

# Chapter 4. Frequency Analysis: The Fourier Series

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# Introduction

- This chapter considers
  - **Spectral representation**: The frequency representation of periodic and aperiodic signals indicates **how their power or energy is allocated to different frequencies**.
    - The spectrum of a periodic signal is discrete, as its power is concentrated at frequencies **multiples of fundamental frequency**.
    - The spectrum of a aperiodic signal is a continuous function of frequency.
  - **Eigenfunctions and Fourier analysis**: **Complex exponentials and sinusoids** are used in the Fourier representation of signals by taking advantage of the **eigenfunction property of LTI systems**.
  - **Steady-state analysis**: **Fourier analysis is in the steady state**, while Laplace analysis considers both transient and steady state.
  - **Application of Fourier analysis**: The frequency representation of signals and systems is extremely **important in signal processing and in communications**.

# Eigenfunctions Revisited

If  $x(t) = e^{j\Omega_0 t}$ ,  $-\infty < t < \infty$  is the input to a **causal** and a **stable system** with **impulse response**  $h(t)$ , the output in the steady state is given by

$$y(t) = e^{j\Omega_0 t} H(j\Omega_0)$$

where

$$H(j\Omega_0) = \int_0^{\infty} h(\tau) e^{-j\Omega_0 \tau} d\tau$$

is the **frequency response** of the system at  $\Omega_0$ .

- The **input signal**  $x(t)$  is a **linear combination of complex exponentials**, with different amplitudes, frequencies, and phases, or

$$x(t) = \sum_k X_k e^{j\Omega_k t} \Rightarrow y(t) = \sum_k X_k e^{j\Omega_k t} H(j\Omega_k)$$

# Eigenfunctions Revisited

- The input signal is an **integral (a sum, after all)** of complex exponentials,

$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\Omega) e^{j\Omega t} d\Omega \Rightarrow y(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\Omega) e^{j\Omega t} H(j\Omega) d\Omega$$

- The **frequency response** is  $H(j\Omega) = H(s)|_{s=j\Omega}$ .
  - Its **magnitude** is an **even function of frequency**, or  $|H(j\Omega)| = |H(-j\Omega)|$ .
  - Its **phase** is an **odd function of frequency**, or  $\angle H(j\Omega) = -\angle H(-j\Omega)$ .
- One application of LTI systems is **filtering**, where one is interested in **preserving desired frequency** and **getting rid of less-desirable components** of a signal.

# Phasor Interpretation of Eigenfunction Property

For a **stable LTI system** with **transfer function**  $H(s)$ , if the input is

$$x(t) = A \cos(\Omega_0 t + \theta) = \operatorname{Re}[X e^{j\Omega_0 t}]$$

where  $X = A e^{j\theta}$  is the phasor of  $x(t)$ , the steady-state output of the system is

$$\begin{aligned} y(t) &= \operatorname{Re}[X H(j\Omega_0) e^{j\Omega_0 t}] = \operatorname{Re}[A H(j\Omega_0) e^{j(\Omega_0 t + \theta)}] \\ &= A |H(j\Omega_0)| \cos(\Omega_0 t + \theta + \angle H(j\Omega_0)) \end{aligned}$$

where the **frequency response** of the system at  $\Omega_0$  is

$$H(j\Omega_0) = H(s) \Big|_{s=j\Omega} = Y/X$$

[Ex 4.1] Consider the RC circuit with  $R = 1\Omega$  and  $C = 1\text{F}$ . Let voltage source be  $v_s(t) = 4 \cos(t + \pi/4)$  volts, find the steady-state voltage across the capacitor.

