Central limit theorem for fluctuations of linear eigenvalue statistics of large random graphs. Diluted regime.

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Abstract

We study the linear eigenvalue statistics of large random graphs in the regimes when the mean number of edges for each vertex tends to infinity. We prove that for a rather wide class of test functions the fluctuations of linear eigenvalue statistics converges in distribution to a Gaussian random variable with zero mean and variance which coincides with "non gaussian" part of the Wigner ensemble variance.

1 Introduction

In this paper we study the spectral properties of ensembles of adjacency matrices of large random graphs. Following Erdős (see, e.g. [6]), we introduce the probability measure considering the set of all graphs with n vertices and set the weight of each graph G as

$$P(G) = (p_n/n)^{e(G)} (1 - p_n/n)^{\binom{n}{2} - e(G)}, \tag{1.1}$$

where e(G) is the number of edges of G and $0 \le p_n \le n$. The set of n-vertices graphs with this measure (usually denoted by $\mathbf{G}(n, p_n/n)$) is one of the classes of the prime reference in the theory of random graphs. Most of the random graphs studies are devoted to the cases where $p_n/n \to 0$, as $n \to \infty$. There are two major asymptotic regimes: $p_n \gg 1$ and $p_n = O(1)$ and corresponding models can be called dilute random graphs and sparse random graphs, respectively.

It is well known that there is one-to-one correspondence between the graphs and their adjacency matrices. For $\mathbf{G}(n, p_n/n)$ the ensemble corresponding to (1.1) consists of random symmetric $n \times n$ adjacency matrices \widetilde{A} is $\widetilde{\mathcal{A}} = \{\widetilde{a}_{ij}\}_{i,j=1}^n$ with $\widetilde{a}_{ii} = 0$, and i.i.d.

$$\widetilde{a}_{ij} = \begin{cases}
1, & \text{with probability } p_n/n, \\
0, & \text{with probability } 1 - p_n/n,
\end{cases}$$
(1.2)

This is a particular case of the random matrix ensemble. Since the pioneering works by Wigner [19] a big part of the random matrix theory is devoted to the limiting transition $n \to \infty$. The results obtained with this limiting transition provide a rather good approximation of the spectral properties of random matrices (or random graphs) of a finite dimensionality.

An important advantage of random matrices (1.2) is that their entries are independent up to the symmetry condition $(a_{ij} = a_{ji})$. This allows one to use the methods of random matrix theory which were developed to study the classical Wigner matrices. Spectral properties of random adjacency matrix (1.4) were examined in the limit $n \to \infty$ both in numerical and theoretical physics studies. The first results on the spectral properties of sparse and dilute random matrices in the physical literature are related with the works [12], [13], [11],

where equations for the limiting density of states of sparse random matrices were derived. In the papers [11] and [5] a number of important results on the universality of the correlation functions and the Anderson localization transition were obtained. Unfortunately, these results were obtained with non rigorous replica and super symmetry methods.

The first result on mathematical level of rigor for the matrices (1.2) was obtained in [2], where the eigenvalue distribution moments of the matrix (1.2) with $p_n = p$ were studied in the limit $n \to \infty$. It was shown that for any fixed natural m there exists nonrandom limiting moment $\lim_{n\to\infty} n^{-1} \operatorname{Tr} A^m$ and these moments can be found from the system of certain recurrent relations. The results of [2] were generalized to the case of weighted random graphs in [8], where the resolvent of the adjacency matrix was studied and equations for the Stieltjes transform g(z) of the limiting eigenvalue distribution were derived rigorously (note, that the same equation for gaussian weights were obtained in [12], [13], [11] by using the replica and the super symmetry approaches). But the limiting eigenvalue distribution, which is an analog of the low of large numbers of the probability theory, is only the first step in studies of linear eigenvalue statistics, corresponding to the test function φ

$$\mathcal{N}_n[\varphi] = \sum \varphi(\lambda_i) = \operatorname{Tr} \varphi(\mathcal{A}).$$
 (1.3)

Here and below $\{\lambda_i\}_{i=1}^n$ are eigenvalues of the matrix \mathcal{A} . The next step is to study the behavior of fluctuations of linear eigenvalue statistics. For the case of sparse random matrices this step was done in [17] with some modification of the method of [8]. It was shown in [17] that the random variable $n^{-1/2}(\mathcal{N}_n[\varphi] - \mathbf{E}\{\mathcal{N}_n[\varphi]\})$ converges in distribution to the gaussian random variable, as $n \to \infty$ (here and below $\mathbf{E}\{...\}$ means the averaging with respect to all $\{a_{ij}\}_{1 \le i < j \le n}$).

The case of diluted matrices $(p_n \to \infty)$ is less complicated technically than that with $p_n = p$. It was shown in [7] that in this case to have finite limits for $\mathbf{E}\{\mathcal{N}_n[\varphi]\}$ one should consider the matrix $\mathcal{A}' = \widetilde{\mathcal{A}}/\sqrt{p_n}$. Then it was proven in [7] that for integrable test functions φ

$$\lim_{n,p_n\to\infty,p_n/n\to 0} \mathbf{E}\{\mathcal{N}_n[\varphi]\} = \frac{1}{2\pi} \int_{-2}^2 \varphi(\lambda) \sqrt{4-\lambda^2} d\lambda,$$

which coincides with the limits for the Wigner model [19]. Let us note that the method, used in [7], is rather similar to that for the Wigner model. But the problem to study the fluctuations of linear eigenvalue statistics usually is much more complicated than the problem to find the limiting eigenvalue distribution of random matrix ensemble. Even for the classical Wigner case the central limit theorem (CLT) for fluctuations of linear eigenvalue statistics was proven only recently in the series of papers with improving results [15, 1, 9, 18].

In the present paper we prove CLT for fluctuations of linear eigenvalue statistics of diluted matrices, more precisely, we prove that the random variable $(p_n/n)^{1/2}(\mathcal{N}_n[\varphi] - \mathbf{E}\{\mathcal{N}_n[\varphi]\})$ in the limit $n, p_n \to \infty, p_n/n \to 0$ converges in distribution to the normal random variable. The method of the paper is a generalization of that of [18]. It allows us to prove CLT under rather weak assumptions on the test function φ (see Theorems 1 and 2 below).

It will be more convenient for us to study the matrix $\mathcal{A} = \widetilde{\mathcal{A}}/\sqrt{p_n} - \mathbf{E}\{\widetilde{\mathcal{A}}/\sqrt{p_n}\}$, where $\mathbf{E}\{...\}$ means averaging with respect to all entries of $\widetilde{\mathcal{A}}$. It is easy to see that \mathcal{A} differs from \mathcal{A}' by the rank one matrix $\mathbf{E}\{\widetilde{\mathcal{A}}/\sqrt{p_n}\}$. So, everywhere below we will assume that the entries a_{ij} of \mathcal{A} are distributed as

$$a_{ij} = \begin{cases} \frac{1}{\sqrt{p_n}} - \frac{\sqrt{p_n}}{n}, & \text{with probability } p_n/n, \\ -\frac{\sqrt{p_n}}{n}, & \text{with probability } 1 - p_n/n, \end{cases}$$
(1.4)

Let us note that the case $p_n \sim \alpha n$ here corresponds to the Wigner ensemble, hence the model (1.4) allows us to make "smooth transition" from the matrix studied in [8] to the Wigner matrix.

Let us set our main notations. For any measurable function f we denote by $\mathbf{E}\{f(\mathcal{A})\}$ the averaging with respect to all random variables $\{a_{ij}\}_{1 \leq i < j \leq n}$ and

$$\mathbf{Var}\{f(\mathcal{A})\} := \mathbf{E}\{|f(\mathcal{A}) - \mathbf{E}\{f(\mathcal{A})\}|^2\}. \tag{1.5}$$

We denote also for any random variable ξ

$$\overset{\circ}{\xi} = \xi^{\circ} = \xi - \mathbf{E}\{\xi\}.$$

Introduce the resolvent of A

$$G_{jk}(z) = (A - z)_{jk}^{-1}, \quad \Im z \neq 0, \quad \gamma_n(z) = \text{Tr } G(z).$$
 (1.6)

In what follows it will be important for us that

$$||G|| \le |\Im z|^{-1}, \quad \sum_{i=1}^{n} |G_{ij}|^2 = (GG^*)_{ii} \le ||G||^2 \le |\Im z|^{-2},$$
 (1.7)

$$\Im(Ge, e)\Im z \ge 0, \quad \forall e \in \mathbb{R}^n.$$
 (1.8)

Here and everywhere below $||\mathcal{A}||$ means the operator norm of the matrix \mathcal{A} .

The main result of the paper is the central limit theorem for the linear eigenvalue statistics of any sufficiently smooth function φ which grows not faster than exponential at infinity. But we prove CLT first for the functions, which are smooth enough and decaying. Set

$$||\varphi||_s^2 = \int (1+2|k|)^{2s} |\widehat{\varphi}(k)|^2 dk, \quad \widehat{\varphi}(k) = \frac{1}{2\pi} \int e^{ikx} \varphi(x) dx$$
 (1.9)

and let \mathcal{H}_s be the space of all function possessing the norm $||.||_s$.

Theorem 1 Consider the adjacency matrix (1.4) with $p_n \to \infty$, $p_n/n \to 0$. Assume that the real valued function $\varphi \in \mathcal{H}_s$ with s > 3/2 and that

$$\int_{-2}^{2} \varphi(\mu) \frac{2 - \mu^2}{\sqrt{4 - \mu^2}} d\mu \neq 0. \tag{1.10}$$

Then the random variable $(p_n/n)^{1/2} \mathring{\mathcal{N}}_n[\varphi]$ converges in distribution to a Gaussian random variable with zero mean and variance

$$V[\varphi] = \frac{1}{2\pi^2} \left(\int_{-2}^2 \varphi(\mu) \frac{2 - \mu^2}{\sqrt{4 - \mu^2}} d\mu \right)^2.$$
 (1.11)

It is interesting to compare (1.11) with that for the Wigner model

$$M = n^{-1/2} \{ w_{ij} \}_{i,j=1}^n, \quad E\{w_{ij}\} = 0, \quad E\{|w_{ij}|^2\} = 1, \quad (i \neq j),$$

we have (see [18])

$$V_{W}[\varphi] = \lim_{n \to \infty} \mathbf{Var} \{ \mathring{\mathcal{N}}_{n}[\varphi] \} = \frac{1}{2\pi^{2}} \int_{-2}^{2} \int_{-2}^{2} \left(\frac{\varphi(\lambda_{1}) - \varphi(\lambda_{2})}{\lambda_{1} - \lambda_{2}} \right)^{2} \frac{(4 - \lambda_{1}\lambda_{2})d\lambda_{1}d\lambda_{2}}{\sqrt{4 - \lambda_{1}^{2}}\sqrt{4 - \lambda_{2}^{2}}} + \frac{\kappa_{4}}{2\pi^{2}} \left(\int_{-2}^{2} \varphi(\mu) \frac{2 - \mu^{2}}{\sqrt{4 - \mu^{2}}} d\mu \right)^{2} + \frac{w_{2} - 2}{4\pi^{2}} \left(\int_{-2}^{2} \frac{\varphi(\mu)\mu d\mu}{\sqrt{4 - \mu^{2}}} \right)^{2},$$

Here $\kappa_4 = n^2(\mathbf{E}\{M_{ij}^4\} - 3\mathbf{E}^2\{M_{ij}^2\}) = \mathbf{E}\{w_{ij}^4\} - 3$, $w_2 = nE\{|M_{ii}|^2\}$. One can see that (1.11) coincides with the term multiplying κ_4 . This can be understood if we recall that in our case $\kappa_4 = n^2(\mathbf{E}\{a_{ij}^4\} - 3\mathbf{E}^2\{a_{ij}^2\}) \sim n/p_n$ and we consider the random variable $(p_n/n)^{1/2} \mathring{\mathcal{N}}_n[\varphi]$, while in the Wigner case one should consider $\mathring{\mathcal{N}}_n[\varphi]$.

