A DISCRETE MODEL FOR NONLOCAL TRANSPORT EQUATIONS WITH FRACTIONAL DISSIPATION

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ABSTRACT. In this note, we propose a discrete model to study one-dimensional transport equations with non-local drift and supercritical dissipation. The inspiration for our model is the equation

$$\theta_t + (H\theta)\theta_x + (-\Delta)^{\alpha}\theta = 0$$

where H is the Hilbert transform. For our discrete model, we present blow-up results that are analogous to the known results for the above equation. In addition, we will prove regularity for our discrete model which suggests supercritical regularity in the range $1/4 < \alpha < 1/2$ in the continuous setting.

1. Introduction

Before we introduce the equation we will primarily be concerned with, we will introduce the more complicated equation we hope our equation can model. Consider the equation

(1)
$$\theta_t(x,t) + (H\theta)(x,t)\theta_x(x,t) + (-\Delta)^{\alpha}\theta(x,t) = 0 \quad (x,t) \in \mathbb{R} \times [0,\infty)$$

where H is the Hilbert transform and $(-\Delta)^{\alpha}$ where $\alpha > 0$ is the fractional Laplacian. When $\alpha = 0$, by convention, we take the dissipation to be absent in the equation. The equation can be viewed as a toy-model for the 2D surface quasi-geostrophic (SQG) equation

(2)
$$\theta_t(x,t) + u \cdot \nabla \theta(x,t) + (-\Delta)^{\alpha} \theta(x,t) = 0, \quad u = \nabla^{\perp} (-\Delta)^{-1/2} \theta, \quad (x,t) \in \mathbb{R}^2 \times [0,\infty)$$

for which the question of regularity in the range $0 \le \alpha < 1/2$ is a longstanding open problem. See [1] and [7] for recent progress and references.

By virtue of (1) satisfying an L^{∞} maximum principle, we can roughly divide the analysis of (1) into three categories according to the value of α . For $\alpha > 1/2$, the equation is subcritical in that one expects the dissipative term to dominate the nonlinearity and one has global regularity from functional analysis methods. For $\alpha = 1/2$, the equation is critical in the sense that dissipative term and nonlinear term are thought to be equally balanced. For $0 \le \alpha < 1/2$, the equation is said to be supercritical in that the nonlinear term is thought to dominate the dynamics due to scaling considerations. It was shown in [3] and [8] that when $0 \le \alpha < 1/4$, solutions can develop finite time blow-up. In [5], it was shown that for $\alpha \ge 1/2$ solutions exist for all time. In the supercritical range $1/4 \le \alpha < 1/2$, the question of global existence versus blow-up is open. It was shown in [4] that supercritical solutions derived in the vanishing viscosity limit, in a sense, become eventually regular. Also, it was shown that solutions supercritical by a logarithm are globally regular. We are not aware of any example where supercritical, with respect to the basic conservation laws, regularity has been proved for a fluid mechanics PDE.

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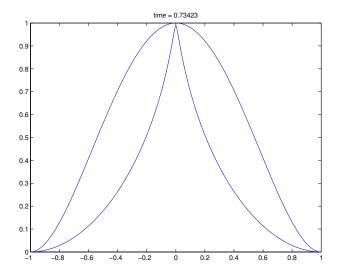


FIGURE 1. Plot of the solution for the inviscid equation with initial data $\theta_0(x) = (1-x^2)^2 \chi_{[-1,1]}(x)$

1.1. Motivation for the model. The proof of blow-up for $0 \le \alpha < 1/4$ relies on the following novel inequality that for $f \in C_c^{\infty}(\mathbb{R}^+)$ and $0 < \delta < 1$, there exists a constant C_{δ} such that

$$-\int_0^\infty \frac{f_x(x)(Hf)(x)}{x^{1+\delta}} dx \ge C_\delta \int_0^\infty \frac{f^2(x)}{x^{2+\delta}} dx.$$

The proof given in [3] uses tools from complex analysis and is used to prove blow-up for even positive initial data with a maximum at 0. In [7], another more elementary proof of blow-up for $0 \le \alpha < 1/4$ was given. The proof also goes by way of (3) but without appealing to complex analysis. In particular, the complicated non-locality of the Hilbert transform is handled by the following key inequality from [7]:

Proposition 1.1. Suppose that the function f(x) is C^1 , even, $f'(x) \ge 0$ for x > 0 and f is bounded on \mathbb{R} with f(0) = 0. Then for 1 < q < 2,

$$Hf(x) \le \log(q-1)(f(qx) - f(q^{-1}x)).$$

Consider initial data for (1) that is C^1 , even, and monotone decaying away from the origin. These properties are preserved by (1). The proposition is then applied to $f(x) = \theta(0,t) - \theta(x,t)$. To derive our model, we will look at dyadic points. Consider the following system of ODEs

(4)
$$a'_k(t) = -(a_k - a_{k-1})^2(t)2^k, \quad a_k(0) = a_k^0, \quad k \ge 1.$$

and set $a_0 = 0$. This system serves as a discrete model for

(5)
$$\theta_t(x,t) = -(H\theta)\theta_x(x,t).$$

One can think of $a_k(t)$ as approximating $\theta(2^{-k},t)$. In particular, one could think that $-(a_k - a_{k-1})^2(t)2^k \approx \theta_x(2^{-k},t)$ and that, using Proposition 1.1 as inspiration, we shall take $(a_k - a_{k-1})$ as our approximation of $H\theta(2^{-k},t)$. The reason for considering a dyadic model is the nature of conjectured blow up for (5). A cusp appears to form in finite time at the origin (see Figure 1) and it is our hope that studying (4) will help us understand such phenomenon.

