

Butterworth Filters

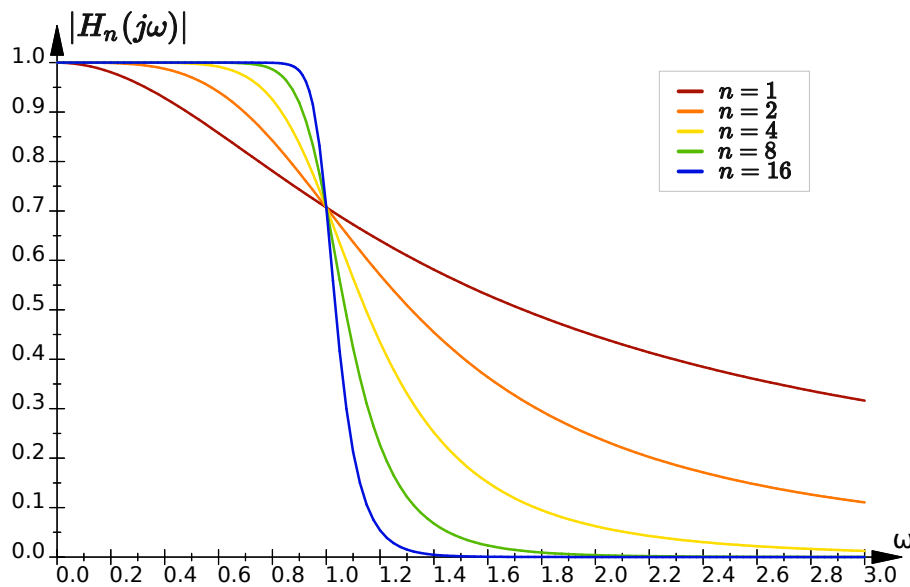
Pieter P

This page will cover the derivation of the transfer functions of low-pass and high-pass [Butterworth filters](#). Butterworth filters are designed to have a very flat frequency response in the passband.

Definition

Normalized Butterworth filters are defined in the frequency domain as follows:

$$|H_n(j\omega)| \triangleq \frac{1}{\sqrt{1 + \omega^{2n}}} \quad (1)$$



In order to determine the transfer function, we'll start from the frequency response squared. We'll assume that the transfer function $H_n(s)$ is a rational function with real coefficients. Therefore, $\overline{H_n(s)} = H_n(\bar{s})$.

$$\begin{aligned} |H_n(j\omega)|^2 &= H_n(j\omega) \overline{H_n(j\omega)} \\ &= H_n(j\omega) H_n(\overline{j\omega}) \\ &= H_n(j\omega) H_n(-j\omega) \\ &= \frac{1}{1 + \omega^{2n}} \end{aligned}$$

We're looking for the transfer function $H_n(s)$, so we'll use the identity $s = j\omega \Leftrightarrow \omega = \frac{s}{j}$.

$$H_n(s)H_n(-s) = \frac{1}{1 + \left(\frac{s}{j}\right)^{2n}}$$

Poles of $H_n(s)H_n(-s)$

The poles of this transfer function are given by:

$$\begin{aligned} \left(\frac{s}{j}\right)^{2n} &= -1 \\ \Leftrightarrow s^{2n} &= -1(j)^{2n} \\ \Leftrightarrow s^{2n} &= -1(-1)^n \\ \Leftrightarrow s^{2n} &= (-1)^{n+1} \\ \Leftrightarrow s^{2n} &= e^{j\pi(n+1)} \end{aligned}$$

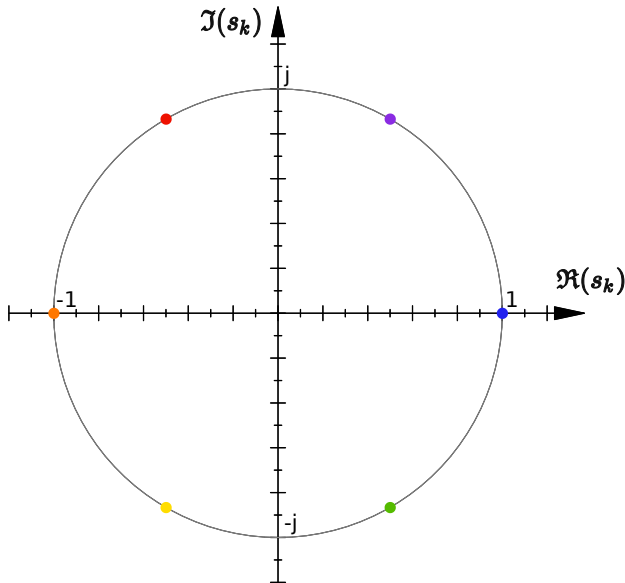
Keep in mind that this is a polynomial of order $2n$, so it has $2n$ complex roots.