# Complex Exponential Fourier Series

- **Fourier series** is a **representation of a periodic signal  $x(t)$**  in terms of complex exponentials or sinusoids of **frequency multiples of its fundamental frequency**.
  - Consider a set of complex functions  $\{\psi_k(t)\}$  defined in  $[a, b]$ . Let's say  $\psi_l(t)$  and  $\psi_m(t)$ ,  $l \neq m$ , is called **orthonormal**, iff
$$\int_a^b \psi_l(t) \psi_m^*(t) dt = \begin{cases} 0 & l \neq m \\ 1 & l = m \end{cases}$$
  - A finite energy signal  $x(t)$  defined in  $[a, b]$  can be approximated by a series

$$\hat{x}(t) = \sum_k a_k \psi_k(t)$$

- We consider a **periodic signal  $x(t)$**  such that
  - It is defined for  $-\infty < t < \infty$ .
  - For any integer  $k$ ,  $x(t + kT_0) = x(t)$ , where  $T_0$  is the **fundamental period**.

# Complex Exponential Fourier Series

The **Fourier series representation** of a periodic signal  $x(t)$  of period  $T_0$  is an **infinite sum of weighted complex exponentials** with **frequency multiples of the signal's fundamental frequency  $\Omega_0 = 2\pi/T_0$** , or

$$x(t) = \sum_{k=-\infty}^{\infty} X_k e^{jk\Omega_0 t} \quad \Omega_0 = \frac{2\pi}{T_0}$$

where the **Fourier coefficients  $X_k$**  are found according to

$$X_k = \frac{1}{T_0} \int_{t_0}^{t_0+T_0} x(t) e^{-jk\Omega_0 t} dt$$

for  $k = 0, \pm 1, \pm 2, \dots$  and any  $t_0$ .



# Parseval's Theorem

- Periodic signals are **infinite-energy finite-power signals**.
  - **Fourier series** provides a way to find **how much of the signal power is in a certain band of frequencies**.

The **power**  $P_x$  of a periodic signal  $x(t)$  of period  $T_0$  can be equivalently calculated in either the time or the frequency domain:

$$P_x = \frac{1}{T_0} \int_{T_0} |x(t)|^2 dt = \sum_k |X_k|^2$$

- The power of the signal is distributed over the **harmonic frequencies**  $\{k\Omega_0\}$ .
  - Given the **discrete nature** of the harmonic frequencies, this plot consists of a line at each frequency, which is called the **power line spectrum**.

# Symmetry of Line Spectra

For a **real-valued periodic signal**  $x(t)$  of period  $T_0$ , the **Fourier coefficients**  $\{X_k = |X_k|e^{j\angle X_k}\}$  at **harmonic frequencies**  $\{k\Omega_0 = 2\pi k/T_0\}$  we have that  $X_k = X_{-k}^*$ . Equivalently,  $|X_k| = |X_{-k}|$  (**even function**) and  $\angle X_k = -\angle X_{-k}$  (**odd function**).

# Trigonometric Fourier Series

The **trigonometric Fourier series** of a **real-valued periodic signal**  $x(t)$  of period  $T_0$ , is given by

$$\begin{aligned} x(t) &= X_0 + 2 \sum_{k=1}^{\infty} |X_k| \cos(k\Omega_0 t + \theta_k) \\ &= X_0 + 2 \sum_{k=1}^{\infty} [c_k \cos(k\Omega_0 t) + d_k \sin(k\Omega_0 t)] \quad \Omega_0 = \frac{2\pi}{T_0} \end{aligned}$$

where  $X_0 = c_0$  is the **DC component** and  $2|X_k| \cos(k\Omega_0 t + \theta_k)$  are the  **$k$ -th harmonics**.

$$\begin{aligned} c_k &= \frac{1}{T_0} \int_{t_0}^{t_0+T_0} x(t) \cos(k\Omega_0 t) dt \quad k = 0, 1, \dots \\ d_k &= \frac{1}{T_0} \int_{t_0}^{t_0+T_0} x(t) \sin(k\Omega_0 t) dt \quad k = 1, 2, \dots \end{aligned}$$

# Example

[Ex 4.4] Find the Fourier series of a raised-cosine signal ( $B \geq A$ ),

$$x(t) = B + A \cos(\Omega_0 t + \theta)$$

which is periodic with period  $T_0$  and fundamental frequency  $\Omega_0 = 2\pi/T_0$ .

Call

$$y(t) = B + A \cos(\Omega_0 t - \pi/2),$$

then find its Fourier series coefficients and compare them to those for  $x(t)$ .