Now, we would like to formulate a discrete version of the fractional Laplacian $(-\Delta)^{\alpha}$. Recall for θ smooth enough,

$$(-\Delta)^{\alpha}\theta(x) = P.V. \int_{-\infty}^{\infty} \frac{\theta(x) - \theta(y)}{|x - y|^{1 + 2\alpha}} dy,$$

see [2] for a derivation. If we take θ to be even then

$$(-\Delta)^{\alpha} \theta(x) = \int_0^{\infty} \left[\frac{1}{|x - y|^{1 + 2\alpha}} + \frac{1}{|x + y|^{1 + 2\alpha}} \right] (\theta(x) - \theta(y)) \, dy$$

Let $x=2^{-k}$ and $y=2^{-n}$. If k < n then $\frac{1}{|x-y|^{1+2\alpha}} + \frac{1}{|x+y|^{1+2\alpha}} \approx 2^{(1+2\alpha)k}$. Similarly, if k > n, then $\frac{1}{|x-y|^{1+2\alpha}} + \frac{1}{|x+y|^{1+2\alpha}} \approx 2^{(1+2\alpha)n}$. Also, $\theta(x) - \theta(y) \approx a_k - a_n$. Combining these observations, we can see that a reasonable model for fractional dissipation would be

(6)
$$(\mathcal{L}^{\alpha}a)_k = \sum_{n=0}^{k-1} (a_k - a_n) 2^{2\alpha n} + \sum_{n=k+1}^{\infty} (a_k - a_n) 2^{2\alpha k} 2^{k-n}.$$

In our discrete formulation, we have ignored the tails of the integral defining $(-\Delta)^{\alpha}$. Adding the tails into (6) (i.e. letting the first sum go to $-\infty$) doesn't change the main results of this note and just adds more complications to the calculations.

Thus, our model for (1) is

(7)
$$a'_{k} = -(a_{k} - a_{k-1})^{2} 2^{k} - (\mathcal{L}^{\alpha} a)_{k}.$$
$$a'_{0} = -(\mathcal{L}^{\alpha} a)_{0}$$

For simplicity of notation, we will sometimes omit the α and simply write \mathcal{L} instead. It should be noted that our model can also serve as discrete model for other non-local transport equations with fractional diffusion such as

$$\theta_t(x,t) = u(x,t)\theta_x(x,t) - (-\Delta)^{\alpha}(x,t)$$

where

$$u(x,t) = \left\{ \begin{array}{ll} \theta(x,t) - \theta(2x,t) & x \geq 0 \\ \theta(2x,t) - \theta(x,t) & x < 0 \end{array} \right.$$

It is unknown whether solutions to this type of equations with general initial data can blow-up or exist globally in time in the supercritical range $0 \le \alpha < 1/2$. In [10], a related kind of "non-local" Burgers type equation was studied and blow-up was observed in the non-viscous case in certain situations.

1.2. Outline of Main Results. In Section 2, we present several properties of solutions to (7) and show that solutions to (4) blow-up, strengthening the analogy between the model and (1). In Section 3, we prove an a-priori bound on solutions akin to a global in-time Hölder 1/2 bound in the continuous setting. In Section 4, we use this bound to prove blow-up for $0 < \alpha < 1/4$ and global regularity for $1/4 \le \alpha < 1/2$, which can be considered the main result of this note. If this result were carried into the continuous setting, it would suggest that solutions to (1) in the supercritical range $1/4 \le \alpha < 1/2$ are globally regular, contrary to natural scaling considerations.

2. Local existence and properties of solutions

Define the space

$$X^{s} = \{\{a_{k}\}_{k=0}^{\infty} : ||a||_{X}^{s} := \sup_{k} |a_{k}| + \sup_{k \ge 1} |a_{k} - a_{k-1}| 2^{sk} < \infty\}.$$

One could think of X^s as being analogous to the Hölder spaces. However, the X^s are made to deal with behavior near the origin. Let $b_{k,s} = (a_k - a_{k-1})2^{sk}$.

Before we show local existence for the full system, we will state some facts about \mathcal{L} and its associated semigroup.

Lemma 2.1. Let $0 < \alpha < 1/2$.

(a) Suppose that for the index k, $b_{k,s} > c_s ||a||_{X^s}$ where

$$c_s = \frac{3}{4}(2^s - 1)^{-1} \left(1 - \frac{1}{2^{s+1} - 1}\right) < 1.$$

Then

$$((\mathcal{L}^{\alpha}a)_k - (\mathcal{L}^{\alpha}a)_{k-1}) 2^{sk} \ge C(\alpha)(2^{2\alpha k} - 2^{2\alpha})b_{k,s}$$

where $C(\alpha)$ is a positive constant only dependent on α .

- (b) The operator \mathcal{L}^{α} generates a contracting semigroup $e^{-t\mathcal{L}^{\alpha}}$ on X^s for all s>0.
- (c) The following identity, which is analogous to $\int (-\Delta)^{\alpha} \theta = 0$, is true

$$\sum_{k=0}^{\infty} (\mathcal{L}^{\alpha} a)_k 2^{-k} = 0.$$

Proof. (a) We will consider k > 1 as the case for k = 1 is similar. By direct computation,

$$(\mathcal{L}a)_k - (\mathcal{L}a)_{k-1} = \sum_{n=0}^{k-2} (a_k - a_{k-1}) 2^{2\alpha n} + 2(a_k - a_{k-1}) 2^{2\alpha(k-1)}$$

$$+ 2^{(1+2\alpha)k} (1 - 2^{-(1+2\alpha)}) \sum_{n=k+1}^{\infty} (a_k - a_n) 2^{-n}$$

$$= (a_k - a_{k-1}) \left(\frac{2^{2\alpha(k-1)} - 1}{2^{2\alpha} - 1}\right) + 2(a_k - a_{k-1}) 2^{2\alpha(k-1)}$$

$$+ 2^{(1+2\alpha)k} \sum_{n=k+1}^{\infty} (a_k - a_n) (1 - 2^{-(1+2\alpha)}) 2^{-n}$$

$$:= I + II + III$$

By hypothesis, we have $c_s b_{j,s} < b_{k,s}$ for j > k. In terms of a, this means, for $n \ge k+1$,

$$(a_n - a_k) = \sum_{j=k+1}^n (a_j - a_{j-1}) \le c_s^{-1} (a_k - a_{k-1}) \sum_{j=1}^{n-k} 2^{-js} = \frac{c_s^{-1}}{2^s - 1} (a_k - a_{k-1}) (1 - 2^{-(n-k)s}).$$