$$s_k = e^{j2\pi \frac{2k+n+1}{4n}} \quad k \in \{0, 1, \dots, 2n-1\}$$

For example, for $n = 3$, the poles are:

$$\begin{aligned} s_0 &= e^{j2\pi \frac{0+3+1}{12}} = e^{j2\pi \frac{2}{6}} \\ s_1 &= e^{j2\pi \frac{2+3+1}{12}} = e^{j2\pi \frac{3}{6}} \\ s_2 &= e^{j2\pi \frac{4+3+1}{12}} = e^{j2\pi \frac{4}{6}} \\ s_3 &= e^{j2\pi \frac{6+3+1}{12}} = e^{j2\pi \frac{5}{6}} \\ s_4 &= e^{j2\pi \frac{8+3+1}{12}} = e^{j2\pi \frac{6}{6}} \\ s_5 &= e^{j2\pi \frac{10+3+1}{12}} = e^{j2\pi \frac{1}{6}} \end{aligned}$$

These are all points on the unit circle, $\pi/3 = 60^\circ$ apart.



The poles are stable if they are in the left half plane, if their complex argument is between 90° and 270° :

$$\begin{aligned} 2\pi \frac{2k+n+1}{4n} &\in \left(\frac{\pi}{2}, \frac{3\pi}{2} \right) \\ \Leftrightarrow 2k+n+1 &\in (n, 3n) \\ \Leftrightarrow k &\in \left(-\frac{1}{2}, n - \frac{1}{2} \right) \\ \Rightarrow k &\in \left(-\frac{1}{2}, n - \frac{1}{2} \right) \cup \{0, 1, \dots, 2n-1\} \\ \Leftrightarrow k &\in \{0, 1, \dots, n-1\} \\ s_{k,stable} &= e^{j2\pi \frac{2k+n+1}{4n}} \quad k \in \{0, 1, \dots, n-1\} \end{aligned} \quad (2)$$

Poles of $H_n(s)$

We want our filter $H_n(s)$ to be stable, so we pick the poles in the left half plane to be the poles of $H_n(s)$. The unstable poles, for $k \in \{n, n+1, \dots, 2n-1\}$ are the poles of $H_n(-s)$. They are the opposites of the poles of $H_n(s)$:

$$\begin{aligned} s_{k,unstable} &= e^{j2\pi \frac{2k+n+1}{4n}} \quad k \in \{n, n+1, \dots, 2n-1\} \\ l &\triangleq k - n \\ &= e^{j2\pi \frac{2(l+n)+n+1}{4n}} \quad l \in \{0, 1, \dots, n-1\} \\ &= e^{j\left(2\pi \frac{2l+n+1}{4n} + \pi\right)} \\ &= e^{j\pi} \cdot e^{j2\pi \frac{2l+n+1}{4n}} \\ &= -1 \cdot e^{j2\pi \frac{2l+n+1}{4n}} \\ &= -s_{l,stable} \end{aligned}$$

Butterworth Polynomials

We'll define the normalized Butterworth polynomial as follows:

$$B_n(s) \triangleq \prod_{k=0}^{n-1} \left(s - e^{j2\pi \frac{2k+n+1}{4n}} \right) \quad (3)$$

We'll rearrange the product to group each pole with its complex conjugate. Then, using the identity $e^{j\theta} + e^{-j\theta} = 2 \cos \theta$, we can further simplify this expression:

Even order n :

