能隙方程的计算

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我们在这里处理能隙方程,这是一个典型的非线性方程,即自治方程,其中问题的难点并不是解方程本身,而是 其中的无量纲化问题.

1 BCS 能隙方程

$$\frac{1}{N(0)V} = \int_0^{\hbar\omega_D} \frac{\tanh\frac{\beta\sqrt{\xi^2 + \Delta^2}}{2}}{\beta\sqrt{\xi^2 + \Delta^2}} d\xi,\tag{1}$$

其中 N(0) 是能态密度,V 是相互吸引作用强度, ω_D 是德拜频率, Δ 是能隙函数.

下面开始做准备工作. 令

$$\delta = \frac{\Delta(T)}{\Delta(0)},\tag{2}$$

$$\tau = \frac{T}{T_c},\tag{3}$$

$$\eta = \frac{1}{N(0)V}. (4)$$

1.1 求 T_c

$$\eta = \int_0^{\hbar\omega_D} \tanh \frac{\xi}{2k_B T_c} \frac{d\xi}{\xi}
= \int_0^{\kappa} \frac{\tanh z}{z} dz \qquad \left(\kappa = \frac{\hbar\omega_D}{2k_B T_c}\right)
= \ln z \tanh z \Big|_0^{\kappa} - \int_0^{\kappa} \frac{\ln z}{\cosh^2 z} dz
\approx \ln \kappa + \ln \frac{4e^{\gamma}}{\pi}
= \ln \frac{4e^{\gamma} \kappa}{\pi},$$
(5)

其中用到了欧拉积分公式

$$\int_0^\infty \frac{\ln x}{\cosh^2 x} dx = -\ln \frac{4e^\gamma}{\pi} \tag{6}$$

其中 γ 为欧拉常数,并有条件 $\hbar\omega_D \gg k_B T_c$. 因此,

$$\frac{\hbar\omega_D}{k_B T_c} = \frac{2e^{\gamma}}{\pi} e^{-\eta} \approx 1.134e^{-\eta}.$$
 (7)

1.2 求零温能隙

$$\eta = \int_0^{\hbar\omega_D} \frac{d\xi}{\sqrt{\xi^2 + \Delta_0^2}},\tag{8}$$

得出

$$\frac{\hbar\omega_D}{\Delta_0} = \sinh\eta,\tag{9}$$

$$\Delta_0 = 2\pi k_B T_c e^{-\gamma} \approx 1.76 k_B T_c,\tag{10}$$

其中用到了 $\hbar\omega_D \gg k_B T_c$.

1.3 无量纲化能隙方程

$$\eta = \int_{0}^{\delta^{-1} \sinh \eta} \tanh \left(\frac{\Delta \sqrt{1+z^{2}}}{2k_{B}T} \right) \frac{dz}{\sqrt{1+z^{2}}} \qquad \left(z = \frac{\xi}{\Delta(T)} \right) \\
= \int_{0}^{\delta^{-1} \sinh \eta} \tanh \left(\frac{\Delta}{\Delta_{0}} \frac{\Delta_{0}}{2k_{B}T} \sqrt{1+z^{2}} \right) \frac{dz}{\sqrt{1+z^{2}}} \\
= \int_{0}^{\delta^{-1} \sinh \eta} \tanh \left(\delta \frac{1.76k_{B}T_{c}}{2k_{B}T} \sqrt{1+z^{2}} \right) \frac{dz}{\sqrt{1+z^{2}}} \\
= \int_{0}^{\delta^{-1} \sinh \eta} \tanh \left(\frac{\delta}{\tau} 0.882\sqrt{1+z^{2}} \right) \frac{dz}{\sqrt{1+z^{2}}}.$$
(11)

1.4 Matlab 代码

```
n = 0.35;
D = zeros(40,1);
t = 0;
for n = 1:40

d0 = [rand];
F = @(d)[
n-integral(@(z) tanh(0.882*d.*sqrt(1+z.^2)./t)./...
sqrt(1+z.^2),0,sinh(n)./d)
];
d = fsolve(F,d0);
t = t + 0.025;
T(n,:) = [t];
D(n,:) = [d];
end
plot(T,D,'p')
```

1.5 一些说明

在解物理中的方程时,无量纲化是很重要的,有时候忘了无量纲化,你选择的参数范围可能是非物理的,所画出来的图自然也是非物理的,即没有任何规律可言.