Then

$$III \geq -\frac{c_s^{-1}}{2^s - 1} 2^{(1+2\alpha)k} (a_k - a_{k-1}) (1 - 2^{-(1+2\alpha)}) \sum_{n=k+1}^{\infty} (1 - 2^{-(n-k)s}) 2^{-n}$$

$$= -\frac{4}{3} (1 - 2^{-(1+2\alpha)}) (a_k - a_{k-1}) 2^{2\alpha k}.$$

From this inequality and using that $0 < \alpha < 1/2$, it is easy to see that $II + III \ge 0$. Then

$$((\mathcal{L}a)_k - (\mathcal{L}a)_{k-1}) \, 2^{sk} \ge I \cdot 2^{sk} = C(\alpha)(2^{2\alpha k} - 2^{2\alpha}) b_{k,s}$$

(b) Consider the system

$$a_k' = -(\mathcal{L}a)_k.$$

Using (a), it's not hard to see that $\frac{d}{dt}||a||_{X^s}(t) < 0$ and so $e^{-t\mathcal{L}}$ is a contracting semigroup.

(c) follows by explicit computation. \Box

We will need the following version of Picard's theorem to prove local existence

Theorem 2.2. (Picard Fixed Point) Let Y be a Banach space and let $\Gamma: Y \times Y \to Y$ be a bilinear operator such that for all $a, b \in Y$,

$$\|\Gamma(a,b)\|_Y \le \eta \|a\|_Y \|b\|_Y.$$

Then for any $a^0 \in Y$ with $4\eta \|a^0\|_Y < 1$, the equation $a = a^0 + \Gamma(a, a)$ has a unique solution $a \in Y$ such that $\|a\|_Y \le 1/2\eta$.

Theorem 2.3. (Local existence) Let $\{a_k(0)\} \in X^s$ where $s \ge 1$. Then there exists $T = T(\|a(0)\|_{X^s})$ such that there is a unique solution $\{a_k(t)\} \in C([0,T],X^s)$.

Proof. The argument is fairly standard and we will provide a sketch. By Lemma 2.1, the operator \mathcal{L} generates a contracting semigroup $e^{-t\mathcal{L}}$ on X^s . Observe that a solution will satisfy

$$a_k(t) = e^{-t\mathcal{L}}a_k^0 - \int_0^t e^{-(t-s)\mathcal{L}} \{2^k(a_k - a_{k-1})^2(s)\} ds$$

In writing $e^{-(t-s)\mathcal{L}}$ in the integral, we have slightly abused notation. The integrand is the kth element of $e^{-(t-s)\mathcal{L}}$ applied to the sequence $\{2^j(a_j-a_{j-1})^2(s)\}_{j=1}^{\infty}$.

Define a bilinear operator $\Gamma: X^s \times X^s \to X^s$ by

$$\Gamma(a,b)_j(t) = -\int_0^t e^{-(t-s)\mathcal{L}} \{ 2^k (a_k - a_{k-1})(s)(b_k - b_{k-1})(s) \} ds$$

Define $\gamma(a,b)(s)$ by

$$\gamma(a,b)_k(s) = -e^{-(t-s)\mathcal{L}} \{ 2^k (a_k - a_{k-1})(s)(b_k - b_{k-1})(s) \}$$

Then after doing basic estimates, one has $\|\gamma(a,b)(s)\|_{X^s} \leq C\|a(s)\|_{X^s}\|b(s)\|_{X^s}$ for $s \geq 1$. From this, we have an estimate on Γ and choosing t small enough, we can apply Picard. \square

Also, we have the following preservation properties:

Lemma 2.4. Let $\{a_k(t)\}$ be a solution of (7) in $C([0,T],X^s)$, s>1. Suppose $\{a_k^0\}$ is non-decreasing in k and is non-negative.

- (a) (Monotonicity and positivity) Then for all $t \leq T$, $\{a_k(t)\}$ is non-decreasing in k and is non-negative.
 - **(b)** (Max stays at 0) For $t \in [0,T]$, we have that

$$\sup_{k} a_k(t) = \lim_{k \to \infty} a_k(t)$$

(c) $(\ell^{\infty} \text{ maximum principle})$ The supremum $\sup_k a_k(t)$ is non-increasing in t and $a_0(t)$ is increasing in t. If $\{a_k(t)\}$ is a solution to (4), then $\sup_k a_k(t)$ is constant in time

Proof.

(a) Set $b_k(t) = (a_k - a_{k-1})(t)2^k$. For a contradiction, suppose there exists j and a time t_0 such that $b_j(t_0) < 0$. Choose a sufficiently small ϵ with $0 < \epsilon < 1$ such that $b_j(t_0) < -\epsilon/t_0$. Define a function f by $f(t) = \frac{-\epsilon}{2t_0 - t}$ for $0 \le t < 2t_0$, so $b_j(t_0) < f(t_0)$.

