PHYS2600J Recitation Class 5 (Final) Week 13

Zixiang Lin

University of Michigan - Shanghai Jiao Tong University Joint Institute

December 9, 2024



UM-SJTU JI

Outline

- Circuits
- EM Waves
- Optics
- Quantum Physics

UM-SJTU JI

Appendix

RLC Circuit

Circuits 000

Differential Equation

$$\frac{\mathrm{d}^2 q(t)}{\mathrm{d}t^2} + \frac{R}{L} \frac{\mathrm{d}q(t)}{\mathrm{d}t} + \frac{1}{LC} q(t) = 0$$

Solutions

■ Underdamped: $R < \sqrt{4L/C}$

$$q(t) = Ae^{-\frac{R}{2L}t}\cos\left(\sqrt{1 - \frac{R^2}{4L/C}}\frac{t}{\sqrt{LC}} + \varphi\right)$$

■ Overdamped: $R > \sqrt{4L/C}$

$$q(t) = Ae^{-(\frac{R}{\sqrt{4L/C}} + \sqrt{\frac{R^2}{4L/C} - 1})\frac{t}{\sqrt{LC}}} + Be^{-(\frac{R}{\sqrt{4L/C}} - \sqrt{\frac{R^2}{4L/C} - 1})\frac{t}{\sqrt{LC}}}$$

■ Critical: $R = \sqrt{4L/C}$

$$q(t) = (A + Bt)e^{-\frac{R}{2L}t}$$

AC Circuit

Circuits

AC Current

$$I(t) = I\cos(\omega t + \phi) = I\operatorname{Re}\left\{e^{j(\omega t + \phi)}\right\}$$
$$I_{arv} = \frac{1}{T} \int_{t_1}^{t_1 + T} |I(t)| dt = \frac{2}{\pi}I$$
$$I_{rms} = \sqrt{\frac{1}{T} \int_{t_1}^{t_1 + T} I^2(t) dt} = \frac{I}{\sqrt{2}}$$

Resonance

$$|Z| = \sqrt{\frac{\frac{1}{j\omega L} \frac{\pi/2}{\pi/2}}{R^2 + (\omega L - \frac{1}{\omega C})^2}}, \quad \omega_0 = \frac{1}{\sqrt{LC}}$$

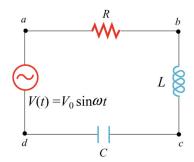
900

Circuits 0000

Suppose an AC generator with $V(t) = (150V)\sin(100t)$ is connected to a

series RLC circuit with $R=40.0\Omega$, L=80.0mH, and $C=50.0\mu F$.

- I Calculate V_{R0} , V_{L0} , and V_{C0} , the maximum of the voltage drops across each circuit element.
- Calculate the maximum potential difference across the inductor and the capacitor between points b and d.



Circuits 0000

1

$$|Z| = \sqrt{R^2 + (\omega L - \frac{1}{\omega C})^2} = 196\Omega$$

$$I_0 = \frac{V_0}{|Z|} = 0.765A$$

$$V_{R0} = I_0 R = 30.6V$$

$$V_{L0} = I_0(\omega L) = 6.12V$$

$$V_{C0} = I_0(\frac{1}{\omega C}) = 153V$$

2

$$|V_{bd} = |V_{L0} - V_{C0}| = 147V$$

UM-SJTU JI

Wave Equation

$$\frac{\partial^2 f}{\partial x^2} - \frac{1}{v^2} \frac{\partial^2 f}{\partial t^2} = 0$$

