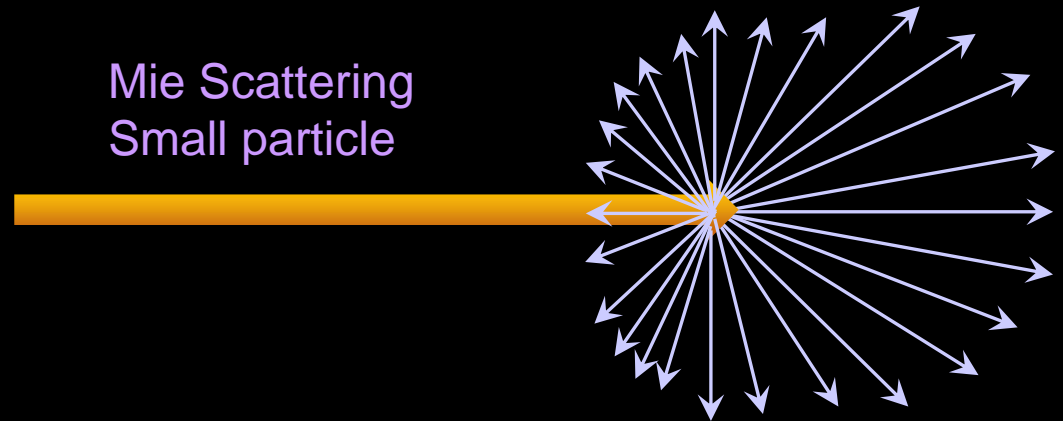
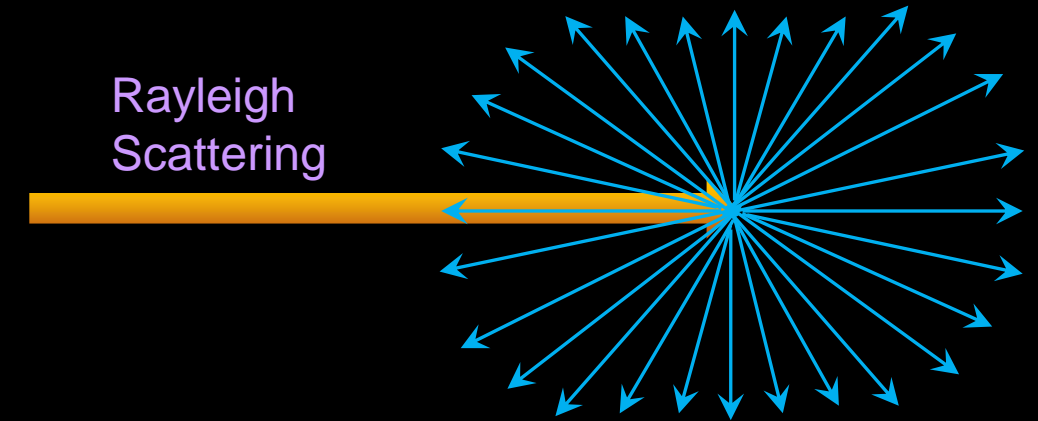
For a given particle, the extinction coefficient is a complex function of particle size, wavelength of light, and complex refractive index of the particle.

Different approximations can be employed for certain particle size ranges.

$2R \ll \lambda/10$ $R < 10$ nm Rayleigh scattering (λ^{-4})

$2R \approx \lambda/10$ $10 < R < 50$ nm Mie scattering

$2R > \lambda$ $R > 800$ nm geometric domain,
classical optics apply, non λ selective
scattering



Clouds are white because they contain bigger drops of water that have a non wavelength selective scattering



Mie's scattering theory (dipolar approximation)

$$\sigma_{ext} = \sigma_{abs} + \sigma_{sca}$$

$$\sigma_{ext} = \frac{24\pi^2 R^3 \epsilon_m^{3/2}}{\lambda} \times \frac{\epsilon_2(\lambda)}{(\epsilon_1(\lambda) + 2\epsilon_m)^2 + \epsilon_2(\lambda)^2}$$

ϵ_1, ϵ_2 are the real (ϵ_1) and imaginary (ϵ_2) parts of the wavelength dependent permittivity function of the material that determine the scattering and absorption, respectively

ϵ_m is the permittivity of the medium (or dielectric function)

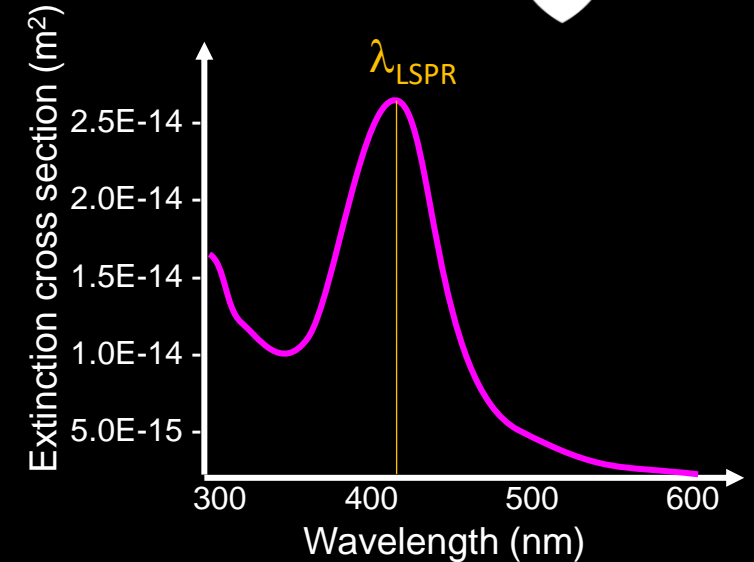
$$\epsilon_m = n^2$$

R radius of the NP

Mie's scattering theory (dipolar approximation)

$$\sigma_{ext} = \sigma_{abs} + \sigma_{sca}$$

$$\sigma_{ext} = \underbrace{\frac{24\pi^2 R^3 \epsilon_m^{3/2}}{\lambda}}_{\text{Linear with } 1/\lambda} \times \underbrace{\frac{\epsilon_2(\lambda)}{(\epsilon_1(\lambda) + 2\epsilon_m)^2 + \epsilon_2(\lambda)^2}}_{\text{Non-linear } \lambda}$$



$$\epsilon_1(\lambda_{LSPR}) = -2\epsilon_m$$

Resonance condition = maximize σ_{ext}
This is verified in some metals depending on the λ

$$\epsilon_m > 1$$

$$\epsilon_1(\lambda_{LSPR}) < 0$$

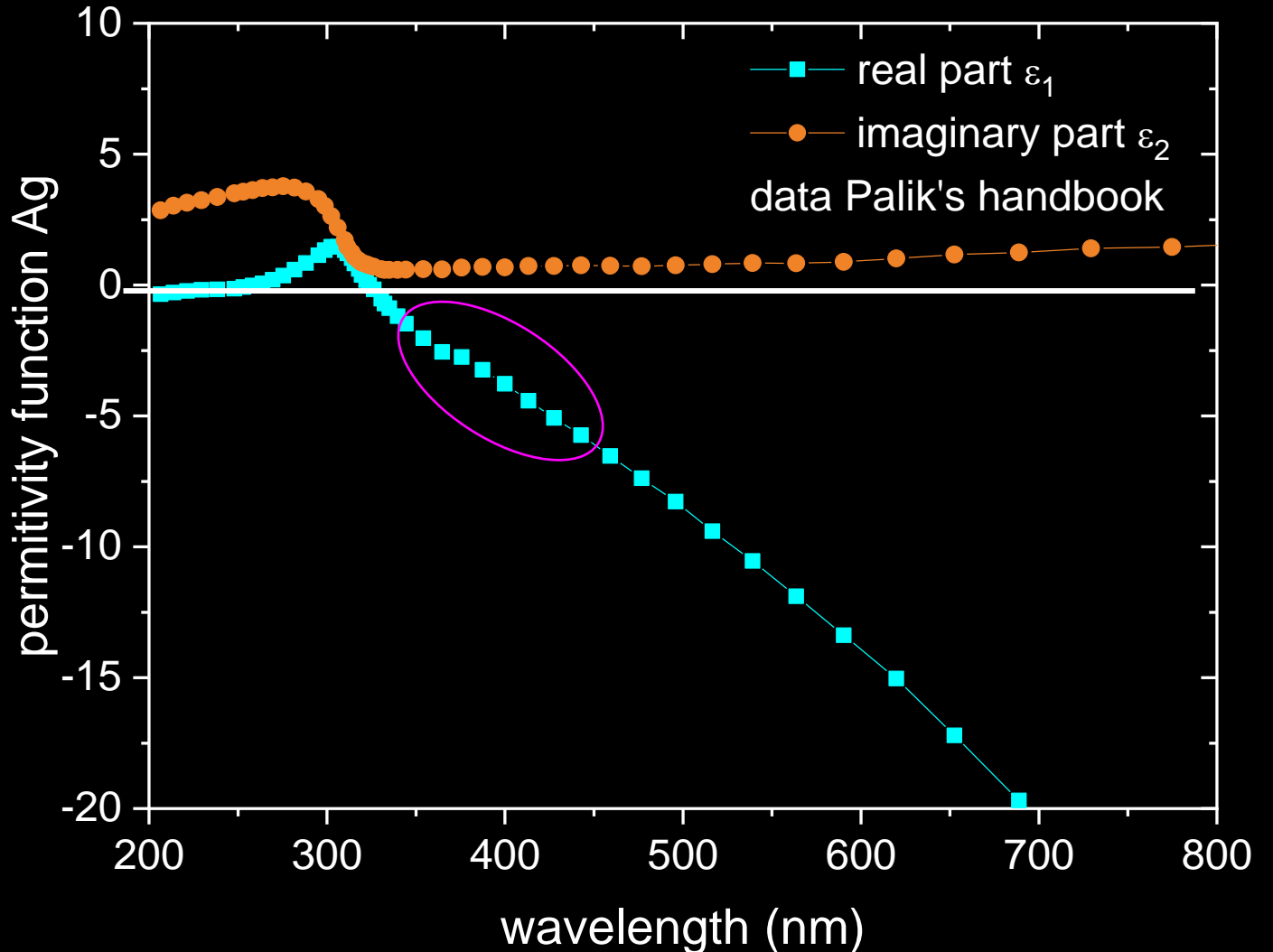
Resonance conditions at optical frequencies for metal

Due to the high **electron density** of metals the resonance conditions for LSP are met at optical frequencies (> 400 nm)

$$\epsilon_1 = -2\epsilon_m$$

$$\epsilon_m = 1-3$$

$$\epsilon_1(\lambda_{LSPR}) < 0$$



Resonance conditions depend on ω_p (electron density)

$$\varepsilon_1(\omega) = 1 - \frac{\omega_p^2}{\omega^2 + \gamma_d^2}$$

ε_1
Real
(scattering)

$\gamma_d = v_F/l$ dumping frequency
(v_F fermi velocity and l is
electron mean free path in
the metal (~30 nm, Au))

ω_p is the frequency of
the bulk plasmon

$$\varepsilon_1(\omega_{LSPR}) = -2\varepsilon_m$$

$$\omega_{LSPR} = \sqrt{\frac{\omega_p^2}{1 + 2\varepsilon_m} - \gamma^2}$$

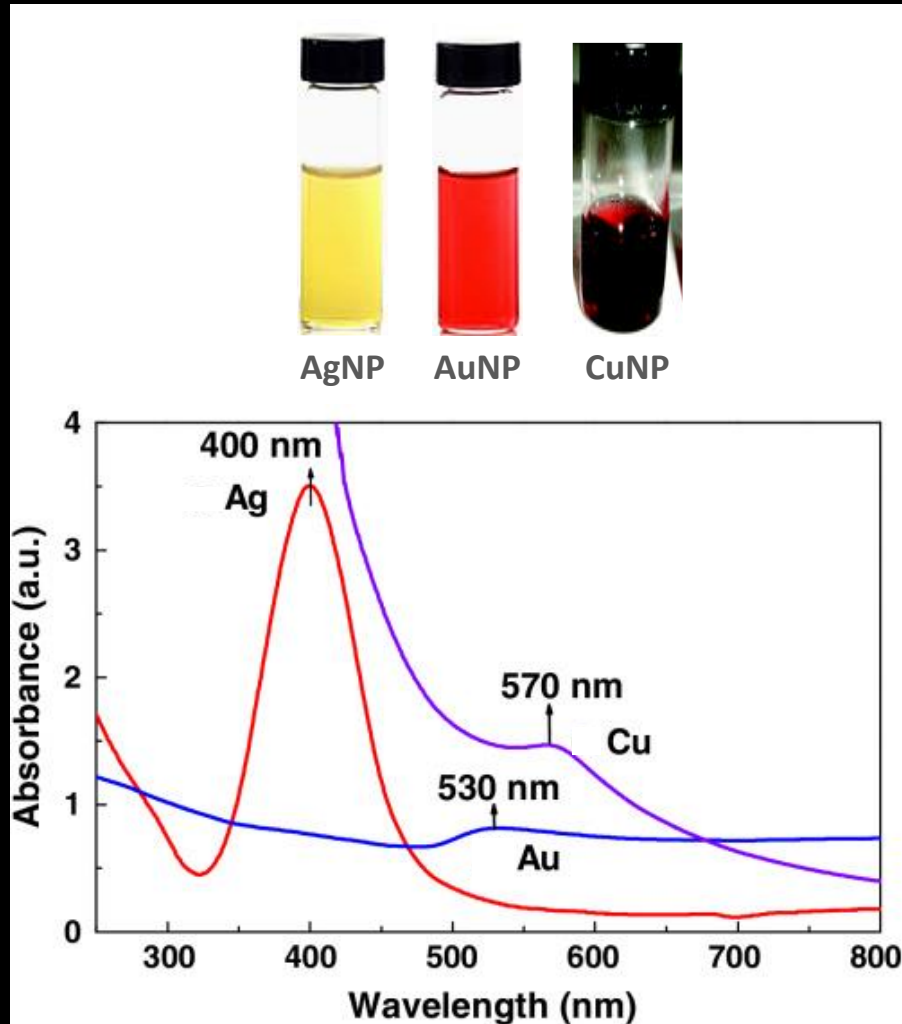
Not an exact answer for LSPR because
 $\varepsilon_{Drude}(\omega)$ is valid for the bulk



$$\omega_p = \sqrt{\frac{N_e e^2}{\varepsilon_0 m_{eff}}}$$

N_e electron density
 e electron charge
 ε_0 vacuum permittivity
 m_{eff} electron mass