Let $g(t) = \inf_k b_k(t)$. Let t_1 be the first time g crosses f so $\inf_k b_k(t_1) = f(t_1)$. Because $b_k(t_1) \to 0$ as $k \to \infty$, the infimum is actually achieved for some index. Let k_1 be the first index for which $b_{k_1}(t_1) = f(t_1)$. By minimality of t_1 , $b'_{k_1}(t_1) \le f'(t_1) = -\frac{\epsilon}{(2t_0 - t)^2}$. However, at time t_1 , b_{k_1} also satisfies

$$b'_{k_1} = -b_{k_1}^2 + 2b_{k_1-1}^2 - ((\mathcal{L}a)_{k_1} - (\mathcal{L}a)_{k_1-1}) 2^{k_1}$$

$$\geq -\frac{\epsilon^2}{(2t_0 - t)^2} - ((\mathcal{L}a)_{k_1} - (\mathcal{L}a)_{k_1-1}) 2^{k_1}$$

By an argument analogous to Lemma 2.1, the contribution from the dissipative terms on the right side of the inequality above is negative. From this, we arrive at a contradiction to the above inequality. Because a_0 is always non-negative, a_k stays non-negative for all time by the monotonicity just proved.

- (b) follows directly from (a).
- (c) By (a), we see that $(\mathcal{L}a)_0(t) < 0$ for all time $t \leq T$ so $a_0(t)$ is increasing in t. Now, for $t \in [0,T]$

$$\frac{d}{dt} \sup_{k \to \infty} a_k(t) = \frac{d}{dt} \lim_{k \to \infty} a_k(t) = \lim_{k \to \infty} \frac{d}{dt} a_k(t) = \lim_{k \to \infty} \left(-(a_k - a_{k-1})^2(t) 2^k - (\mathcal{L}^{\alpha} a)_k(t) \right)$$

Since we are on a compact time interval, the convergence above is uniform and so we are justified in interchanging limit and derivative. The limit of the first term on the far right is zero. Using that $\{a_k(t)\}\in X^s$ where $2\alpha<1< s$,

$$\lim_{k \to \infty} (\mathcal{L}^{\alpha} a)_k(t) = \lim_{k \to \infty} \left(\sum_{n=0}^{k-1} (a_k - a_n)(t) 2^{2\alpha n} + \sum_{n=k+1}^{\infty} (a_k - a_n)(t) 2^{2\alpha k} 2^{k-n} \right) > 0$$

because the limit of the first term is positive by monotonicity and the limit of the second term is zero. This completes the proof.

In the spirit of [7], we show that solutions of (4) blow up in finite time.

Theorem 2.5. Let $\{a_k\}_{k=0}^{\infty}$ be a solution of (4) in X^s , s > 1 with initial data that is non-negative and non-decreasing in k. Then $\{a_k\}$ develops blow up in X^s for every s > 1 in finite time.

Define

$$J(t) = \sum_{k=1}^{\infty} (a - a_k(t)) 2^{k\delta}$$

where $0 < \delta < 1$. Observe that, $J \le ||a||_{X^1} \sum_{k=1}^{\infty} 2^{k(\delta-1)}$. To prove we have blow up, we will show that J must become infinite in finite time. We will need the following lemma

Lemma 2.6. Let $0 < \delta < 1$. Suppose that $\{a_k\} \in X^s$, s > 1, non-decreasing in k, and $a := \lim_{k \to \infty} a_k < \infty$. Then

(8)
$$\sum_{k=1}^{\infty} (a_k(t) - a_{k-1}(t))^2 2^{k(\delta+1)} \ge C_0(\delta) \sum_{k=1}^{\infty} (a - a_k(t))^2 2^{k(\delta+1)}$$

where $C_0(\delta)$ is constant depending only on δ .

Proof. Choose c > 0 such that $(c+1)^{-2}2^{\delta+1} > 1$. We call k "good" if $a_k - a_{k-1} \ge c(a - a_k)$ and bad otherwise.

Claim: If the set of all good k are finite, then both sides of (8) are infinite.

If the set of good k is finite, then there exists K such that for k > K, $a_k - a_{k-1} \le c(a - a_k)$ or equivalently $(a - a_k) \ge (c + 1)^{-1}(a - a_{k-1})$. Then

$$(a-a_k)^2 2^{k(\delta+1)} \ge (c+1)^{-2(k-K+1)} (a-a_K)^2 2^{k(\delta+1)} \to \infty$$

as $k \to \infty$ by the choice of c, so the right side of (8) diverges.

Define c_n by $a_k - a_{k-1} = c_k(a - a_k)$. Then $(a - a_{k-1}) = (1 + c_k)(a - a_k)$. Since $\lim_{k \to \infty} (a - a_k) = 0$, $\prod_{k=K}^{\infty} (1 - \frac{c_k}{c_k + 1}) = \prod_{k=K}^{\infty} \frac{1}{1 + c_k} = 0$, so $\sum_k c_k = \infty$. For the bad k,

$$(a_k - a_{k-1})^2 2^{k(\delta+1)} = c_k^2 (a - a_k)^2 2^{k(\delta+1)}$$

from which it is not hard to see that the left side of (8) will be infinite as well. The claim is proven.

Now suppose k_{j-1} , k_j are good such that for $k_{j-1} < k < k_j$, k is bad. Then $(a - a_{k-1}) \le (1+c)(a-a_k)$, which implies

$$\sum_{k=k_{j-1}+1}^{k_j} (a-a_k)^2 2^{k(\delta+1)} \leq \sum_{k=k_{j-1}+1}^{k_j} (1+c)^{2(k_j-k)} (a-a_{k_j})^2 2^{k(\delta+1)}$$

$$= (a-a_{k_j})^2 2^{k_j(\delta+1)} \sum_{k_{j-1}+1}^{k_j} (1+c)^{2(k_j-k)} 2^{(k-k_j)(\delta+1)}$$

$$= (a-a_{k_j})^2 2^{k_j(\delta+1)} \sum_{n=0}^{k_j-k_{j-1}+1} (1+c)^{2n} 2^{-n(\delta+1)}$$

$$\leq C(\delta)(a-a_{k_j})^2 2^{k_j(\delta+1)}$$

where the last inequality comes from the choice of c. Treating all bad k this way, the inequality (8) follows. \square