$$\begin{aligned} B_n(s) &= \prod_{k=0}^{n-1} \left(s - e^{j2\pi \frac{2k+n+1}{4n}} \right) \\ &= \prod_{k=0}^{\frac{n}{2}-1} \left(s - e^{j2\pi \frac{2k+n+1}{4n}} \right) \prod_{l=\frac{n}{2}}^{n-1} \left(s - e^{j2\pi \frac{2l+n+1}{4n}} \right) \\ &\quad l = n - k - 1 \\ &= \prod_{k=0}^{\frac{n}{2}-1} \left(s - e^{j2\pi \frac{2k+n+1}{4n}} \right) \left(s - e^{j2\pi \frac{2(n-k-1)+n+1}{4n}} \right) \\ &= \prod_{k=0}^{\frac{n}{2}-1} \left(s - e^{j2\pi \frac{2k+n+1}{4n}} \right) \left(s - \frac{e^{j2\pi \frac{-2k+3n-1}{4n}}}{1} \right) \\ &= \prod_{k=0}^{\frac{n}{2}-1} \left(s - e^{j2\pi \frac{2k+n+1}{4n}} \right) \left(s - \frac{e^{j2\pi \frac{-2k+3n-1}{4n}}}{e^{j2\pi}} \right) \\ &= \prod_{k=0}^{\frac{n}{2}-1} \left(s - e^{j2\pi \frac{2k+n+1}{4n}} \right) \left(s - e^{j2\pi \left(\frac{-2k+3n-1}{4n} - 1 \right)} \right) \\ &= \prod_{k=0}^{\frac{n}{2}-1} \left(s - e^{j2\pi \frac{2k+n+1}{4n}} \right) \left(s - e^{j2\pi \frac{-2k+3n-4n-1}{4n}} \right) \\ &= \prod_{k=0}^{\frac{n}{2}-1} \left(s - e^{j2\pi \frac{2k+n+1}{4n}} \right) \left(s - e^{j2\pi \frac{-2k-n-1}{4n}} \right) \\ &= \prod_{k=0}^{\frac{n}{2}-1} \left(s - e^{j2\pi \frac{2k+n+1}{4n}} \right) \left(s - e^{-j2\pi \frac{2k+n+1}{4n}} \right) \\ &= \prod_{k=0}^{\frac{n}{2}-1} \left(s^2 - s e^{j2\pi \frac{2k+n+1}{4n}} - s e^{-j2\pi \frac{2k+n+1}{4n}} + 1 \right) \\ &= \prod_{k=0}^{\frac{n}{2}-1} \left(s^2 - 2 \cos \left(2\pi \frac{2k+n+1}{4n} \right) s + 1 \right) \end{aligned}$$

Odd order n :

In this case, $n - 1$ is even, and you get a special pole for $k = \frac{n-1}{2}$:

$$\begin{aligned} s_{\frac{n-1}{2}} &= e^{j2\pi \frac{2 \cdot \frac{n-1}{2} + n + 1}{4n}} \\ &= e^{j2\pi \frac{2n}{4n}} \\ &= e^{j\pi} \\ &= -1 \end{aligned}$$

After isolating this pole, we're left with an even number of complex conjugate poles, just like in the case where n was even.

In conclusion, the normalized Butterworth polynomial of degree n is given by:

$$B_n(s) = \begin{cases} \prod_{k=0}^{\frac{n}{2}-1} \left(s^2 - 2 \cos \left(2\pi \frac{2k+n+1}{4n} \right) s + 1 \right) & \text{even } n \\ (s + 1) \prod_{k=0}^{\frac{n-1}{2}-1} \left(s^2 - 2 \cos \left(2\pi \frac{2k+n+1}{4n} \right) s + 1 \right) & \text{odd } n \end{cases} \quad (4)$$

Butterworth Transfer Function $H_n(s)$

The transfer function $H_n(s)$ has no zeros, so the numerator is a constant. The poles of $H_n(s)$ are given by Equation (2), so the denominator is given by Equation (3).

$$H_n(s) = \frac{c}{B_n(s)}$$

We wanted a DC gain of 1 ($= 0dB$) for $\omega = 0$:

$$\begin{aligned}
 |H_n(0j)| &= 1 \\
 \Leftrightarrow \left| \frac{c}{B_n(0)} \right| &= 1 \\
 \Leftrightarrow \left| \frac{c}{\prod_{k=0}^{n-1} \left(0 - e^{j2\pi \frac{2k+n+1}{4n}} \right)} \right| &= 1 \\
 \Leftrightarrow \frac{|c|}{\prod_{k=0}^{n-1} \left| -e^{j2\pi \frac{2k+n+1}{4n}} \right|} &= 1 \\
 \Leftrightarrow \frac{|c|}{1} &= 1
 \end{aligned}$$

If we want no phase offset for low frequencies, we can postulate that $\angle H_n(0j) = 0$:

$$\begin{aligned}
 \angle H_n(0j) &= 0 \\
 \Leftrightarrow \angle \left(\frac{c}{B_n(0)} \right) &= 0 \\
 \Leftrightarrow \angle c - \angle \left(\prod_{k=0}^{\frac{n}{2}-1} \left(0^2 - 2 \cos \left(2\pi \frac{2k+n+1}{4n} \right) \cdot 0 + 1 \right) \right) &= 0 \\
 \Leftrightarrow \angle c - \angle 1 &= 0
 \end{aligned}$$

The derivation is analogous for odd n .