One more interesting question is what is happening if the l.h.s. of (1.10) is zero. It is easy to guess that in this case one have to change the normalization factor in front of $\mathcal{N}_n[\varphi]$. But it could happen that the new expression for the limiting variance in this case will depend on the rate of convergence of $p_n/n \to 0$. We are going to study this situation in the future works.

Consider the set $\mathcal{H}_s^{(c)}$ of the functions, represented in the form

$$\varphi(\lambda) = \cosh(c\lambda)\,\widetilde{\varphi}(\lambda), \quad \widetilde{\varphi} \in \mathcal{H}_s.$$
 (1.12)

Theorem 2 Consider the adjacency matrix (1.4) with $p_n \to \infty$, $p_n/n \to 0$. Assume that the real valued function $\varphi \in \mathcal{H}_s^{(c)}$ with some c > 0, s > 3/2 and (1.10) is satisfied. Then the random variable $(p_n/n)^{1/2} \mathring{\mathcal{N}}_n[\varphi]$ converges in distribution to a Gaussian random variable with zero mean and variance (1.11).

2 Proofs

The proof follows the strategy developed in citeS:10 for the Wigner model. We start from the lemma.

Lemma 1 Let $\gamma_n(z)$ be defined by (1.6). Then for any $z: \Im z > 0$ there exists a constant C such that

$$\frac{p_n}{n} \mathbf{Var}\{\gamma_n(z)\} \le C/|\Im z|^4, \quad \left(\frac{p_n}{n}\right)^2 \mathbf{E}\{|\gamma_n^{\circ}(z)|^4\} \le C/|\Im z|^{12}. \tag{2.1}$$

Moreover, for any $\varepsilon > 0$ we have

$$\frac{p_n}{n} \mathbf{Var}\{\gamma_n(z)\} \le C \mathbf{E}\{|G_{11}|^{1+\varepsilon}\}/|\Im z|^{3+\varepsilon}, \tag{2.2}$$

and for any smooth function F and any $z:\Im z>a$

$$\mathbf{Var}\Big\{n^{-1}\sum_{j=1}^{n} F(G_{jj}(z))\Big\} \le n^{-1} \sup_{\zeta:0<\Im\zeta,|\zeta|< a^{-1}} |F'(\zeta)|^{2}.$$
(2.3)

Proof of Lemma 1 To prove (2.1) we use the following proposition proven in [4]

Proposition 1 Let ξ_{α} , $\alpha = 1, ..., \nu$ be independent random variables, assuming values in $\mathbb{R}^{m_{\alpha}}$ and having probability laws P_{α} , $\alpha = 1, ..., \nu$ and let $\Phi : \mathbb{R}^{m_1} \times \cdots \times \mathbb{R}^{m_{\nu}} \to \mathbb{C}$ be a Borelian function. Set

$$\Phi_{\alpha}(\xi_1, \dots, \xi_{\alpha}) = \int \Phi(\xi_1, \dots, \xi_{\alpha}, \xi_{\alpha+1}, \dots, \xi_{\nu}) P_{\alpha+1}(d\xi_{\alpha+1}) \dots P_{\nu}(d\xi_{\nu})$$
 (2.4)

so that $\Phi_{\nu} = \Phi$, $\Phi_0 = \mathbf{E}\{\Phi\}$, where $\mathbf{E}\{\dots\}$ denotes the expectation with respect to the product measure $P_1 \dots P_{\nu}$.

Then for any positive $p \geq 1$ there exists C'_p , independent of ν and such that

$$\mathbf{E}\{|\Phi - \mathbf{E}\{\Phi\}|^{2p}\} \le C_p' \nu^{p-1} \sum_{\alpha=1}^{\nu} \mathbf{E}\{|\Phi_{\alpha} - \Phi_{\alpha-1}|^{2p}\}.$$
(2.5)

Let $\Phi = \gamma_n(z)$, $\xi_\alpha = \{a_{\alpha j}\}_{j \leq \alpha}$. Denote also $E_\alpha\{.\}$ the averaging with respect to the random variables $\{a_{\alpha j}\}_{j=1}^n$. Then it is easy to see that

$$\Phi_{\alpha} = E_{\alpha+1} \dots E_n \Phi$$

and by the Hölder inequality

$$\mathbf{E}\{|\Phi_{\alpha} - \Phi_{\alpha-1}|^{2p}\} = \mathbf{E}\{|E_{\alpha+1} \dots E_n(\Phi - E_{\alpha}\{\Phi)\}|^{2p}\} \le \mathbf{E}\{|\Phi - \mathbf{E}_{\alpha}\{\Phi\}|^{2p}\} = \mathbf{E}\{|\Phi - \mathbf{E}_1\{\Phi\}|^{2p}\}$$
(2.6)

Define $\mathcal{A}^{(1)}$ as a $(n-1) \times (n-1)$ matrix which can be obtained from \mathcal{A} if we remove from \mathcal{A} the first line and the first column. Set also

$$G^{(1)}(z) = (\mathcal{A}^{(1)} - z)^{-1}, \quad \gamma_n^{(1)}(z) = \sum_{i=2}^n G_{ii}^{(1)}(z), \quad a^{(1)} = (a_{12}, \dots, a_{1n}). \tag{2.7}$$

We use the representations:

$$G_{11}(z) = -(z + (G^{(1)}a^{(1)}, a^{(1)}))^{-1},$$

$$G_{ii}(z) = G_{ii}^{(1)}(z) - \frac{(G^{(1)}a^{(1)})_i(G^{(1)}a^{(1)})_i}{z + (G^{(1)}a^{(1)}, a^{(1)})}, \quad i \neq 1.$$
(2.8)

Since $G^{(1)}$ does not depend on $a^{(1)}$, we have

$$\gamma_{n} - \mathbf{E}_{1}\{\gamma_{n}\} = -\frac{1 + (G^{(1)}G^{(1)}a^{(1)}, a^{(1)})}{z + (G^{(1)}a^{(1)}, a^{(1)})} + \mathbf{E}_{1}\left\{\frac{1 + (G^{(1)}G^{(1)}a^{(1)}, a^{(1)})}{z + (G^{(1)}a^{(1)}, a^{(1)})}\right\}
:= -\frac{1 + B(z)}{A(z)} + \mathbf{E}_{1}\left\{\frac{1 + B(z)}{A(z)}\right\}.$$
(2.9)

Hence, it suffices to estimate $\mathbf{E}\{|B/A - \mathbf{E}_1\{B/A\}|^2\}$ and $\mathbf{E}\{|A^{-1} - \mathbf{E}_1\{A^{-1}\}|^2\}$. We show how to estimate the first expression. The second one can be estimated similarly. Denote by $\xi_1^{\circ} = \xi - \mathbf{E}_1\{\xi\}$ for any random variable ξ . Note that since for any a

$$\mathbf{E}_1\{|\xi - a|^2\} = \mathbf{E}_1\{|\xi_1^{\circ}|^2\} + |a - \mathbf{E}_1\{\xi\}|^2 \Rightarrow \mathbf{E}_1\{|\xi_1^{\circ}|^2\} \le \mathbf{E}_1\{|\xi - a|^2\}$$
 (2.10)

it suffices to estimate $\mathbf{E}\{|B/A - \mathbf{E}_1\{B\}/\mathbf{E}_1\{A\}|^2\}$ instead $\mathbf{E}\{|B/A - \mathbf{E}_1\{B/A\}|^2\}$. Then it is easy to see that

$$\left| \frac{B}{A} - \frac{\mathbf{E}_1 \{B\}}{\mathbf{E}_1 \{A\}} \right| = \left| \frac{B_1^{\circ}}{\mathbf{E}_1 \{A\}} - \frac{A_1^{\circ}}{\mathbf{E}_1 \{A\}} \frac{B}{A} \right| \le \left| \frac{B_1^{\circ}}{\mathbf{E}_1 \{A\}} \right| + \left| \frac{A_1^{\circ}}{\Im z \mathbf{E}_1 \{A\}} \right|. \tag{2.11}$$

Here we used the relations that follow from the spectral theorem

$$\Im(G^{(1)}a^{(1)}, a^{(1)}) = \Im z(G^{(1)}a^{(1)}, G^{(1)}a^{(1)}), \quad \Im \operatorname{Tr} G^{(1)} = \Im z \operatorname{Tr} (G^{(1)}G^{(1)*})$$

$$\Rightarrow \frac{(G^{(1)}a^{(1)}, G^{(1)}a^{(1)})}{|z + (G^{(1)}a^{(1)}, a^{(1)})|} \le |\Im z|^{-1}, \quad \frac{n^{-1}\operatorname{Tr} (G^{(1)}G^{(1)*})}{|z + n^{-1}\operatorname{Tr} G^{(1)}|} \le |\Im z|^{-1}. \tag{2.12}$$

The first relation yields, in particular, that $|B/A| \leq |\Im z|^{-1}$. It is evident that

$$A_1^{\circ} = \sum_{i \neq j} G_{ij}^{(1)} a_{1i} a_{1j} + \sum_i G_{ii}^{(1)} (a_{1i}^2)^{\circ}, \tag{2.13}$$

$$\mathbf{E}_1\{|A_1^{\circ}|^2\} \le C(p_n n)^{-1} \mathrm{Tr}\,(G^{(1)} G^{(1)*}).$$

In view of (2.12) and (1.7) we have

$$\frac{n^{-1}\operatorname{Tr}(G^{(1)}G^{(1)*})}{|z+n^{-1}\operatorname{Tr}G^{(1)}|} \le \frac{|\Im z|^{-1+\varepsilon}|n^{-1}\operatorname{Tr}(G^{(1)}G^{(1)*})|^{\varepsilon}}{|z+n^{-1}\operatorname{Tr}G^{(1)}|^{\varepsilon}} \le C\frac{|\Im z|^{-1-\varepsilon}}{|\mathbf{E}_{1}\{A\}|^{\varepsilon}}.$$
(2.14)