Proof of Theorem 2.5 Using the Lemma 2.4(c) and 2.6 as well as Holder's inequality, we have

$$\frac{d}{dt}J(t) = \sum_{k=1}^{\infty} (a_k(t) - a_{k-1}(t))^2 2^{k(\delta+1)} \ge C_0 \sum_{k=1}^{\infty} (a - a_k(t))^2 2^{k(\delta+1)}
\ge C_0 \left(\sum_{k=1}^{\infty} 2^{k(\delta-1)}\right)^{-1} \left(\sum_{k=1}^{\infty} (a - a_k(t)) 2^{k\delta}\right)^2 = C_1(\delta)J(t)^2$$

The result follows from Gronwall's inequality. \square

3. A-PRIORI HÖLDER-1/2 BOUND

The purpose of this section is to prove an a-priori bound for solutions of (7) from which the results of next section will follow. If this result of this section were to be carried to the continuous setting, it would mean that solutions to (1) are bounded in the Hölder class $C^{1/2}$. This regularization effect has recently been conjectured in [9] for the vanishing viscosity approximation. If such a bound were to hold, it would mean that $\alpha = 1/4$ is the true critical power for regularity and would answer the open question regarding (1) stated in the introduction. For the SQG equation (2), the $C^{1-2\alpha}$ norm is critical [6]: weak solutions that are bounded in time in $C^{1-2\alpha}$ are classical solutions.

First, we will prove bounds on (4), the model without dissipation, as it is more elementary. Then we will generalize to the full model (7) with dissipation.

We rewrite (4) in a different form, which allows us to give a more detailed picture of how blow-up can occur. Let $b_k = (a_k - a_{k-1})2^k$. Then the b_k 's satisfy

$$b_k'(t) = -b_k^2 + 2b_{k-1}^2$$

with corresponding initial data

$$b_k(0) = -(a_k^0 - a_{k-1}^0)2^k := b_k^0$$

By convention, we set $b_1' = -b_1^2$ and $b_0 = 0$. We will take $b_k(0) > 0$ for all k so b_k will remain pnon-negative by Lemma 2.4(a). The fact that the a_k 's blow up in X^s , s > 1, means that there exists T > 0 such that $\lim_{t \to T} \sup_k b_k(t) 2^{k(s-1)} = \infty$. In what follows, we work with X^s , s > 1 solutions. We will need the following group of lemmas.

Lemma 3.1. Suppose that for $T_1 \le t \le T_2$, $b_{k-1}(t) < \sqrt{2}b_{k-2}(t)$ (or >) and $b_k(T_1) < \sqrt{2}b_{k-1}(T_1)$ (or >). Then for $T_1 < t < T_2$, we have $b_k(t) < \sqrt{2}b_{k-1}(t)$ (or >).

Proof

For a contradiction suppose the set $A := \{T_2 \ge t > T_1 : b_k(t) \ge \sqrt{2}b_{k-1}(t)\}$ is non-empty. By continuity, A has an minimum call it t_0 . Then $b_k(t_0) = \sqrt{2}b_{k-1}(t_0)$, which implies

$$(b_k - \sqrt{2}b_{k-1})'(t_0) = -b_k^2(t_0) + 2b_{k-1}^2(t_0) + \sqrt{2}b_{k-1}^2(t_0) - 2\sqrt{2}b_{k-2}^2(t_0)$$
$$= \sqrt{2}(b_{k-1}^2(t_0) - 2b_{k-2}^2(t_0)) < 0.$$

From this, we can arrive at a contradiction to t_0 being the minimum of A. The proof for the opposite inequality is similar \square

The proof of the following lemma can be done in a similar way.

Lemma 3.2. Suppose that for $T_1 \le t \le T_2$, $b_{k-1}(t) < b_{k-2}(t)$ (or >) and $b_k(T_1) < b_{k-1}(T_1)$ (or >). Then for $T_1 \le t < T_2$, we have $b_k(t) < b_{k-1}(t)$ (or >).

Now, we analyze the following scenario. Suppose there exists K_0 such that the initial data $\{b_k^0\}$ satisfies

(9)
$$b_k^0 < b_{k-1}^0, \text{ for } k > K_0$$

(10)
$$b_{k-1}^0 < b_k^0 < \sqrt{2}b_{k-1}^0, \text{ for } k \le K_0$$

and consider solutions $\{b_k\}$ with this type of initial data. Observe that $\max_k b_k^0 = b_{K_0}$. As we let

time evolve, due to the control given by the above lemmas, the solution will have a similar structure as the initial data. That is, at anytime t < T, there exists K_t such that

$$b_k(t) < b_{k-1}(t), \text{ for } k > K_t$$

 $b_{k-1}(t) < b_k(t) < \sqrt{2}b_{k-1}(t), \text{ for } k \le K_t.$

At this point, we have $\max_k b_k(t) = b_{K_t}(t)$ and $K_t \geq K_0$. Thinking in terms of $a_k \approx \theta(2^{-k}, t)$ and b_k being the modulus of the slope of the line between $(2^{-k}, a_k)$ and $(2^{-k+1}, a_{k-1})$, we are approximating a solution θ of (1) by a piecewise linear function taking the value a_k at 2^{-k} (see figure 2).

The picture is of a front propagating towards the origin attempting to form a cusp. Once the front hits the origin (i.e. the time when $\sup_k b_k = \lim_{k \to \infty} b_k$), blow-up occurs. The index K_t can be thought of as the position of the inflection point of the function that sequence $\{b_k\}$ is approximating. One can hypothesize that $b_k(t) \approx \sqrt{2}b_{k-1}(t)$ as $t \to T$. In fact, the following result suggests that this may be the case.