Use symbolic MATLAB to compute the Fourier series of  $y(t) = 1 + \sin(100t)$ , find and plot its magnitude and phase line spectra.

# Fourier Coefficients from Laplace

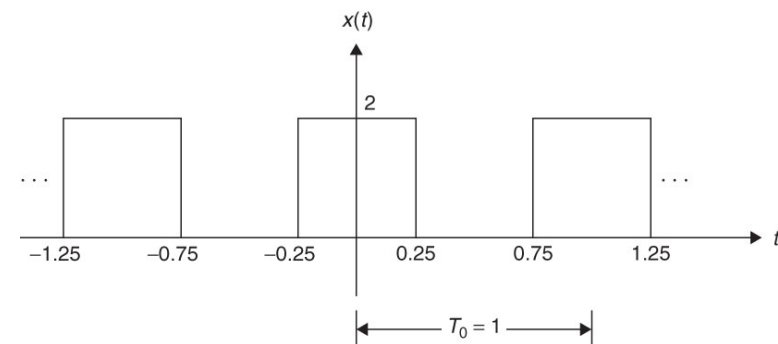
For a **periodic signal**  $x(t)$  of period  $T_0$ , if we know or easily **compute the Laplace transform of a period of  $x(t)$**

$$x_1(t) = x(t)[u(t - t_0) - u(t - t_0 - T_0)]$$

Then the **Fourier coefficients of  $x(t)$**  are given by

$$X_k = \frac{1}{T_0} \mathcal{L}[x_1(t)]_{s=jk\Omega_0} \quad \Omega_0 = \frac{2\pi}{T_0}$$

[Ex 4.5] Consider the periodic pulse train  $x(t)$  of period  $T_0 = 1$ , shown below. Find its Fourier series.



# Reflection and Even/Odd Periodic Signals

The **Fourier coefficients** of  $x(-t)$  are  $\{X_{-k}\}$ .

The **Fourier coefficients**  $X_k$  of even periodic signal  $x(t)$  are **real**, and its **trigonometric Fourier series** is

$$x(t) = X_0 + 2 \sum_{k=1}^{\infty} X_k \cos(k\Omega_0 t)$$

The **Fourier coefficients**  $X_k$  of odd periodic signal  $x(t)$  are **imaginary** and its **trigonometric Fourier series** is

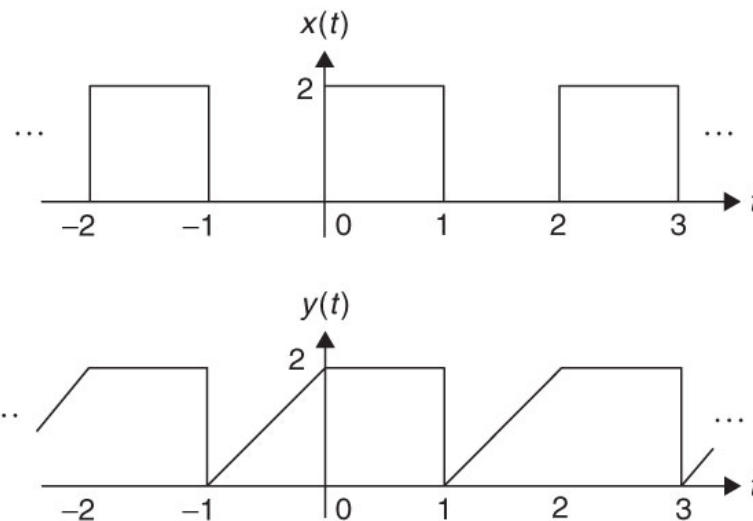
$$x(t) = 2 \sum_{k=1}^{\infty} jX_k \sin(k\Omega_0 t)$$