Composition Effect on LSP



LSPR of metal nanostructures appear in the visible

LSPR are easier to observe for coinage metals because they do not oxidize so easily

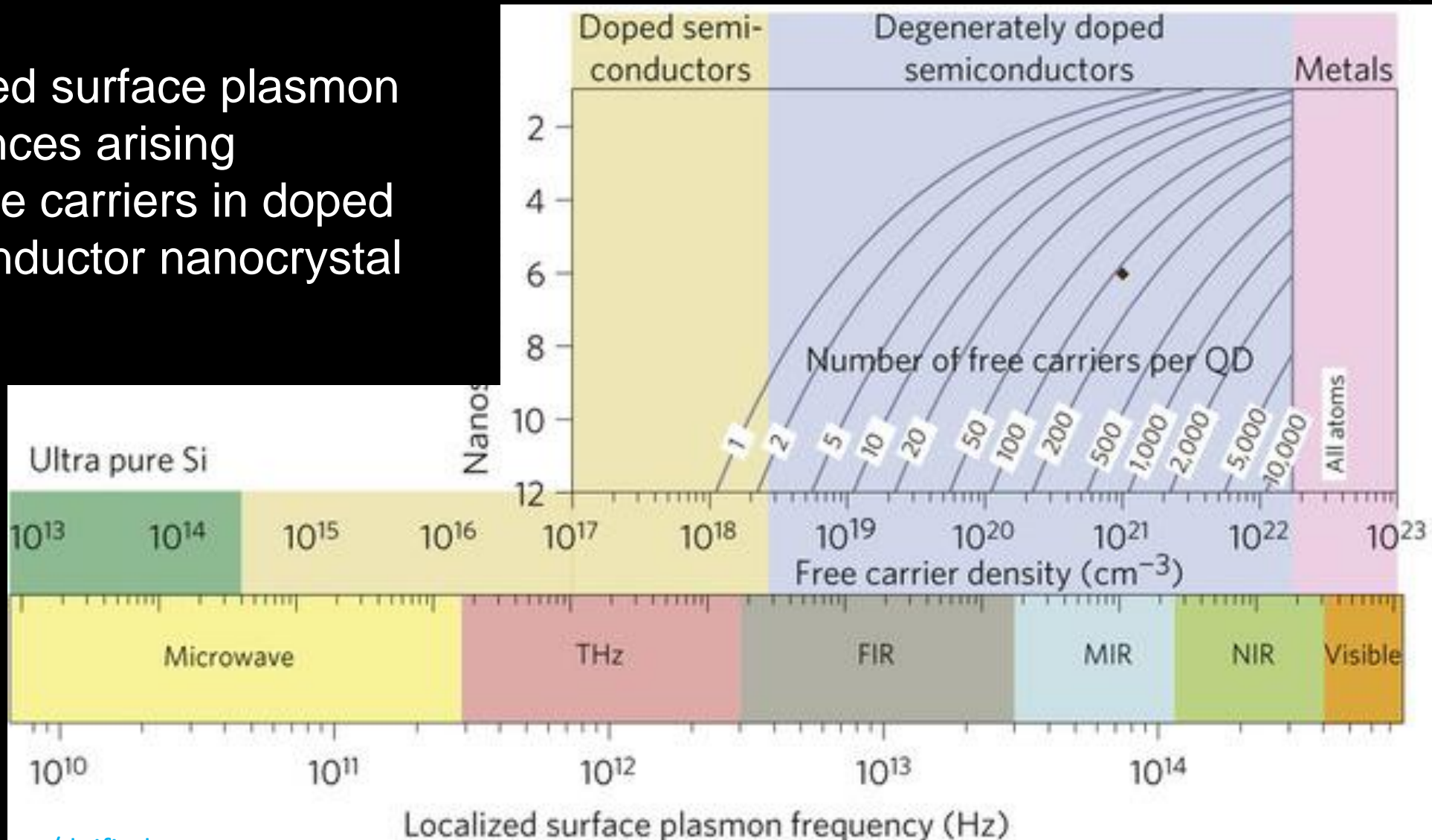
Element	Z	Valence configuration	Ne density 10 ²⁸ e/m ³
Cu	29	4s ¹ 3d ¹⁰	8.5x10 ²⁸
Ag	47	5s ¹ 4d ¹⁰	5.8x10 ²⁸
Au	79	6s ¹ 5d ¹⁰ 4f ¹⁴	5.9x10 ²⁸

$$\omega_p = \sqrt{\frac{N_e e^2}{\epsilon_0 m_{eff}}}$$

N_e electron density
 e electron charge
 ϵ_0 vacuum permittivity
 m_{eff} electron mass

Composition Effect on LSP

Localized surface plasmon resonances arising from free carriers in doped semiconductor nanocrystal



<http://www.nature.com/doi/10.1038/nmat3004>

Size effect (R)

$$\sigma_{ext} = \frac{24\pi^2 R^3 \epsilon_m^{3/2}}{\lambda} \times \frac{\epsilon_2(\lambda)}{(\epsilon_1(\lambda) + 2\epsilon_m)^2 + \epsilon_2(\lambda)^2}$$

For particle $R < 10$ nm the λ_{SPR} is size independent

$R < 10$ nm Abs dominates with a single peak at 355 nm, intensity increases with R

$R > 10$ Scattering increases and band broadening occurs due to higher multipoles

Larger particle no longer behave like simple dipole such that high-multipolar distributions can be induced. The spectrum broadens and band shape becomes asymmetric

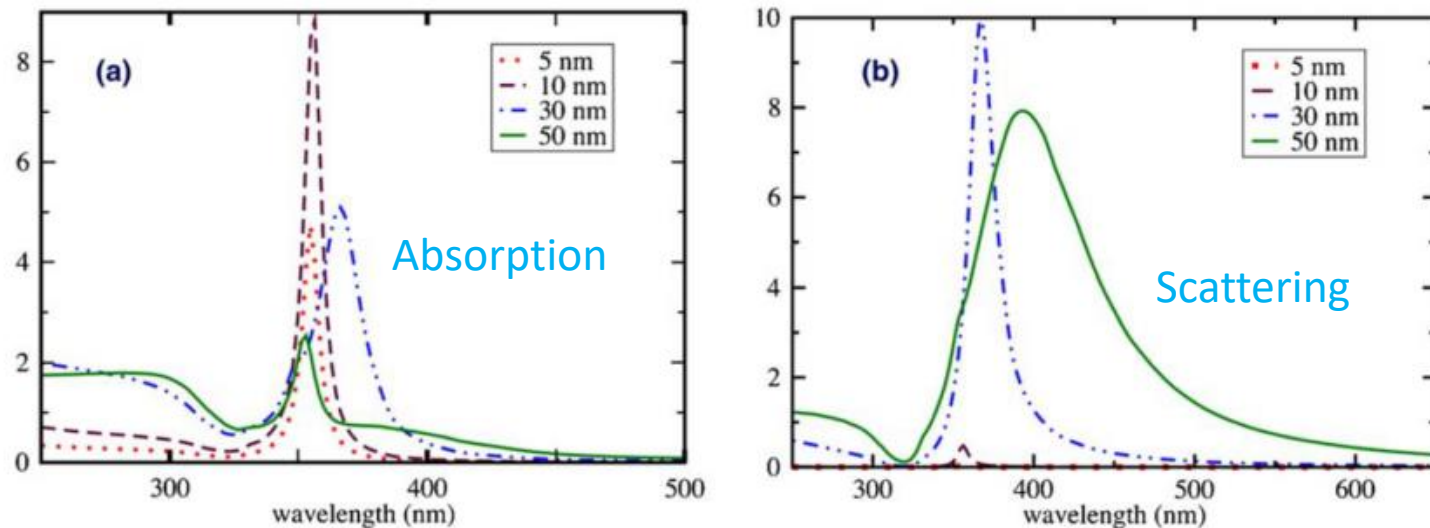


Fig. 1. (a) Absorption and (b) scattering efficiencies for spherical nanoparticles of different radii.

C. Noguez, Optical Materials 27 (2005) 1204–1211

Size effect

For spherical particles, the effect of particle size on the extinction spectra is small (almost no color change is noticed for AuNPs or AgNPs)



5-100 nm AgNP



5-100 nm AuNP

Effect of the medium (ϵ_m)

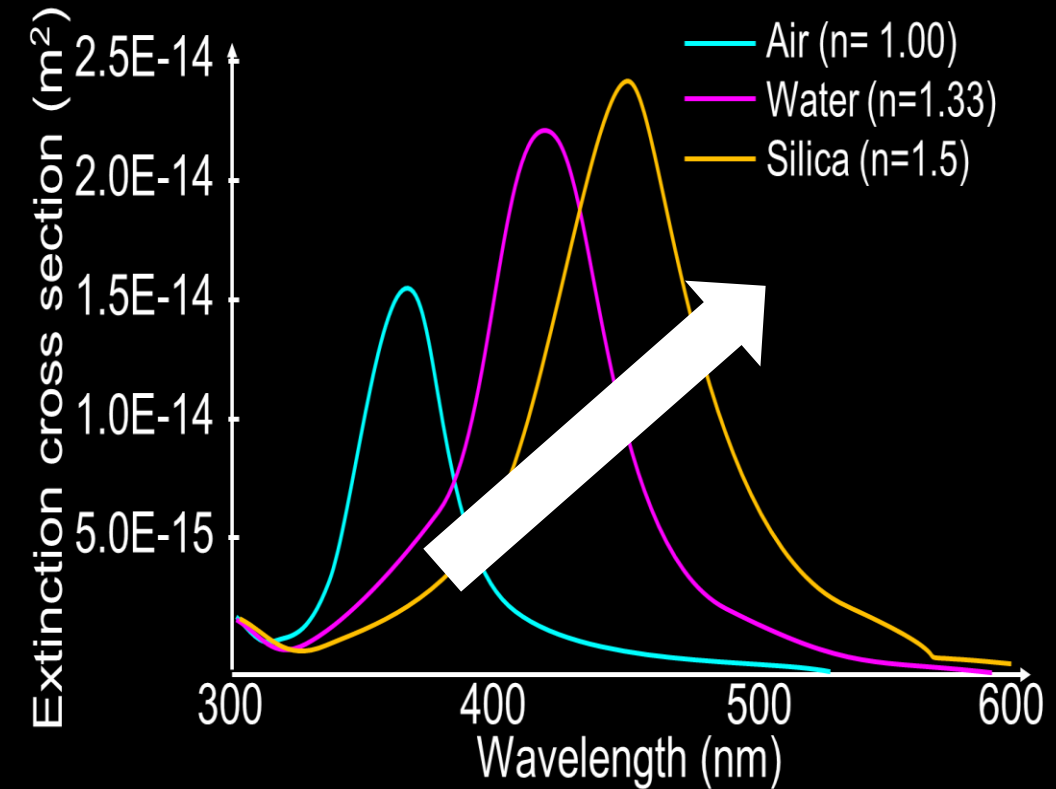
$$\sigma_{ext} = \frac{24\pi^2 R^3 \epsilon_m^{3/2}}{\lambda} \times \frac{\epsilon_2(\lambda)}{(\epsilon_1(\lambda) + 2\epsilon_m)^2 + \epsilon_2(\lambda)^2}$$

$$\epsilon_m = n^2$$

$$\omega_{LSPR} = \frac{\omega_p}{\sqrt{1 + 2n^2}}$$

Approximated linear relationship between λ_{SPR} and environment refractive index (n_m)