Theorem 3.3. Suppose we have initial data satisfying (9) and (10). Then for $0 < \beta \le 1/2$,

$$\sup_{t < T, k} b_k 2^{k(\beta - 1)} < \infty$$

Proof. Follows from Lemma 3.1 and 3.2. \square

3.1. With Dissipation. We will consider a class of initial data slightly different from the previous section. The blow-up picture becomes less clear, but we still are able to have the analogue to Theorem 3.3. We still consider initial data $\{a_k^0\}_{k=0}^{\infty} \in X^s$, s > 1, that is positive and non-decreasing in k. As before we let $b_k = (a_k - a_{k-1})2^k$ and we start with the following configuration

(11)
$$b_k^0 < b_{k-1}^0, \text{ for } k > K_0$$

(12)
$$b_{k}^{0} < \sqrt{2}b_{k-1}^{0}, \text{ for } k \leq K_{0}$$
(13)
$$\sup_{k} b_{k}^{0} = b_{K_{0}}^{0}$$

$$\sup_{k} b_k^0 = b_{K_0}^0$$

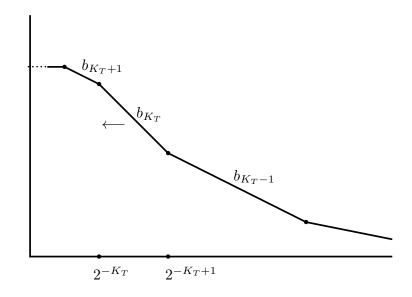


Figure 2.

for some positive integer K_0 . The following lemma is an analogue to the observation made with the model without dissipation.

Lemma 3.4. Let $0 < \alpha < 1/2$. Suppose we have initial data as in (11)-(13). For all t such that the solution is smooth, there exists a positive integer K_t such that

$$(14) b_k < b_{k-1}, for k > K_t$$

$$(15) b_k < \sqrt{2}b_{k-1}, for k \le K_t$$

In addition, $K_t \geq K_0$.

Proof. The result will quickly follow from the following two claims.

Claim 1: Fix t. Suppose we have a j such that $b_j(t) < b_{j-1}(t) < b_{j-2}(t)$. For later times r > t, as long as $b_{j-1}(r) < b_{j-2}(r)$, we have $b_j(r) < b_{j-1}(r)$.

Claim 2: Fix t. Suppose we have a j such that $b_j(t) < \sqrt{2}b_{j-1}(t)$ and $b_{j-1}(t) < \sqrt{2}b_{j-2}(t)$. For later times r > t, as long as $b_{j-1}(r) < \sqrt{2}b_{j-2}(r)$, we have $b_j(r) < \sqrt{2}b_{j-1}(r)$.

Proof of Claim 1:

For a contradiction, suppose there exists a minimal time $t_0 > t$ such that $b_j(t_0) = b_{j-1}(t_0)$ or equivalently $(a_j - a_{j-1})(t_0) = \frac{1}{2}(a_{j-1} - a_{j-2})(t_0)$. First,

$$(b_{j} - b_{j-1})'(t_{0}) = -b_{j}^{2} + b_{j-1}^{2} + 2(b_{j-1}^{2} - b_{j-2}^{2}) + 2^{j} ((\mathcal{L}a)_{j-1} - (\mathcal{L}a)_{j}) + 2^{j-1} ((\mathcal{L}a)_{j-2} - (\mathcal{L}a)_{j-1})$$

$$= 2(b_{j-1}^{2} - b_{j-2}^{2}) - 2^{j} ((\mathcal{L}a)_{j} - (\mathcal{L}a)_{j-1}) + 2^{j-1} ((\mathcal{L}a)_{j-1} - (\mathcal{L}a)_{j-2})$$

where everything from now on will be evaluated at $t = t_0$. By hypothesis, the first term on the right side is negative. If we can show

$$(\mathcal{L}a)_j - (\mathcal{L}a)_{j-1} > \frac{1}{2} ((\mathcal{L}a)_{j-1} - (\mathcal{L}a)_{j-2})$$

we will have a contradiction. From now on, we are left with lengthy algebraic computations. By direct calculation,

(16)
$$(\mathcal{L}a)_{j} - (\mathcal{L}a)_{j-1} = \sum_{n=0}^{j-2} (a_{j} - a_{j-1}) 2^{2\alpha n} + 2(a_{j} - a_{j-1}) 2^{2\alpha(j-1)}$$

$$+ 2^{(1+2\alpha)j} (1 - 2^{-(1+2\alpha)}) \sum_{n=j+1}^{\infty} (a_{j} - a_{n}) 2^{-n}$$

Also, we have

$$\frac{1}{2} [(\mathcal{L}a)_{j-1} - (\mathcal{L}a)_{j-2}] = \sum_{n=0}^{j-3} \frac{1}{2} (a_{j-1} - a_{j-2}) 2^{2\alpha n} + (a_{j-1} - a_{j-2}) 2^{2\alpha(j-2)}
+ 2^{(1+2\alpha)(j-1)-1} (1 - 2^{-(1+2\alpha)}) \sum_{n=j}^{\infty} (a_{j-1} - a_n) 2^{-n}
= \sum_{n=0}^{j-3} (a_j - a_{j-1}) 2^{2\alpha n} + 2(a_j - a_{j-1}) 2^{2\alpha(j-2)}
+ 2^{(1+2\alpha)(j-1)-1} (1 - 2^{-(1+2\alpha)}) \sum_{n=j}^{\infty} (a_{j-1} - a_n) 2^{-n}$$

where we have used $b_k = b_{k-1}$ in the last equality. Then

$$(\mathcal{L}a)_{j} - (\mathcal{L}a)_{j-1} - \frac{1}{2} \left[(\mathcal{L}a)_{j-1} - (\mathcal{L}a)_{j-2} \right] = (2 - 2^{-2\alpha})(a_{j} - a_{j-1})2^{2\alpha(j-1)}$$

$$+ \frac{1}{4}(1 - 2^{-(1+2\alpha)})(a_{j} - a_{j-1})2^{2\alpha(j-1)}$$

$$+ 2^{(1+2\alpha)j}(1 - 2^{-(1+2\alpha)}) \sum_{n=j+1}^{\infty} (a_{j} - a_{n})2^{-n}$$

