

Coupled ODEs

$$\dot{\mathbf{y}} = \mathbf{A}\mathbf{y} + \mathbf{r}$$

$$\ddot{y} + p(t)y + q(t) = r(t) \longrightarrow \begin{pmatrix} y \\ \dot{y} \end{pmatrix} = \dot{\mathbf{y}} = \begin{pmatrix} 0 & 1 \\ -q(t) & -p(t) \end{pmatrix} \mathbf{y} + \begin{pmatrix} 0 \\ r(t) \end{pmatrix}$$

The eigenvalues of \mathbf{A} found by solving $\det(\mathbf{A} - \lambda\mathbf{I}) = 0$ gives the same characteristic equation for the uncoupled system. General solutions for a second-order system with real constant coefficients, with eigenvalues λ_1, λ_2 and eigenvectors $\mathbf{x}_1, \mathbf{x}_2$, are given by:

$$\lambda_1 \neq \lambda_2: \mathbf{y} = c_1 \mathbf{x}_1 e^{\lambda_1 t} + c_2 \mathbf{x}_2 e^{\lambda_2 t}$$

$$\lambda_1, \lambda_2 \in \mathbb{C}: \mathbf{y} = c_1 \operatorname{Re}(\mathbf{x}_1 e^{\lambda_1 t}) + c_2 \operatorname{Im}(\mathbf{x}_1 e^{\lambda_1 t})$$

$$\lambda_1 = \lambda_2: \mathbf{y} = c_1 \mathbf{x} e^{\lambda t} + c_2 (t\mathbf{x} + \mathbf{p}) e^{\lambda t}, \text{ where } (\mathbf{A} - \lambda\mathbf{I})\mathbf{p} = \mathbf{x}.$$

Phase Portraits of 2D Autonomous Systems

Let $\dot{\mathbf{y}} = \mathbf{f}(\mathbf{y})$ with $\mathbf{y} \in \mathbb{R}^2$, $\mathbf{f}: \mathbb{R}^2 \rightarrow \mathbb{R}^2$, where \mathbf{f} does not depend on t . \mathbf{f} defines a direction field in the y_1 - y_2 plane, and a trajectory of the system is a curve following the field. \mathbf{f} being continuous with continuous partials implies existence and uniqueness for these trajectories, so they do not cross. By the chain rule, the slope of a trajectory is

$$\frac{dy_2}{dy_1} = \frac{\dot{y}_2}{\dot{y}_1} = \frac{f_2}{f_1}.$$

A nullcline is a curve in phase space where the slope of the trajectory is either 0 or ∞ . Set $\dot{y}_2 = 0$ for horizontal nullclines and $\dot{y}_1 = 0$ for vertical nullclines, and solve for y_2 . Solving for the trajectories themselves is possible in specific cases, usually when the ODE in y_2 from the slope formula is separable.

For the lines in phase space in the direction of the eigenvectors (called eigendirections), there are straight line solutions either flowing towards or away from the origin, determined by the sign of the corresponding eigenvalue (− stable, + unstable).

$\lambda_1 \neq \lambda_2$ and both have the same sign:

- Negative λ gives a *stable improper node*, where trajectories curve like a parabola and approach the eigendirection with smallest absolute value, since that exponential term dominates at large t .
- Positive λ gives an *unstable improper node*, where trajectories approach eigendirection with largest absolute value.

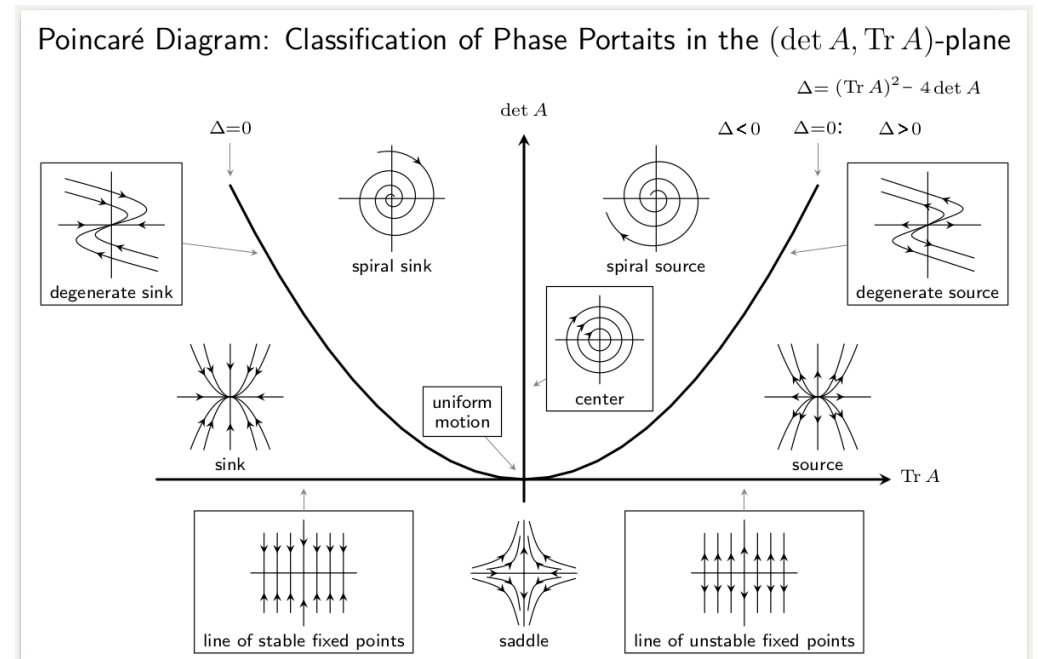
$\lambda_1 > 0, \lambda_2 < 0$:

- This gives a *saddle point*. Trajectories come in parallel to the stable eigendirection then turn around and become parallel to the unstable direction.

$\lambda_1 = \lambda_2$:

- If there are 2 linearly independent eigenvectors (eigenspace of λ has dimension 2), all trajectories travel in straight lines either towards or away from the origin, called a *proper node*. For a 2D system, this is only possible if \mathbf{A} is some multiple of the identity matrix, i.e. $\mathbf{A} = k\mathbf{I}$ for $k \in \mathbb{R}$.
 - If there is only 1 eigenvector, trajectories come tangent to the eigendirection near the origin and do a U-turn away, called a *degenerate node*.
- $\lambda = \alpha \pm \beta i$:
- If $\alpha = 0$, trajectories are ellipses around the origin, and the origin is called a *centre*. To determine the direction of flow, consider the sign of \dot{y}_1 at a point in the positive quadrant.
 - If $\alpha \neq 0$, trajectories trace out a logarithmic spiral $r(t) = e^{\alpha t}$ shown by converting to polar coordinates. This is called a *spiral*. Plotting nullclines helps with sketching trajectories.

The following stability chart allows for the system to be classified given values of $\operatorname{tr} \mathbf{A}$ and $\det \mathbf{A}$ [1]. The discriminant Δ is defined as $\Delta = (\operatorname{tr} \mathbf{A})^2 - 4 \det \mathbf{A}$.



Non-Linear or Inhomogeneous Systems

A system $\dot{\mathbf{y}} = \mathbf{f}(\mathbf{y})$ has a critical point at \mathbf{y}^* if $\mathbf{f}(\mathbf{y}^*) = \mathbf{0}$. Linear systems always have one at $\mathbf{0}$, and only one unless $\det \mathbf{A} = 0$. An inhomogeneous linear system $\dot{\mathbf{y}} = \mathbf{A}\mathbf{y} + \mathbf{p}$ is a linear translation away from a homogeneous system. Shifting the critical point at \mathbf{p} to $\mathbf{0}$ via the change of variables $\mathbf{x} = \mathbf{y} - \mathbf{p}$ gives the homogeneous $\dot{\mathbf{x}} = \mathbf{A}\mathbf{x}$, so techniques discussed earlier are applicable.

For a nonlinear system $\dot{\mathbf{y}} = \mathbf{f}(\mathbf{y})$, the nature of the system around the critical point is approximated via the Jacobian matrix \mathbf{Df} . The stability diagram can be used to classify these critical points by evaluating \mathbf{Df} at each \mathbf{y}^* . The linearised system near \mathbf{y}^* is then

$$\dot{\mathbf{y}} = \mathbf{Df}|_{\mathbf{y}^*}(\mathbf{y} - \mathbf{y}^*), \quad \mathbf{Df} = \begin{pmatrix} \frac{\partial f_1}{\partial y_1} & \frac{\partial f_1}{\partial y_2} \\ \frac{\partial f_2}{\partial y_1} & \frac{\partial f_2}{\partial y_2} \end{pmatrix}.$$

Special Functions

The *Dirac Delta function* $\delta(t)$ is an “impulse” with infinite value at $t = 0$ and zero value everywhere else, satisfying $\int_0^\infty \delta(t) dt = 1$. $\delta(t - a)$ represents an impulse at $t = a$.

The *unit step function* $u(t)$ is defined by

$$u(t) = \begin{cases} 0, & t < 0 \\ 1, & t \geq 0. \end{cases}$$

To write a piecewise-continuous function in terms of u , if $f(t)$ changes from $g_1(t)$ to $g_2(t)$ at $t = c$, add a $u(t - c)(g_2(t) - g_1(t))$ term to the new expression of f [2].