Here in the first inequality the numerator N and the denominator D are just written as $N = N^{\varepsilon}N^{1-\varepsilon}$, $D = D^{\varepsilon}D^{1-\varepsilon}$, then for $(N/D)^{1-\varepsilon}$ the second inequality of (2.12) is used, and then for N^{ε} the inequality (1.7) is used. Thus, in view of the second line of (2.13)

$$\mathbf{E}_1 \left\{ \left| \frac{A_1^{\circ}}{\mathbf{E}_1 \{A\}} \right|^2 \right\} \le C \frac{(p_n n)^{-1} \mathrm{Tr} \left(G^{(1)} G^{(1)*} \right)}{|z + n^{-1} \mathrm{Tr} G^{(1)}|^2} \le C \frac{|\Im z|^{-1 - \varepsilon}}{p_n |\mathbf{E}_1 \{A\}|^{1 + \varepsilon}}.$$

Similarly

$$\mathbf{E}_1 \left\{ \left| \frac{B_1^{\circ}}{\mathbf{E}_1 \{B\}} \right|^2 \right\} \leq C \frac{\operatorname{Tr} \left(G^{(1)} G^{(1)} G^{(1)*} G^{(1)*} \right)}{p_n n |z + n^{-1} \operatorname{Tr} G^{(1)}|^2} \leq C \frac{p_n^{-1} n^{-1} \operatorname{Tr} \left(G^{(1)} G^{(1)*} \right)}{|\Im z|^2 |z + n^{-1} \operatorname{Tr} G^{(1)}|^2} \leq C \frac{p_n^{-1} |\Im z|^{-3 - \varepsilon}}{|\mathbf{E}_1 \{A\}|^{1 + \varepsilon}},$$

because, using the averaging with respect to $\{a_{1j}\}$, we obtain for $E_1\{|B^{\circ}|^2\}$ the same bound as in the second line of (2.13), but with $G^{(1)}$ replaced by $(G^{(1)})^2$. This gives the first inequality above. Then we use that $\operatorname{Tr}(G^{(1)}G^{(1)}G^{(1)*}G^{(1)*}) \leq |\Im z|^{-2}\operatorname{Tr}(G^{(1)}G^{(1)*})$ (since $|G^{(1)}|^2 \leq |\Im z|^{-2}$ and finally use (2.14).

Then, the Jensen inequality $|\mathbf{E}_1\{A\}|^{-1} \leq \mathbf{E}_1\{|A|^{-1}\}$, and the relation $A^{-1} = -G_{11}(z)$ yield

$$\mathbf{E}\{|(\gamma_n(z))_1^{\circ}|^2\} \le \frac{C\mathbf{E}\{|G_{11}(z)|^{1+\varepsilon}\}}{p_n|\Im z|^{3+\varepsilon}}.$$

Then (2.5) for p = 1 implies (2.2). Putting here $\varepsilon = 0$ we get (2.1).

To prove the second inequality of (2.1), we use (2.5) for p=2. In view of (2.9) it is enough to check that

$$\mathbf{E}_1\{|A_1^{\circ}|^4\} \le Cp_n^{-2}|\Im z|^{-4}, \quad \mathbf{E}_1\{|B_1^{\circ}|^4\} \le Cp_n^{-2}|\Im z|^{-8}. \tag{2.15}$$

The first relation here evidently follow from (2.13), if we take the fourth power and average with respect to $\{a_{1i}\}$. The second one can be obtained similarly.

To prove (2.3) we note first that (2.5) and (2.6) for $\Phi = n^{-1} \sum F(G_{ij})$ yield

$$\mathbf{Var}\{\Phi\} \le n\mathbf{E}\Big\{ \Big| \Phi - \mathbf{E}_1\{\Phi\} \Big|^2 \Big\} \le n^{-1}\mathbf{E}\Big\{ \Big| \sum_{j} \left(F(G_{jj}) - F(G_{jj}^{(1)}) \right) \Big|^2 \Big\}$$

$$\le n^{-1} \sup_{\zeta:0<\Im\zeta, |\zeta|$$

if we take into account that $n^{-1} \sum F(G_{jj}^{(1)})$ does not depend on a_{1j} and hence may play the role of a in the inequality (2.10). Moreover, using (2.8) and (2.12), we get

$$\sum_{j} \left| G_{jj} - G_{jj}^{(1)} \right| \le \frac{1 + (G^{(1)}a^{(1)}, G^{(1)}a^{(1)})}{|z + G^{(1)}a^{(1)}, a^{(1)})|} \le |\Im z|^{-1}.$$

The above two bounds prove (2.3). \square

Lemma 1 gives the bound for the variance of the linear eigenvalue statistics for the functions $\varphi(\lambda) = (\lambda - z)^{-1}$. Now we are going to extend the bound to a wider class of test functions. For this aim we use Proposition below. We formulate it for the variance of linear eigenvalue statistics, but one can see easily that it can be applied also to a more general case even without reference to random matrix, see e.g. [14]. Proposition was proven in [18], but for the completeness we give its proof here. We also would like to thank Prof. A.Soshnikov for the fruitful discussion on the proposition, which allows us to make proof the proof more simple.

Proposition 2 Let A be any random $n \times n$ matrix, $\mathcal{N}_n[\varphi]$ be its linear eigenvalue statistic (1.3), and $\gamma_n(z)$ be defined by (1.6). Then

$$\mathbf{Var}\{\mathcal{N}_n[\varphi]\} \le C_s ||\varphi||_s^2 \int_0^\infty dy e^{-y} y^{2s-1} \int_{-\infty}^\infty \mathbf{Var}\{\gamma_n(x+iy)\} dx$$
 (2.16)

where $||\varphi||_s$ is defined in (1.9).

Remark 1 If the integral in the r.h.s. is equal to infinity, then the inequality is not interesting, hence we will assume that this integral is finite.

Proof. Consider the operators \mathcal{D}_s , \mathcal{V} defined in the space of the Fourier transforms of the functions of the standard $L_2(\mathbb{R})$:

$$\widehat{\mathcal{D}_s f}(k) = (1 + 2|k|)^s \widehat{f}(k),$$

$$\widehat{\mathcal{V}}f(k) = \int dk' \widehat{\mathcal{V}}(k, k') \widehat{f}(k'), \quad \widehat{\mathcal{V}}(k_1, k_2) = \mathbf{Cov}\{\mathrm{Tr}e^{ik_1\mathcal{A}}, \mathrm{Tr}e^{ik_2\mathcal{A}}\}.$$
(2.17)

It is easy to see that if we introduce the operator $K := \mathcal{D}_s^{-1} \mathcal{V} \mathcal{D}_s^{-1}$ then

$$\mathbf{Var}\{\mathcal{N}_n[\varphi]\} = (2\pi)^{-2} (\mathcal{V}\varphi, \varphi) = (2\pi)^{-2} (K\mathcal{D}_s\varphi, \mathcal{D}_s\varphi)$$

$$\leq (2\pi)^{-2} ||K|| \cdot ||\mathcal{D}_s\varphi||^2 \leq (2\pi)^{-2} ||\varphi||_s^2 \operatorname{Tr} K$$
(2.18)

Let us check that the operator K is indeed of the trace class in $L_2(\mathbb{R})$. Note first that in the Fourier space his kernel has the form

$$\widehat{K}(k_1, k_2) = (1 + 2|k_1|)^{-s} \widehat{\mathcal{V}}(k_1, k_2) (1 + 2|k_1|)^{-s},$$

with $\widehat{\mathcal{V}}(k_1, k_2)$ of (2.17). It is evident that $K \geq 0$ in the operator sense, and $K(k_1, k_2)$ is a continuous function of k_1, k_2 , since $\widehat{\mathcal{V}}(k_1, k_2)$ can be written as a finite sum of the products of Fourier transforms of the positive unit measures, which are the distributions of eigenvalues of \mathcal{A} . Moreover, we will prove below that

$$\int \widehat{K}(k,k)dk < \infty. \tag{2.19}$$

Then it follows from the inequality $|\widehat{K}(k_1, k_2)|^2 \leq \widehat{K}(k_1, k_1)\widehat{K}(k_2, k_2)$ (which is valid for any continuous positive definite kernels) that \widehat{K} belongs to the Hilbert-Schmidt class, and therefore \widehat{K} has a basis $\{\phi_j(k)\}_{j=1}^{\infty}$, which is made from the continuous eigenfunctions with corresponding eigenvalues $\lambda_j > 0$. Then if we consider a finite rank operator with the kernel $\widehat{K}_N(k_1, k_2) = \sum_{j=1}^N \lambda_j \phi_j(k_1) \overline{\phi}_j(k_2)$, we have $K - K_N \geq 0$ in the operator sense, and $\widehat{K}(k_1, k_2) - \widehat{K}_N(k_1, k_2)$ is continuous, thus $\widehat{K}(k, k) \geq \widehat{K}_N(k, k)$ and

$$\sum_{j=1}^{N} \lambda_j = \int \widehat{K}_N(k,k) dk \le \int \widehat{K}(k,k) dk.$$

Since N is arbitrary and $\lambda_i > 0$ we obtain that K is a trace class operator.

We are left to prove the inequality of (2.19). We have

$$\begin{split} \int \widehat{K}(k,k)dk &= \int (1+2|k|)^{-2s}\widehat{\mathcal{V}}(k,k)dk \\ &= \frac{1}{\Gamma(2s)} \int_0^\infty dy e^{-y} y^{2s-1} \int e^{-2|k|y} \widehat{\mathcal{V}}(k,k)dk \\ &= \frac{1}{\Gamma(2s)} \int_0^\infty dy e^{-y} y^{2s-1} \int dx \int \int dk_1 dk_2 e^{i(k_1-k_2)x} \widehat{\mathcal{V}}(k_1,k_2) e^{-|k_1|y-|k_2|y} \\ &= \frac{1}{\Gamma(2s)} \int_0^\infty dy e^{-y} y^{2s-1} \int dx \mathbf{Var} \{ \mathcal{N}_n[P_y(x-.)] \} \\ &= \frac{1}{\Gamma(2s)} \int_0^\infty dy e^{-y} y^{2s-1} \int dx \mathbf{Var} \{ \Im \gamma_n(x+iy) \}, \end{split}$$

where P_y is the Poisson kernel

$$P_y(x) = \frac{y}{\pi(x^2 + y^2)}. (2.20)$$

and we used that

$$\int P_y(x-\lambda)e^{ik\lambda}d\lambda = e^{ikx-|k|y}.$$

This relation combined with (2.18) proves (2.16).