$$- 2^{(1+2\alpha)(j-1)-1}(1 - 2^{-(1+2\alpha)}) \sum_{n=j+1}^{\infty} (a_{j-1} - a_{n})2^{-n}$$

$$= (2 + 2^{-2} - 2^{-2\alpha} - 2^{-(3+2\alpha)})(a_{j} - a_{j-1})2^{2\alpha(j-1)}$$

$$+ (1 - 2^{-(1+2\alpha)})2^{(1+2\alpha)(j-1)-1} \sum_{n=j+1}^{\infty} (a_{j} - a_{j-1})2^{-n}$$

$$+ (1 - 2^{-(1+2\alpha)})(1 - 2^{-(2+2\alpha)})2^{(1+2\alpha)j} \sum_{n=j+1}^{\infty} (a_{j} - a_{n})2^{-n} .$$

Using that $b_{k+1} < b_k$ for $k \ge j$, we have that $\sum_{n=j+1}^{\infty} (a_j - a_n) 2^{-n} \ge -\frac{2}{3} (a_j - a_{j-1}) 2^{-j}$. Inserting

this into the above estimate we get that the above expression is bounded below by

$$\left(\frac{5}{2} - 5(2^{-(2+2\alpha)}) - \frac{2^{1+2\alpha}}{3}(1 - 2^{-(1+2\alpha)})(1 - 2^{-(2+2\alpha)})\right)(a_j - a_{j-1})2^{2\alpha(j-1)}.$$

Since $0 < \alpha < 1$, one can see that the expression is indeed positive.

<u>Proof of Claim 2:</u> The proof is done similarly to Claim 1, but we will provide the details for completeness. Assume there exists a time t_0 such that $b_j(t_0) = \sqrt{2}b_{j-1}(t_0)$. Also, at this time, $b_k(t_0) \leq \sqrt{2}b_{k-1}(t_0)$ for all $k \neq j$. If we can show $(\sqrt{2}b_{j-1}-b_j)'(t_0) > 0$ then we have a contradiction and the claim is proven. From now on all function will be evaluated at t_0 . By direct calculation,

$$(\sqrt{2}b_{j-1} - b_j)' = (b_j^2 - \sqrt{2}b_{j-1}^2) - 2(b_{j-1}^2 - \sqrt{2}b_{j-2}^2) + [((\mathcal{L}a)_j - (\mathcal{L}a)_{j-1}) - \frac{1}{\sqrt{2}}((\mathcal{L}a)_{j-1} - (\mathcal{L}a)_{j-2})]2^j$$

The contribution on the first line is positive. It suffices to show

$$D_j := ((\mathcal{L}a)_j - (\mathcal{L}a)_{j-1}) - \frac{1}{\sqrt{2}}((\mathcal{L}a)_{j-1} - (\mathcal{L}a)_{j-2}) > 0$$

After similar algebraic manipulations, using that $b_j = \sqrt{2}b_{j-1}$ as in Claim 1 we get

$$D_{j} = (1 - 2^{-(1+2\alpha)}) \left[\left(\frac{2\sqrt{2} + 1}{\sqrt{2}} \right) (a_{j} - a_{j-1}) 2^{2\alpha(j-1)} + (1 - 2^{-(3/2+2\alpha)}) 2^{(1+2\alpha)j} \sum_{n=j+1}^{\infty} (a_{j} - a_{n}) 2^{-n} \right]$$

Using $b_k \le \sqrt{2}b_{k-1}$ for $k \ne j$,

$$\sum_{n=j+1}^{\infty} (a_j - a_n) 2^{-n} \ge -\frac{2}{2\sqrt{2} - 1} (a_j - a_{j-1}) 2^{-j}$$

Inserting this bound into the expression for D_j , we see that D_j is indeed positive for $0 < \alpha < 1/2$.

Thus, by Claims 1 and 2, we know the solution must be of the form given by (13) and (14). At time t=0, we have control over all b_k 's except for b_{K_0+1} . Though we initially have $b_{K_0+1}(0) < b_{K_0}(0)$, there may exist a time t_1 such that $b_{K_0}(t_1) \le b_{K_0+1}(t_1) < \sqrt{2}b_{K_0}(t_1)$. By Claim 2, we know for $t > t_1$, $b_{K_0+1}(t) < \sqrt{2}b_{K_0}(t)$. In addition, by Claim 1, at time t_1 , our solution is of the form (13) and (14) with $K_{t_1} = K_0 + 1$. If there exists no time t_1 , then $K_t = K_0$. Continuing with this line of reasoning, the result follows. \square

Theorem 3.5. (A-priori Hölder-1/2 bound) Let $\{a_k\}$ be a solution to (7) with initial data $\{a_k^0\} \in X^s$, s > 1, that is increasing in k and positive with $b_j^0 < \sqrt{2}b_{j+1}^0$ for $j \ge 0$. Then

$$\sup_{t>0,k} (a_k(t) - a_{k-1}(t))2^{k/2} = \sup_{t>0,k} b_k(t)2^{-k/2} < \infty$$

Proof. By Claim 2 in the proof of the previous Lemma, $b_k(t) < \sqrt{2}b_{k+1}(t)$ for all t and $k \ge 0$. Then for all k, $b_k(t) < 2^{\frac{k-1}{2}}b_1(t) < C2^{k/2}$. \square

4. Regularity for $1/4 < \alpha \le 1/2$ and Blow-up for $\alpha < 1/4$

The previous theorem will allow us to adapt the argument of Theorem 2.5 and prove blow-up for (7) for $0 < \alpha < 1/4$.

Theorem 4.1. Let $0 < \alpha < 1/4$ Suppose $\{a_k^0\}$ is increasing in k and is non-negative. Then there exists initial datum in X^s , s > 1, such that solutions blow-up in finite time.