The *Gamma function* $\Gamma(z)$ generalises the factorial by $\Gamma(n + 1) = n!$ for $n \in \mathbb{N}$, and is defined for all positive z by

$$\Gamma(z) = \int_0^\infty e^{-x} x^{z-1} dx,$$

In this course the *error function* is defined as

$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-u^2} du.$$

It is odd and monotone increasing. The following formulas may be useful for computation:

$$\int_0^\infty e^{-u^2} du = \int_{-\infty}^0 e^{-u^2} du = \frac{\sqrt{\pi}}{2}, \quad \frac{1}{\sqrt{\pi}} \int_a^b e^{-u^2} du = \operatorname{erf}(b) - \operatorname{erf}(a).$$

For the heat equation with parameter c , the *Gaussian function* $G(x, t)$ is defined as

$$G(x, t) = \frac{1}{\sqrt{4\pi c^2 t}} \exp\left(\frac{-x^2}{4c^2 t}\right).$$

$\int_{-\infty}^\infty G(x, t) dx = 1$ for $t > 0$, $\lim_{t \rightarrow 0^+} G(x, t) = 0$ for $x \neq 0$ and $\lim_{t \rightarrow 0^+} G(0, t) = +\infty$.

Laplace Transforms

The following formulas are given in the formula sheet for the exam, but have been added here for completeness.

$$\mathcal{L}\{f(t)\} = F(s) = \int_0^\infty f(t)e^{-st} dt$$

$f(t)$	$F(s)$	$f(t)$	$F(s)$
c	$\frac{c}{s}$	$ag(t) + bh(t)$	$aG(s) + bH(s)$
t^n	$\frac{n!}{s^{n+1}}$	$e^{at}f(t)$	$F(s - a)$
t^a	$\frac{\Gamma(a + 1)}{s^{a+1}}$	$f(t - k)u(t - k)$	$e^{-ks}F(s)$
$\cos \alpha t$	$\frac{s}{s^2 + \alpha^2}$	$f'(t)$	$sF(s) - f(0)$
$\sin \alpha t$	$\frac{\alpha}{s^2 + \alpha^2}$	$f''(t)$	$s^2F(s) - sf(0) - f'(0)$
$\delta(t - k)$	e^{-ks}	$tf(t)$	$-F'(s)$

The convolution theorem states that:

$$\mathcal{L}^{-1}\{F(s)G(s)\} = \int_0^t f(\tau)g(t - \tau) d\tau = (f * g)(t).$$

If $f(t)$ has period p , then the Laplace integral can be simplified:

$$\mathcal{L}\{f(t)\} = \frac{1}{1 - e^{-sp}} \int_0^p f(t)e^{-st} dt.$$

The following formulas are not included in the formula sheet, but may be helpful to know.

$$\begin{aligned} \mathcal{L}^{-1}\left\{\frac{s}{(s^2 + \alpha^2)^2}\right\} &= \frac{1}{2\alpha} t \sin \alpha t \\ \mathcal{L}\{f(t)u(t - k)\} &= e^{-ks} \mathcal{L}\{f(t + k)\} \\ \mathcal{L}\{f^{(n)}(t)\} &= s^n F(s) - \sum_{k=1}^n s^{n-k} f^{(k-1)}(0) \\ \mathcal{L}\{t^n f(t)\} &= (-1)^n F^{(n)}(s). \end{aligned} \quad [2]$$

Differential Operators and Change of Variables

A homogeneous linear differential equation has the form $L(u) = 0$ where L is a *linear differential operator* satisfying $L(u + v) = L(u) + L(v)$ and $L(cu) = cL(u)$ for $c \in \mathbb{R}$ and functions u, v . This can be used to verify a DE is linear using the linearity of the derivative. A common differential operator is the del operator ∇ , defined in 3 dimensions by $\nabla = (\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z})$. It can be applied to a scalar field f to get f 's gradient field, or be dotted or crossed with a vector field \mathbf{v} to get its *divergence* or *curl*:

$$\text{grad } f = \nabla f = \left(\frac{\partial f}{\partial x}, \frac{\partial f}{\partial y}, \frac{\partial f}{\partial z} \right), \quad \text{div } \mathbf{v} = \nabla \cdot \mathbf{v} = \frac{\partial v_1}{\partial x} + \frac{\partial v_2}{\partial y} + \frac{\partial v_3}{\partial z},$$

$$\text{curl } \mathbf{v} = \nabla \times \mathbf{v} = \begin{vmatrix} \hat{\mathbf{i}} & \hat{\mathbf{j}} & \hat{\mathbf{k}} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ v_1 & v_2 & v_3 \end{vmatrix}.$$

