03_04_05_Two_Dimensional_Systems

1. Graphical Representations of Two-Dimensional Systems

Diagrams

2. Linear Time Invariant Systems in Modal Coordinates

Change of Coordinates

3. Qualitative Behavior of Linear Time Invariant Systems

Two Real Eigenvalues

Complex Eigenvalues

4. Classification of Nonlinear Equilibrium Points

Concepts

5. Periodic Orbits and Limit Cycles

Concepts

Existence of Periodic Orbits

Absence of Periodic Orbits

5. Bifurcations

Conceptions

Pitchfork Bifurcations

Transcritical and Hopf Bifurcations

Finding Bifurcation Points

Useful Result From Linear Systems Theory

Summary

1. Graphical Representations of Two-Dimensional Systems

Diagrams

Definition: Vector Field Diagram

Consider the autonomous two-dimensional system

$$\left[egin{array}{c} \dot{x}_1 \ \dot{x}_2 \end{array}
ight] = f\left(\left[egin{array}{c} x_1 \ x_2 \end{array}
ight]
ight)$$

A <u>vector field diagram</u> is obtained when, for each x on a grid in the $x_1 - x_2$ plane, an **arrow is drawn from** x **to** x + cf(x), where c > 0 is a scale factor such that the diagram is easy to interpret.

• Sometimes, the length of the arrows is also set to a fixed value

Definition: Trajectories

Consider the autonomous two-dimensional system

$$\dot{x}_{1}=f_{1}\left(x_{1},x_{2}
ight),\quad \dot{x}_{2}=f_{2}\left(x_{1},x_{2}
ight),\quad \left[egin{array}{c}x_{1}\left(t_{0}
ight)\ x_{2}\left(t_{0}
ight)\end{array}
ight]=x_{0}$$

The curve that results from plotting $x_2(t)$ over $x_1(t)$ for $t \geq 0$ is called **trajectory (or orbit)** from x_0 .

Definition: Phase Portrait

The **phase portrait** of an autonomous two-dimensional system

$$\left[egin{array}{c} \dot{x}_1\ \dot{x}_2 \end{array}
ight]=f\left(\left[egin{array}{c} x_1\ x_2 \end{array}
ight]
ight)$$

is obtained if the **trajectories for many different initial points** x_0 **are shown in a single plot**. Sketching a phase portait becomes easier when it is drawn on top of a vector field diagram.

Definition: Isoclines

The **isocline** of an autonomous two-dimensional system

$$\dot{x}_1 = f_1(x_1, x_2), \quad \dot{x}_2 = f_2(x_1, x_2),$$

with slope $s \in \mathbb{R}$ is the curve of all $x \in \mathbb{R}^2$ that satisfy

$$f_2(x_1,x_2)=sf_1(x_1,x_2).$$

The <u>nullclines</u> are the two curves given by $f_1=0$ and $f_2=0$.

Note:

- Isocline means equal slope
- The nullclines divide the phase plane into regions in which the signs of f_1 and f_2 are constant (if f is continuous). In these regions, the arrows on an isocline all point into the same direction
- The intersections of the nullclines are the equilibria of the system

Method: Method of Drawing Isoclines

- 1. **Sketch the nullclines** $f_1=0$ and $f_2=0$. Indicate the slope of f on the nullclines using equal-length arrows
- 2. **Sketch a few isoclines** $f_2=sf_1$ for different slopes s. Indicate the slope of f on the isoclines using equal-length arrows
- 3. Sketch trajectories and add equilibria.

2. Linear Time Invariant Systems in Modal Coordinates

If the eigenvalues of A are distinct and different from zero, there exists a real invertible matrix M such that $M^{-1}AM = J_r$, where

$$J_r = \left[egin{array}{cc} \lambda_1 & 0 \ 0 & \lambda_2 \end{array}
ight]$$

if the eigenvalues λ_1,λ_2 are both real. Otherwise, $\lambda_1=lpha+eta j.\lambda_2=lpha-eta j$

$$J_r = \left[egin{array}{cc} lpha & -eta \ eta & lpha \end{array}
ight]$$

with α , β real. (This is the **real Jordan form**.)