Proof.

As before we let

$$J(t) = \sum_{k=0}^{\infty} (a_{\infty}(t) - a_k(t)) 2^{k\delta}.$$

for $0 < \delta < 1 - 4\alpha$ where $a_{\infty}(t) = \lim_{k \to \infty} a_k(t)$. Then computing and using Lemma 2.6 we get

$$\frac{d}{dt}J(t) = \sum_{k=1}^{\infty} (a_k(t) - a_{k-1}(t))^2 2^{k(\delta+1)} - \sum_{k=1}^{\infty} ((\mathcal{L}a)_{\infty} - (\mathcal{L}a)_k) 2^{k\delta}
\geq C_0 \sum_{k=1}^{\infty} (a_{\infty}(t) - a_k(t))^2 2^{k(\delta+1)} - \sum_{k=1}^{\infty} ((\mathcal{L}a)_{\infty} - (\mathcal{L}a)_k) 2^{k\delta}$$

where

$$(\mathcal{L}a)_{\infty} = \lim_{k \to \infty} (\mathcal{L}a)_k = \sum_{n=0}^{\infty} (a_{\infty} - a_n) 2^{2\alpha n}$$

Now,

$$(\mathcal{L}a)_{\infty} - (\mathcal{L}a)_k = \sum_{n=0}^{k-1} (a_{\infty} - a_k) 2^{2\alpha n} + \sum_{n=k}^{\infty} (a_{\infty} - a_n) 2^{2\alpha n} - \sum_{n=k+1}^{\infty} (a_k - a_n) 2^{(1+2\alpha)k} 2^{-n}$$

We analyze each of the above three terms separately.

$$\sum_{n=0}^{k-1} (a_{\infty} - a_k) 2^{2\alpha n} \le C(a_{\infty} - a_k) 2^{2\alpha k}$$

By Theorem 3.5,

$$\sum_{n=k}^{\infty} (a_{\infty} - a_n) 2^{2\alpha n} \le C \sum_{n=k}^{\infty} 2^{(2\alpha - 1/2)n} \le C 2^{(2\alpha - 1/2)(k-1)}$$

For the last term we have

$$\sum_{n=k+1}^{\infty} (a_n - a_k) 2^{(1+2\alpha)k} 2^{-n} \le (a_{\infty} - a_k) 2^{(1+2\alpha)k} \sum_{n=k+1}^{\infty} 2^{-n} \le C(a_{\infty} - a_k) 2^{2\alpha k}.$$

Then using $\alpha < 1/4$ and $0 < \delta < 1 - 4\alpha$,

$$\sum_{k=1}^{\infty} ((\mathcal{L}a)_{\infty} - (\mathcal{L}a)_k) \, 2^{k\delta} \le C_1 + C_2 \sum_{k=1}^{\infty} (a_{\infty} - a_k) 2^{(2\alpha + \delta)k}$$

For every $\epsilon > 0$, by an application of Holder's inequality and using that $\alpha < 1/4$,

$$\sum_{k=1}^{\infty} (a_{\infty} - a_k) 2^{(2\alpha + \delta)k} \le \epsilon \sum_{k=1}^{\infty} (a_{\infty} - a_k)^2 2^{(\delta + 1)k} + C_{\epsilon} ||a||_{\ell^{\infty}}$$

Given these bounds and following the proof of Theorem 2.6 we can show

$$J'(t) \ge C(\delta)J(t)^2 - C(1 + ||a||_{\ell^{\infty}}).$$

By choosing initial data appropriately, the inequality above leads to blow-up. \Box

Now, we will move towards proving regularity for $1/4 < \alpha < 1/2$.

Theorem 4.2. Let $1/4 < \alpha < 1/2$. Suppose we have the same hypotheses as in Theorem 3.5. Then solutions exist for all time, in particular, for s > 1

$$\sup_{t>0} \|a(t)\|_{X^s} < \infty.$$

Proof.

After a computation,

$$(b_{k,s})'(t) = -b_k \cdot b_{k,s}(t) + 2^s b_{k-1} \cdot b_{k-1,s}(t) - ((\mathcal{L}a)_k - (\mathcal{L}a)_{k-1})(t) 2^{sk}$$

From Theorem 3.5, $b_{k-1}(t) \leq C_1 2^{k/2}$ for all $t \in [t_0, t_1]$ where the constant C_1 is independent of the time interval. Then there exists K' such that if $k \geq K'$ and $b_{k,s}(t) > c_s ||a||_{X^s}(t)$ where c_s is the constant from Lemma 2.1,

$$(b_{k,s})'(t) \le b_{k,s} \left(C_1 2^{k/2+s} - C(\alpha)(2^{2\alpha k} - 2^{2\alpha}) \right) < 0.$$

In the above inequality, we use $0 < \alpha < 1/4$. This implies that for any T > 0, $\lim_{t \to T, k \to \infty} b_{k,s}(t) \neq \infty$. Therefore, the solution exists globally in time. \square

Remark: It is unclear whether the hypothesis that $b_j^0 \leq \sqrt{2}b_{j+1}^0$ can be weakened as it is crucial to the proof of the a-priori bound. We conjecture that any reasonable initial data will eventually satisfy such a condition at least for large enough indexes.

Summary: We have shown that when $\alpha > 1/4$, under mild assumptions on the initial data, it is possible to have global solutions: the dissipation terms win over the nonlinearity. In order to have blow-up, energy initially present in the lower modes must reach the higher modes. By making use of a-priori Holder-type control and an estimate on the dissipative terms, we have shown that such an energy transfer must stop at some time.

It is important to note that our conditions on the initial data are not a condition of "smallness" with regard to some norm. Our results of global regularity can hold for large initial data. Also, one can easily find initial data for which our model blows up for $0 < \alpha < 1/4$ and is regular for $\alpha > 1/4$.

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