Now we are ready to prove the bound for the variance of linear eigenvalue statistics for a rather wide class of the test functions

Lemma 2 If $||\varphi||_{3/2+\alpha} \leq \infty$, with any $\alpha > 0$, then

$$\frac{p_n}{n} \operatorname{Var} \{ \mathcal{N}_n[\varphi] \} \le C_\alpha ||\varphi||_{3/2 + \alpha}^2$$
(2.21)

Proof. In view of Proposition 2 we need to estimate

$$I(y) = \int_{-\infty}^{\infty} \mathbf{Var} \{ \gamma_n(x+iy) \} dx$$

Take in (2.2) $\varepsilon = \alpha/2$. Then we need to estimate

$$\int_{-\infty}^{\infty} \mathbf{E}\{|G_{11}(x+iy)|^{1+\alpha/2}\}dx.$$

Use the spectral representation

$$G_{11} = \int \frac{N_{11}(d\lambda)}{\lambda - x - iy}, \text{ were } N_{11}(\Delta) = \sum_{k=1}^{n} |\psi_1^{(k)}|^2 \mathbf{1}_{\Delta}(\lambda_k)$$

with $\psi^{(k)} = (\psi_1^{(k)}, \dots, \psi_n^{(k)})$ being an eigenvector of \mathcal{A} , corresponding the eigenvalue λ_k , i.e. $\mathcal{A}\psi^{(k)} = \lambda_k \psi^{(k)}$. Then the Jensen inequality with respect to $N_{11}(d\lambda)$ yields

$$\int_{-\infty}^{\infty} |G|_{11}^{1+\alpha/2}(x+iy)dx \le \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} \frac{N_{11}(d\lambda)}{(|x-\lambda|^2+y^2)^{(1+\alpha/2)/2}} \le C|y|^{-\alpha/2}.$$

Taking $s = 3/2 + \alpha$ in (2.16) we get

$$\frac{p_n}{n} \mathbf{Var} \{ \mathcal{N}_n[\varphi] \} \le ||\varphi||_{3/2 + \alpha}^2 C \int_0^\infty e^{-y} y^{2 + 2\alpha} y^{-3 - \alpha} dy \le C ||\varphi||_{3/2 + \alpha}^2.$$

The next lemma is technical one. We accumulate relations which we need to prove CLT.

Lemma 3 Using notations of (2.9) we have uniformly in $z_1, z_2 : \Im z_{1,2} > a$ with any a > 0:

$$\mathbf{E}_1\{|A^{\circ}|^4\} = O(p_n^{-2}), \quad \mathbf{E}_1\{|B^{\circ}|^4\} = O(p_n^{-2}),$$
 (2.22)

$$\left(\mathbf{E}_{1}\{A^{-1}\}\right)^{\circ} = -\left(1 + O(p_{n}^{-1}) + O(p_{n}/n)\right) \frac{n^{-1}(\gamma_{n}^{(1)})^{\circ}}{\mathbf{E}^{2}\{A\}} + r,\tag{2.23}$$

with $E\{|r^{\circ}|^2\} \le C/n^2 + C/p_n^2 n$,

$$p_n \mathbf{E}_1 \{ A^{\circ}(z_1) A^{\circ}(z_2) \} = \frac{1}{n} \sum_{i} G_{ii}^{(1)}(z_1) G_{ii}^{(1)}(z_2) + p_n \mathring{\gamma}_n^{(1)}(z_1) \mathring{\gamma}_n^{(1)}(z_2) / n^2, \tag{2.24}$$

$$p_n \mathbf{E}_1 \{ A^{\circ}(z_1) B^{\circ}(z_2) \} = p_n \frac{d}{dz_2} \mathbf{E}_1 \{ A^{\circ}(z_1) A^{\circ}(z_2) \}, \tag{2.25}$$

$$\mathbf{Var}\{p_n\mathbf{E}_1\{A^{\circ}(z_1)A^{\circ}(z_2)\}\} = O(n^{-1}), \quad \mathbf{Var}\{p_n\mathbf{E}_1\{A^{\circ}(z_1)B^{\circ}(z_2)\}\} = O(n^{-1}), \quad (2.26)$$

$$\mathbf{E}\{|\overset{\circ}{\gamma}_{n}^{(1)}(z) - \overset{\circ}{\gamma}_{n}(z)|^{2}\} = O(p_{n}^{-1}). \tag{2.27}$$

Moreover,

$$\mathbf{Var}\{G_{ii}^{(1)}(z_1)\} = O(p_n^{-1}), \quad |\mathbf{E}\{G_{ii}^{(1)}(z_1)\} - \mathbf{E}\{G_{ii}(z_1)\}| = O(p_n^{-1}), \tag{2.28}$$

$$|\mathbf{E}\{\gamma_n^{(1)}(z)\}/n - f(z)| = O(n^{-1}), \quad |\mathbf{E}^{-1}\{A(z)\} + f(z)| = O(p_n^{-1}),$$
 (2.29)

where

$$f(z) = \frac{1}{2}(\sqrt{z^2 - 4} - z). \tag{2.30}$$

Proof. Note that since $\Im z \Im (G^{(1)}m,m) \geq 0$, we can use the bound

$$|\Im A| \ge |\Im z| \Rightarrow |A^{-1}| \le |\Im z|^{-1} \le a^{-1}.$$
 (2.31)

Relations (2.22), (2.24), and (2.25) follow from the representation

$$A^{\circ} = \sum_{i \neq j} G_{ij}^{(1)} a_{1i} a_{1j} + \sum_{i} G_{ii}^{(1)} (a_{1i}^{2})^{\circ} + n^{-1} \mathring{\gamma}_{n}^{(1)}(z) = A_{1}^{\circ} + n^{-1} \mathring{\gamma}_{n}^{(1)}(z), \qquad (2.32)$$

$$B^{\circ} = \sum_{i \neq j} (G^{(1)} G^{(1)})_{ij} a_{1i} a_{1j} + \sum_{i} (G^{(1)} G^{(1)})_{ii} (a_{1i}^{2})^{\circ} + n^{-1} \frac{d}{dz} \mathring{\gamma}_{n}^{(1)}(z)$$

$$= B_{1}^{\circ} + n^{-1} \frac{d}{dz} \mathring{\gamma}_{n}^{(1)}(z),$$

and Lemma 1 (see (2.1) and (2.15)).

The first bound of (2.26) follows from (2.24) and (2.3) for $F(z) = z^2$. The second bound of (2.26) follows from (2.25) and the first relation of (2.26), if we use the fact that the variance of the derivative of an analytic function by the Cauchy theorem can be bounded by the variance of the initial function.

Relations (2.27) follow from the representation (see (2.8))

$$\mathring{\gamma}_{n}^{(1)}(z) - \mathring{\gamma}_{n}(z) = (A^{-1})^{\circ} + (BA^{-1})^{\circ}$$

and (2.34). The first relation of (2.28) is the analog of the relation

$$\mathbf{Var}\{G_{ii}(z_1)\} = \mathbf{Var}\{G_{11}(z_1)\} = O(p_n^{-1})$$
(2.33)

if in the latter we replace the matrix \mathcal{A} by $\mathcal{A}^{(1)}$. But since $G_{11}(z_1) = -A^{-1}(z_1)$, (2.33) follows from (2.22) and (2.31). The second relation of (2.28) follows from the symmetry of the problem and (2.9)

$$\mathbf{E}\{G_{ii}^{(1)}(z_1)\} - \mathbf{E}\{G_{ii}(z_1)\} = \frac{1}{n-1}\mathbf{E}\{\gamma_n - \gamma_n^{(1)} - G_{11}\} = \frac{1}{n-1}\mathbf{E}\{B/A\} = O(n^{-1}).$$

The first relation of (2.29) follows from the above bound for $n^{-1}\mathbf{E}\{\gamma_n - \gamma_n^{(1)}\}$ and the estimate (see [7])

$$n^{-1}\mathbf{E}\{\gamma_n\} - f(z) = O(p_n^{-1}).$$

The second relation of (2.29) is the corollary of the above estimate and the representation

$$A^{-1} = \mathbf{E}^{-1}\{A\} - A^{\circ}\mathbf{E}^{-2}\{A\} + (A^{\circ})^{2}A^{-1}\mathbf{E}^{-2}\{A\}, \tag{2.34}$$

which implies

$$\mathbf{E}\{A(z)\}^{-1} = \mathbf{E}\{A(z)^{-1}\} + O(\mathbf{Var}\{A(z)\}) = -\mathbf{E}\{G_{11}\} + O(p_n^{-1}) = -n^{-1}\mathbf{E}\{\gamma_n\} + O(p_n^{-1}).$$

We are left to prove (2.23). Set

$$\widetilde{A} = z + \sum_{i} G_{ii}^{(1)} a_{1i}^2. \tag{2.35}$$

Using the analog of (2.34) for A and \widetilde{A} , we write first

$$A^{-1} = \widetilde{A}^{-1} - \widetilde{A}^{-2}(A - \widetilde{A}) + r_1, \quad r_1 = \widetilde{A}^{-2}A^{-1}(A - \widetilde{A})^2.$$

We have

$$A - \widetilde{A} = \sum_{i \neq j} G_{ij}^{(1)} a_{1i} a_{1j}, \quad \mathbf{E}\{|r_1|^2\} \le \frac{1}{|\Im z|^6} \mathbf{E}\{|A - \widetilde{A}|^4\} \le C/n^2 + C/np_n^2.$$

Moreover, the analog of (2.34) for \widetilde{A} yields

$$\begin{aligned} \mathbf{E}_{1}\{\widetilde{A}^{-2}(A-\widetilde{A})\} &= \mathbf{E}_{1}^{-2}\{A\}\mathbf{E}_{1}\{(A-\widetilde{A})\} \\ &- 2\mathbf{E}_{1}^{-3}\{A\}\mathbf{E}_{1}\{(A-\widetilde{A})(\widetilde{A}-\mathbf{E}_{1}\{A\})\} + r_{2} = r_{2} \\ &r_{2} &= \mathbf{E}_{1}\Big\{\mathbf{E}_{1}^{-2}\{A\}\Big(\widetilde{A}^{-2} + 2\mathbf{E}_{1}^{-1}\{A\}\widetilde{A}^{-1}\Big)\Big(A-\widetilde{A}\Big)\Big(\widetilde{A}-\mathbf{E}_{1}\{A\}\Big)^{2}\Big\}. \end{aligned}$$

Since $\widetilde{A} - \mathbf{E}_1\{A\} = \sum_{i} G_{ii}^{(1)}(a_{1i}^2 - \mathbf{E}_1\{a_{1i}^2\})$ we have

$$\mathbf{E}\{r_2^2\} \le C\mathbf{E}\{|A - \widetilde{A}|^2 |\widetilde{A} - \mathbf{E}_1\{A\}|^4\} \le C/np_n^2.$$

Hence we have proved that

$$\mathbf{E}_1\{A^{-1}\} = \mathbf{E}_1\{\widetilde{A}^{-1}\} + \widetilde{r}, \quad E\{|\widetilde{r}|^2\} \le C/n^2 + C/np_n^2. \tag{2.36}$$