For any periodic signal  $x(t) = x_e(t) + x_o(t)$ , then

$$X_k = X_{ek} + X_{ok} = 0.5[X_k + X_{-k}] + 0.5[X_k - X_{-k}]$$

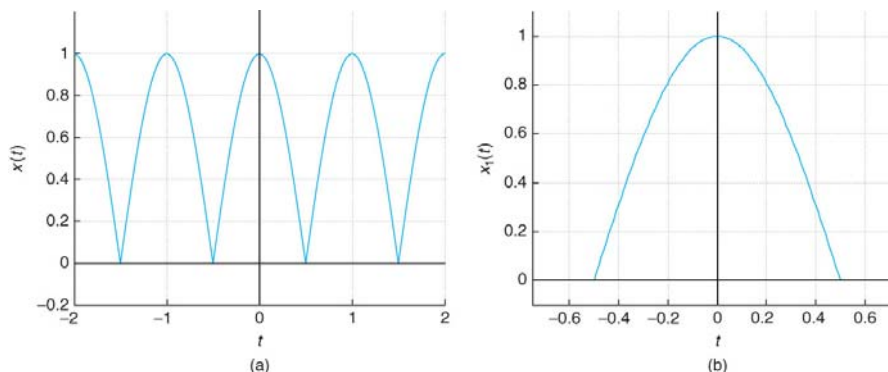
# Examples

[Ex 4.6] Consider the periodic signals  $x(t)$  and  $y(t)$  shown in below. Determine their Fourier coefficients by using the symmetry conditions and the even-odd decomposition.

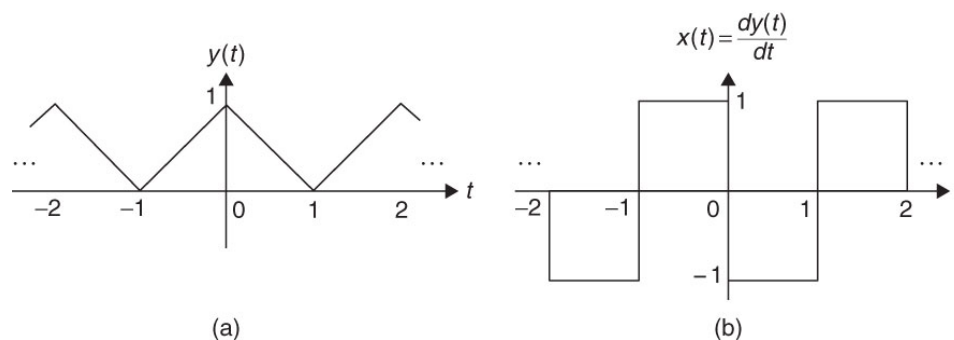


# Examples

[Ex 4.7] Find the Fourier series of full-wave rectified signal  $x(t) = |\cos \pi t|$  shown below.



[Ex 4.8] Consider a train of triangular pulses  $y(t)$  with fundamental period  $T_0 = 2$ . Let  $x(t) = dy(t)/dt$ . Find its Fourier series and compare  $|X_k|$  and  $|Y_k|$  to determine which of the signals is smoother.





# Convergence of Fourier Series

- A signal  $x(t)$  is said to be **piecewise smooth** if it has a **finite number of discontinuities**, while a **smooth signal** has a **derivative** that changes continuously.

The Fourier series of a **piecewise smooth periodic signal**  $x(t)$  converges for all values of  $t$ . The **sufficient condition of the Fourier series convergence** is (over a one period)

- Be **absolutely integrable**
- Have a **finite number of maxima, minima, and discontinuities**.

The infinite series **equals**  $x(t)$  **at every continuity point** and **equals the average**  $0.5[x(t + 0^+) + x(t + 0^-)]$  of the right limit  $x(t + 0^+)$  and the left limit  $x(t + 0^-)$  **at every discontinuity point**.

# Gibb's Phenomenon

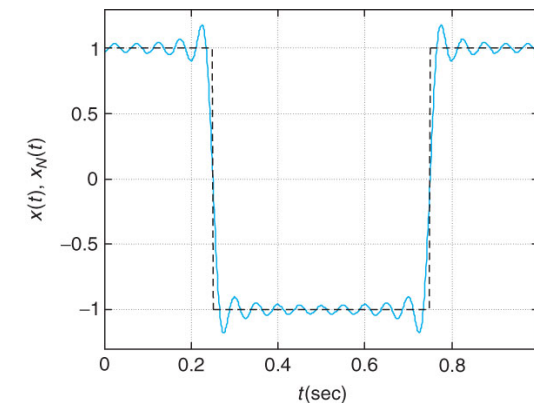
- Although the Fourier series converges to **the arithmetic average at discontinuities**, there is a **ringing** before and after the discontinuity point, called **Gibb's phenomenon**.