Sinusoidal Wave Function

- $f(x,t) = A\cos(k(x-vt)+\varphi)$
- Complex Notation: $\tilde{f}(x,t) = \tilde{A}e^{i(kx-\omega t)}$
- Parameters:

$$\begin{array}{ll} k & \text{wave number} \\ k(x-vt)+\varphi & \text{phase} \\ \varphi & \text{phase constant (}0 \leq \varphi < 2\pi\text{)} \\ \lambda = k/2\pi & \text{wave length} \end{array}$$

amplitude

 $T = 2\pi/kv$ period $\nu = 1/T$ frequency

 $\omega = 2\pi\nu$

angular frequency

EM Waves

in Vacuum

$$\nabla^2 \left(\vec{\mathbf{E}}, \vec{\mathbf{B}} \right) - \mu_0 \varepsilon_0 \frac{\partial^2 (\vec{\mathbf{E}}, \vec{\mathbf{B}})}{\partial t^2} = 0$$

in Conducting Medium

$$\nabla^2 \Big(\vec{\mathbf{E}}, \vec{\mathbf{H}} \Big) - \sigma \mu \frac{\partial (\vec{\mathbf{E}}, \vec{\mathbf{H}})}{\partial t} - \mu \varepsilon \frac{\partial^2 (\vec{\mathbf{E}}, \vec{\mathbf{H}})}{\partial t^2} = 0$$

Complex Notation

$$\nabla = ik, \quad \frac{\partial}{\partial t} = -i\omega$$

Poynting Vector

$$\vec{\mathbf{S}} = \vec{\mathbf{E}} \times \vec{\mathbf{H}}$$
$$\nabla \cdot \vec{\mathbf{S}} = -\frac{\partial u}{\partial t}$$

Exercise 2

- \blacksquare Consider a long straight cylindrical wire of electrical conductivity σ and radius a carrying a uniform axial current of density J. Calculate the magnitude and direction of the Poynting vector at the surface of the wire.
- 2 Consider a thick conducting slab (conductivity σ) exposed to a plane EM wave with peak amplitudes E_0, B_0 . Calculate the Poynting vector within the slab, averaged in time over a wave period. Consider σ large, i.e. $\sigma \gg \omega \varepsilon_0$.
- \blacksquare In part 2, if σ is infinite, what is the value of the average Poynting vector everywhere in space?

Solution 2

$$\vec{\mathbf{E}} = \frac{J}{\sigma} \mathbf{e_z}, \quad \vec{\mathbf{B}} = \frac{\mu_0 J a}{2} \mathbf{e_\theta}$$
$$\vec{\mathbf{S}} = \frac{1}{\mu_0} \vec{\mathbf{E}} \times \vec{\mathbf{B}} = -\frac{J^2 a}{2\sigma} \mathbf{e_r}$$

 $\vec{\mathbf{k}} = (\beta + i\alpha)\mathbf{e_z}, \quad \alpha = \beta \approx \sqrt{\frac{\omega\mu_0\sigma}{2}}$ (from HW Problems)

2

$$\vec{\mathbf{E}}(\vec{\mathbf{r}},t) = \vec{\mathbf{E}}_0(\vec{\mathbf{r}})e^{-\alpha z}e^{i(\beta z - \omega t)}$$

$$\vec{\mathbf{S}} = \vec{\mathbf{E}} \times \vec{\mathbf{H}} = \vec{\mathbf{E}} \times \frac{1}{\omega\mu_0}(\vec{\mathbf{k}} \times \vec{\mathbf{E}}) = \sqrt{\frac{\sigma}{\omega\mu_0}}e^{i\frac{\pi}{4}}E^2\mathbf{e_z}$$

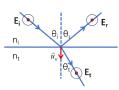
3

$$\sigma \to \infty, \alpha \to \infty, \sqrt{\sigma}e^{-2\alpha z} \to 0, \quad \langle \vec{\mathbf{S}} \rangle \to 0$$

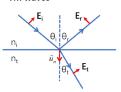
 $\langle \vec{\mathbf{S}} \rangle = \frac{\sqrt{2}}{4} \sqrt{\frac{\sigma}{\omega \mu_0}} E_0^2 e^{-2\alpha z} \mathbf{e_z}$

Fresnel Equations

TE waves



TM waves



Boundary Condition: $\hat{m{u}}_{m{n}} imes (m{E_i} + m{E_r}) = \hat{m{u}}_{m{n}} imes m{E_t}$

$$r_{TE} = \frac{n_i \cos \theta_i - n_t \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t} \quad r_{TM} = \frac{n_t \cos \theta_i - n_i \cos \theta_t}{n_t \cos \theta_i + n_i \cos \theta_t}$$