Using (1.4) we can write

$$\begin{split} i\mathbf{E}_{1}\{\widetilde{A}^{-1}\} &= \int_{0}^{\infty} dt e^{izt} \prod \mathbf{E}_{1}\{e^{itG_{jj}^{(1)}a_{1j}^{2}}\} \\ &= \int_{0}^{\infty} dt e^{izt} e^{itp_{n}\gamma_{n}^{(1)}/n^{2}} \prod \left(1 + \frac{p_{n}}{n} \left(e^{itG_{ii}^{(1)}(1/p_{n}-2/n)} - 1\right)\right) \\ &= \int_{0}^{\infty} dt e^{izt} e^{itp_{n}\gamma_{n}^{(1)}/n^{2}} \exp\left\{\frac{p_{n}}{n} \sum_{i} \left(e^{itG_{ii}^{(1)}(1/p_{n}-2/n)} - 1\right)\right\} + O(n^{-1}) \\ &= \int_{0}^{\infty} dt e^{izt} e^{it\gamma_{n}^{(1)}(1-p_{n}/n)/n} \exp\left\{in^{-1} \sum_{i} F(G_{ii}^{(1)}, t)\right\} + O(n^{-1}), \end{split}$$

where

$$F(x,t) = p_n \Big(e^{itx(1/p_n - 2/n)} - 1 - itx(1/p_n - 2/n) \Big),$$

Then in view of (2.3), since $\sup_{\Im x>0}|F'_x(x,t)|\leq C|t|p_n^{-1},$ we obtain

$$i\mathbf{E}_{1}\{\widetilde{A}^{-1}\} = \int_{0}^{\infty} dt e^{izt} e^{it\gamma_{n}^{(1)}(1-p_{n}/n)/n} \exp\left\{in^{-1}\sum \mathbf{E}\{F(G_{ii}^{(1)},t)\}\right\}$$

$$\cdot \left(1 + O\left(n^{-1}\sum F^{\circ}(G_{ii}^{(1)},t)\right)\right) + O(n^{-1})$$

$$= \int_{0}^{\infty} dt e^{izt} e^{it\gamma_{n}^{(1)}(1-p_{n}/n)/n} \exp\left\{in^{-1}\sum \mathbf{E}\{F(G_{ii}^{(1)},t)\}\right\} + r',$$

$$\mathbf{E}\{|r'|^{2}\} \leq C/n^{2} + C/np_{n}^{2}.$$

Finally, replacing similarly to the above $\gamma_n^{(1)}$ by $\mathbf{E}\{\gamma_n^{(1)}\}$ in the exponent, we get

$$i\mathbf{E}_{1}\{\widetilde{A}^{-1}\} = \int_{0}^{\infty} dt e^{izt} e^{it\mathbf{E}\{\gamma_{n}^{(1)}\}(1-p_{n}/n)/n} \exp\left\{in^{-1}\sum_{n} \mathbf{E}\{F(G_{ii}^{(1)},t)\}\right\}$$

$$\left(1 + itn^{-1}(\gamma_{n}^{(1)})^{\circ}(1-p_{n}/n) + O((n^{-1}(\gamma_{n}^{(1)})^{\circ})^{2}) + r'.$$

Taking $(\mathbf{E}_1\{\widetilde{A}^{-1}\})^{\circ}$, we can see that the term which corresponds to 1 in the r.h.s. disappears, and since $\mathbf{E}\{F(G_{ii},t)\}=O(p_n^{-1})$, the coefficient in front of $(\gamma_n^{(1)})^{\circ}$ equals

$$\int_0^\infty ite^{izt}e^{it\mathbf{E}\{\gamma_n^{(1)}\}(1-p_n/n)/n}\exp\Big\{in^{-1}\sum\mathbf{E}\{F(G_{ii}^{(1)},t)\}\Big\}dt$$

$$=i(z+\mathbf{E}\{\gamma_n^{(1)}\}/n)^{-2}(1+O(p_n^{-1})+O(p_n/n))$$

$$=-i\mathbf{E}^{-2}\{A\}\left(1+O(p_n^{-1})+O(p_n/n)\right).$$

In view of (2.1) and (2.36) we obtain (2.23).

Proof of Theorem 1. We prove first Theorem 1 for the function φ_{η} of the form

$$\varphi_{\eta} = P_{\eta} * \varphi_{0}, \quad \int |\varphi_{0}(\lambda)| d\lambda \le C < \infty,$$
(2.37)

where P_{η} is the Poisson kernel (see (2.20)) and φ_0 is a real valued function from $L_1(\mathbb{R})$. One can see easily that

$$\mathcal{N}_{n}^{\circ}[\varphi_{\eta}] = \left(\operatorname{Tr}\varphi_{\eta}(\mathcal{A})\right)^{\circ} = \frac{1}{\pi} \int \varphi_{0}(\mu) \Im\left(\operatorname{Tr}G(\mu + i\eta)\right)^{\circ} d\mu$$
$$= \frac{1}{2\pi i} \int \varphi_{0}(\mu) (\gamma_{n}^{\circ}(z_{\mu}) - \gamma_{n}^{\circ}(\overline{z}_{\mu})) d\mu, \quad z_{\mu} = \mu + i\eta. \tag{2.38}$$

Set

$$d_n = \left(\frac{n}{p_n}\right)^{1/2}, \quad Z_n(x) = \mathbf{E}\left\{e^{ix\mathcal{N}_n^{\circ}[\varphi_{\eta}]/d_n}\right\},$$

$$e(x) = e^{ix\mathcal{N}_n^{\circ}[\varphi_{\eta}]/d_n}, \quad Y_n(z, x) = d_n^{-1}\mathbf{E}\left\{\operatorname{Tr}G(z)e^{\circ}(x)\right\}.$$
(2.39)

Then it is easy to see that

$$\frac{d}{dx}Z_n(x) = \frac{1}{2\pi} \int \varphi_0(\mu) (Y_n(z_\mu, x) - Y_n(\overline{z}_\mu, x)) d\mu. \tag{2.40}$$

On the other hand, using the symmetry of the problem and the notations of (2.9), we have

$$Y_n(z,x) = d_n^{-1} \mathbf{E} \{ \operatorname{Tr} G(z) e^{\circ}(x) \} = n d_n^{-1} \mathbf{E} \{ G_{11}(z) e^{\circ}(x) \}$$

$$= -n d_n^{-1} \mathbf{E} \{ (A^{-1})^{\circ} e_1(x) \} - n d_n^{-1} \mathbf{E} \{ (A^{-1})^{\circ} (e(x) - e_1(x)) \} = T_1 + T_2,$$
(2.41)

where

$$e_1(x) = e^{ix(\mathcal{N}_{n-1}^{(1)}[\varphi_{\eta}])^{\circ}/d_n}, \quad (\mathcal{N}_{n-1}^{(1)}[\varphi_{\eta}])^{\circ} = (\operatorname{Tr}\varphi_{\eta}(\mathcal{A}^{(1)}))^{\circ} = \frac{1}{\pi} \int d\mu \, \varphi_0(\mu) \Im_{\eta}^{\circ(1)}(z_{\mu}).$$

Since $e_1(x)$ does not depend on $\{a_{1i}\}$, using that $\mathbf{E}\{...\} = \mathbf{E}\{\mathbf{E}_1\{...\}\}$, we obtain in view of the above representation and (2.23)

$$T_1 = d_n^{-1} \mathbf{E} \{ (\gamma_n^{(1)}(z))^{\circ} e_1(x) \} / \mathbf{E}^2 \{ A \} \left(1 + O(d_n^{-1}) + O(p_n^{-1}) \right) + O(d_n^{-1}).$$

Write

$$e(x) - e_1(x) = \frac{ix}{\pi d_n} \int \varphi_0(\mu) \left(\Im\left(\gamma_n^{\circ} - \gamma_n^{\circ}\right)^{(1)}\right) e_1(x) + O\left(\left(\gamma_n^{\circ} - \gamma_n^{\circ}\right)^{(1)}\right)^2\right) ix/d_n d\mu. \tag{2.42}$$

Then (2.27), the relations $|e(x)| = |e_1(x)| = 1$, and (2.1) yield

$$d_n^{-1}|\mathbf{E}\{(\gamma_n^{(1)})^{\circ}e_1(x)\} - \mathbf{E}\{\gamma_n^{\circ}e(x)\}| \le Cd_n^{-1}\mathbf{E}\{|(\gamma_n^{(1)})^{\circ} - \gamma_n^{\circ}|(1 + |x||\gamma_n^{\circ}|d_n^{-1})\}$$

$$\le Cd_n^{-1}\mathbf{E}^{1/2}\{|(\gamma_n^{(1)})^{\circ} - \gamma_n^{\circ}|^2\}\left(1 + |x|d_n^{-1}\mathbf{E}^{1/2}\{|\gamma_n^{\circ}|^2\}\right) = O(d_n^{-1}p_n^{-1/2}).$$

Hence we obtain

$$T_1 = Y_n(z, x) / \mathbf{E}^2 \{A\} \left(1 + O(d_n^{-1}) + O(p_n^{-1}) \right) + O(d_n^{-1} p_n^{-1/2}). \tag{2.43}$$

To compute T_2 we use (2.42). Then, taking into account (2.27), we conclude that the term $O(nd_n^{-3}(\gamma_n^{\circ} - (\gamma_n^{(1)})^{\circ})^2)$ gives the contribution $O(d_n^{-1})$. Then, since $e_1(x)$ does not depend on $\{a_{1i}\}$, we average first with respect to $\{a_{1i}\}$ and obtain in view of (2.9)

$$T_{2} = -\frac{ixn}{d_{n}^{2}\pi} \int d\mu \varphi_{0}(\mu) \mathbf{E} \left\{ (A^{-1})^{\circ}(z) e_{1}(x) \Im \left(\gamma_{n}^{\circ}(z_{\mu}) - (\gamma_{n}^{(1)}(z_{\mu}))^{\circ} \right) \right\} + O(d_{n}^{-1})$$

$$= -\frac{ixp_{n}}{\pi} \int d\mu \varphi_{0}(\mu) \mathbf{E} \left\{ e_{1}(x) \mathbf{E}_{1} \left\{ (A^{-1})^{\circ}(z) \Im \left(\gamma_{n}^{\circ}(z_{\mu}) - (\gamma_{n}^{(1)}(z_{\mu}))^{\circ} \right) \right\} \right\} + O(d_{n}^{-1})$$

$$= \frac{ixp_{n}}{\pi} \int d\mu \varphi_{0}(\mu) \mathbf{E} \left\{ e_{1}(x) \mathbf{E}_{1} \left\{ (A^{-1})^{\circ}(z) \Im \left((1 + B(z_{\mu})) A^{-1}(z_{\mu}) \right)^{\circ} \right\} \right\} + O(d_{n}^{-1}).$$