- **The smoother the signal  $x(t)$  is, the easier it is to approximate it** with a Fourier series with a finite number of terms.
- The  $N$ -th order approximation of a periodic signal  $x(t)$  is

$$X_N(t) = \sum_{k=-N}^N X_k e^{jk\Omega_0 t}$$

- The average quadratic error over a period is

$$E_N = \frac{1}{T_0} \int_{T_0} [x(t) - x_N(t)]^2 dt$$

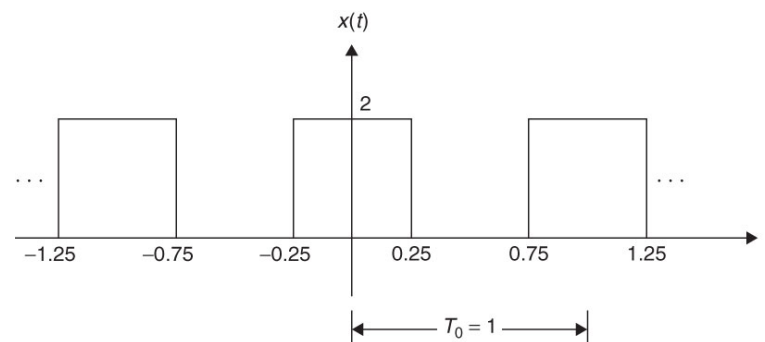


$N = 20$

# Examples

[Ex 4.11] We wish to approximate  $x(t)$  below by  $x_2(t) = \alpha + 2\beta \cos \Omega_0 t$ . Compute  $\alpha$  and  $\beta$  to minimize the mean-square error

$$E_2 = \frac{1}{T_0} \int_{T_0} [x(t) - x_2(t)]^2 dt$$



# Time and Frequency Shifting

**Time shifting** causes **only a change in phase**.

- A periodic signal  $x(t)$  of period  $T_0$  remains periodic of the same period when shifted in time. If  $X_k$  are the Fourier coefficients of  $x(t)$ , the **Fourier coefficients for  $x(t - t_0)$**  are

$$\{X_k e^{-jk\Omega_0 t_0} = |X_k| e^{j(\angle X_k - k\Omega_0 t_0)}\}$$

When a periodic signal  $x(t)$  of period  $T_0$  **modulates a complex exponential  $e^{j\Omega_1 t}$  (frequency shifting)**.

- If  $\Omega_1 = M\Omega_0$  for an integer  $M \geq 1$ , the modulated signal  $x(t)e^{j\Omega_1 t}$  is **periodic with period  $T_0$** .
- The **Fourier coefficients  $X_k$**  are shifted to frequencies  $k\Omega_0 + \Omega_1$ .
- The modulated signal is real-valued by  $x(t) \cos \Omega_1 t$ .

# Response of LTI Systems to Periodic Signals

If the input  $x(t)$  of a causal and stable LTI system, with impulse response  $h(t)$ , is periodic of period  $T_0$  and has the Fourier series

$$x(t) = X_0 + 2 \sum_{k=1}^{\infty} |X_k| \cos(k\Omega_0 t + \angle X_k) \quad \Omega_0 = \frac{2\pi}{T_0}$$

the **steady-state response** of the system is

$$y(t) = X_0 |H(j0)| \cos(\angle H(j0)) + 2 \sum_{k=1}^{\infty} |X_k| |H(jk\Omega_0)| \cos(k\Omega_0 t + \angle X_k + \angle H(jk\Omega_0))$$

where

$$H(jk\Omega_0) = \int_0^{\infty} h(\tau) e^{-jk\Omega_0 \tau} d\tau$$

is the frequency response of the system at  $k\Omega_0$ .