Change of Coordinates

$$\dot{z}(t)=J_rz(t),\quad z(t):=M^{-1}\chi(t).$$

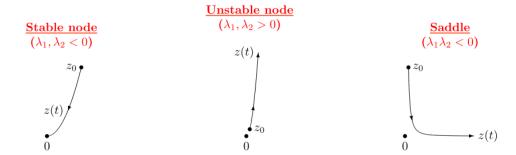
Then we will have

$$\dot{x}(t) = Ax(t) = MJ_rM^{-1}x(t)$$

3. Qualitative Behavior of Linear Time Invariant Systems

Two Real Eigenvalues

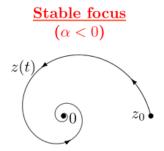
$$egin{aligned} \dot{Z}(t) = \left[egin{array}{cc} \lambda_1 & 0 \ 0 & \lambda_2 \end{array}
ight] z(t) \Rightarrow \dot{Z}_1 = \lambda_1 Z_1, \dot{z}_2 = \lambda_2 Z_2 \ z_1(t) = z_1^0 e^{\lambda_1 t}, \quad z_2(t) = z_2^0 e^{\lambda_2 t} \end{aligned}$$

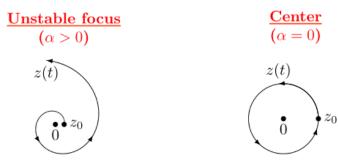


Complex Eigenvalues

We consider it in polar representation

$$egin{aligned} \dot{r} &= lpha r, \quad \dot{ heta} &= eta. \ \ z(t) &= r_0 e^{lpha t} \left[egin{array}{c} \cos \left(heta_0 + eta t
ight) \ \sin \left(heta_0 + eta t
ight) \end{array}
ight] \end{aligned}$$





4. Classification of Nonlinear Equilibrium Points

Concepts

Definition: Hyperbolic Equilibrium Points for Linear Systems

The equilibrium point $x^* = 0$ of the linear system $\dot{x}(t) = Ax(t)$ is called **Hyperbolic** if no eigenvalue of A are zero real part.

- Hyperbolic equilibrium points do not change their type in the perturbed system $\dot{x}=(A+\Delta A)x$ if ΔA is small enough
- They are thus said to be **structurally stable**
- Nodes, focuses and saddles are hyperbolic, centers are not

Definition: Hyperbolic Equilibrium Points for Nonlinear Systems

The equilibrium point x^* is called <u>hyperbolic</u> if the linearized system has a hyperbolic equilibrium point at zero.

Note:

- If an equilibrium point of $x^st = f(x)$ is hyperbolic, the system will behave locally like its linearization around that point
- This is not always true for non-hyperbolic equilibrium points
- If the linearization has a center, nonlinear analysis is required

Theorem

An equilibrium point x^* of the nonlinear system $\dot{x}=f(x)$ is called <u>(un)stable node / (un)stable focus / saddle</u> if the linearized system

$$\dot{y}=Ay,\quad A=\left \lceil rac{\partial f\left (x
ight)}{\partial x}
ight
ceil$$

has a (un)stable node / (un)stable focus / saddle at zero.

5. Periodic Orbits and Limit Cycles

Concepts

Definition: Oscillates and Periodic Orbits

A system **oscillates** if it has a nonconstant periodic solution

$$x(t) = x(t+T) \quad orall t \geq t_0, \quad x(t)
ot\equiv ext{const.}$$

These are usually called periodic or closed orbits

Definition: Harmonic Oscillators

system has a equilibrium point as a center \Rightarrow all trajectories with $z(t_0)
eq 0$ are periodic orbits

Note:

A harmonic oscillator is useless in practice: even very small pertubations can turn it into a focus