$$\lambda_{SPR} = \frac{2\pi c}{\omega_p} \sqrt{2n_m^2 + 1}$$



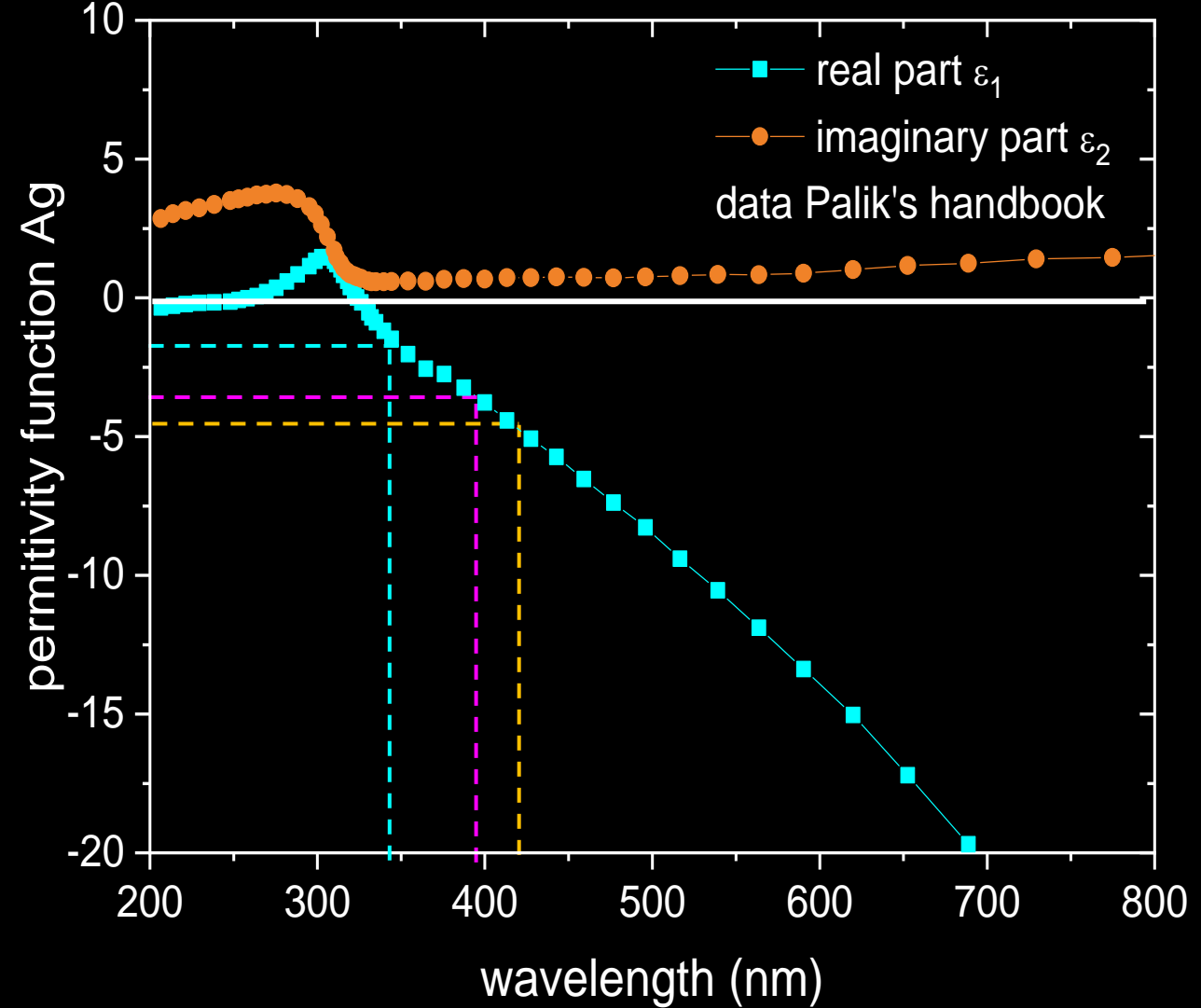
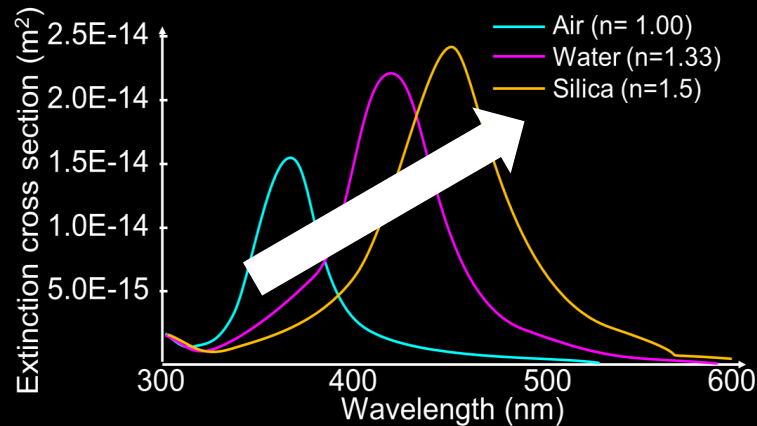
As the refractive index of the medium increases ($\epsilon_m = n^2$), the nanoparticle spectrum shifts to longer wavelengths.

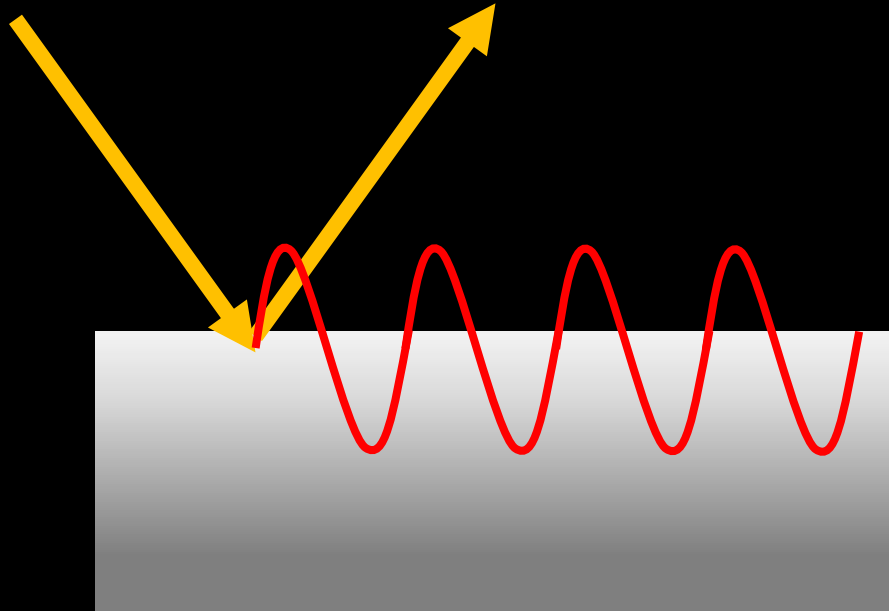
Effect of the medium (ϵ_m)

$$\lambda_{SPR} = ???$$

$$n_m^2 = \epsilon_m$$

$$\epsilon_1 = -2\epsilon_m$$





Plasmon polariton are collective oscillation of electrons at the interface of conductor materials and a dielectric

The evanescent wave propagates only very short distances in the dielectric for LSPR (20-40 nm)

Strong coupling of light with SPP require a change in permittivity sign between the NP material and the medium

$$\varepsilon_1 = -2\varepsilon_m$$

Resonance at optical frequencies only occurs for noble metals and some semiconductors with a given electron density

Optical properties of Metal NPs

depend on:

Permittivity of media ϵ_m

Composition ($\epsilon(\lambda)$, $\epsilon(\omega)$)

Size (R or V)

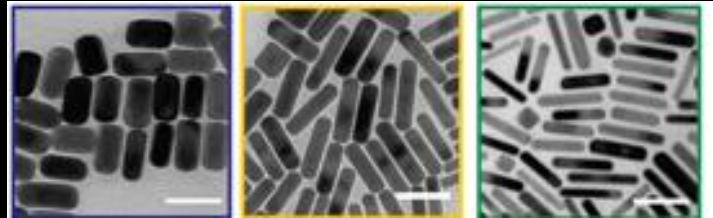
Shape ???

$$C_{ext} = \frac{24\pi^2 R^3 \epsilon_m^{3/2}}{\lambda} \times \frac{\epsilon_2(\lambda)}{(\epsilon_1(\lambda) + 2\epsilon_m)^2 + \epsilon_2(\lambda)^2}$$

$$\epsilon_{Drude}(\omega) = 1 - \frac{\omega_p^2}{\omega^2 + \gamma_d^2} + i \frac{\gamma_d \omega_p^2}{\omega(\omega^2 + \gamma_d^2)}$$

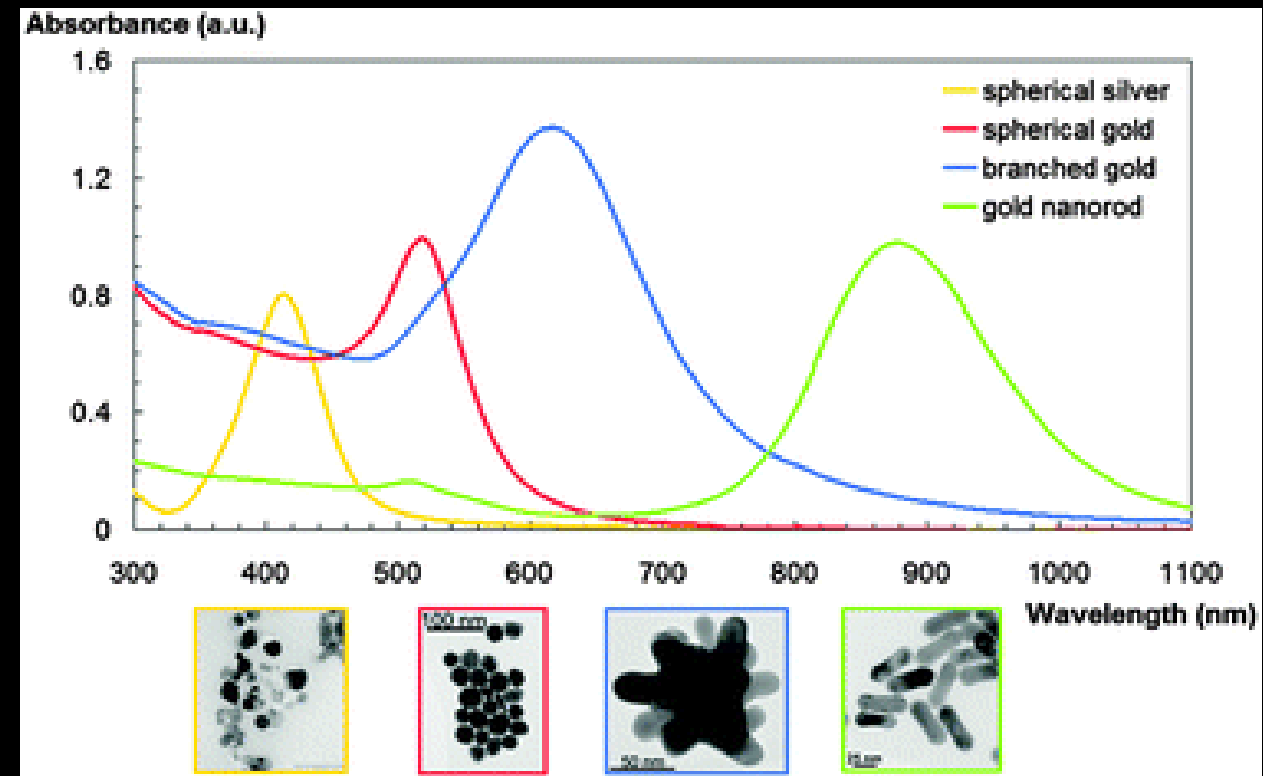
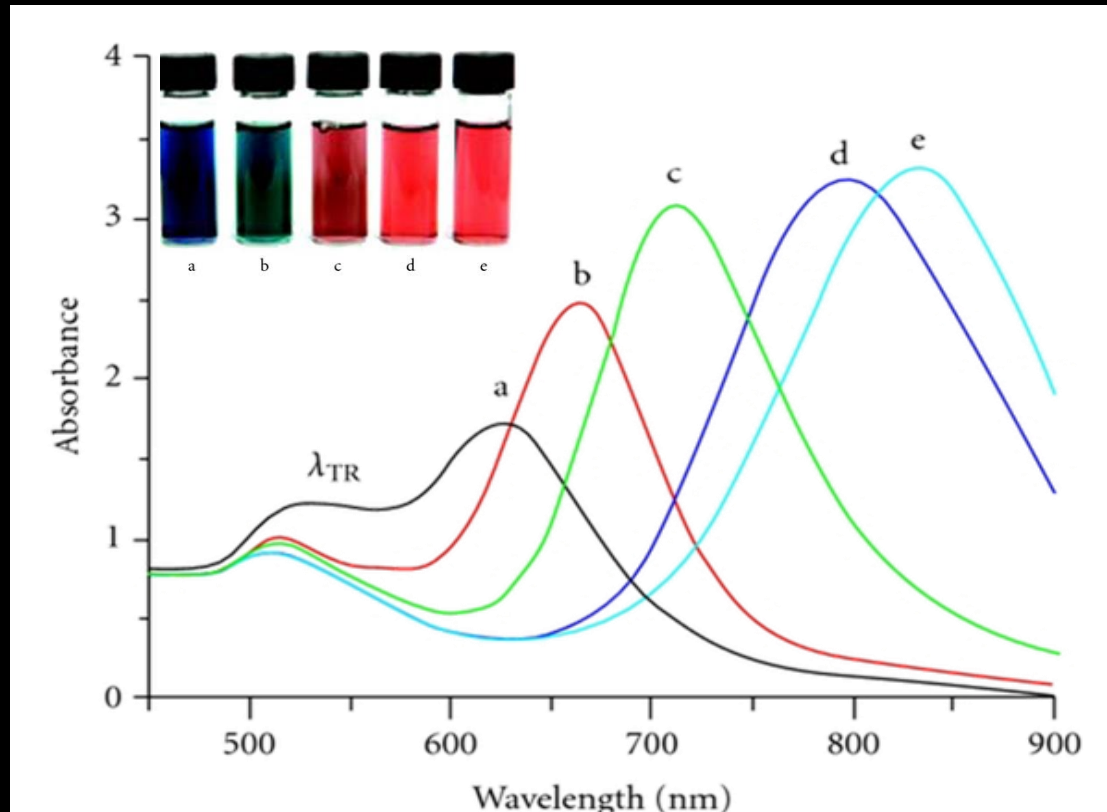
$$\omega = \frac{2\pi c}{\lambda}$$