Optics •0000

$$t_{TE} = \frac{2n_i \cos \theta_i}{n_i \cos \theta_i + n_t \cos \theta_t} \quad t_{TM} = \frac{2n_i \cos \theta_i}{n_t \cos \theta_i + n_i \cos \theta_t}$$

 π Phase Shift: $r_{TE}, r_{TM} < 0 (n_i < n_t)$ or $r_{TE} > 0, r_{TM} < 0 (n_i > n_t)$ Brewster's Angle / Polarization Angle: $r_{TM} = 0, \ \theta_p = \arctan(n_t/n_i)$

Interference

Young's Double-Slit Experiment

$$\frac{I}{I_0} = \left[\cos(\delta/2)\right]^2, \delta = \frac{2\pi d \sin \theta}{\lambda}$$

Optics 00000

Newton's Rings

$$r_N = \begin{cases} \sqrt{\lambda RN} & \text{dark rings} \\ \sqrt{\lambda R(N-\frac{1}{2})} & \text{bright rings} \end{cases}$$

Interference in Thin Films

$$\frac{I_T}{I_0} = \frac{(1-R)^2}{(1-R)^2 + 4R\sin^2(\delta/2)}, \quad \frac{I_R}{I_0} = \frac{4R\sin^2(\delta/2)}{(1-R)^2 + 4R\sin^2(\delta/2)}$$
$$\delta = \frac{4\pi nh\cos i}{2\pi nh\cos i}$$

Diffraction

Fraunhofer Diffraction

Single Slit:

$$\frac{I}{I_0} = \left[\frac{\sin(\alpha)}{\alpha}\right]^2, \alpha = \frac{\pi a \sin \theta}{\lambda}$$

Rectangular Aperture:

$$\frac{I}{I_0} = [\frac{\sin(\alpha)}{\alpha}]^2 [\frac{\sin(\beta)}{\beta}]^2, \alpha = \frac{\pi a \sin \theta}{\lambda}, \beta = \frac{\pi b \sin \theta}{\lambda}$$

Circular Aperture:

$$\frac{I}{I_0} = \left[\frac{2J_1(x)}{x}\right]^2, x = \frac{2\pi a \sin \theta}{\lambda}$$
$$\Delta \theta = 0.61 \frac{\lambda}{a} = 1.22 \frac{\lambda}{D}$$

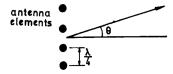
Diffraction Grating

$$\frac{I}{I_0} = \left[\frac{\sin(\alpha)}{\alpha}\right]^2 \left[\frac{\sin(N\beta)}{\beta}\right]^2, \alpha = \frac{\pi a \sin \theta}{\lambda}, \beta = \frac{\pi d \sin \theta}{\lambda} = \frac{\pi (a+b) \sin \theta}{\lambda}$$

Exercise 3

There is a phased array of radar antenna elements spaced at 1/4 wavelength. The transmitting phase is shifted by steps of $\pi/6$ from element to element, i.e., element 0 has 0 phase, element 1 has $\pi/6$ phase shift, element 2 has $2\pi/6$, element 3 has $3\pi/6$, etc.

- ${\rm 1\! I \! I}$ At what angle θ is the zero order constructive interference transmission lobe?
- 2 Are there any secondary lobes?



Solution 3

$$\delta = \frac{2\pi}{\lambda} \frac{\lambda}{4} \sin \theta + \frac{\pi}{6} = \frac{\pi}{2} (\sin \theta + \frac{1}{3})$$

$$\tilde{A}_{tot} = \sum_{k=0}^{N-1} \tilde{A} e^{ik\delta}, \quad I \sim \left[\frac{\sin(N\delta/2)}{\sin(\delta/2)}\right]^2$$

1

$$\delta/2 = 0 \to \theta = -\arcsin\frac{1}{3}$$

2

$$\delta/2 = \pm \pi \rightarrow \sin \theta = -\frac{1}{3} \pm 4$$

which means no secondary lobe can occur.