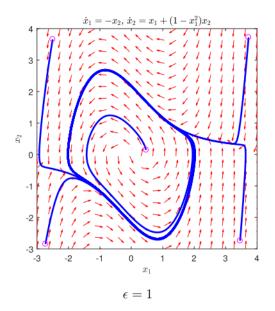
Definition: Limit Cycle

An isolated closed orbit is called a limit cycle

- In the case where all the neighboring trajectories approach the limit cycle as time approaches infinity, it is called a **stable or attractive limit cycle** (ω-limit cycle).
- If instead, all neighboring trajectories approach it as time approaches negative infinity, then it is an <u>unstable limit</u> <u>cycle</u> (α-limit cycle).
- If there is a neighboring trajectory which spirals into the limit cycle as time approaches infinity, and another one which spirals into it as time approaches negative infinity, then it is a **semi-stable limit cycle.**

Example: Van der Pol Oscillator

$$\dot{x}_1=x_2,\dot{x}_2=-x_1+\epsilon\left(1-x_1^2
ight)x_2,\quad\epsilon>0$$



Existence of Periodic Orbits

Lemma: Trapping Trajectories

We consider the system

$$\dot{x}=f(x), \quad f:\mathbb{R}^2 o\mathbb{R}^2 ext{ continuously differentiable.}$$

Let $V:\mathbb{R}^2 o \mathbb{R}$ be continuously differentiable such that

$$f(x) \cdot \nabla V(x) \le 0$$
 whenever $V(x) = c$

for some constant c>0. Then, any trajectory starting in

$$M:=\left\{x\in\mathbb{R}^2:V(x)\leq c
ight\},$$

will stay in M. That is, $x\left(t_{0}
ight)\in M$ implies $x(t)\in M$ for all $t\geq t_{0}$.

Note:

$$f(x) \cdot \nabla V(x) \leq 0$$
 whenever $V(x) = c$

Guarantees that f(x) is pushed to the opposite direction of V(x) increasing

• The lemma still holds if the condition is replaced with $f(x)\cdot
abla V(x)\geq 0$ and the set is changed to $M':=\{x\in \mathbb{R}^2: V(x)\geq c\}.$

The condition means x change opposite to the V(x) change

Theorem: Poincare-Bendixson Criterion

We consider the autonomous second-order system

$$\dot{x}_{1}=f_{1}\left(x_{1},x_{2}
ight),\quad \dot{x}_{2}=f_{2}\left(x_{1},x_{2}
ight),\quad x\left(t_{0}
ight)=x_{0},$$

where f is continuously differentiable. Let M be a closed bounded subset of the plane such that

- ullet M contains **at most one equilibrium point** that is either an unstable focus or an unstable node
- ullet Every solution with $x_0 \in M$ satisfies $x(t) \in M$ for all $t \geq t_0$

Then, there exists an initial condition $x_0 \in M$ and a period T>0 such that the solution of the system is not constant and satisfies

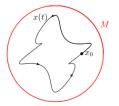
$$x(t+T)=x(t) \quad \forall t \geq t_0.$$

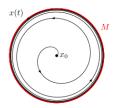
That is, M contains at least one periodic orbit

Intuition Behind The Poincarè-Bendixson Criterion

If there is no equilibrium point in M:

- x(t) cannot cross itself (true for any trajectory)
- x(t) cannot converge (that would be an equilibrium)
- x(t) also cannot leave M (by assumption)
- x(t) thus either arrives somewhere it has already been, or it gets closer and closer to some closed curve (e.g., the boundary of M)

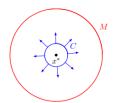




Intuition Behind The Poincarè-Bendixson Criterion

If there is one unstable focus or node x^* in M:

- Choose a curve C around x^* such that f always points into M on C
- Trajectories starting outside C will never intersect it
- Remove the area enclosed by C from M and call the result M'
- We can now apply the "no equilibrium" case to M'



Absence of Periodic Orbits

Theorem: Bendixson Criterion

We consider the system

$$\dot{x} = f(x), \quad f: \mathbb{R}^2 \to \mathbb{R}^2$$
 continuously differentiable.

Let $D \subseteq \mathbb{R}^2$ be a simply connected domain. If either

$$abla \cdot f(x) > 0 \quad orall x \in D \quad ext{ or } \quad
abla \cdot f(x) < 0 \quad orall x \in D,$$

then there are no periodic orbits that lie entirely in D.