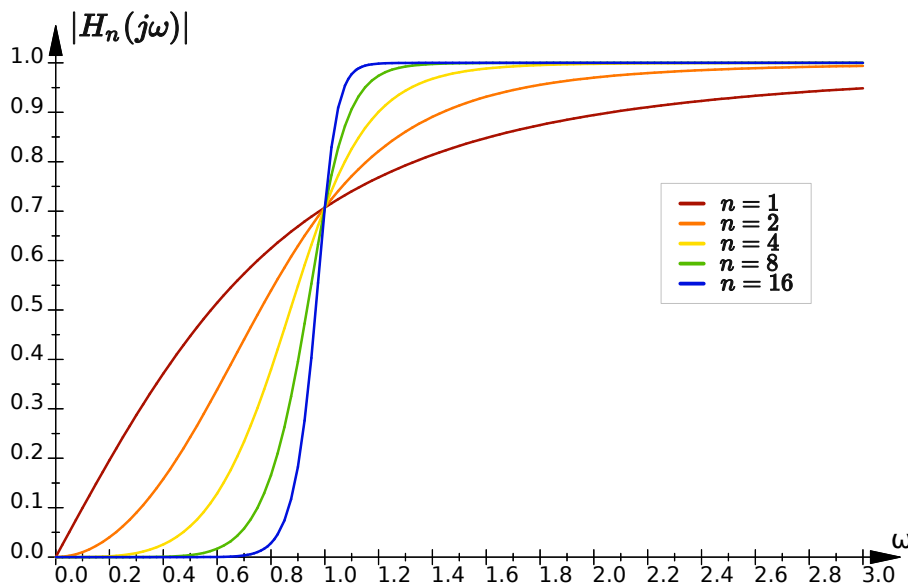
Therefore, $c = 1$, and we've eliminated all unknown parameters from the transfer function:

$$H_n(s) = \frac{1}{B_n(s)} \quad (5)$$

High-Pass Butterworth filters

Up until now, we only looked at the low-pass Butterworth filter. There's also a high-pass version:

$$|H_{n,hp}(j\omega)| \triangleq \frac{1}{\sqrt{1 + \omega^{-2n}}} \quad (6)$$



We can just multiply the numerator and the denominator by ω^n to get a more familiar form:

$$|H_{n,hp}(j\omega)| = \frac{\omega^n}{\sqrt{1 + \omega^{2n}}}$$

As you can see, the poles will be the same as for the low-pass version. On top of that, there now are n zeros for $s = 0$. So the transfer function becomes:

$$H_{n,hp}(s) = \frac{s^n}{B_n(s)} \quad (7)$$

Non-normalized Butterworth Filters

Up until now, we only looked at normalized Butterworth filters, that have a corner frequency of 1 rad/s . To get a specific corner frequency ω_c , we can just scale ω , so the definitions become:

$$|H_{n,lp}(j\omega)| \triangleq \frac{1}{\sqrt{1 + \left(\frac{\omega}{\omega_c}\right)^{2n}}} \quad (8)$$

$$|H_{n,hp}(j\omega)| \triangleq \frac{1}{\sqrt{1 + \left(\frac{\omega_c}{\omega}\right)^{2n}}} \quad (9)$$

If you start recalculating the transfer functions, you'll quickly realize that this just scales everything by a factor of ω_c . The poles no longer lie on the unit circle, but on a circle with radius $|s_k| = \omega_c$.

This results in the following transfer functions:

$$H_{n,lp}(s) = \frac{1}{B_n\left(\frac{s}{\omega_c}\right)} \quad (10)$$

$$H_{n,hp}(s) = \frac{s^n}{\omega_c^n B_n\left(\frac{s}{\omega_c}\right)} \quad (11)$$

The gain at the corner frequency can easily be determined from the definitions:

$$\begin{aligned} |H_{n,lp}(j\omega_c)| &= |H_{n,hp}(j\omega_c)| = \frac{1}{\sqrt{1 + \left(\frac{\omega_c}{\omega_c}\right)^{2n}}} \\ &= \frac{1}{\sqrt{2}} \\ &= \frac{\sqrt{2}}{2} \\ &\approx 0.707 \\ 20 \log_{10}|H_n(j\omega_c)| &= 20 \log_{10}\left(\frac{\sqrt{2}}{2}\right) \\ &= 10 \log_{10}\left(\frac{1}{2}\right) \\ &\approx -3.01 \text{ dB} \end{aligned}$$

This is often called the -3 dB -point or the half-power point, because a sinusoidal input signal at that frequency will result in an output signal that has only half of the power of the input signal: $|H_n(j\omega_c)|^2 = \frac{1}{2}$.