Shape effect

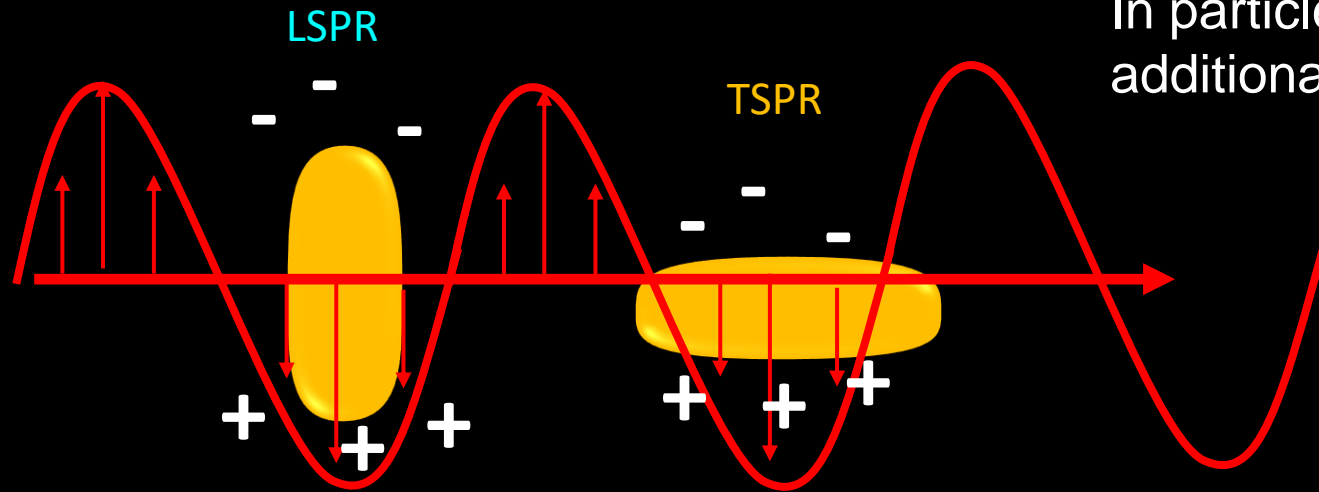


**Au
nanorods**

nanoparticles of different shapes can
show very different colors



Shape effect: NanoRods



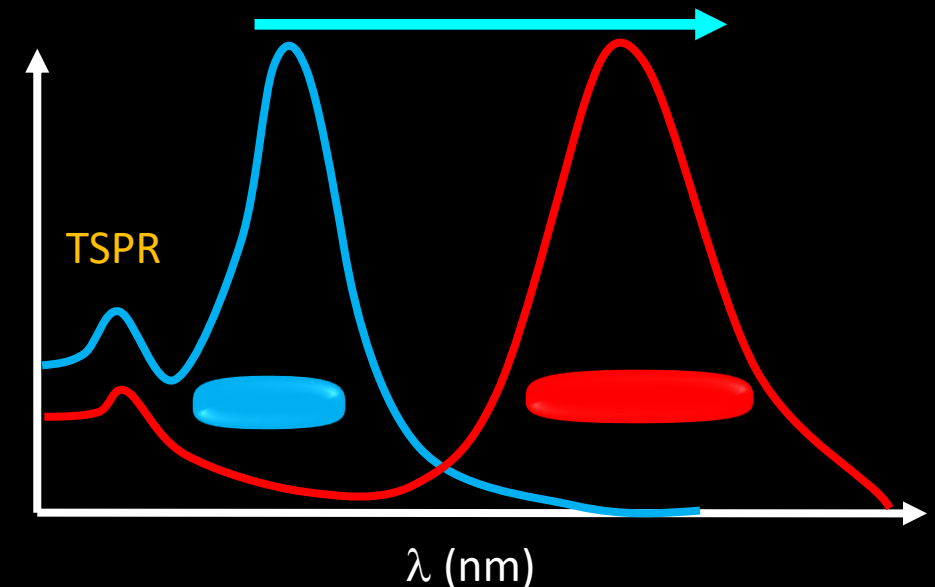
In particles without spherical symmetry there are additional modes of plasmon resonance (*Gans theory*)

Metal nanorods present two SPR modes: **longitudinal** and **transversal**

LSPR increases with aspect ratio (L/W)

Metal nanorods are more easily polarized longitudinally: LSPR at lower energies, longer λ

Increasing the aspect ratio (*length/width*), the LSPR can be tuned from 550 to over 2000 nm, while the TSPR remains relatively constant at ~ 520 nm



Shape effect: NanoRods

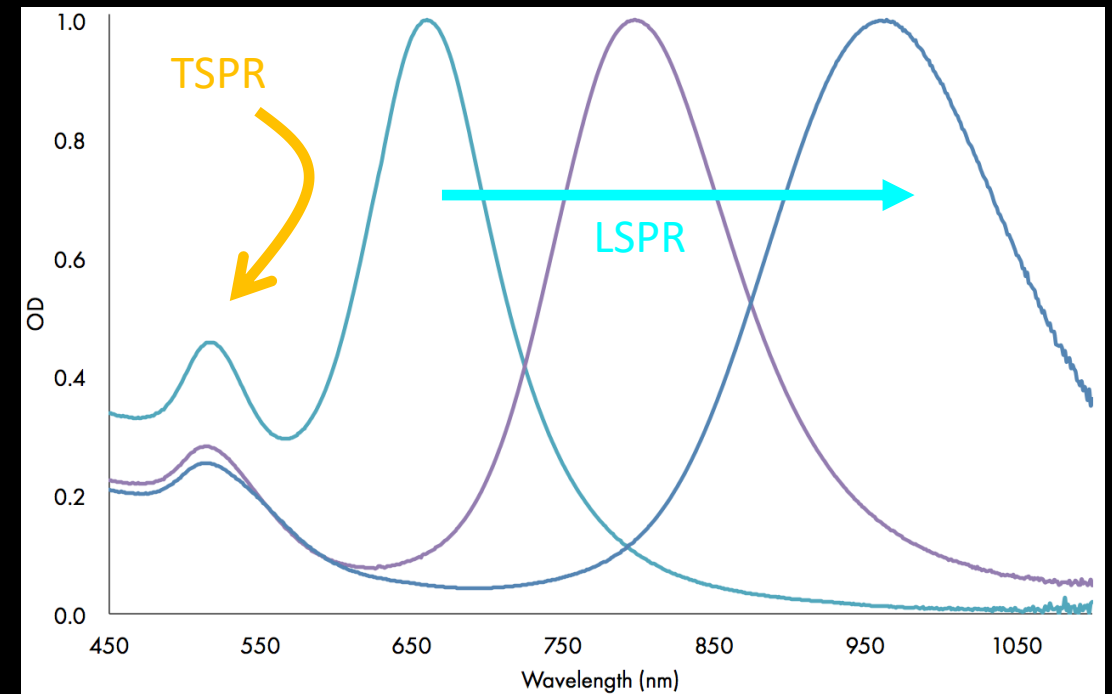
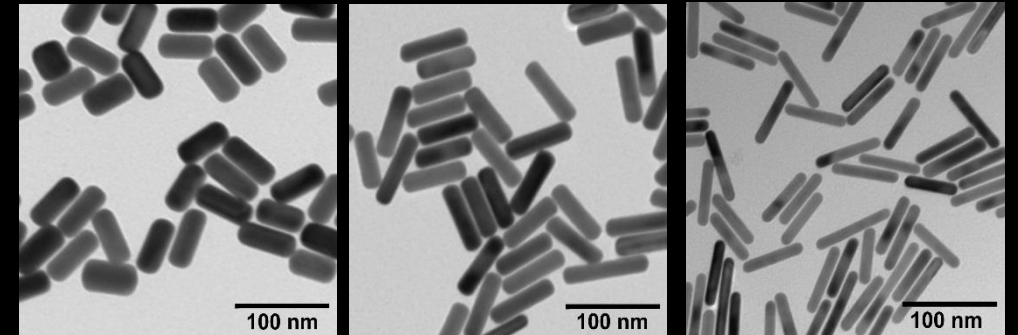
Variation of *longitudinal* and *transverse* SPR modes of gold nanorods (AuNR) with aspect ratio

Resonance	Length	Width	Aspect Ratio
660 nm	50 nm	19 nm	2.7
800 nm	70 nm	19 nm	3.6
980 nm	70 nm	12 nm	6.1

Empiric expression

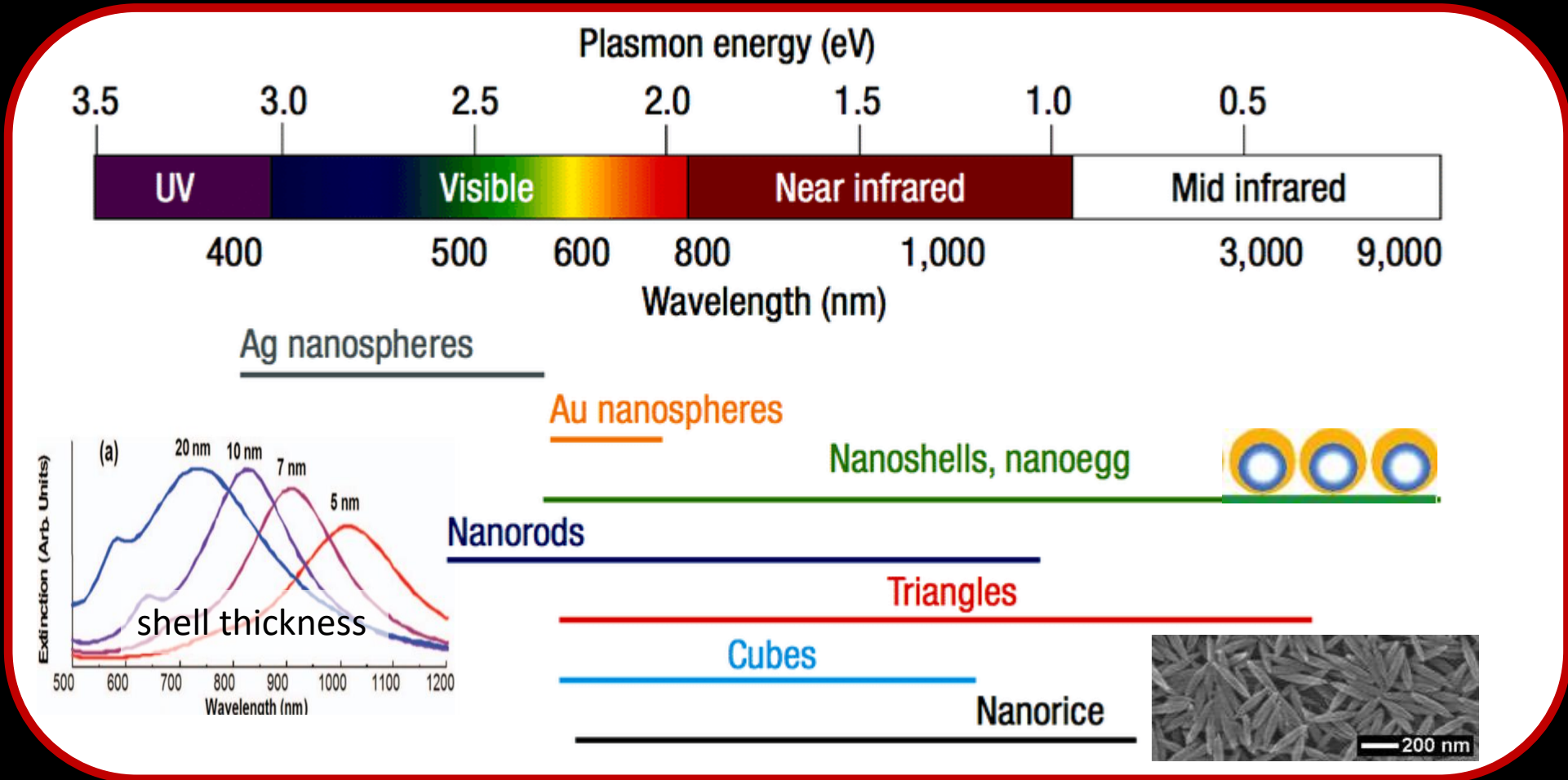
$$\lambda_{LSPR} = (K_1 D - K_2) \varepsilon_m + K_3$$