2 Slave-boson mean-field self-consistent equations

这其实是求解非线性方程组,其中的非线性方程组是从 salve-boson 平均场方法通过求解 t-J 模型得到的,下面的方程取自 Physical Review B **70**,054504(2004):

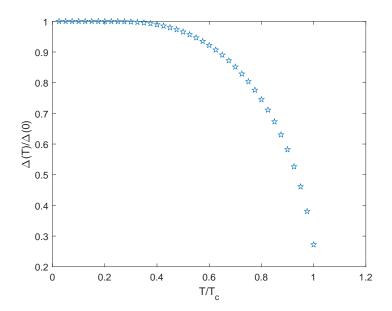


Figure 1: 能隙随温度的变化

$$\frac{1}{J} = \frac{3}{8} \int \frac{d^2k}{(2\pi)^2} \frac{(\cos k_x - \cos k_y)^2}{E_k},\tag{12}$$

$$\chi = -\frac{1}{4} \int \frac{d^2k}{(2\pi)^2} \left(\frac{\xi_{\mathbf{k}}}{E_{\mathbf{k}}}\right) \left(\cos k_x + \cos k_y\right),\tag{13}$$

$$x = \int \frac{d^2k}{(2\pi)^2} \left(\frac{\xi_k}{E_k}\right),\tag{14}$$

其中 x 为掺杂浓度.

$$\xi_{\mathbf{k}} = \epsilon_{\mathbf{k}} - \mu$$

$$E_{\mathbf{k}} = \sqrt{\xi_{\mathbf{k}}^2 + \Delta_{\mathbf{k}}^2}$$

$$\epsilon_{\mathbf{k}} = -2\left(xt + \frac{3J\chi}{4}\right)\left(\cos k_x + \cos k_y\right) + 4xt'\cos k_x\cos k_y, t' = t/4$$

$$\Delta_{\mathbf{k}} = \frac{\Delta^{sb}\left(\cos k_x - \cos k_y\right)}{2} = \frac{3J|\Delta|\left(\cos k_x - \cos k_y\right)}{2}$$

2.1 以 J 标度使方程无量纲化

$$\tilde{\epsilon_k} = \frac{\epsilon_k}{J} = -2\left(x\frac{t}{J} + \frac{3\chi}{4}\right)\left(\cos k_x + \cos k_y\right) + x\frac{t}{J}\cos k_x\cos k_y$$

$$\tilde{\mu} = \frac{\mu}{J}$$

$$\tilde{\Delta_k} = \frac{\Delta_k}{J} = 1.5|\Delta|\left(\cos k_x - \cos k_y\right)$$

$$\tilde{\xi_k} = \frac{\epsilon_k}{J} = -2\left(x\frac{t}{J} + \frac{3\chi}{4}\right)\left(\cos k_x + \cos k_y\right) + x\frac{t}{J}\cos k_x\cos k_y - \tilde{\mu}$$

$$\tilde{E_k} = \sqrt{\left\{2\left(x\frac{t}{J} + \frac{3\chi}{4}\right)\left(\cos k_x + \cos k_y\right) - x\frac{t}{J}\cos k_x\cos k_y + \tilde{\mu}\right\}^2 + 2.25|\Delta|^2\left(\cos k_x - \cos k_y\right)^2}$$

因此当 t/J=3 时,我们有下面这一组无量纲化的方程:

$$0 = 1 - \frac{3}{8} \int \frac{d^2k}{(2\pi)^2} \frac{\left(\cos k_x - \cos k_y\right)^2}{\sqrt{\left\{\left(6x + \frac{3\chi}{2}\right)\left(\cos k_x + \cos k_y\right) - 3x\cos k_x\cos k_y + \tilde{\mu}\right\}^2 + 2.25|\Delta|^2 \left(\cos k_x - \cos k_y\right)^2}},$$
 (15)

$$0 = \chi - \frac{1}{4} \int \frac{d^2k}{(2\pi)^2} \left(\frac{\left(6x + \frac{3\chi}{2}\right) \left(\cos k_x + \cos k_y\right) - 3x \cos k_x \cos k_y + \tilde{\mu}}{\sqrt{\left\{\left(6x + \frac{3\chi}{2}\right) \left(\cos k_x + \cos k_y\right) - 3x \cos k_x \cos k_y + \tilde{\mu}\right\}^2 + 2.25|\Delta|^2 \left(\cos k_x - \cos k_y\right)^2}} \right) \left(\cos k_x + \cos k_y\right),$$

$$(16)$$

$$0 = x + \int \frac{d^2k}{(2\pi)^2} \frac{\left(6x + \frac{3\chi}{2}\right)\left(\cos k_x + \cos k_y\right) - 3x\cos k_x\cos k_y + \tilde{\mu}}{\sqrt{\left\{\left(6x + \frac{3\chi}{2}\right)\left(\cos k_x + \cos k_y\right) - 3x\cos k_x\cos k_y + \tilde{\mu}\right\}^2 + 2.25|\Delta|^2\left(\cos k_x - \cos k_y\right)^2}}.$$
 (17)