The condition means, we cannot find a trajectory in which f(x) is periodic, that is increasing of f(x) along the closed trajectory is zero

We use a proof by contradiction. Assume that a periodic orbit $\gamma := \{x(t) : t \ge t_0\}$ would exist. Since

$$\left[\begin{array}{c} \frac{dx_1}{dt} \\ \frac{dx_2}{dt} \end{array}\right] = \dot{x} = f(x) \quad \Longrightarrow \quad \frac{dx_2}{dx_1} = \frac{\dot{x}_2}{\dot{x}_1} = \frac{f_2}{f_1} \quad \Longrightarrow \quad f_1 dx_2 = f_2 dx_1,$$

we find that

$$\oint_{\mathbb{R}} f_2(x_1, x_2) dx_1 - f_1(x_1, x_2) dx_2 = 0.$$

Let S denote the interior of γ . By Green's theorem, we would thus have

$$\iint_{S} \left(\frac{\partial f_{1}}{\partial x_{1}} + \frac{\partial f_{2}}{\partial x_{2}} \right) dx_{1} dx_{2} = 0.$$

However, our assumptions imply that either

$$\frac{\partial f_1}{\partial x_1} + \frac{\partial f_2}{\partial x_2} > 0 \quad \forall x \in D \qquad \text{or} \qquad \frac{\partial f_1}{\partial x_1} + \frac{\partial f_2}{\partial x_2} < 0 \quad \forall x \in D$$

everywhere in S. Hence, it must be

$$\iint_{S} \left(\frac{\partial f_1}{\partial x_1} + \frac{\partial f_2}{\partial x_2} \right) dx_1 dx_2 \neq 0,$$

and there can be no periodic orbit in D.

Green's Theorem

Let C be a positively oriented, piecewise smooth, simple closed curve in a plane, and let D be the region bounded by C. If L and M are functions of (x,y) defined on an open region containing D and having continuous partial derivatives there, then

$$\oint_C (Ldx + Mdy) = \iint_D \left(rac{\partial M}{\partial x} - rac{\partial L}{\partial y}
ight) dx dy$$

where the path of integration along ${\cal C}$ is **anticlockwise**

5. Bifurcations

For bifurcations, we consider parameterized system model:

$$\left[egin{array}{c} \dot{x}_1 \ \dot{x}_2 \end{array}
ight] = f\left(\left[egin{array}{c} x_1 \ x_2 \end{array}
ight], \mu
ight), \quad \mu \in \mathbb{R}, f ext{ continuously differentiable}$$

We care interested in the **influence of small changes in the parameter** μ **on changes in system behavior**

Conceptions

Definition: Bifurcations

The system $\dot{x} = f(x, \mu)$ <u>bifurcates</u> at the point $\mu = \mu_0$ if the phase portrait is **qualitatively different** for $\mu < \mu_0$ and $\mu > \mu_0$.

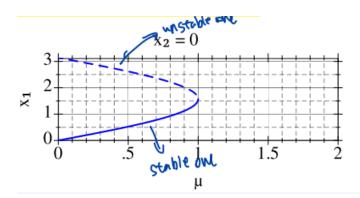
Bifurcation Diagram

A plot with the axes μ, x_1 and x_2 that shows the locations of the involved equilibrium points/limit cycles.

- Straight lines indicate stable equilibria (i.e., stable nodes, focuses or centers) or limit cycles;
- Dashed lines indicate unstable equilibria (i.e., unstable nodes, focuses or saddles) or limit cycles.