D = aspect ratio = main/minor axis



Shape effect

Range of SPR modes in metal nanoparticles with different morphology

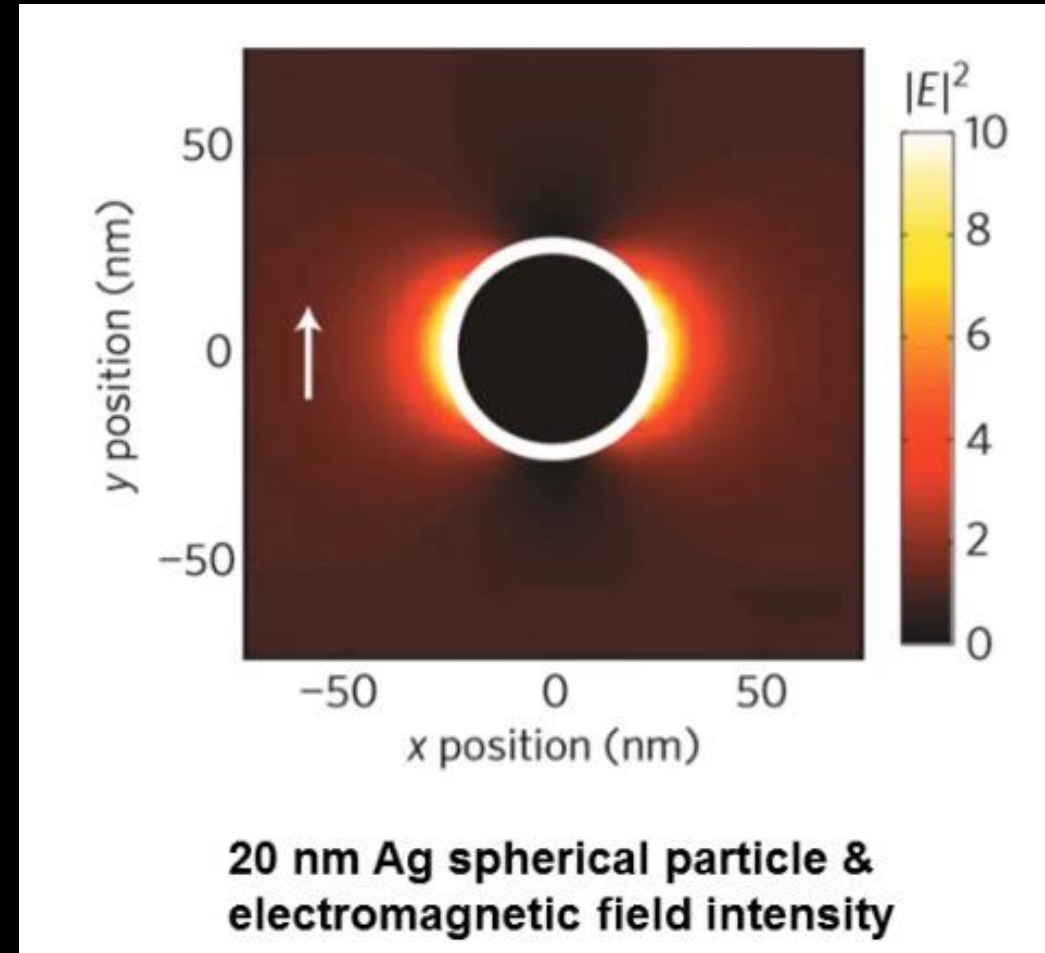


Field Enhancement

In NPs the plasmon field can couple with the incident electric field resulting in field enhancement

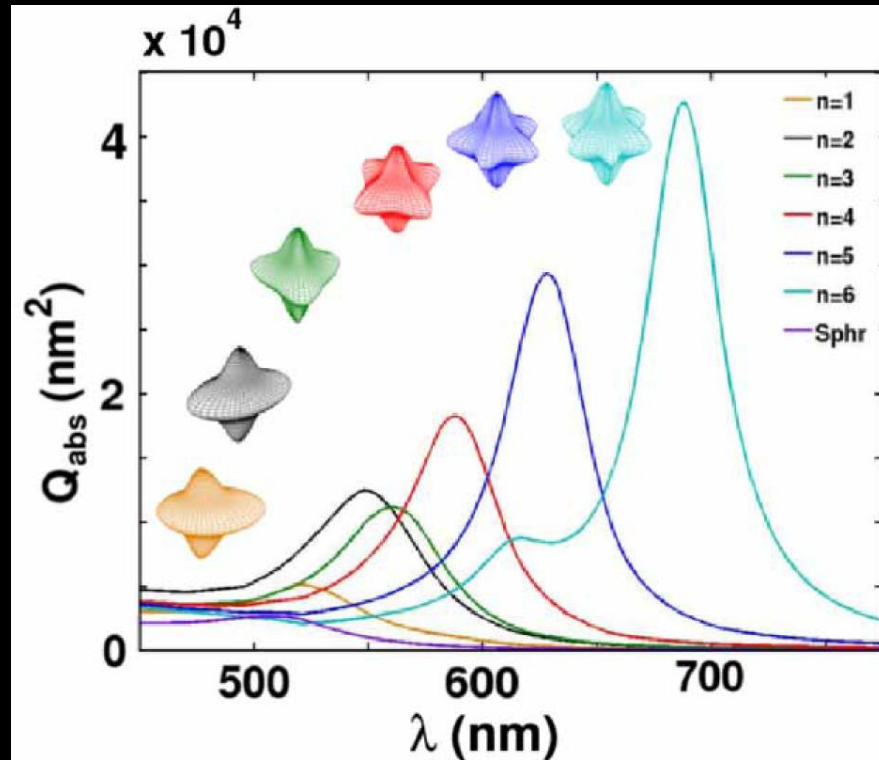
$$\frac{|E_{sphere}|}{|E_0|} = 1 + \left| \frac{\varepsilon_1 - \varepsilon_m}{\varepsilon_1 + 2\varepsilon_m} \right|$$

field enhancement is maximized when the real part of the denominator vanishes

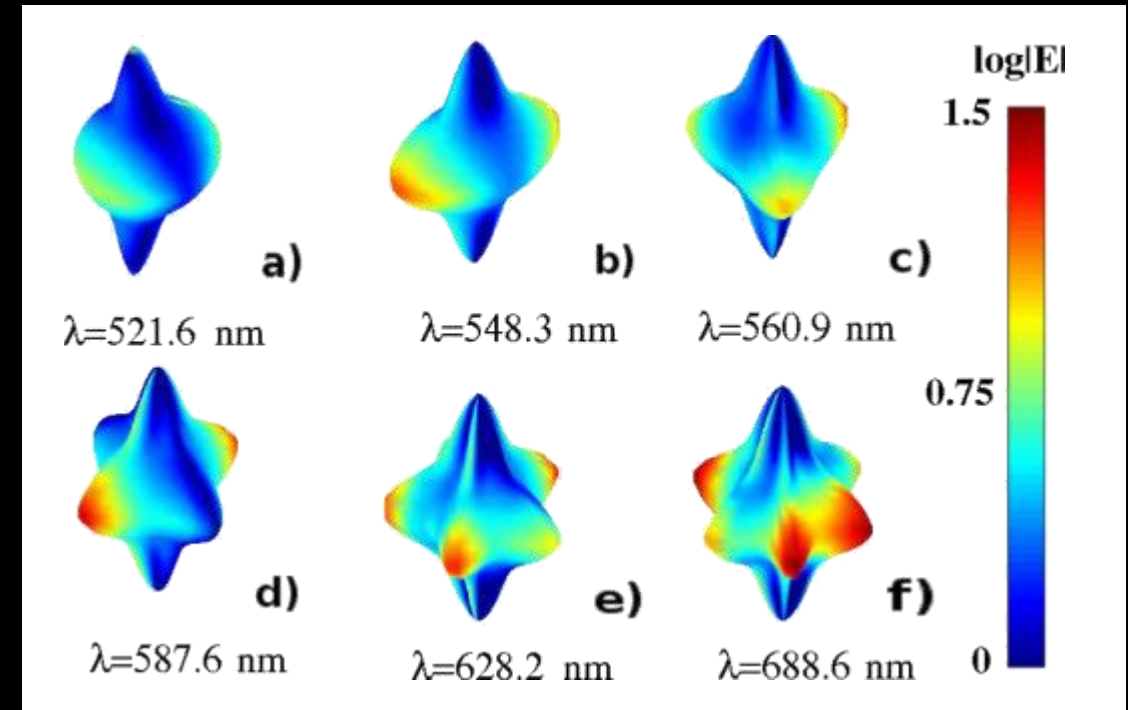


Field Enhancement: hotspots

Metal nanoparticles present surface electrical field enhancement due to the SPR. At high curvature points there are **hotspots** with higher reactivity and field enhancement



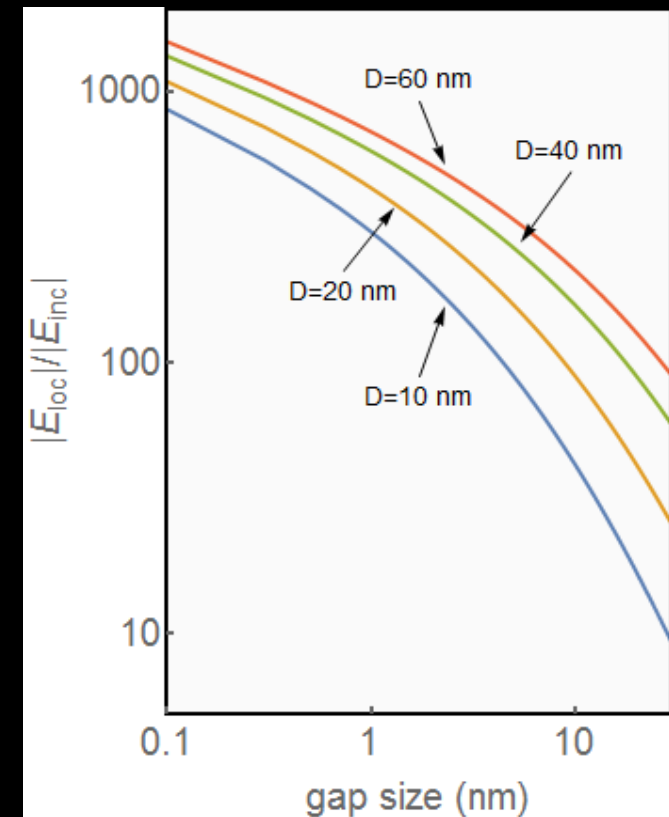
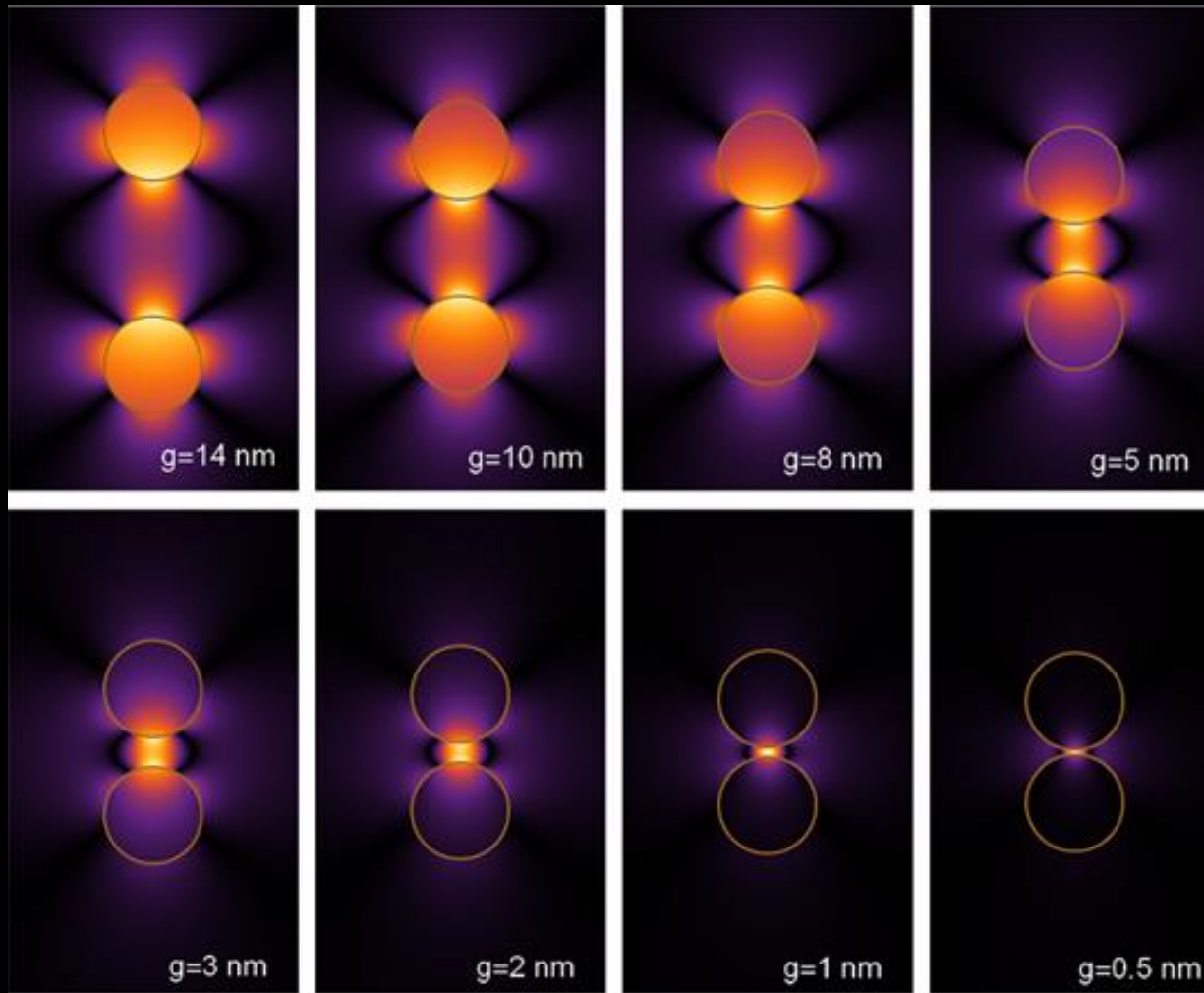
Variation of SPR modes of **gold nanostars** (AuNS) with morphology



Distribution of electric field amplitudes on the surface of AuNSs, with corresponding LSPR wavelengths

Field Enhancement: coupled NPs

When NPs interact, the surface plasmon resonances couple, shifting the resonance frequencies and changing the field characteristics.