2.1.1 以 t 标度使方程无量纲化

$$\begin{split} \tilde{\epsilon_{k}} &= \frac{\epsilon_{k}}{t} = -2\left(x + \frac{3J\chi}{4t}\right)\left(\cos k_{x} + \cos k_{y}\right) + x\cos k_{x}\cos k_{y} \\ \tilde{\mu} &= \frac{\mu}{t} \\ \tilde{\Delta_{k}} &= \frac{\Delta_{k}}{t} = \frac{3J|\Delta|\left(\cos k_{x} - \cos k_{y}\right)}{2t} \end{split}$$

$$\tilde{E}_{k} = \sqrt{\left\{2\left(x + \frac{3J\chi}{4t}\right)\left(\cos k_x + \cos k_y\right) - x\cos k_x\cos k_y + \tilde{\mu}\right\}^2 + 2.25\left(\frac{J}{t}\right)^2 |\Delta|^2 \left(\cos k_x - \cos k_y\right)^2}$$

因此当 t/J=3 时,我们有下面这一组无量纲化的方程:

$$0 = 1 - \frac{1}{8} \int \frac{d^2k}{(2\pi)^2} \frac{\left(\cos k_x - \cos k_y\right)^2}{\sqrt{\left\{\left(2x + \frac{\chi}{2}\right)\left(\cos k_x + \cos k_y\right) - x\cos k_x\cos k_y + \tilde{\mu}\right\}^2 + 0.25|\Delta|^2 \left(\cos k_x - \cos k_y\right)^2}},$$
 (18)

$$0 = \chi - \frac{1}{4} \int \frac{d^2k}{(2\pi)^2} \left(\frac{\left(2x + \frac{\chi}{2}\right) \left(\cos k_x + \cos k_y\right) - x \cos k_x \cos k_y + \tilde{\mu}}{\sqrt{\left\{\left(2x + \frac{\chi}{2}\right) \left(\cos k_x + \cos k_y\right) - x \cos k_x \cos k_y + \tilde{\mu}\right\}^2 + 0.25|\Delta|^2 \left(\cos k_x - \cos k_y\right)^2}} \right) \left(\cos k_x + \cos k_y\right),$$

$$(19)$$

$$0 = x + \int \frac{d^2k}{(2\pi)^2} \frac{\left(2x + \frac{\chi}{2}\right)(\cos k_x + \cos k_y) - x\cos k_x\cos k_y + \tilde{\mu}}{\sqrt{\left\{\left(2x + \frac{\chi}{2}\right)(\cos k_x + \cos k_y) - x\cos k_x\cos k_y + \tilde{\mu}\right\}^2 + 0.25|\Delta|^2 \left(\cos k_x - \cos k_y\right)^2}}.$$
 (20)

2.2 Matlab 代码

P = zeros(20,3);

```
 \begin{aligned} &d=0.01;\\ &p0=\left[-0.01; rand; 1\right];\\ &for\ n=\ 1:20\\ &F=@(p)[\\ &1-((1/8)*(1/(4*pi*pi))*(integral2(@(x,y)\ (cos(x)-cos(y)).^2./...\\ &sqrt\left(((2*d+0.5*p(2)).*(cos(x)+cos(y))-d*cos(x).*cos(y)+p(1)).^2+...\\ &0.25*abs(p(3)).^2.*(cos(x)-cos(y)).^2),-pi,pi,-pi,pi)));\\ &p(2)-((1/4)*(1/(4*pi*pi))*(integral2(@(x,y)\ ...\\ &((2*d+0.5*p(2)).*(cos(x)+cos(y))-d*cos(x).*cos(y)+p(1)).*...\\ &(cos(x)+cos(y))./...\\ &sqrt\left(((2*d+0.5*p(2)).*(cos(x)+cos(y))-d*cos(x).*cos(y)+p(1)).^2+...\\ &0.25*abs(p(3)).^2.*(cos(x)-cos(y)).^2),-pi,pi,-pi,pi)));\\ &d+((1/(4*pi*pi))*(integral2(@(x,y)\ ...) \end{aligned}
```

```
 \begin{split} &((2*d+0.5*p(2)).*(\cos(x)+\cos(y))-d*\cos(x).*\cos(y)+p(1))./... \\ & \text{sqrt}\left(((2*d+0.5*p(2)).*(\cos(x)+\cos(y))-d*\cos(x).*\cos(y)+p(1)).^2+... \\ & 0.25*abs(p(3)).^2.*(\cos(x)-\cos(y)).^2),-pi,pi,-pi,pi))) \\ &]; \\ & p = fsolve(F,p0); \\ & d = d+0.02; \\ & P(n,:) = [p(1),p(2),p(3)]; \\ & p0 = P(n,:); \\ & end \\ & m = 0.01:0.02:0.4; \\ & plot(m,P(:,3),'p') \\ & hold on \\ & plot(m,P(:,2),'*') \\ \end{split}
```

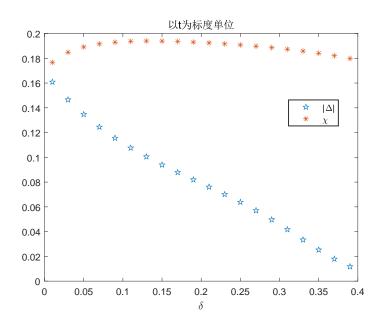


Figure 2: 能隙随掺杂的变化