Additionally, the *Laplacian* is defined as $\nabla^2 = \nabla \cdot \nabla$. For scalar fields,

$$\nabla^2 f = \nabla \cdot \nabla f = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2},$$

and for a vector field \mathbf{v} in cartesian coordinates, the Laplacian is just defined elementwise, so $\nabla^2 \mathbf{v} = (\nabla^2 v_1, \nabla^2 v_2, \nabla^2 v_3)$. Equivalently, it can be defined as

$$\nabla^2 \mathbf{v} = \nabla(\nabla \cdot \mathbf{v}) - \nabla \times (\nabla \times \mathbf{v}).$$

The above formulas only apply for *cartesian coordinates*; other coordinate systems have different expressions. Importantly, the Laplacian in polar coordinates is given by

$$\nabla^2 = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2}.$$

Changing variables that a PDE is expressed in may simplify the equation and give different types of solutions not apparent from the initial coordinate system. The goal is to express derivatives in the *initial* coordinates in terms of derivatives in the *new* coordinates. If $f(x, y)$ is changed to $f(u, v)$ via $u(x, y)$ and $v(x, y)$, then by the chain rule,

$$\frac{\partial f}{\partial x} = \frac{\partial f}{\partial u} \frac{\partial u}{\partial x} + \frac{\partial f}{\partial v} \frac{\partial v}{\partial x},$$

$$\frac{\partial f}{\partial y} = \frac{\partial f}{\partial u} \frac{\partial u}{\partial y} + \frac{\partial f}{\partial v} \frac{\partial v}{\partial y}.$$

Note that f can be removed from the above equations to get a relationship purely involving the differential operators $\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial u}$ and $\frac{\partial}{\partial v}$. This can then be substituted into the original PDE to get a new equation.

An example is with the PDE $u_t + cu_x = 0$ for $u(x, t)$. Changing variables to $x' = x - ct$ and $t' = t$ allows the equation to be written as $u_{t'} = 0$, thus $u(x', t') = f(x')$ and $u(x, t) = f(x - ct)$ for differentiable f . The solution is defined on a plane in x - t space.

Fourier Series

If a function $f(x)$ has period $2L$, i.e. $f(x + 2L) = f(x)$ for all x , then it can be written as a *Fourier series*:

$$f(x) = a_0 + \sum_{n=1}^{\infty} \left(a_n \cos\left(\frac{n\pi x}{L}\right) + b_n \sin\left(\frac{n\pi x}{L}\right) \right),$$

with Fourier coefficients given by

$$a_n = \frac{1}{L} \int_{-L}^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx, \quad b_n = \frac{1}{L} \int_{-L}^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx,$$

$$a_0 = \frac{1}{2L} \int_{-L}^L f(x) dx.$$

In the case that f is *even* ($f(-x) = f(x)$) or *odd* ($f(-x) = -f(x)$), the formulas for the coefficients can be simplified, only changing the fraction and integral bounds:

$$a_0 = \frac{1}{L} \int_0^L (\dots) dx, \quad a_n = \frac{2}{L} \int_0^L (\dots) dx, \quad b_n = \frac{2}{L} \int_0^L (\dots) dx.$$

Odd f makes $a_n = 0$ for all n , and even f makes $b_n = 0$ for all n .

These simplified formulas can also be used when extending f only defined on $[0, L]$ to be odd or even periodic over period $2L$. This is called a *half-range expansion* of f ; odd periodic extension makes $a_n = 0$ and even period extension makes $b_n = 0$ for all n .

The following identities may be useful for computation.

$$2 \sin A \sin B = \cos(A - B) - \cos(A + B)$$

$$2 \cos A \cos B = \cos(A - B) + \cos(A + B)$$

$$2 \sin A \cos B = \sin(A - B) + \sin(A + B)$$

If f has period $2L$ and is piecewise continuous within that period, then its Fourier series converges to $\frac{1}{2}(f(x^+) + f(x^-))$ at every x where f has left and right derivatives.

To derive formulas for the coefficients, define functions f_1 and f_2 to be orthogonal if $\int_{-L}^L f_1(x) f_2(x) dx = 0$. It can be shown that all functions $\{1, \cos(\frac{n\pi x}{L}), \sin(\frac{n\pi x}{L})\}$ are mutually orthogonal by integration by parts. Both sides of the Fourier Series equation can then be integrated and solved for the coefficients, where orthogonal terms vanish. This orthogonality is further discussed in MATH2001.

Important PDEs

Name	1D	2D	General
Wave Equation	$u_{tt} = c^2 u_{xx}$	$u_{tt} = c^2(u_{xx} + u_{yy})$	$u_{tt} = c^2 \nabla^2 u$
Heat Equation	$u_t = c^2 u_{xx}$	$u_t = c^2(u_{xx} + u_{yy})$	$u_t = c^2 \nabla^2 u$
Laplace's Equation	$u_{xx} = 0$	$u_{xx} + u_{yy} = 0$	$\nabla^2 u = 0$

These equations are all linear and homogeneous.