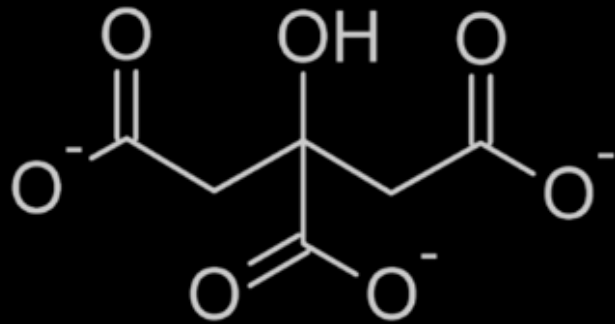




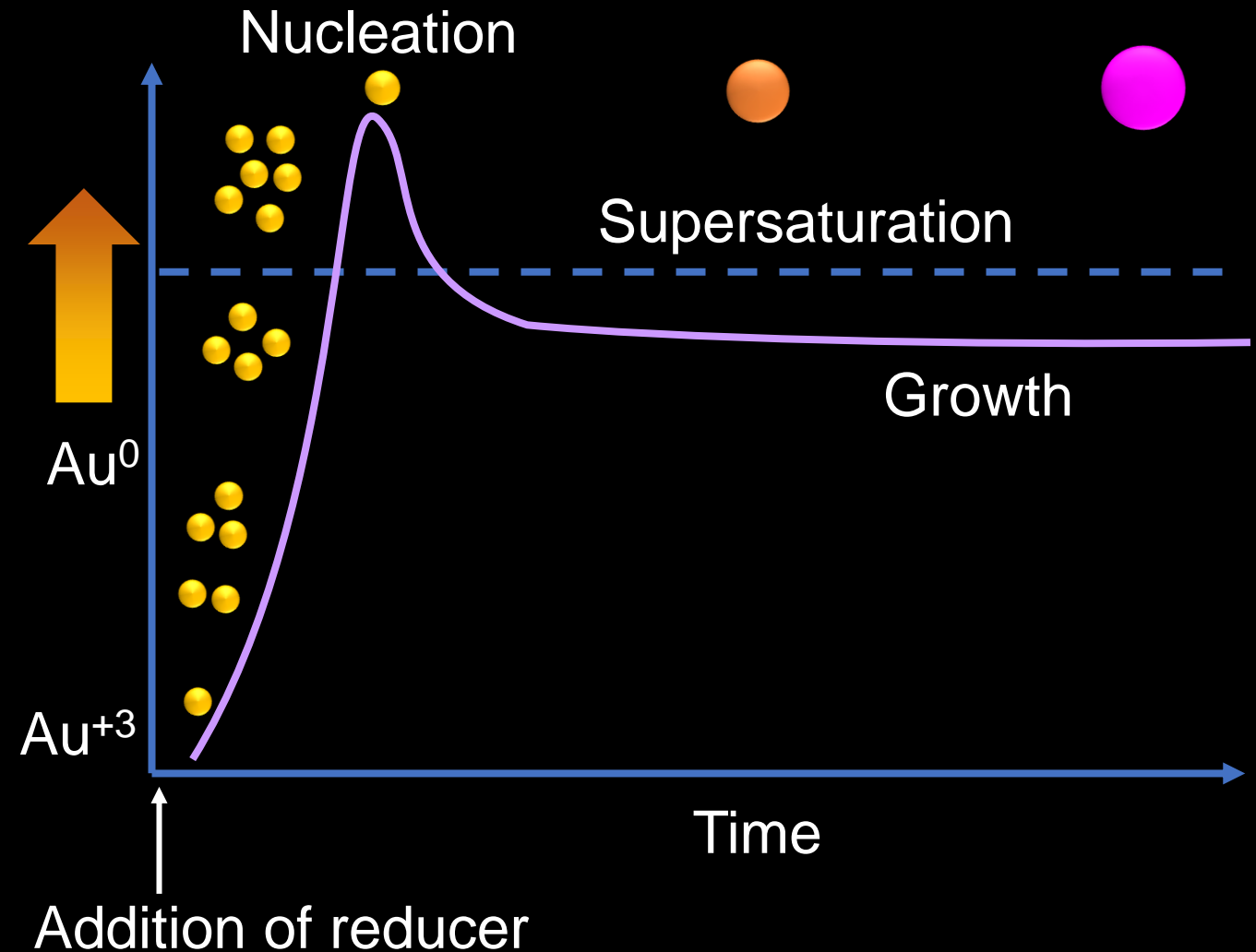
Metal nanoparticles: preparation

Metal NPs

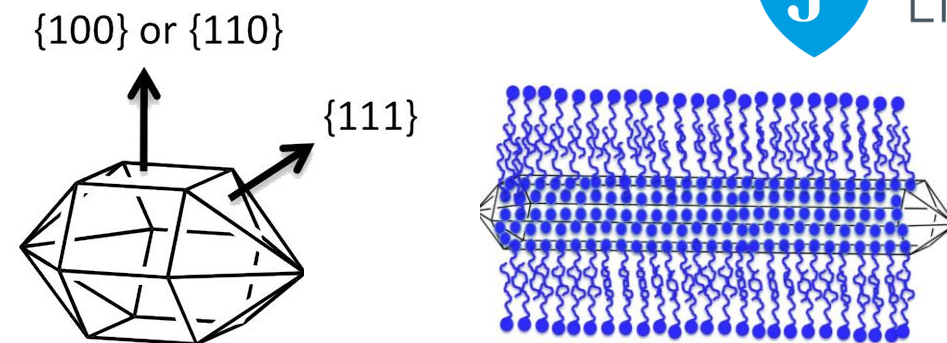
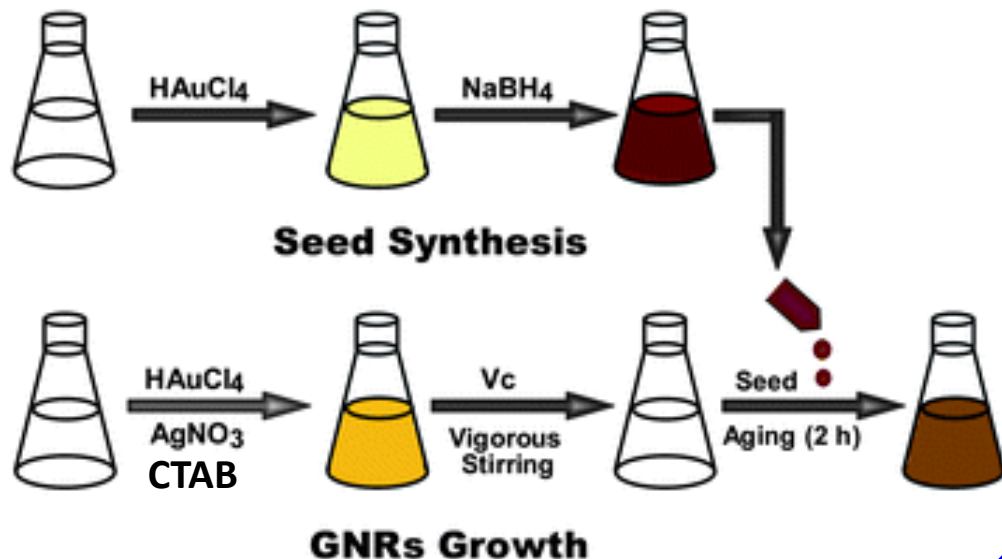
Au and Ag colloidal particles usually prepared in aqueous dispersion from HAuCl_4 and AgNO_3 , in the presence of a reducer, for example of citrate



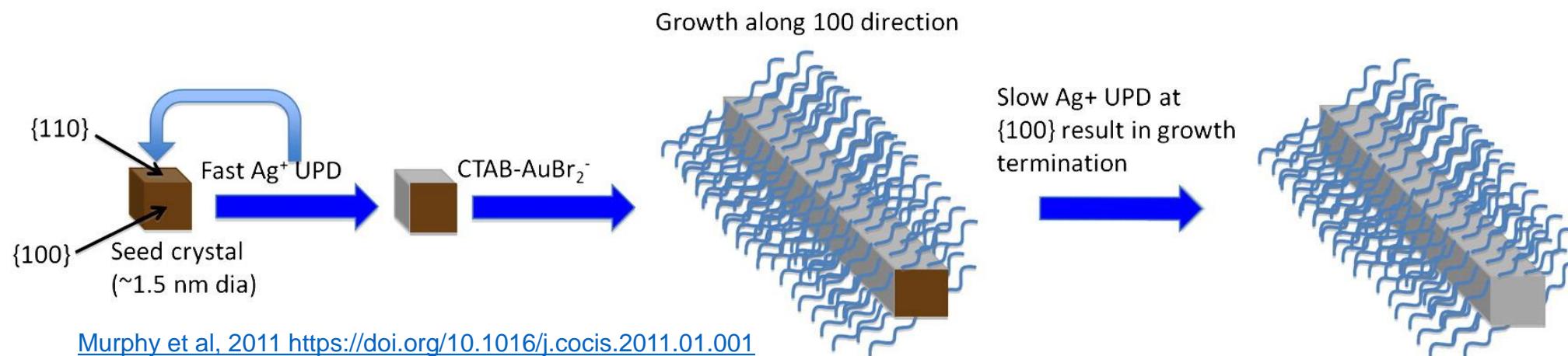
ascorbic acid,
hydroxylamine, sodium
borohydride (NaBH_4)



Nanorods seed-mediated growth



Bilayer packing density is lower near the tip due to its curvature
The probability for micellized gold ions to approach the tips is higher compared to the sides



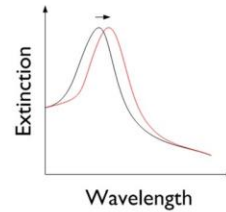
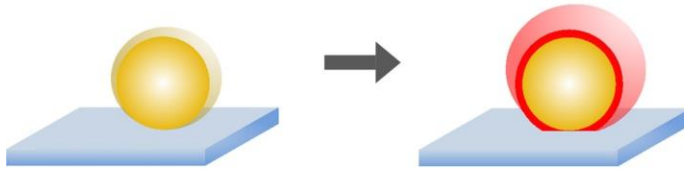
Murphy et al, 2011 <https://doi.org/10.1016/j.cocis.2011.01.001>



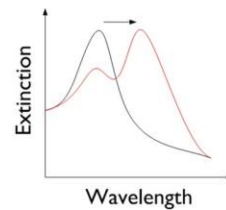
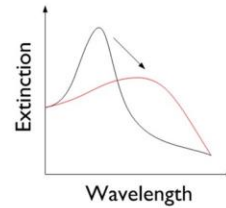
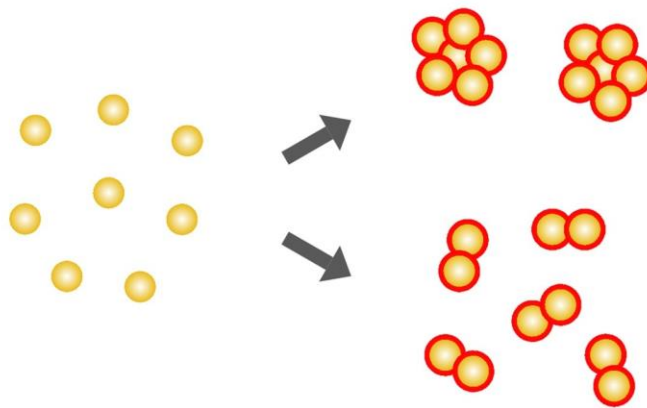
Metal nanoparticles: Application

Sensing

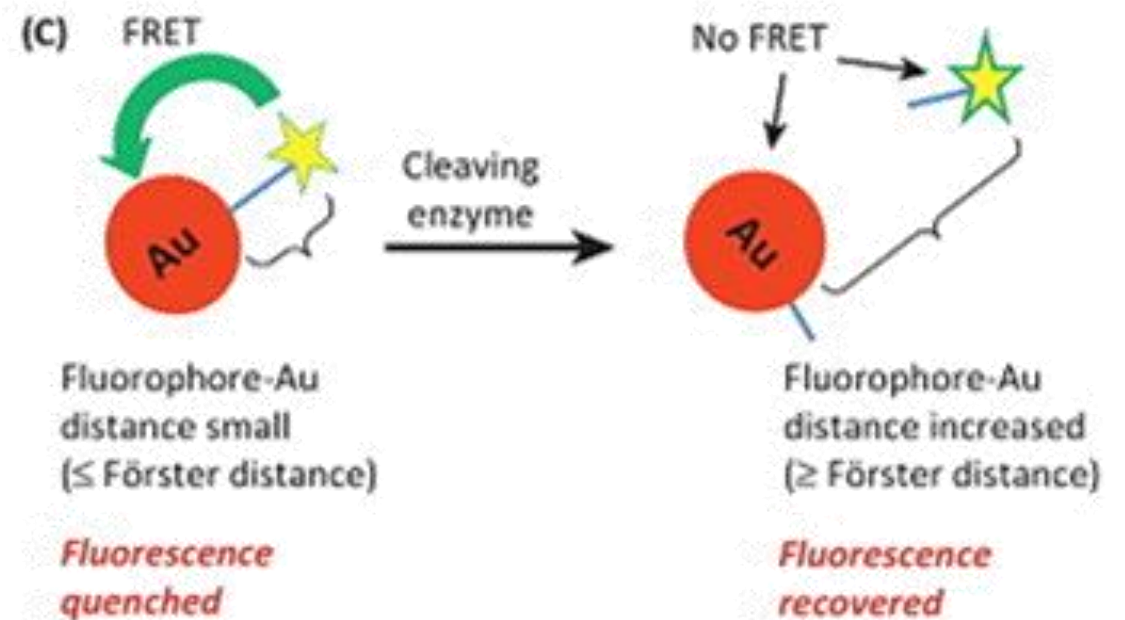
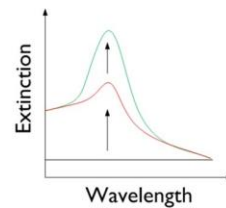
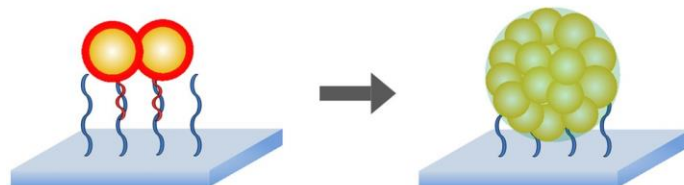
Refractive Index Change