## Sum of Periodic Signals

If  $x(t)$  and  $y(t)$  are periodic signals with the same fundamental frequency  $\Omega_0$ . The **Fourier series coefficients of  $z(t) = \alpha x(t) + \beta y(t)$**  are  **$Z_k = \alpha X_k + \beta Y_k$** .

If  $x(t)$  is periodic of period  $T_1$  and  $y(t)$  is periodic of period  $T_2$  such that  $T_2/T_1 = N/M$  for non-divisible integers  $N$  and  $M$ . The **Fourier series coefficients of  $z(t) = \alpha x(t) + \beta y(t)$**  are periodic of period  **$T_0 = MT_2 = NT_1$** , are  **$Z_k = \alpha X_{k/N} + \beta Y_{k/M}$**  for **integers  $k$  such that  $k/N$ , and  $k/M$  are integers**.

[Ex 4.15] Consider the sum  $z(t)$  of a periodic signal  $x(t)$  of period  $T_1 = 2$ , with a periodic signal  $y(t)$  with period  $T_2 = 0.2$ . Find the Fourier coefficients  $Z_k$  of  $z(t)$  in terms of the Fourier coefficients  $X_k$  and  $Y_k$ .

# Multiplication of Periodic Signals

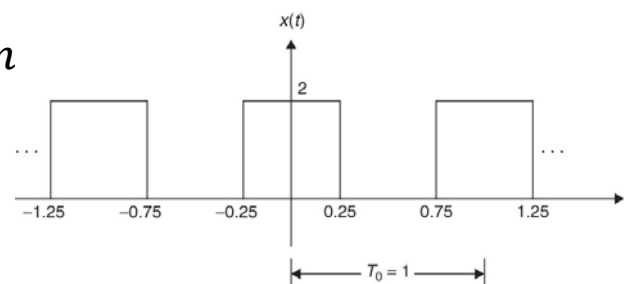
If  $x(t)$  and  $y(t)$  are periodic signals of the same period  $T_0$ , then their product  $z(t) = x(t)y(t)$  is also periodic of period  $T_0$ , and with Fourier coefficients that are the convolution sum of the Fourier coefficients of  $x(t)$  and  $y(t)$ :

$$Z_k = \sum_l X_l Y_{k-l}$$

[Ex 4.16] Consider the train of rectangular pulses  $x(t)$  shown below. Let  $z(t) = 0.25 x^2(t)$ . Use the Fourier series  $z(t)$  to show that

$$X_k = \alpha \sum_m X_m X_{k-m}$$

for some constant  $\alpha$ . Determine  $\alpha$ .



# Derivatives and Integrals of Periodic Signals

The derivative  $dx(t)/dt$  of periodic signal  $x(t)$ , of period  $T_0$ , is periodic of the same period  $T_0$ . If  $X_k$  are the **coefficients of the Fourier series of  $x(t)$** , the **Fourier coefficients of  $dx(t)/dt$**  are  $jk\Omega_0 X_k$ .

For a **zero-mean**, periodic signal  $y(t)$ , of period  $T_0$ , the **integral**

$$z(t) = \int_{-\infty}^t y(\tau) d\tau$$

is periodic of the same period as  $y(t)$ , with **Fourier coefficients**

$$Z_k = \frac{Y_k}{jk\Omega_0} \quad k \neq 0, \quad Z_0 = - \sum_{m \neq 0} Y_m \frac{1}{jm\Omega_0} \quad \Omega_0 = \frac{2\pi}{T_0}$$



# Examples

[Ex 4.17] Let  $g(t)$  be the derivative of a triangular train of pulses  $f(t)$ , of period  $T_0 = 1$ . The period of  $f(t)$ ,  $0 \leq t \leq 1$ , is

$$f_1(t) = 2r(t) - 4r(t - 0.5) + 2r(t - 1)$$

Use the Fourier series of  $g(t)$  to find the Fourier series of  $f(t)$ .

[Ex 4.18] Find the Fourier coefficients  $Z_k$  of  $z(t)$

$$z(t) = \int_{-\infty}^t g(\tau) d\tau$$

where  $g(t)$  is defined in [Ex 4.17]



*Thank You*