1D Wave, Infinite Domain

The general solution to the 1D wave equation $u_{tt} = c^2 u_{xx}$ can be found with the change of variables $v = x + ct$ and $w = x - ct$, giving $utt - c^2 u_{xx} = u_{vw} = 0$. This results in $u(x, t) = F(x + ct) + G(x - ct)$ for arbitrary twice-differentiable F, G .

From here, consider an infinite string with $u(x, 0) = f(x)$ as initial deflection and $u_t(x, 0) = g(x)$ as initial velocity profile. From the general solution derived earlier, taking derivatives gives $u_t = cF'(x + ct) - cG'(x - ct)$ where the prime denotes a derivative with respect to the entire argument. After substituting initial conditions and solving for F and G , the *D'Alembert's solution* to the IVP is

$$u(x, t) = \frac{1}{2} (f(x + ct) + f(x - ct)) + \frac{1}{2C} \int_{x-ct}^{x+ct} g(s) ds.$$

A *plucked* string is modelled with $g(x) = 0$ and $f(x)$ arbitrary. A *struck* string is modelled with $f(x) = 0$ and $g(x) = Ae^{-x^2}$ for $A \in \mathbb{R}$. The resulting integral can be written in terms of the error function,

$$u(x, t) = B \operatorname{erf}(x + ct) - B \operatorname{erf}(x - ct), \quad B = \frac{A\sqrt{\pi}}{4c}.$$

1D Wave, Finite Domain

A finite string is modelled with $u_{tt} - c^2 u_{xx} = 0$ on $x \in [0, L]$ for $t \geq 0$, and again has initial deflection $u(x, 0) = f(x)$ and initial velocity profile $u_t(x, 0) = g(x)$. Boundary conditions are $u(0, t) = u(L, t) = 0$.

By Fourier's method (discussed in a later section), the non-trivial solution to this IVP is

$$u(x, t) = \sum_{n=1}^{\infty} \left(A_n \cos\left(\frac{n\pi ct}{L}\right) + B_n \sin\left(\frac{n\pi ct}{L}\right) \right) \sin\left(\frac{n\pi x}{L}\right)$$

$$A_n = \frac{2}{L} \int_0^L \sin\left(\frac{n\pi x}{L}\right) f(x) dx \quad \left(\frac{1}{2}\text{-range sin}\right)$$

$$B_n = \frac{2}{n\pi c} \int_0^L \sin\left(\frac{n\pi x}{L}\right) g(x) dx.$$

Note that A_n are just the coefficients of the half-range sin series of f .

1D Heat, Finite Domain

Consider the 1D heat equation $u_t = c^2 u_{xx}$ on $x \in [0, L]$ for $t \geq 0$, with ICs $u(x, 0) = f(x)$ and BCs $u_x(0, t) = u_x(L, t) = 0$, indicating that the endpoints are insulated. Again using Fourier's method, one possible solution is given by $u(x, t) = A_0 \in \mathbb{R}$. A non-constant solution is given by the following formula. A_n are the half-range cos coefficients of f .

$$u(x, t) = \sum_{n=0}^{\infty} A_n e^{-\left(\frac{n\pi c}{L}\right)^2 t} \cos\left(\frac{n\pi x}{L}\right)$$

$$\left. \begin{aligned} A_0 &= \frac{1}{L} \int_0^L f(x) dx \\ A_n &= \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx \end{aligned} \right\} \quad \left(\frac{1}{2}\text{-range cos}\right)$$

For the case that BCs are $u(0, t) = u(L, t) = 0$ involving u instead of u_x , a half-range sin series for coefficients B_n is used instead, and $u(x, t) = \sum_{n=1}^{\infty} B_n \exp\left(-\left(\frac{n\pi c}{L}\right)^2 t\right) \sin\left(\frac{n\pi x}{L}\right)$. This was verified in an assignment question.

1D Heat, Infinite Domain

For $u_t = c^2 u_{xx} = 0$ on $x \in \mathbb{R}$ for $t \geq 0$, with IC $u(x, 0) = f(x)$ and no boundary conditions, the Fourier method will not be useful. Instead, the Gaussian function $G(x, t)$ is a special solution to the equation, and so is $G(x - y, t)$ with $y \in \mathbb{R}$ as a parameter. It represents an injection of 1 unit of heat energy at $x = y, t = 0$. Due to the superposition principle, the following formula can be verified as the solution to the IVP by distributing partials into the integral and taking the limit as $t \rightarrow 0$ for the IC.

$$u(x, t) = \int_{-\infty}^{\infty} G(x - y, t) f(y) dy = \frac{1}{\sqrt{4\pi c^2 t}} \int_{-\infty}^{\infty} \exp\left(-\frac{(x - y)^2}{4c^2 t}\right) f(y) dy.$$