Using (2.34) and (2.22), we conclude that only linear terms with respect to B° and A° give non vanishing contribution, hence we obtain

$$D_{n}(z, z_{\mu}) := p_{n} \mathbf{E}_{1} \Big\{ (A^{-1})^{\circ}(z) \Big((1 + B(z_{\mu})) A^{-1}(z_{\mu}) \Big)^{\circ} \Big\}$$

$$= p_{n} \mathbf{E}^{-2} \{ A(z) \} \mathbf{E}^{-2} \{ A(z_{\mu}) \} \Big(1 + \mathbf{E} \{ B(z_{\mu}) \} \Big) \mathbf{E}_{1} \{ A^{\circ}(z) A^{\circ}(z_{\mu}) \}$$

$$- p_{n} \mathbf{E}^{-2} \{ A(z) \} \mathbf{E}^{-1} \{ A(z_{\mu}) \} \mathbf{E}_{1} \{ A^{\circ}(z) B^{\circ}(z_{\mu}) \} + O(p_{n}^{-1/2})$$

$$= f^{3}(z) f^{3}(z_{\mu}) (1 + f'(z_{\mu})) + f^{3}(z) f(z_{\mu}) f'(z_{\mu}) + O(p_{n}^{-1/2}).$$

Here we used first (2.24) and (2.25) to express $\mathbf{E}_1\{A^{\circ}(z)A^{\circ}(z_{\mu})\}$ and $\mathbf{E}_1\{A^{\circ}(z)B^{\circ}(z_{\mu})\}$ in terms of $G_{ii}^{(1)}(z)$, and $\frac{d}{dz_{\mu}}G_{ii}^{(1)}(z_{\mu})$, and then 2.28) combined with (2.29) to replace $\mathbf{E}_1\{A^{\circ}(z)A^{\circ}(z_{\mu})\}$

by $f(z)f(z_{\mu})$ and $\mathbf{E}_1\{A^{\circ}(z)B^{\circ}(z_{\mu})\}$ by $f(z)f'(z_{\mu})$. Moreover, we used (2.29) to replace $\mathbf{E}^{-1}\{A(z)\}$ by -f(z). Hence

$$D_n(z, z_\mu) = \left(f^3(z)f^3(z_\mu)(1 + f'(z_\mu)) + f^3(z)f(z_\mu)f'(z_\mu)\right) + O(p_n^{-1/2}). \tag{2.44}$$

In addition, similarly to (2.42) we have

$$\mathbf{E}\{e_1(x)\} = Z_n(x) + O(d_n^{-1}).$$

Hence, relations (2.41)–(2.44) imply

$$Y_{n}(z,x) = f^{2}(z)Y_{n}(z,x) + ixZ_{n}(x) \int d\mu \varphi_{0}(\mu) \frac{D_{n}(z,z_{\mu}) - D_{n}(z,\overline{z_{\mu}})}{2i\pi} + O(p_{n}^{-1/2}) + O(d_{n}^{-1}),$$

$$Y_{n}(z,x) = ixZ_{n}(x) \int d\mu \varphi_{0}(\mu) \frac{C_{n}(z,z_{\mu}) - C_{n}(z,\overline{z_{\mu}})}{2i\pi} + O(p_{n}^{-1/2}) + O(d_{n}^{-1}), \qquad (2.45)$$

$$C_{n}(z,z_{\mu}) := \frac{D_{n}(z,z_{\mu})}{1 - f^{2}(z)}.$$

Using that

$$f(z)(f'(z)+1) = \frac{f(z)}{1-f^2(z)} = -\frac{1}{\sqrt{z^2-4}}, \quad f' = -\frac{f(z)}{\sqrt{z^2-4}},$$

we can transform $C_n(z, z_{\mu})$ to the form

$$C_n(z, z_{\mu}) = C(z, z_{\mu}) + O(p_n^{-1/2}) + O(d_n^{-1/2})$$

$$C(z, z_{\mu}) := 2 \frac{f^2(z) f^2(z_{\mu})}{(z^2 - 4)^{1/2} (z_{\mu}^2 - 4)^{1/2}}.$$
(2.46)

Taking into account (2.40), (2.45), and (2.46), we obtain the equation

$$\frac{d}{dx}Z_{n}(x) = -xV[\varphi_{\eta}]Z_{n}(x) + O(p_{n}^{-1/2}) + O(d_{n}^{-1})
V[\varphi_{\eta}] = -\frac{1}{4\pi^{2}} \int \int \varphi_{0}(\mu_{1})\varphi_{0}(\mu_{2}) \Big(C(z_{\mu_{1}}, z_{\mu_{2}}) - C(z_{\mu_{1}}, \overline{z_{\mu_{2}}}) - C(\overline{z_{\mu_{1}}}, z_{\mu_{2}})
+ C(\overline{z_{\mu_{1}}}, \overline{z_{\mu_{2}}})\Big) d\mu_{1} d\mu_{2}.$$
(2.47)

Formulas (2.46) and (2.47) imply that

$$\begin{split} V[\varphi_{\eta}] &= \frac{2}{\pi^2} \bigg(\int \varphi_0(\mu) \Im \frac{f^2(z_{\mu})}{(z_{\mu}^2 - 4)^{1/2}} d\mu \bigg)^2 = \frac{1}{2\pi^2} \bigg(\int \varphi_0(\mu) \Im \bigg(\frac{z_{\mu}^2 - 2}{(z_{\mu}^2 - 4)^{1/2}} - z_{\mu} \bigg) d\mu \bigg)^2 \\ &= \frac{1}{2\pi^2} \bigg(\int d\mu \varphi_0(\mu) \Im \bigg(\frac{1}{\pi} \int_{-2}^2 \frac{(\lambda^2 - 2) d\lambda}{(\mu + i\eta - \lambda)\sqrt{4 - \lambda^2}} \bigg) \bigg)^2 = \frac{1}{2\pi^2} \bigg(\int_{-2}^2 d\lambda \varphi_{\eta}(\lambda) \frac{(\lambda^2 - 2)}{\sqrt{4 - \lambda^2}} \bigg)^2, \end{split}$$

where we used also the well known relations

$$\frac{1}{\pi} \int_{-2}^{2} \frac{d\lambda}{(z-\lambda)\sqrt{4-\lambda^{2}}} = \frac{1}{(z^{2}-4)^{1/2}}, \quad \frac{1}{\pi} \int_{-2}^{2} \frac{d\lambda}{\sqrt{4-\lambda^{2}}} = 1.$$

Now if we consider

$$\widetilde{Z}_n(x) = e^{x^2 V[\varphi_\eta]/2} Z_n(x),$$

then (2.47) yields that for any $|x| \leq C$

$$\frac{d}{dx}\widetilde{Z}_n(x) = O(p_n^{-1/2}) + O(d_n^{-1}),$$

and since $\widetilde{Z}_n(0) = Z_n(0) = 1$, we obtain uniformly in $x \leq C$

$$\widetilde{Z}_n(x) = 1 + O(p_n^{-1/2}) + O(d_n^{-1})
\Rightarrow Z_n(x) = e^{-x^2 V[\varphi_\eta]/2} + O(p_n^{-1/2}) + O(d_n^{-1}).$$
(2.48)

Thus, we have proved CLT for the functions of the form (2.37). To extend CLT to a wider class of functions we use

Proposition 3 Let $\{\xi_l^{(n)}\}_{l=1}^n$ be a triangular array of random variables, $\mathcal{N}_n[\varphi] = \sum_{l=1}^n \varphi(\xi_l^{(n)})$ be its linear statistics, corresponding to a test function $\varphi : \mathbb{R} \to \mathbb{R}$, and

$$V_n[\varphi] = \mathbf{Var}\{d_n^{-1}\mathcal{N}_n[\varphi]\}$$

be the variance of $\mathcal{N}_n[\varphi]$, where $\{d_n\}_{n=1}^{\infty}$ is some bounded from above sequence of numbers. Assume that

(a) there exists a vector space \mathcal{L} endowed with a norm ||...|| and such that V_n is defined on \mathcal{L} and admits the bound

$$V_n[\varphi] \le C||\varphi||^2, \quad \forall \varphi \in \mathcal{L},$$
 (2.49)

where C does not depend on n;

(b) there exists a dense linear manifold $\mathcal{L}_1 \subset \mathcal{L}$ such that the Central Limit Theorem is valid for $\mathcal{N}_n[\varphi]$, $\varphi \in \mathcal{L}_1$, i.e., if $Z_n[x\varphi] = \mathbf{E}\{e^{ixd_n^{-1}\mathring{\mathcal{N}}_n[\varphi]}\}$ is the characteristic function of $d_n^{-1/2}\mathring{\mathcal{N}}_n[\varphi]$, then there exists a continuous quadratic functional $V: \mathcal{L}_1 \to \mathbb{R}_+$ such that we have uniformly in x, varying on any compact interval

$$\lim_{n \to \infty} Z_n[x\varphi] = e^{-x^2 V[\varphi]/2}, \quad \forall \varphi \in \mathcal{L}_1;$$
(2.50)

Then V admits a continuous extension to \mathcal{L} and Central Limit Theorem is valid for all $\mathcal{N}_n[\varphi]$, $\varphi \in \mathcal{L}$.

Proof. Let $\{\varphi_k\}$ be a sequence of elements of \mathcal{L}_1 converging to $\varphi \in \mathcal{L}$. We have then in view of the inequality $|e^{ia} - e^{ib}| \leq |a - b|$, the linearity of $\mathring{\mathcal{N}}_n[\varphi]$ in φ , the Schwarz inequality, and (2.49):

$$\left| Z_n(x\varphi) - Z_n(x\varphi) \right|_{\varphi = \varphi_k} \le |x| \mathbf{E} \left\{ \left| d_n^{-1} \mathring{\mathcal{N}}_n[\varphi] - d_n^{-1} \mathring{\mathcal{N}}_n[\varphi_k] \right| \right\}$$

$$\leq |x| \mathbf{Var}^{1/2} \{ d_n^{-1} \mathcal{N}_n[\varphi - \varphi_k] \} \leq C|x| \quad ||\varphi - \varphi_k||.$$
(2.51)

Now, passing first to the limit $n \to \infty$ and then $k \to \infty$, we obtain the assertion of the proposition. \square

Let us show now that hypothesis (a) and (b) of Proposition 3 are fulfilled in some vector space. Consider the space \mathcal{H}_s of all functions with the norm (1.9) and set $\mathcal{L} = \mathcal{H}_s \cap L_1(\mathbb{R})$ and

$$||\varphi|| = \int |\varphi(\lambda)|d\lambda + ||\varphi||_s = ||\varphi||_{L_1(\mathbb{R})} + ||\varphi||_s. \tag{2.52}$$

Then for s > 3/2 Lemma 2 guarantees that assumption (a) of Proposition 3 is fulfilled. Moreover, the Lebesgue theorem about the dominated convergence yields that

$$||\varphi - \varphi * P_{\eta}||_{s}^{2} \le C \int |1 - e^{-\eta |k|}|^{2} (1 + 2|k|)^{2s} |\widehat{\varphi}(k)|^{2} dk \to 0, \quad \eta \to 0.$$

Hence the set of the functions $\varphi * P_{\eta}$ is dense in \mathcal{L} with respect to the norm $||.||_s$. Thus, if we prove that the set of the functions $\varphi * P_{\eta}$ is dense in \mathcal{L} with respect to the norm $||.||_{L_1(\mathbb{R})}$, then (2.48) will imply assumption (b) of Proposition 3.