Colorimetric Sensing via LSPR Coupling



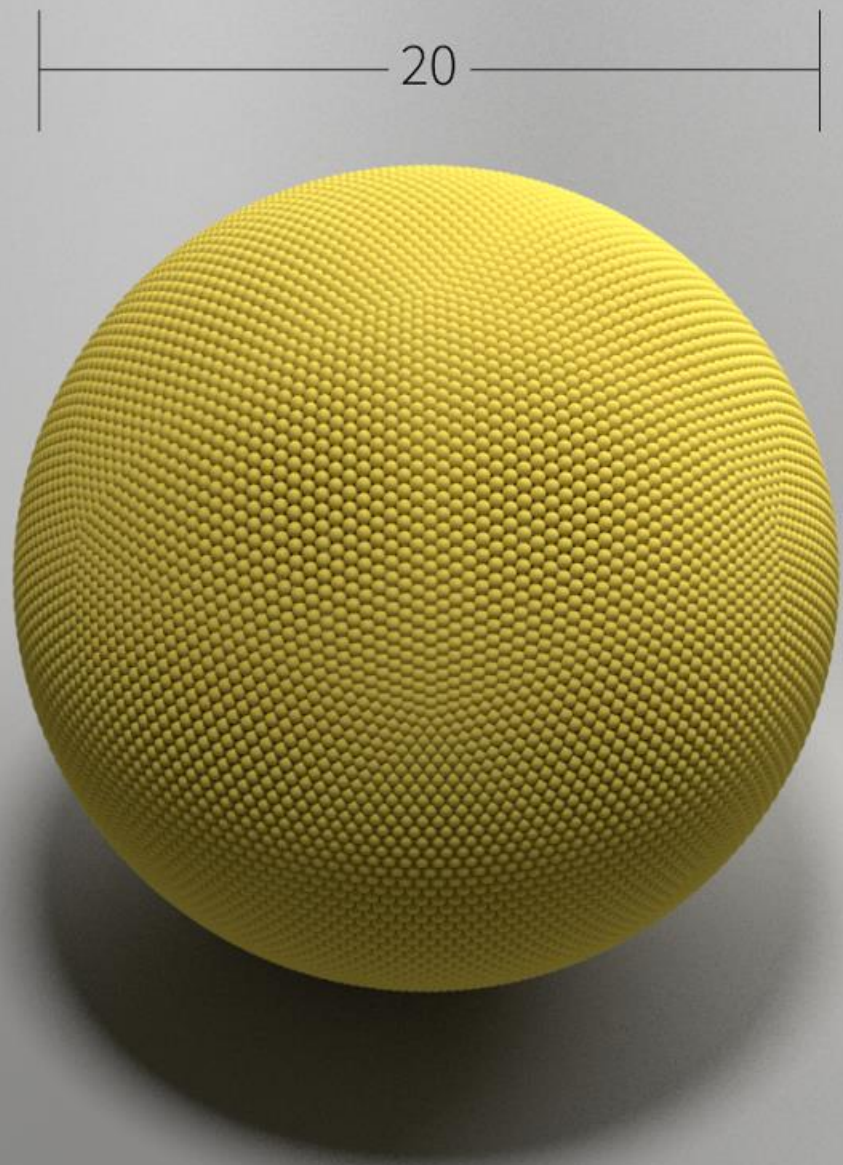
Signal Amplification by Nanoparticle Growth



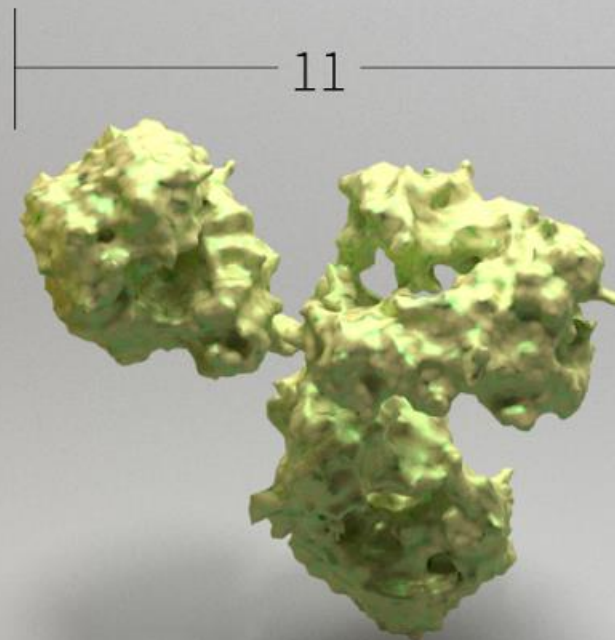
Distance between fluorophore and NP:
< 5 nm quenching effect
10-30 nm fluorescence enhancement

Trends in Pharmacological Sciences 2013, 34, 497

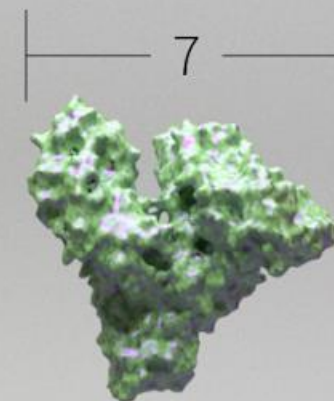
Nanotoday, 2015, <https://doi.org/10.1016/j.nantod.2015.02.007>