1D Heat, Positive Domain

For $u_t = c^2 u_{xx}$ on $x, t \geq 0$, with IC $u(x, 0) = F(x)$ for $x \geq 0$ and BC $u(0, t) = 0$, this is converted into a problem over all of \mathbb{R} considering values $x < 0$ to be *unphysical*. Define an odd extension of F :

$$f(x) = \begin{cases} F(x), & x \geq 0 \\ -F(-x), & x < 0. \end{cases}$$

The solution to the infinite domain problem with $u(x, 0) = f(x)$ on \mathbb{R} is a solution to the original problem for $x \geq 0$. The BC is satisfied since f odd $\implies f(0) = 0$. Using such symmetric extensions is called the *method of images*. The solution given above for the infinite case can be simplified to the following formula:

$$u(x, t) = \int_0^{\infty} [G(x - y, t) - G(x + y, t)] F(y) dy.$$

1D Heat, Positive Domain (cont.)

For the case that an *insulation condition* $u_x(0, t) = 0$ is given instead of $u(0, t) = 0$, extend F with even symmetry:

$$f(x) = \begin{cases} F(x), & x \geq 0 \\ F(-x), & x < 0. \end{cases}$$

By a similar derivation as before, the solution is given by

$$u(x, t) = \int_0^\infty [G(x - y, t) + G(x + y, t)] F(y) dy.$$

Note the $+$ instead of $-$ in the square brackets.

For a non-zero boundary condition $u(0, t) = u_0 \neq 0$, the problem can be modified to have the solution $\hat{u}(x, t) = u(x, t) - u_0$, which also satisfies $\hat{u}_t = c^2 \hat{u}_{xx}$ and has BC $\hat{u}(0, t) = u(0, t) - u_0 = 0$, so the previous formulas can be used to obtain $\hat{u}(x, t)$ using

$$\hat{F}(x) = F(x) - u_0.$$

This can then be transformed back to the original problem by adding u_0 to $\hat{u}(x, t)$.

1D Heat, Finite, Sinusoidal BC

Consider $u_t = c^2 u_{xx}$ on $x, t \geq 0$ with BC $u(0, t) = u_0 \cos(\omega t)$ and no ICs. This can model temperature variation in soil due to periodic variation of surface temperature. The equation can be solved by considering a complex function $v(x, t) = X(x)e^{i\omega t}$ which solves the PDE, and representing the BC as a complex exponential $v(0, t) = u_0 e^{i\omega t}$. The desired solution is then given by $u(x, t) = \text{Re}(v(x, t))$.

Finding derivatives and substituting into the PDE gives $i\omega X e^{i\omega t} = c^2 X'' e^{i\omega t}$ so $X'' = \frac{i\omega}{c^2} X = \alpha^2 X$ with $\alpha = \pm \sqrt{\frac{i\omega}{c^2}} = \pm \sqrt{\frac{\omega}{2c^2}}(1 + i)$. This gives

$$X(x) = A \exp\left(\sqrt{\frac{\omega}{2c^2}}(1 + i)x\right) + B \exp\left(-\sqrt{\frac{\omega}{2c^2}}(1 + i)x\right)$$

$A = 0$ since the solution should be bounded as $x \rightarrow \infty$. Applying the BC gives $v(0, t) = u_0 e^{i\omega t} = B e^{i\omega t}$ so $B = u_0$. Thus,

$$v(x, t) = X(x)e^{i\omega t} = u_0 \exp\left(-\sqrt{\frac{\omega}{2c^2}}x\right) e^{i\omega t},$$

and the desired solution is a sinusoid with exponential decay,

$$u(x, t) = u_0 e^{-\sqrt{\frac{\omega}{2c^2}}x} \cos\left(\omega t - \sqrt{\frac{\omega}{2c^2}}x\right).$$

1D Heat, Finite, Inhomogeneous, General Case $q = q(x, t)$

If there is a heat source present, the 1D heat equation is modified to $u_t - c^2 u_{xx} = q(x, t)$. Consider this PDE on $x \in [0, L]$ for $t \geq 0$ with IC $u(x, 0) = f(x)$ and BCs $u(0, t) = u(L, t) = 0$. Assume a solution of the form

$$u(x, t) = \sum_{n=1}^{\infty} B_n(t) \sin\left(\frac{n\pi x}{L}\right).$$

The IC gives $u(0, t) = f(x) = \sum_{n=1}^{\infty} B_n(0) \sin\left(\frac{n\pi x}{L}\right)$ so using a half-range sin series,

$$B_n(0) = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx.$$

q can also be written as a half-range sin series:

$$q(x, t) = \sum_{n=1}^{\infty} Q_n(t) \sin\left(\frac{n\pi x}{L}\right), \quad Q_n(t) = \frac{2}{L} \int_0^L q(x, t) \sin\left(\frac{n\pi x}{L}\right) dx.$$

Substituting the half-range sums into the PDE and taking all terms to one side gives

$$\sum_{n=1}^{\infty} \left[B'_n(t) + \left(\frac{n\pi c}{L}\right)^2 B_n(t) - Q_n(t) \right] \sin\left(\frac{n\pi x}{L}\right) = 0.$$

For every n , since $\sin\left(\frac{n\pi x}{L}\right)$ are orthogonal, then

$$B'_n(t) + \left(\frac{n\pi c}{L}\right)^2 B_n(t) - Q_n(t) = 0 \quad (*)$$

This is a 1st-order linear ODE with constant coefficients for $B_n(t)$.