Introduction to Quantum Physics

Wave Function

- Normalization: $\iiint_{\mathbb{D}^3} |\Psi(\vec{\mathbf{r}},t)|^2 d^3 r = 1$
- Probability Density: $Pr(\vec{\mathbf{r}},t) = |\Psi(\vec{\mathbf{r}},t)|^2$
- Stationary State: $|\Psi(\vec{\mathbf{r}},t)|^2 = \Phi(\vec{\mathbf{r}})$
- Inner Product: $\langle \cdot, \cdot \rangle \in \mathcal{L}^2(\mathbb{R}^3), \langle \Psi, \Phi \rangle = \iiint_{\mathbb{R}^3} \Psi^*(\vec{\mathbf{r}}, t) \Phi(\vec{\mathbf{r}}, t) d^3r$

Operators

- Position Operator: $\hat{x} = x$
- Momentum Operator: $\hat{p} = -i\hbar \frac{\partial}{\partial x}$
- Hamiltonian Operator: $\hat{H} = \hat{K} + \hat{V} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x)$
- Energy Operator: $\hat{E} = i\hbar \frac{\partial}{\partial x}$

Hermitian

- Hermitian Conjugate: $\langle \hat{A}^{\dagger} \Psi, \Phi \rangle = \langle \Psi, \hat{A} \Phi \rangle$
- Hermitian Operator: $\hat{A}^{\dagger} = \hat{A}$

Physical Quantity

- Average: $\langle A \rangle_{\Psi} = \langle \Psi, \hat{A}\Psi \rangle = \int_{-\infty}^{\infty} A |\Psi(x,t)|^2 dx$
- Variance: $(\Delta_A^{(\Psi)})^2 = \langle (A \langle A \rangle_{\Psi})^2 \rangle_{\Psi} = \langle A^2 \rangle_{\Psi} \langle A \rangle_{\Psi}^2$
- Standard Deviation: $\Delta_A^{(\Psi)} = \sqrt{\langle A^2 \rangle_\Psi \langle A \rangle_\Psi^2}$

Commutator

$$[\hat{A}, \hat{B}] = \hat{A}\hat{B} - \hat{B}\hat{A}$$

- $[\hat{A}, \hat{B}] = -[\hat{B}, \hat{A}]$
- $\ \, [\hat{A},\hat{B}\hat{C}] = \hat{B}[\hat{A},\hat{C}] + [\hat{A},\hat{B}]\hat{C}$
- $\hat{A} [\hat{A}, \hat{B} + \hat{C}] = [\hat{A}, \hat{B}] + [\hat{A}, \hat{C}]$
- $\mathbf{4} \ [\hat{A}, \gamma] = 0$
- $[\hat{A}, f(\hat{A})] = 0$

Introduction to Quantum Physics

Schrödinger-Robertson Uncertainty Principle

$$\Delta_A^{(\Psi)} \Delta_B^{(\Psi)} \geq \frac{1}{2} |\langle \Psi, i [\hat{A}, \hat{B}] \Psi \rangle|$$

Heisenberg Uncertainty Principle

$$\Delta_x^{(\Psi)} \Delta_p^{(\Psi)} \ge \frac{\hbar}{2}$$

1D Schrödinger Equation

$$\hat{H}\Psi = \hat{E}\Psi : -\frac{\hbar^2}{2m}\frac{\partial^2}{\partial x^2}\Psi(x,t) + V(x)\Psi(x,t) = i\hbar\frac{\partial}{\partial t}\Psi(x,t)$$

1D Stationary Schrödinger Equation

$$\hat{H}\psi = E\psi : -\frac{\hbar^2}{2m}\frac{\mathrm{d}^2}{\mathrm{d}x^2}\psi(x) + V(x)\psi(x) = E\psi(x)$$

Introduction to Quantum Physics

Solution of 1D Schrödinger Equation

$$\Psi(x,t) = \sum_{n} c_n e^{-\frac{i}{\hbar}E_n t} \psi_n(x)$$

Infinite Potential Well

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

$$\psi_n(x) = \begin{cases} \sqrt{\frac{2}{L}} \sin(\frac{n\pi}{L}x) & 0 \le x \le L\\ 0 & \text{otherwise} \end{cases}$$