It is easy to see that the set of all functions with finite supports, possessing the norm (2.52), is dense in \mathcal{L} with respect to this norm. Hence we need only to prove that if $\varphi \in \mathcal{H}_s$ and has a finite support [-A, A], then

$$\int |\varphi(\lambda) - \varphi * P_{\eta}(\lambda)| d\lambda \to 0, \quad \eta \to 0.$$

But

$$\int |\varphi(\lambda) - \varphi * P_{\eta}(\lambda)| d\lambda = \left(\int_{|\lambda| \le A+1} + \int_{|\lambda| \ge A+1} \right) |\varphi(\lambda) - \varphi * P_{\eta}(\lambda)| d\lambda = I_1 + I_2.$$

We have for I_2

$$I_2 \le \frac{\eta}{\pi} \int_{|\lambda| \ge A+1} d\mu \int_{|\lambda| \le A} \frac{|\varphi(\lambda)| d\lambda}{(\lambda - \mu)^2 + \eta^2} \le C\eta ||\varphi||_{L_1(\mathbb{R})},\tag{2.53}$$

and for I_1 we use the inequalities:

$$I_1 \le 2|A+1| \sup_{|\lambda| \le A+1} |\varphi(\lambda) - \varphi * P_{\eta}(\lambda)| \le C|A+1| \int |1 - e^{-\eta|k|} ||\widehat{\varphi}(k)| dk.$$
 (2.54)

But since

$$\int |\widehat{\varphi}(k)| dk = \int \frac{|\widehat{\varphi}(k)|(1+2|k|)^s}{(1+2|k|)^s} dk \le ||\varphi||_s \left(\int \frac{dk}{(1+2|k|)^{2s}}\right)^{1/2} \le C||\varphi||_s,$$

(2.54) the Lebesgue theorem on the dominated convergence implies that $I_1 \to 0$, as $\eta \to 0$. Combining this with (2.54) we get that the set of all functions with finite supports, possessing the norm (2.52) is dense in \mathcal{L} . As it was mentioned above this implies that the set of the functions $\varphi * P_{\eta}$ with $\varphi \in \mathcal{L}$ is dense in \mathcal{L} with respect to the norm $||.||_{L_1(\mathbb{R})}$ and in view of (2.48) proves assumption (b) of the proposition. \square

Proof of Theorem 2. Let us note first that in the case when

$$p_n > C_* \log^{1/3} n, \quad n \to \infty, \tag{2.55}$$

the proof of Theorem 2 is rather simple. By the method of [3] one can prove the estimate

$$n^{-1}\mathbf{E}\{\operatorname{Tr} \mathcal{A}^{2m}\} \le C^2(1+\frac{m^3}{p_n})n^{-1}\mathbf{E}\{\operatorname{Tr} \mathcal{A}^{2m-2}\}.$$

Then under condition (2.55) it is easy to get the bound, valid for sufficiently big K:

$$\mathbf{Prob}\{||\mathcal{A}|| \geq K\} \leq \inf_{m} n\mathbf{E}\{\mathrm{Tr}\,(\mathcal{A}/K)^{2m}\} \leq \exp\{-p_n^{1/3}\log(K/2C) + \log n\} \to 0, n \to \infty.$$

Then, for any $\varphi \in \mathcal{H}_s^{(c)}$, if we consider a smooth function $\varphi^{(K)} \in \mathcal{H}_s$ with a finite support and such that $\varphi^{(K)}(\lambda) = \varphi(\lambda)$, $|\lambda| \leq K$, then evidently

$$|\mathbf{E}\{e^{ix\mathcal{N}_n^{\circ}[\varphi]/d_n}\} - \mathbf{E}\{e^{ix\mathcal{N}_n^{\circ}[\varphi^{(K)}]/d_n}\}| \leq \mathbf{Prob}\{||\mathcal{A}||K\} + d_n^{-1}|\mathbf{E}\{\mathcal{N}_n[\varphi]\} - \mathbf{E}\{\mathcal{N}_n[\varphi^{(K)}]\}| \to 0, \quad n \to \infty.$$

Thus, we can derive Theorem 2 from Theorem 1 almost immediately.

But if the inequality (2.55) is not fulfilled, then the proof of Theorem 2 is more complicated. It is based on the bound which is the analog of (2.16)

$$\frac{p_n}{n} \mathbf{Var}\{\mathcal{N}_n[\varphi]\} \le C(c) ||\widetilde{\varphi}||_s, \quad \varphi \in \mathcal{H}_s^{(c)}, \tag{2.56}$$

where $\widetilde{\varphi}(\lambda) = \varphi(\lambda) \cosh^{-1}(c\lambda)$. The main step here is the lemma, which is the generalization of Lemma 1

Lemma 4 Denote by $\gamma_n^{(c)} = \operatorname{Tr} G(z) e^{cA}$. Then for any $1 > \varepsilon > 0$

$$\frac{p_n}{n} \mathbf{Var}\{\gamma_n^{(c)}\} \le C(c, \varepsilon) \mathbf{E}\{|G_{11}|^{1+\varepsilon}\}/|\Im z|^{3+\varepsilon}. \tag{2.57}$$

Proof. According to Proposition 1 it is enough to prove that

$$\mathbf{E}\{|\gamma_n^{(c)} - \mathbf{E}_1\{\gamma_n^{(c)}\}|^2\} \le C(c,\varepsilon)\mathbf{E}\{|G_{11}|^{1+\varepsilon}\}/|\Im z|^{3+\varepsilon}p_n. \tag{2.58}$$

Let us set

$$G^{(1)}(z) = (\mathcal{A}^{(1)} - z)^{-1}, \quad \gamma_n^{(1c)} = \operatorname{Tr} G^{(1)}(z) e^{c\mathcal{A}^{(1)}}.$$

Note that differently from the proof of Proposition 1 here and below we denote by $\mathcal{A}^{(1)}$ the $n \times n$ matrix whose first line and column are zero and the other entries coincide with those of \mathcal{A} . We also denote $a^{(1)} = (0, a_{12}, \dots, a_{1n})$. Then we can write

$$\begin{split} \gamma_n^{(c)} - \mathbf{E}_1 \{ \gamma_n^{(c)} \} = & \gamma_n^{(c)} - \gamma_n^{(1c)} - \mathbf{E}_1 \{ \gamma_n^{(c)} - \gamma_n^{(1c)} \}, \\ \gamma_n^{(c)} - \gamma_n^{(1c)} = & \operatorname{Tr} \left(G(z) - G^{(1)}(z) \right) e^{c\mathcal{A}^{(1)}} + \operatorname{Tr} G^{(1)}(z) (e^{c\mathcal{A}} - e^{c\mathcal{A}^{(1)}}) \\ & + \operatorname{Tr} \left(G(z) - G^{(1)}(z) \right) (e^{c\mathcal{A}} - e^{c\mathcal{A}^{(1)}}) = I + II + III. \end{split}$$

Let us use the formulas

$$(G(z) - G^{(1)}(z))_{11} = z^{-1} - A^{-1}, \quad (G(z) - G^{(1)}(z))_{1i} = -A^{-1}(G^{(1)}a^{(1)})_{i},$$

$$(G(z) - G^{(1)}(z))_{ij} = -A^{-1}(G^{(1)}a^{(1)})_{i}(G^{(1)}a^{(1)})_{j}, \quad i, j \geq 2,$$

$$(e^{cA} - e^{cA^{(1)}})_{11} = (e^{cA})_{11} - 1, \quad (e^{cA} - e^{cA^{(1)}})_{1i} = c \int_{0}^{1} dt (e^{c(1-t)A})_{11} (e^{ctA^{(1)}}a^{(1)})_{i},$$

$$(e^{cA} - e^{cA^{(1)}})_{ij} = c^{2} \int_{0}^{1} dt \int_{0}^{1-t} d\tau (e^{ctA^{(1)}}a^{(1)})_{i} (e^{c\tau A^{(1)}}a^{(1)})_{j} (e^{c(1-t-\tau)A})_{11}, \quad i, j \geq 2.$$

Here A is defined in (2.9) and to obtain the last two lines we have used the Duhamel formula, valid for any matrices \mathcal{M} and $\mathcal{M}^{(1)}$:

$$e^{c\mathcal{M}} - e^{c\mathcal{M}^{(1)}} = c \int_0^1 dt e^{c\mathcal{M}^{(1)}t} (\mathcal{M} - \mathcal{M}^{(1)}) e^{c\mathcal{M}(1-t)}.$$

Moreover, we have taken into account that

$$\mathcal{A}_{11}^{(1)} = \mathcal{A}_{i1}^{(1)} = G_{i1}^{(1)} = (e^{t\mathcal{A}^{(1)}})_{i1} = 0, \quad \mathcal{A}_{11}^{(1)} = 0, \quad G_{11}^{(1)} = -z^{-1}, \quad (e^{t\mathcal{A}^{(1)}})_{11} = 1.$$

Hence we have

$$\begin{split} I = &z^{-1} - A^{-1} - A^{-1}(e^{c\mathcal{A}^{(1)}}(G^{(1)})^2 a^{(1)}, a^{(1)}), \\ II = &-z^{-1}((e^{c\mathcal{A}})_{11} - 1) + c^2 \int_0^1 dt \int_0^{1-t} d\tau (e^{c(1-t-\tau)\mathcal{A}})_{11}(G^{(1)}e^{ct\mathcal{A}^{(1)}}a^{(1)}, e^{c\tau\mathcal{A}^{(1)}}a^{(1)}) \\ = &-z^{-1}((e^{c\mathcal{A}})_{11} - 1) + c^2 \int_0^1 s(e^{cs\mathcal{A}^{(1)}}G^{(1)}a^{(1)}, a^{(1)})(e^{c(1-s)\mathcal{A}})_{11}ds, \\ III = &((e^{c\mathcal{A}})_{11} - 1)(z^{-1} - A^{-1}) - 2cA^{-1} \int_0^1 dt(e^{ct\mathcal{A}^{(1)}}G^{(1)}a^{(1)}, a^{(1)})(e^{c(1-t)\mathcal{A}})_{11} \\ &-c^2A^{-1} \int_0^1 dt \int_0^{1-t} d\tau (e^{ct\mathcal{A}^{(1)}}G^{(1)}a^{(1)}, a^{(1)})(e^{c\tau\mathcal{A}^{(1)}}G^{(1)}a^{(1)}, a^{(1)})(e^{c(1-t-\tau)\mathcal{A}})_{11}. \end{split}$$