Here are the steps required to solve the PDE:

1. Calculate IC $B_n(0)$ with half-range sin
2. Calculate $Q_n(t)$ with half-range sin
3. Find the homogeneous solution of $(*)$, $B_{n,H}$
4. Guess a particular solution $B_{n,P}$ to $(*)$ (undetermined coefficients)
5. $B_n(t) = B_{n,H} + B_{n,P}$
6. Substitute the calculated $B_n(0)$ and find unknown coefficients
7. $u(x, t) = \sum_{n=1}^{\infty} B_n(t) \sin\left(\frac{n\pi x}{L}\right)$.

1D Heat, Finite, Inhomogeneous, Special Case $q = q(x)$

For the case that the inhomogeneous term q is only a function of x , so $u_t - c^2 u_{xx} = q(x)$, then a simpler process can be used to solve the PDE. BCs and ICs are same as in the general case. Assume that a solution $u_s(x)$ is only a function of x , called the “steady-state” solution. It can be thought of as $u(x, t)$ as $t \rightarrow \infty$. Substituting into the PDE gives

$$(u_s)_t - c^2(u_s)_{xx} = -c^2(u_s)_{xx} = q(x).$$

Integrate twice to get u_s , and apply BCs to solve for constants of integration.

For the “transient” or non-steady-state part, $\hat{u} = u - u_s$ can be shown to satisfy the homogeneous heat equation and preserve the BCs $\hat{u}(0, t) = \hat{u}(L, t) = 0$. The new IC is $\hat{u}(x, 0) = \hat{f}(x) = f(x) - u_s(x)$. The solution to this homogeneous equation, as discussed earlier, is

$$\hat{u}(x, t) = \sum_{n=1}^{\infty} B_n e^{-\left(\frac{n\pi c}{L}\right)^2 t} \sin\left(\frac{n\pi x}{L}\right), \quad B_n = \frac{2}{L} \int_0^L \hat{f}(x) \sin\left(\frac{n\pi x}{L}\right) dx.$$

The total solution is then $u(x, t) = u_s(x) + \hat{u}(x, t)$.

1D Wave, Finite, Inhomogeneous

Consider $u_{tt} - c^2 u_{xx} = F(x, t)$ on $x \in [0, L]$ and $t \geq 0$, where F represents the force per unit mass on the finite string. BCs are $u(0, t) = u(L, t) = 0$ and ICs are $u(x, 0) = f(x)$, $u_t(x, 0) = g(x)$.

If F is a function of x and t , assume that the solution has the form $u(x, t) = \sum_{n=1}^{\infty} B_n(t) \sin\left(\frac{n\pi x}{L}\right)$ and convert F to a half-range sin series. Proceed as in the 1D heat inhomogeneous equation and obtain a second-order ODE in $B_n(t)$:

$$B_n''(t) + \left(\frac{n\pi c}{L}\right)^2 B_n(t) - K_n(t) = 0,$$

where $K_n(t)$, $B_n(0)$ and $B_n'(0)$ are the half-range sin coefficients of $F(x, t)$, $f(x)$ and $g(x)$ respectively. (This formula was not fully derived in lectures.)

If F is only a function of x , the steady-state technique discussed in the 1D heat case can be used to simplify the process, assuming a solution $u_s(x)$ independent of t and then calculating the transient part $\hat{u} = u - u_s$. This can be shown to solve $\hat{u}_{tt} - c^2 \hat{u}_{xx} = 0$ and preserve BCs. New ICs are

$$\hat{u}(x, 0) = \hat{f}(x) = f(x) - u_s(x), \quad \hat{u}_t(x, 0) = g(x).$$

\hat{u} can be found via formulas given earlier, reproduced here for convenience:

$$\hat{u}(x, t) = \sum_{n=1}^{\infty} \left(A_n \cos\left(\frac{n\pi c t}{L}\right) + B_n \sin\left(\frac{n\pi c t}{L}\right) \right) \sin\left(\frac{n\pi x}{L}\right)$$

$$A_n = \frac{2}{L} \int_0^L \sin\left(\frac{n\pi x}{L}\right) \hat{f}(x) dx, \quad B_n = \frac{2}{n\pi c} \int_0^L \sin\left(\frac{n\pi x}{L}\right) g(x) dx.$$