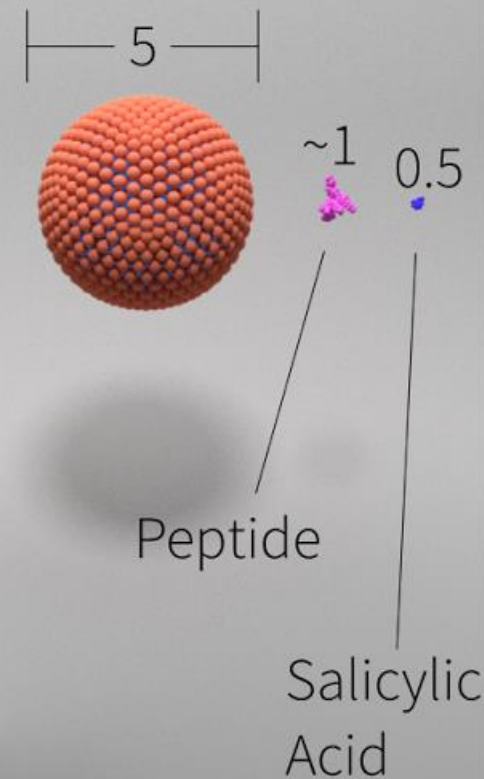
Gold Nanoparticle



Antibody



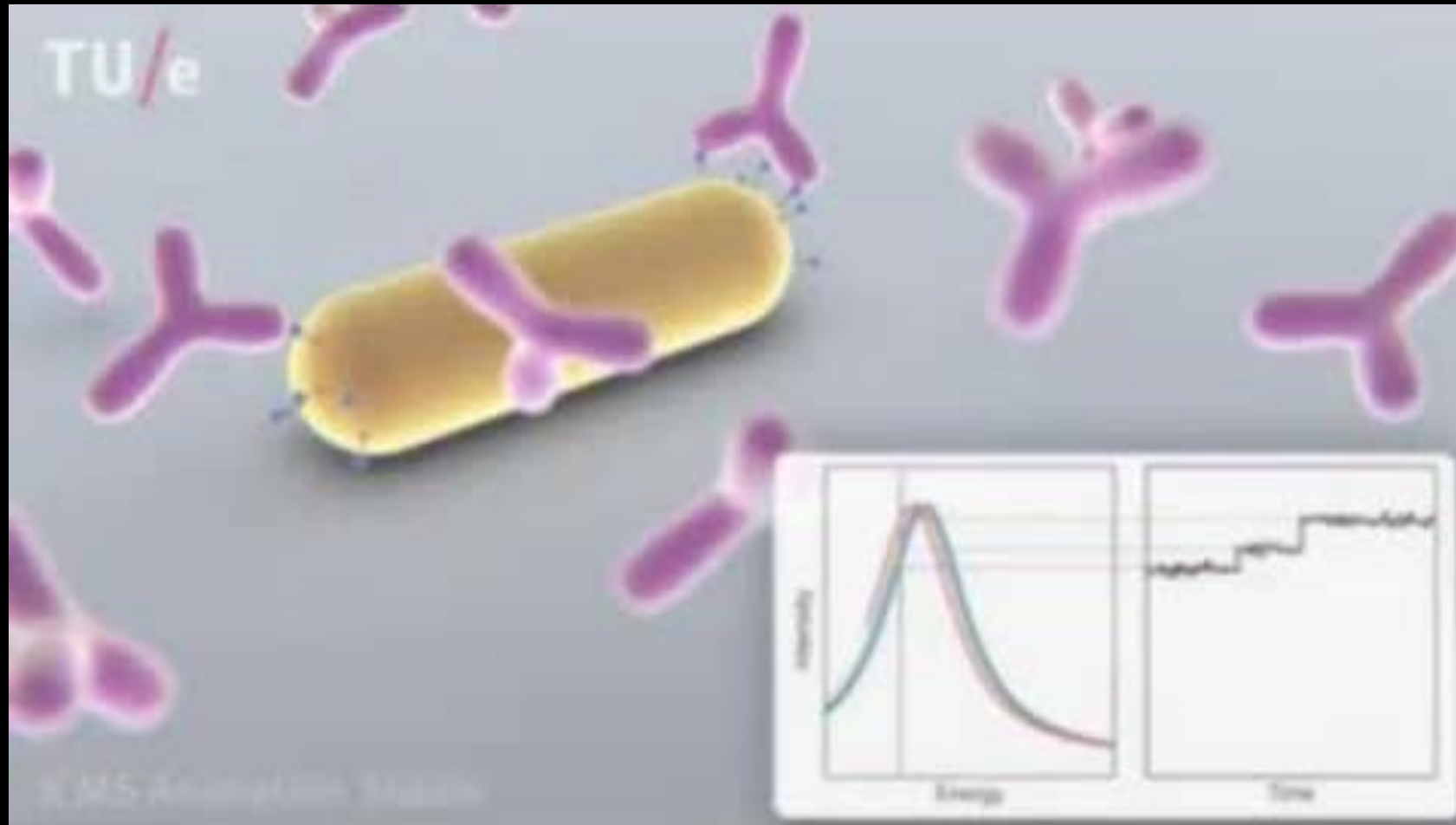
Albumin



Iron Oxide
Nanoparticle

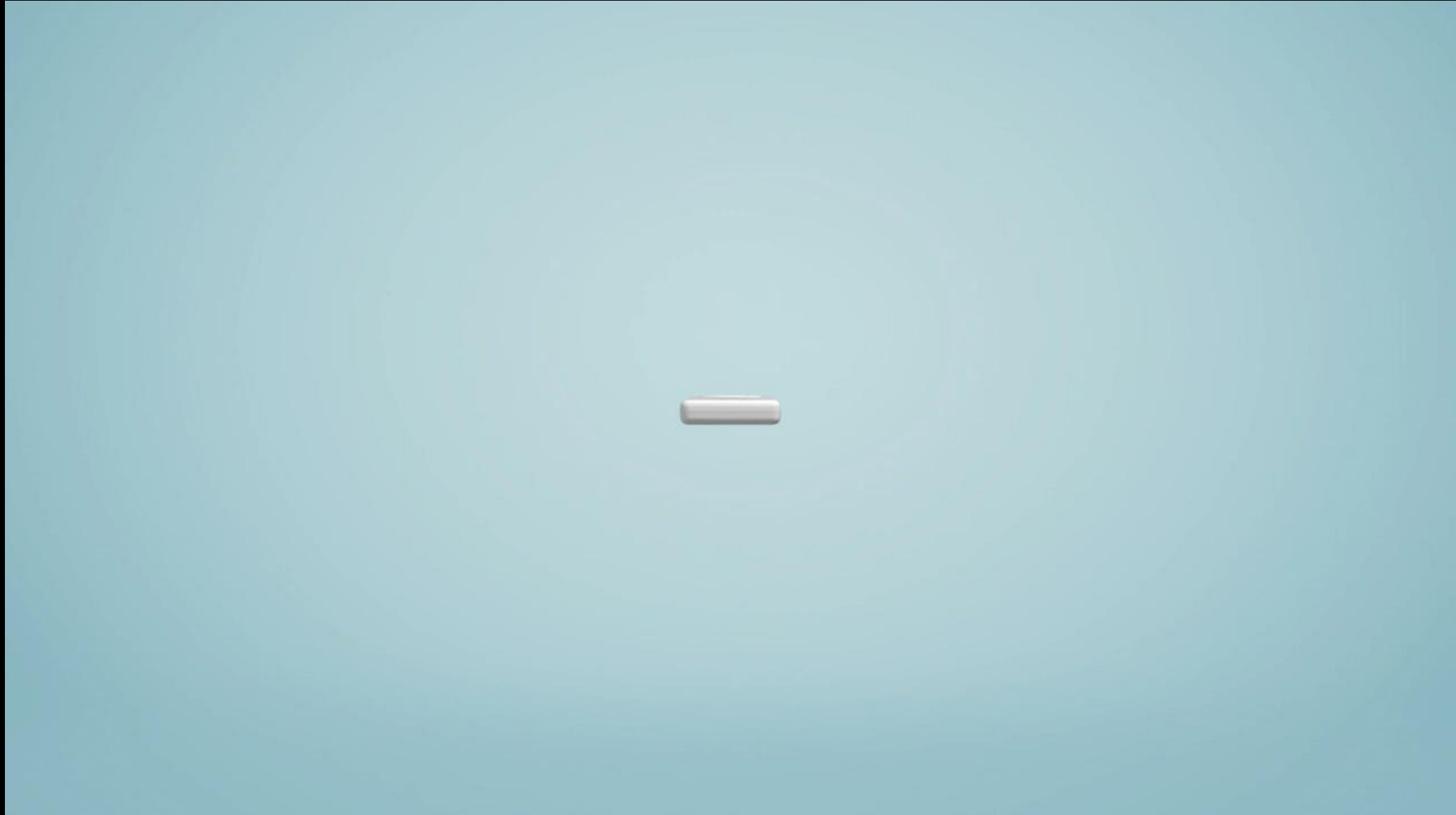
Single-Molecule Detection using Plasmons in Metal Nanoparticles

<https://youtu.be/HCdMXhvRD9A>



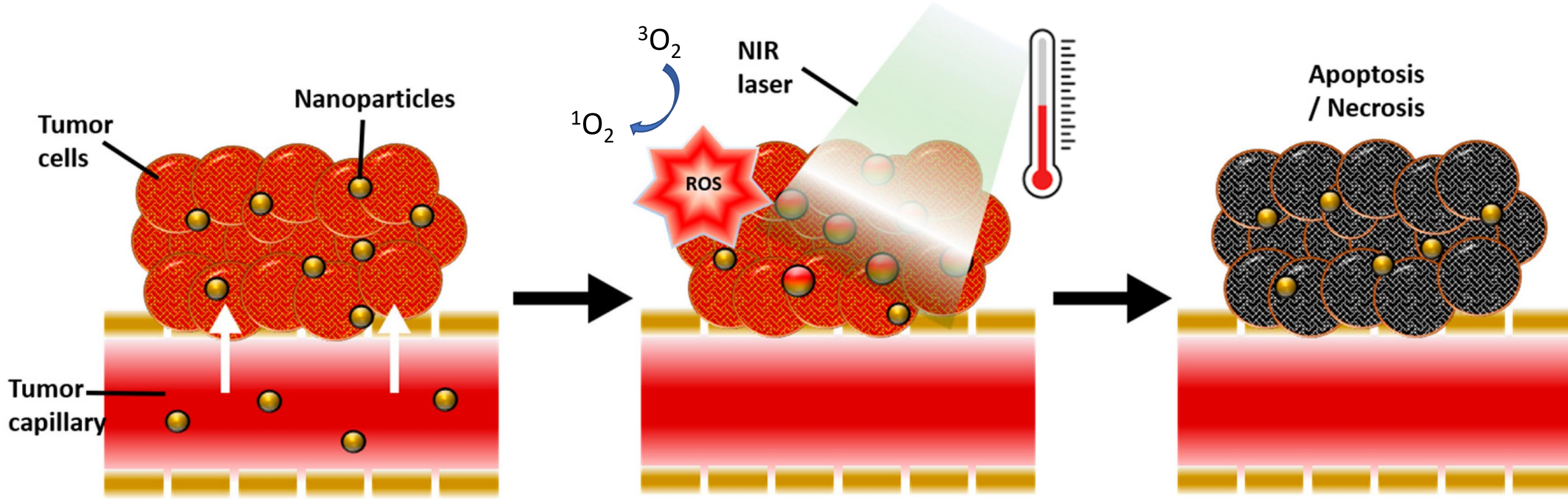
Lateral Flow Immunoassay

<https://youtu.be/z07CK-4JoFo>



Photothermal and photodynamic therapy

NPs with absorption in the NIR (typically NanoRods)
Most important effect comes from absorption of light

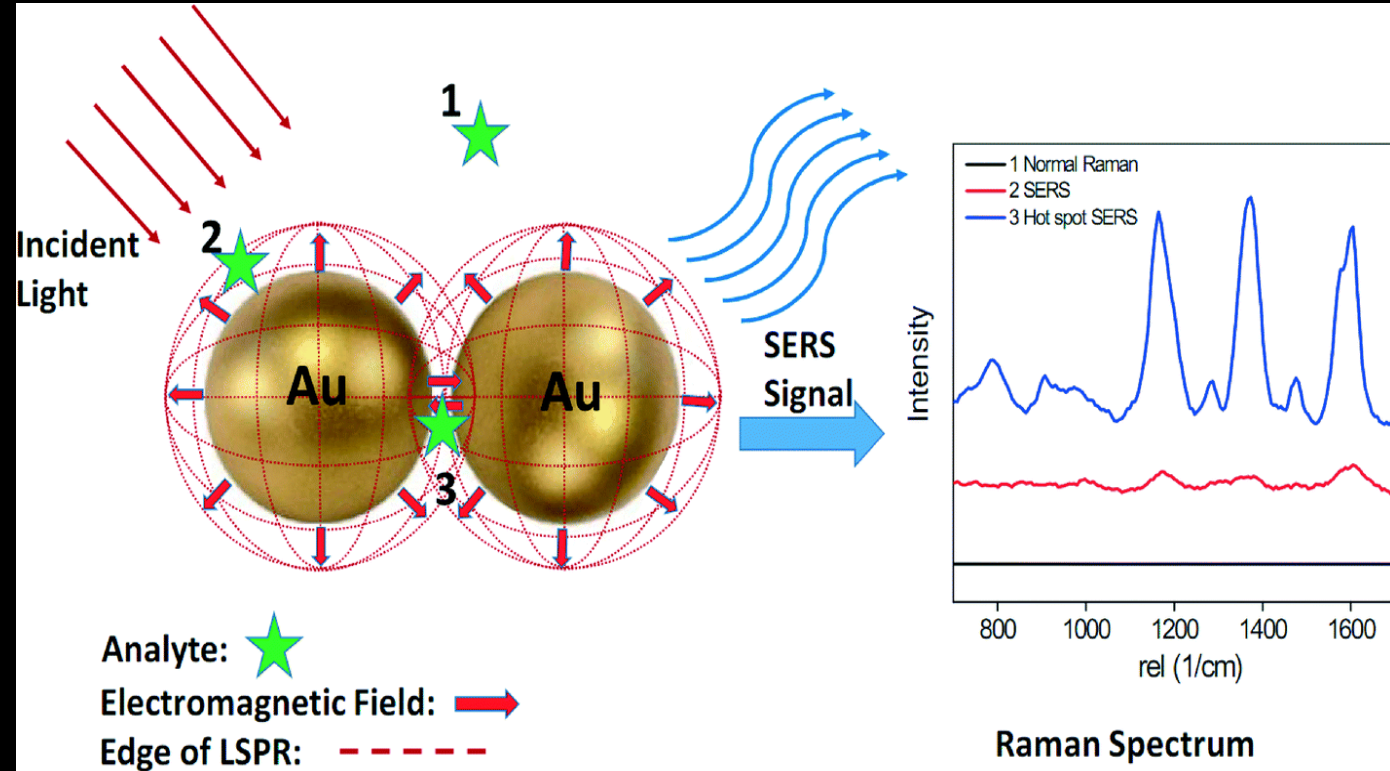


SERS (Surface Enhanced Raman Scattering) TÉCNICO LISBOA

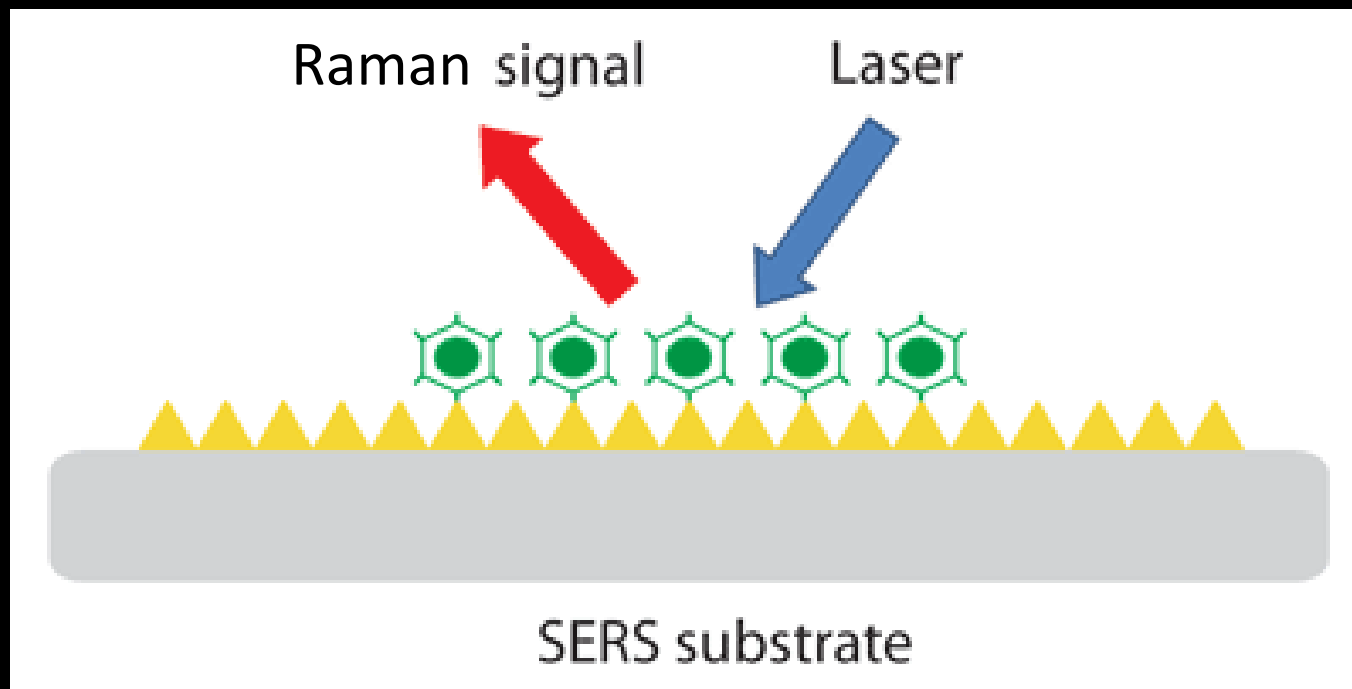
Raman spectroscopy

SERS

- $10^6 - 10^{14}$ amplification
- single molecule detection

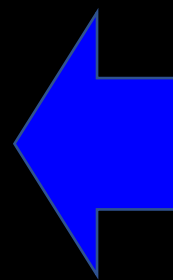


Raman scattering cross-section is extremely small compared with fluorescence process, which is usually 15 orders of magnitude smaller.
SP is ideal for enhancing Raman scattering processes.



SERS

- $10^6 - 10^{14}$ amplification
- single molecule detection



Coupling of LSPR with incident light (E_0)

$$\frac{|E_{sphere}|}{|E_0|} = 1 + \left| \frac{\varepsilon_1(\lambda_L) - \varepsilon_m}{\varepsilon_1(\lambda_L) + 2\varepsilon_m} \right|$$

Coupling of LSPR with the Raman scattered light (E_{mol})

$$\frac{|E_{sphere}|}{|E_0|} = 1 + \left| \frac{\varepsilon_1(\lambda_{Raman}) - \varepsilon_m}{\varepsilon_1(\lambda_{Raman}) + 2\varepsilon_m} \right|$$

λ_{Raman} and λ_L should both match the λ_{LSPR}

Applications

	Size (nm)						shape
	5	10	20	40	60	80	
Lateral flow			X	X	X	X	sphere
Western blot			X	X	X	X	Sphere
Dark filed microscopy				X	X	X	Sphere
Electron microscopy	X	X	X				sphere
Metal enhanced fluorescence			X	X			Sphere/spiky surface
SERS				X	X		Non-spherical/spiky
FRET quenching	X	X	X	X			sphere
Photothermal therapy				X			Rod/hollow sphere
Photodynamic therapy				X			Rod/hollow sphere
Drug delivery	X	X	X	X			Rod/sphere

The background features a central, glowing, hourglass-shaped nanostructure with a bright yellow-orange light source at its narrowest point. This central structure is surrounded by several smaller, spherical, textured particles in shades of blue and teal, all set against a dark, gradient background.

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ermelinda.macoas@tecnico.ulisboa.pt