



# COMPUTATIONAL ASTROPHYSICS

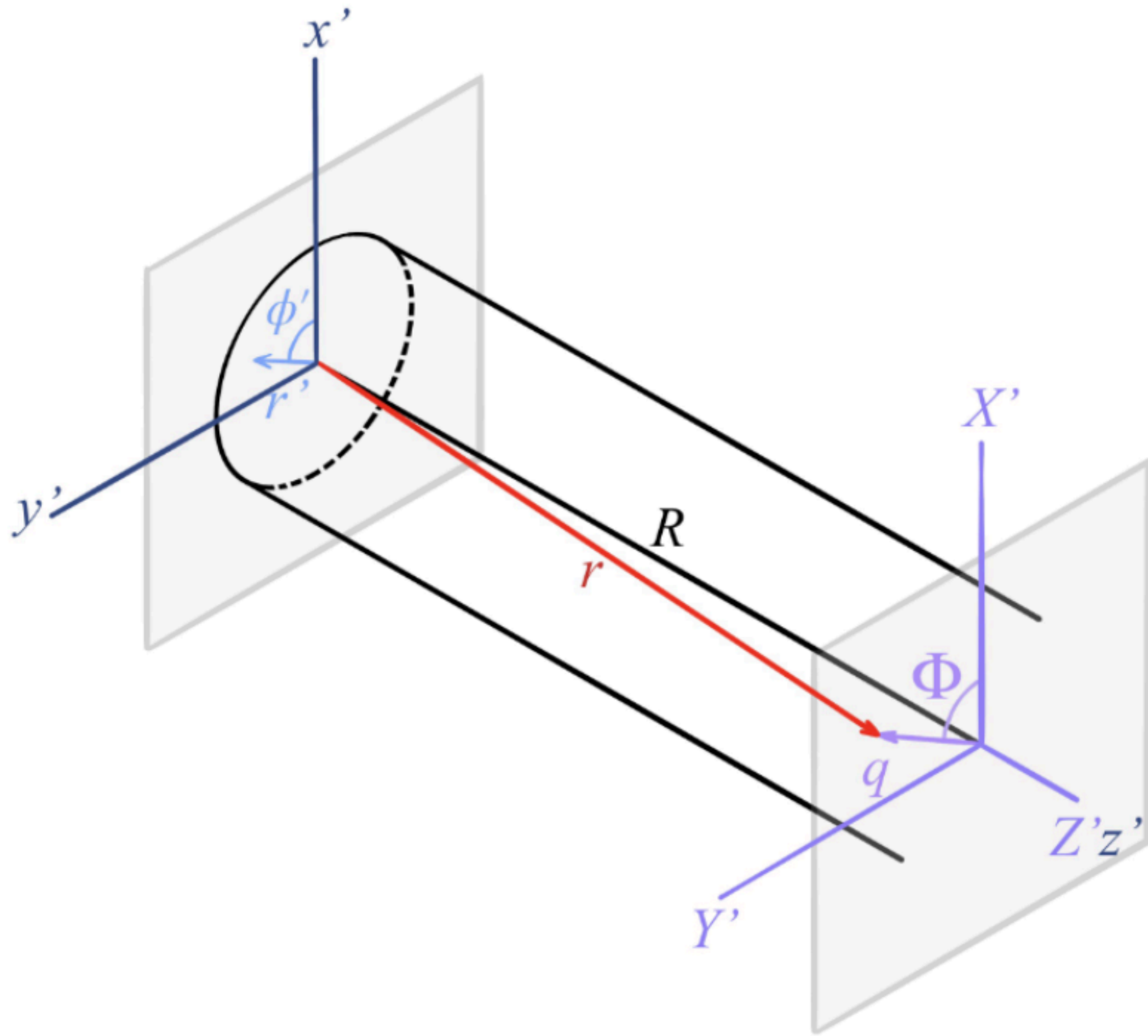
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# Computational Astrophysics 2022

## 03. Airy Diffraction

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# Diffraction



- Plane wave incident into a circular hole with radius  $a$
- We want to calculate the intensity of light registered in a plane screen located at a distance  $R \gg a$
- The electric field at the circular aperture is constant,

$$E(t, x', y') = E_0 e^{-i\omega t}$$

The Fourier transform of the electric field gives

$$E(t, k_x, k_y) = \frac{e^{ikr}}{\lambda r} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} E(t, x', y') e^{-i(k_x x' + k_y y')} dx' dy'$$

Using polar coordinates  $(r', \phi')$  we have

$$E(t, k_x, k_y) = \frac{e^{ikr}}{\lambda r} \int_0^a \int_0^{2\pi} E(t, r', \phi') e^{-i(k_x r' \cos \phi' + k_y r' \sin \phi')} r' dr' d\phi'$$

Introducing the coordinates  $(X, Y)$  at the screen plane and approximating

$$k_x \approx k \frac{X}{r}$$

$$k_y \approx k \frac{Y}{r}$$

and the polar coordinates  $(q, \Phi)$  at the screen plane, gives the relation

$$E(t, q) = \frac{e^{i(kr - \omega t)}}{\lambda r} \int_0^a \int_0^{2\pi} E_0 \exp \left[ -\frac{ikr'q}{r} \cos(\phi' - \Phi) \right] r' dr' d\phi'$$

Due to the azimuthal symmetry, we can take  $\Phi = 0$ , without losing generality, giving

$$E(t, q) = \frac{e^{i(kr - \omega t)}}{\lambda r} \int_0^a \int_0^{2\pi} E_0 \exp \left[ -\frac{ikr'q}{r} \cos(\phi') \right] r' dr' d\phi'$$

Introducing the Bessel functions of the first kind,

$$J_m(u) = \frac{1}{2\pi} \int_0^{2\pi} e^{i(mv + u \cos v)} dv$$

We recognize, in the expression for the electric field, the  $m = 0$  Bessel function,

$$\mathcal{J} = \int_0^{2\pi} \exp \left[ -\frac{ikr'q}{r} \cos(\phi') \right] d\phi' = 2\pi J_0(u) ,$$

$$\text{if } u = \frac{kr'q}{r}.$$

Hence, we can write

$$E(t, q) = \frac{2\pi E_0 e^{i(kr - \omega t)}}{\lambda r} \int_0^a J_0 \left( \frac{kr'q}{r} \right) r' dr'$$

where

$$J_0 \left( \frac{kr'q}{r} \right) = \frac{1}{2\pi} \mathcal{J}.$$



Using the recurrence relation

$$\frac{d}{du} \left[ u^m J_m(u) \right] = u^m J_{m-1}(u)$$

The integration gives

$$E(t, \theta) = \frac{AE_0 e^{i(kr - \omega t)}}{\lambda r} \frac{2J_1(ka\theta)}{ka\theta}$$

where  $\theta = \frac{q}{r}$  and  $A = \pi a^2$ .

The intensity of the radiation is given by

$$I(ka\theta) = \frac{c}{8\pi} E(t, \theta) E^*(t, \theta)$$

which gives

$$I(k\rho) = \frac{cA^2E_0^2}{8\pi\lambda^2r^2} \left[ \frac{2J_1(k\rho)}{k\rho} \right]^2$$

where  $\rho = a\theta$  .

Knowing that

$$\left. \frac{dJ_1}{du} \right|_{u=0} = \frac{1}{2}$$

We can write the final result as

$$I(k\rho) = I_0 \left[ \frac{2J_1(k\rho)}{k\rho} \right]^2$$

where  $I_0 = I(0) = \frac{cA^2E_0^2}{8\pi\lambda^2r^2}$  .

This equation gives the Airy's diffraction pattern for circular diffraction.