Note:

If all equilibrium points satisfy either $x_1=0$ or $x_2=0$, the bifurcation diagram becomes two dimensional



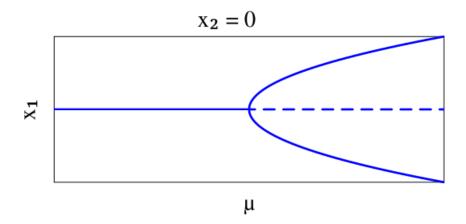
Type 1: Saddle-node Bifurcation

An unstable and a stable equilibrium collide and cancel

Pitchfork Bifurcations

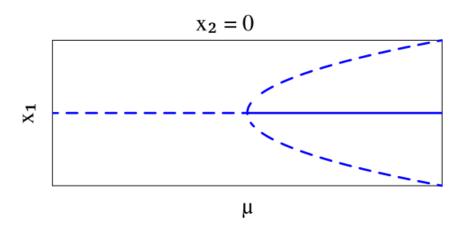
Type 2: Supercritical Pitchfork Bifurcation

A stable equilibrium becomes unstable at the bifurcation point and spawns two new stable equilibria



Type 3: Subcritical Pitchfork Bifurcation

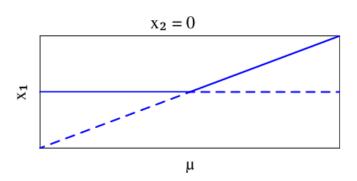
An unstable equilibrium point becomes stable at the bifurcation point and spawns two new unstable equilibria



Transcritical and Hopf Bifurcations

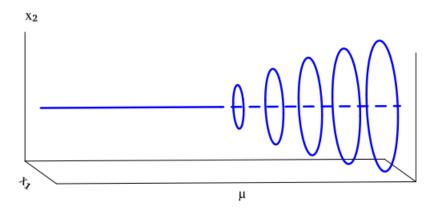
Type 4: Transcritical Bifurcation

A stable equilibrium and a non-stable equilibrium collide at the bifurcation point and exchange stability properties.



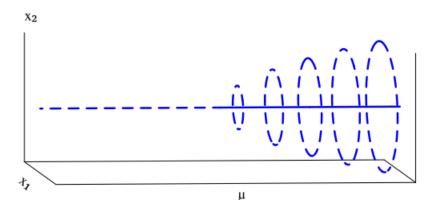
Type 5: Supercritical Hopf bifurcation

A stable focus turns into an unstable focus and spawns a stable limit cycle.



Type 6: Subcritical Hopf bifurcation

An unstable focus turns into a stable focus and spawns an unstable limit cycle



Finding Bifurcation Points

Assume $f(x,\mu)$ is continuously differentiable

Let \bar{x} be the equilibrium of the system $\dot{x}=f(x,\mu)$ at the bifurcation point $\mu=\mu^*$. Then \bar{x} is not hyperbolic. In particular:

- the equilibrium is a center for Hopf bifurcations
- at least one of the eigenvalues of the linearization is zero for the other bifurcation types discussed today.

Useful Result From Linear Systems Theory

Theorem: Routh-Hurwitz Stability Criterion for 2×2 Matrix

The eigenvalues of a matrix $A \in \mathbb{R}^{2 imes 2}$ are all in the open left half-plane $\mathbf{s}(\mathbf{s}) < 0$ if and only if

$$Tr(A) < 0$$
 and $det(A) > 0$.

Summary

- Investigated two-dimensional systems $\dot{x}_1 = f_1(x_1, x_2)$, $\dot{x}_2 = f_2(x_1, x_2)$
- Graphical representations: Vector field diagrams, phase portraits
- Method of isoclines for sketching phase portraits
- Investigated the equilibrium point $x^* = 0$ for a large class of linear time invariant systems (distinct non-zero eigenvalues)
- The equilibrium point can have on of the following types:
 - Stable / unstable node, saddle
 - Stable / unstable focus, center
- An equilibrium point x^* of the nonlinear system $\dot{x} = f(x)$ is hyperbolic of the same type if its linearization is hyperbolic at zero
- If the linearization has non-hyperbolic equilibrium at the origin, a nonlinear analysis has to be carried out
- Limit cycle = isolated periodic orbit
- Poincaré-Bendixson criterion → existence of periodic orbits
- Bendixson criterion

 non-existence of periodic orbits
 - Considered parameter-dependent systems $\dot{x} = f(x, \mu)$
 - At a bifurcation point $\mu = \mu_0$, the phase portrait changes qualitatively
 - The corresponding equilibrium is not hyperbolic at μ_0
 - Different types:
 - pitchfork
 - transcritical
 - Hopf