Measurement Postulate

$$\psi = \sum_{n} c_n \psi_n$$
, $\Pr(A = a_n) = |c_n|^2$

Quantum Harmonic Oscillator

$$E_n = (n + \frac{1}{2})\hbar\omega$$

Exercise 4

A particle with mass m is confined in a one-dimensional box of length l

$$\begin{cases} V = \infty, & x < 0 \\ V = 0, & 0 < x < l \\ V = \infty, & x > l \end{cases}$$

At time t = 0, the particle wave function is

$$\begin{cases} \psi = \sqrt{\frac{30}{l^5}} x(l-x), & 0 < x < l \\ \psi = 0, & x > l \text{ or } x < 0 \end{cases}$$

Find the series representation and expression for the series coefficients of $\psi(x,t>0)$.

Solution 4

$$\psi_n = \sqrt{\frac{2}{l}} \sin\left(\frac{n\pi}{l}x\right), \quad E_n = \frac{n^2 \pi^2 \hbar^2}{2ml^2}$$

$$c_n = \int_0^l \sqrt{\frac{2}{l}} \sin\left(\frac{n\pi}{l}x\right) \sqrt{\frac{30}{l^5}} x(l-x) dx = 4\sqrt{15} \frac{1 - (-1)^n}{n^3 \pi^3}$$

$$\Psi(x,t) = \sum_n c_n e^{-\frac{i}{\hbar} E_n t} \psi_n = 8 \sum_{n=1}^{\infty} \sqrt{\frac{30}{l^5}} \frac{\sin\left(\frac{2n+1}{l} \pi x\right) e^{-i\frac{\hbar}{2m}\left(\frac{2n+1}{l} \pi^2\right)^2 t}}{(2n+1)^3 \pi^3}$$

At some instant of time the state of a particle moving in a harmonic potential well $V(x)=m\omega^2x^2/2$ is described by the (normalized) wave function

$$\psi = \sqrt{\frac{1}{3}}\psi_1 + \sqrt{\frac{2}{3}}\psi_3$$

where ψ_n is the n-th eigenstate of the Hamiltonian.

The total energy of the particle is measured.

- What are possible outcomes of the measurement?
- 2 What is the probability that the measurement yields $3\hbar\omega/2, 5\hbar\omega/2, 7\hbar\omega/2$?
- If the first measurement yields $3\hbar\omega/2$, what is the probability that the subsequent measurement yields $3\hbar\omega/2$?

Appendix

Quantum Physics

Solution 5

1

$$E_1 = \hbar\omega(1 + \frac{1}{2}) = \frac{3}{2}\hbar\omega$$

$$E_3 = \hbar\omega(3 + \frac{1}{2}) = \frac{7}{2}\hbar\omega$$

2

$$\Pr(E = E_1) = |c_1|^2 = \frac{1}{3}$$

 $\Pr(E = E_2) = |c_2|^2 = 0$

$$\Pr(E = E_3) = |c_3|^2 = \frac{2}{3}$$

3 Wave function collapse

$$\Pr(E = E_1) = 1$$

UM-SJTU JI

Tips for the Final

- Check whether the dimensions of the answers are correct.
- Check whether the answers are as expected and make physical sense.
- List relative equations if you don't know how to solve the problem.
- Review HW problems and check whether you can solve them easily.
- Review the lecture notes carefully, especially the unfamiliar concepts!

Optics 00000 Quantum Physics

EM Waves

Thanks for listening!

Appendix O●O

References



D. J. Griffiths.

Introduction to Electrodynamics.
Cambridge University Press, 5 edition, 2023.



M. Krzyzosiak.

Lecture slides, 2024.

PHYS3900J, Modern Physics, UM-SJTU JI.



Z. Qu.

Lecture notes, 2024. PHYS2600J, Honors Physics II, UM-SJTU JI.



J. Wu.

Lecture notes, 2024.

ECE3340J, Principles of Optics, UM-SJTU JI.