2D Laplace, Dirichlet BVP, Rectangular Region

Consider the 2D Laplace equation $\nabla^2 u = u_{xx} + u_{yy} = 0$, which can be viewed as a steady-state temperature distribution since $u_t = 0$. For a rectangular region $x \in [0, a]$, $y \in [0, b]$, define the following boundary conditions:

$$\begin{cases} u(0, y) = 0 \\ u(a, y) = 0 \\ u(x, 0) = 0 \\ u(x, b) = f(x). \end{cases}$$

In other words, u is 0 everywhere on the boundary except for the top side, where it has value $f(x)$. Specifying the values of u on a boundary is known as a Dirichlet boundary value problem (BVP). No ICs are specified.

The Fourier method is used to find solutions in separable form $u(x, t) = F(x)G(y)$, giving

$$u(x, y) = \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{a}\right) \sinh\left(\frac{n\pi y}{a}\right)$$

$$B_n = \frac{2}{a \sinh\left(\frac{n\pi b}{a}\right)} \int_0^a f(x) \sin\left(\frac{n\pi x}{a}\right) dx.$$

This is found using the following formula for arbitrary $h(x)$ and $k, c_1, c_2 \in \mathbb{R}$, where the imaginary exponentials are written in terms of sinh and cosh:

$$h''(x) = k^2 h(x) \implies h(x) = c_1 \cosh(kx) + c_2 \sinh(kx).$$

For a more general problem with functions specifying the entire boundary, $\nabla^2 u = 0$ is linear so solutions considering one side at a time can be added together.

Harmonic Functions

A function $u(x, y)$ satisfying the Laplace equation $\nabla^2 u = 0$ throughout a region $D \subseteq \mathbb{R}^2$ is called *harmonic* on D . Harmonic functions come in pairs, known as *harmonic conjugates*, such that conjugates $u(x, y), v(x, y)$ satisfy the *Cauchy-Riemann Equations*:

$$u_x = v_y, \quad u_y = -v_x. \quad (\text{on } D)$$

This uniquely determines v from u up to a constant, and v can be found by partially integrating with respect to x and y .

For harmonic conjugates u and v , $\nabla u = (u_x, u_y)$ and $\nabla v = (v_x, v_y) = (-u_y, u_x)$ so $\nabla u \cdot \nabla v = -u_x u_y + u_y u_x = 0$. Thus, $\nabla u \perp \nabla v$, but since ∇u and ∇v are perpendicular to u and v 's contour lines, u and v 's contours must then be perpendicular.

In the context of heat where u is temperature, contours of u are lines of constant temperature, and contours of v indicate the flow of heat. v is called the *heat stream function*.

Fourier Method

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About

This is a summary document I made for the 2021 offering of MATH2100 Applied Mathematical Analysis taught in The University of Queensland, which I made to contain the most essential information covered in the course. Most worked examples covered in lectures have been excluded, and have been replaced by either derivations or processes to solve various types of questions. The last part of the course, applications of Laplace equation, covers point sources/sinks in 2D. That section has been omitted from the summary as I felt it was not that important compared to the rest of the content, and would not likely be assessed on exams.

All of the material in this document is either from the course notes [3] and [4], or from sources referenced below.

The source code for this document can be found here: <https://github.com/DukMastaaa/beans/blob/main/uni/MATH2100/MATH2100Summary.tex>

This PDF can be downloaded from here: <https://github.com/DukMastaaa/beans/blob/main/uni/MATH2100/output/MATH2100Summary.pdf>

Useful Resources

Phase plane plotter: <https://www.wolframalpha.com/widgets/view.jsp?id=9298fea31cf266903b3df7174b95ddd7>

Laplace transform worksheet: <https://sites.math.washington.edu/~aloveles/Math307Fall2019/m307LaplacePractice.pdf>

Inverse Laplace worksheet: <https://sites.math.washington.edu/~aloveles/Math307Spring2015/m307InverseLaplacePractice.pdf>

References

- [1] Freesodas, *Stability Diagram*, 2018. [Online]. Available: https://commons.wikimedia.org/wiki/File:Stability_Diagram.png.
- [2] Andrew D. Loveless. “Laplace Transform Practice Problems.” (2019), [Online]. Available: <https://sites.math.washington.edu/~aloveles/Math307Fall2019/m307LaplacePractice.pdf>.
- [3] School of Mathematics and Physics, *MATH2010-MATH2100-MATH7100 Analysis of Ordinary Differential Equations Workbook*, 2021.
- [4] —, *MATH2011 The PDE Section of MATH2100 Lecture Workbook*, 2021.