Thus, denoting $B^{(c)} := (e^{c\mathcal{A}^{(1)}}G^{(1)}a^{(1)}, G^{(1)}a^{(1)})$ and using the Schwarz inequality and (2.10)-(2.11), we get for I

$$\mathbf{E}\{|I - \mathbf{E}_1\{I\}|^2\} \le 3\mathbf{E}\{|A^{-1} - \mathbf{E}_1\{A^{-1}\}|^2\} + 3\mathbf{E}\left\{\left|\frac{B_1^{(c)\circ}}{\mathbf{E}_1\{A\}}\right|^2\right\} + 3\mathbf{E}\left\{\left|\frac{B^{(c)}A_1^{\circ}}{A\mathbf{E}_1\{A\}}\right|^2\right\}.$$

Averaging with respect to $\{a_{1i}\}$ and then using the Hölder inequality, we get

$$\begin{split} \mathbf{E}_{1} \bigg\{ \bigg| \frac{B_{1}^{(c) \circ}}{\mathbf{E}_{1} \{A\}} \bigg|^{2} \bigg\} &\leq C \frac{n^{-1} \mathrm{Tr} \, |G^{(1)}|^{4} e^{2c \mathcal{A}^{(1)}}}{p_{n} |\mathbf{E}_{1} \{A\}|^{2}} \leq C \frac{n^{-1} \mathrm{Tr} \, |G^{(1)}|^{2} e^{2c \mathcal{A}^{(1)}}}{p_{n} |\Im z|^{2} |\mathbf{E}_{1} \{A\}|^{2}} \\ &\leq C \frac{(n^{-1} \mathrm{Tr} \, |G^{(1)}|^{2})^{1-\varepsilon} (n^{-1} \mathrm{Tr} \, |G^{(1)}|^{2} e^{2c \mathcal{A}^{(1)}/\varepsilon})^{\varepsilon}}{p_{n} |\Im z|^{2} |\mathbf{E}_{1} \{A\}|^{2}} \leq \frac{C (n^{-1} \mathrm{Tr} \, e^{2c \mathcal{A}^{(1)}/\varepsilon})^{\varepsilon}}{p_{n} |\Im z|^{3+\varepsilon} |\mathbf{E}_{1} \{A\}|^{1+\varepsilon}}, \end{split}$$

where $|G^{(1)}|^2 = G^{(1)*}G^{(1)}$. Similarly, using that $\Im A = \Im z(1 + (|G^{(1)}|^2 a^{(1)}, a^{(1)}))$ (see (2.12)), we obtain

$$\begin{split} &\mathbf{E}_{1}\bigg\{\bigg|\frac{B^{(c)}A_{1}^{\circ}}{A\mathbf{E}_{1}\{A\}}\bigg|^{2}\bigg\} \leq \mathbf{E}_{1}\bigg\{\bigg(\frac{(|G^{(1)}|^{2}a^{(1)},a^{(1)})^{1-\varepsilon}(e^{2c\mathcal{A}^{(1)}/\varepsilon}|G^{(1)}|^{2}a^{(1)},a^{(1)})^{\varepsilon}|A_{1}^{\circ}|}{|\Im z|(1+(|G^{(1)}|^{2}a^{(1)},a^{(1)}))|\mathbf{E}_{1}\{A\}|}\bigg)^{2}\bigg\} \\ &\leq \frac{\mathbf{E}_{1}\big\{(e^{2c\mathcal{A}^{(1)}/\varepsilon}a^{(1)},a^{(1)})^{2\varepsilon}|A_{1}^{\circ}|^{2}\big\}}{|\Im z|^{2}|\mathbf{E}_{1}\{A\}|^{2}} \leq \frac{\mathbf{E}_{1}^{1-2\varepsilon}\{|A_{1}^{\circ}|^{2}\}\mathbf{E}_{1}^{2\varepsilon}\{(e^{2c\mathcal{A}^{(1)}/\varepsilon}a^{(1)},a^{(1)})|A_{1}^{\circ}|^{2}\}}{|\Im z|^{2}|\mathbf{E}_{1}\{A\}|^{2}} \\ &\leq C\frac{(n^{-1}\mathrm{Tr}\,|G^{(1)}|^{2})^{1-2\varepsilon}(n^{-1}||G^{(1)}||^{2}\mathrm{Tr}\,e^{2c\mathcal{A}^{(1)}/\varepsilon})^{2\varepsilon}}{p_{n}|\Im z|^{2}|\mathbf{E}_{1}\{A\}|^{2}} \leq C\frac{(n^{-1}\mathrm{Tr}\,e^{2c\mathcal{A}^{(1)}/\varepsilon})^{2\varepsilon}}{p_{n}|\Im z|^{3+2\varepsilon}|\mathbf{E}_{1}\{A\}|^{1+2\varepsilon}}. \end{split}$$

The terms with II and III can be estimated similarly, if we use also the bound

$$\mathbf{E}_{1}\{((e^{c\mathcal{A}})_{11} - \mathbf{E}_{1}\{(e^{c\mathcal{A}})_{11}\})^{2}\} \le Cp_{n}^{-1}\mathbf{E}_{1}^{1/2}\{(e^{8|c|\mathcal{A}})_{11}\}. \tag{2.59}$$

To prove (2.59), we prove first that

$$\mathbf{E}_{1}\{((\mathcal{A}^{m})_{11} - \mathbf{E}_{1}\{(\mathcal{A}^{m})_{11}\})^{2}\} \le Cm^{2}2^{m}\mathbf{E}_{1}^{1/2}\{(\mathcal{A}^{4m-2})_{11}\}/p_{n}. \tag{2.60}$$

It is easy to see that

$$(\mathcal{A}^m)_{11} = \sum_{k=1}^{\lfloor m/2 \rfloor} \sum_{l_1 + \dots + l_k = m-2k} \Sigma^{(l_1)} \dots \Sigma^{(l_k)}, \quad \Sigma^{(l_p)} := (\mathcal{A}^{(1)l_p} a^{(1)}, a^{(1)})$$
 (2.61)

with $\mathcal{A}^{(1)}$ and $a^{(1)}$ of (2.7). Thus, using that for all l $\mathbf{E}\{a_{ij}^l\} \geq 0$, we have

$$\begin{split} &\mathbf{E}_{1} \Big\{ \Big(\Sigma^{(l_{1})} \dots \Sigma^{(l_{k})} - \mathbf{E}_{1} \{ \Sigma^{(l_{1})} \} \dots \mathbf{E}_{1} \{ \Sigma^{(l_{k})} \} \Big)^{2} \Big\} \\ &\leq k \sum_{j=1}^{k} \mathbf{E}_{1}^{1/2} \Big\{ \Big(\Sigma^{(l_{j})} - \mathbf{E}_{1} \{ \Sigma^{(l_{j})} \} \Big)^{4} \Big\} \mathbf{E}_{1}^{1/2} \Big\{ \prod_{i \neq j} (\Sigma^{(l_{i})})^{4} \Big\} \\ &\leq C k p_{n}^{-1} \sum_{j=1}^{k} \mathbf{E}_{1}^{1/2} \{ \Sigma^{(4l_{j})} \} \mathbf{E}_{1}^{1/2} \Big\{ \prod_{i \neq j} \Sigma^{(4l_{i})} \Big\} \leq C k^{2} p_{n}^{-1} \mathbf{E}_{1}^{1/2} \Big\{ \prod_{i} \Sigma^{(4l_{i})} \Big\}. \end{split}$$

Taking the sum as in (2.61) and using the Schwarz inequality, we obtain (2.60). The Taylor expansion, the Schwarz inequality, and (2.60) imply (2.59):

$$\begin{aligned} &\mathbf{E}\{(e^{c\mathcal{A}})_{11} - \mathbf{E}_{1}\{(e^{c\mathcal{A}})_{11}\})^{2}\} \leq p_{n}^{-1} \sum_{m=1}^{\infty} \frac{|2c^{2}|^{m}}{(m!)^{2}} \mathbf{E}\left\{\left((\mathcal{A}^{m})_{11} - \mathbf{E}_{1}\{(\mathcal{A}^{m})_{11}\}\right)^{2}\right\} \\ &\leq Cp_{n}^{-1} \sum_{m=1}^{\infty} \frac{m^{2}|2c|^{2m}}{(2m)!} \mathbf{E}_{1}^{1/2}\{(\mathcal{A}^{4m-2})_{11}\} \leq Cp_{n}^{-1} \mathbf{E}_{1}^{1/2}\left\{\sum_{m=1}^{\infty} \frac{|8c|^{4m}}{(4m)!}(\mathcal{A}^{4m-2})_{11}\right\} \\ &\leq Cp_{n}^{-1} \mathbf{E}_{1}^{1/2}\left\{(e^{8|c|\mathcal{A}})_{11}\right\}. \end{aligned}$$

Lemma 4 is proven. \square

The next step is the analog of Proposition 2

Proposition 4 For any $\varphi \in \mathcal{H}_s^{(c)}$

$$\mathbf{Var}\{\mathcal{N}_{n}[\varphi]\} \le C_{s} ||\varphi^{(c)}||_{s}^{2} \int_{0}^{\infty} dy e^{-y} y^{2s-1} \int_{-\infty}^{\infty} \mathbf{Var}\{\gamma_{n}^{(c)}(x+iy)\} dx, \tag{2.62}$$

where $\widetilde{\varphi}(\lambda) = \varphi(\lambda) \cosh^{-1}(c\lambda)$ and $||\widetilde{\varphi}||_s$ is defined in (1.9).

The proof of Proposition 4 coincides with that of Proposition 2, if we replace the operator \mathcal{V} of (2.18) by the operator $\mathcal{V}^{(c)}$ whose Fourier transform has the kernel

$$\widehat{\mathcal{V}^{(c)}}(k_1, k_2) = \mathbf{Cov}\{\operatorname{Tr}\cosh(c\mathcal{A})e^{ik_1\mathcal{A}}, \operatorname{Tr}\cosh(c\mathcal{A})e^{ik_2\mathcal{A}}\}, \quad \mathbf{Var}\{\mathcal{N}_n[\varphi]\} = (\mathcal{V}^{(c)}\widetilde{\varphi}, \widetilde{\varphi}).$$

Now one can derive (2.56) from Lemma 4 and Proposition 4 by the same argument that we used in Lemma 2 to derive (2.21) from Lemma 1 and Proposition 2.

Having in mind the bound (2.56), we can derive Theorem 2 from Proposition 3, if we are able to prove CLT for some dense subset of $\mathcal{H}_s^{(c)}$, e.g., for φ with finite supports. But if φ has a finite support and belongs to $\mathcal{H}_s^{(c)}$, it belongs also automatically to \mathcal{H}_s , thus we can apply Theorem 1 to it. This completes the proof of Theorem 2.

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