

# Eddie Tang LLP Report

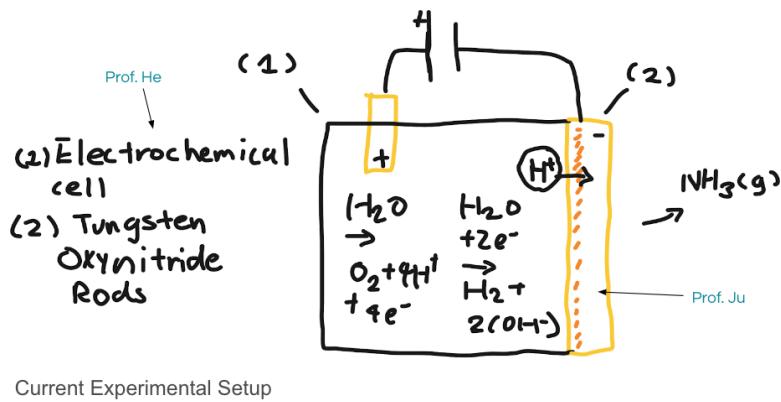
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## 1. Introduction

Green energy is vital in the future of technological advancements and in combating climate change. The hydrogen molecule ( $H_2$ ) serves as an effective energy source that satisfies the requirement for green energy chiefly from its ability to combust cleanly by only forming water and its high gravimetric energy density. However, hydrogen transportation has its challenges and setbacks. Ammonia ( $NH_3$ ) is a possible solution since it has a high energy density, is a carrier of  $H_2$ , and is easier to store at room temperature.

Ammonia is synthesized on a large scale from hydrogen and nitrogen through the Haber-Bosch process. Moreover, since ammonia production requires a large amount of energy to break the nitrogen ( $N_2$ ) triple bond in the  $NH_3$  synthesis reaction and is characterized by centralized facilities to carry out this process, it accounts for 2% of the global energy consumption.

Current developments<sup>1</sup> in combating these downfalls have found that combining electrochemical cells with tungsten oxynitride rods ( $WO_xN_y$ ) and plasma-assisted nitrogen synthesis yielded promising results. Its advantages are that it is a distributed method of synthesizing  $NH_3$  by being compact and relatively cheap, it has a higher yield of up to 200% compared to current  $NH_3$  fuel cells, and it can use surplus energy that usually goes to waste, effectively becoming an electrified method in producing  $NH_3$ .



However, even with experimental results, there is a lack of understanding of the elementary reactions that dictate the many processes in our experiments. Namely, the process of creating  $WO_xN_y$  rods by blasting tungsten oxide ( $WO_3$ ) with hydrogen plasma and then  $N_2$  plasma needs to be better understood from purely experimental analysis, which is the basis of

<sup>1</sup> Our Research Project

question one. Moreover, question two is whether the O-vacancies in the  $\text{WO}_x\text{N}_y$  rods would prefer water ( $\text{H}_2\text{O}$ ) or  $\text{N}_2$  needs to be clearly understood.

Here, the objective of this project is to investigate and study these elementary reactions through the use of density functional theory (DFT) calculations. These calculations will be used to calculate the absorption energies of  $\text{H}$ ,  $\text{N}_2$ ,  $\text{H}_2\text{O}$ , and  $\text{H}_2$  adsorbates to determine chemically and thermodynamically favorable reaction pathways that explain the processes mentioned above.

Lastly, another project that is underway is to study the interaction of nitrogen species and  $\text{WO}_3$  surface vacancies through a machine-learned potential energy surface.

## 2. DFT Methods

### 2.1. Atomic models

A four-layer  $\text{WO}_3$  periodic slab ( $S$ ,  $1 \times 1$ ) of the lowest-energy surface (001) and a 10 Å vacuum layer on top of the slab surface were generated. Therefore, the optimized supercell dimensions were  $7.4984 \text{ \AA} \times 7.4984 \text{ \AA} \times 27 \text{ \AA}$ . Calculating the in-plane ( $\alpha$ ) and out-of-plane ( $\beta$ ) angles I get  $120^\circ$  and  $90^\circ$ , respectively. For gaseous species, namely  $\text{H}$ ,  $\text{N}_2$ ,  $\text{H}_2\text{O}$ , and  $\text{H}_2$ , a  $10 \text{ \AA} \times 11 \text{ \AA} \times 12 \text{ \AA}$  periodic cubic box was used when sampling their free energies for calculations. Additionally, for simulations not involving vacancies, the adsorbates were placed 1.5 Å above the surface. For simulations involving vacancies, the adsorbates were placed 0.5 Å above the surface.

To generate the initial configurations, a custom Python script<sup>2</sup> was programmed using the Atomic Simulation Environment (ASE) Python package and Pymatgen (Python Materials Genomics). This script can add our gaseous species on top of a set location, generate vacancies on the  $\text{WO}_3$  surface, generate gaseous species in the set vacuums, and generate KPOINTS files.

To study the effects further, a  $2 \times 2$  supercell slab was generated with Open Visualization Tool (OVITO) to generate a larger slab ( $LS$ ). This slab has supercell dimensions of  $14.9968 \text{ \AA} \times 14.9968 \text{ \AA} \times 27 \text{ \AA}$ . Furthermore,  $S$  was replicated in OVITO once in the  $y$  direction to generate a medium slab ( $MS$ ,  $1 \times 2$ ). This slab has supercell dimensions of  $7.498 \text{ \AA} \times 14.9968 \text{ \AA} \times 27 \text{ \AA}$  after relaxation.

Both relaxed surfaces were used to study interactions between  $\text{N}_2$ ,  $\text{H}_2\text{O}$ , and O vacancies on the  $\text{WO}_3$  surface. Note that this structure is based on the  $\text{WO}_3$  hexagonal crystal structure. For  $\text{WO}_3$  (001), I used the O-terminated ( $O_t$ ) side for surface simulations since it is more stable than the W-terminated side.

### 2.2. Electronic structure theory

DFT calculations were performed using the Vienna Ab-initio Simulation Package (VASP) version 6.3.1. In these calculations, Perdew-Burke-Ernzerhof (PBE) exchange-correlation (XC) functional were used. The wave functions were expanded in a plane wave basis set with a kinetic energy cutoff of 500 eV (ENCUT=500). The Brillouin zone (BZ) was sampled so that the sampling grid preserves the symmetry in the reciprocal space, specifically, for  $S$ , by  $4 \times 4 \times 1$ , for

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<sup>2</sup> <https://github.com/EDED2314/VASP-scripts>

*LS*, by  $2 \times 2 \times 1$ , and for *MS*, by  $4 \times 2 \times 1$  Gamma-centered meshes. The dipole moment will be calculated only parallel to the direction of the third lattice vector in our simulations (IDIPOL=3). Simulations were run involving H adsorbates with spin-polarization to account for the free electrons around H. Other simulations were run without spin-polarization. Forces and stress tensors were calculated to allow the  $\text{WO}_3$  structure's position degree of freedom to change (ISIF=2). The Gaussian method was used to smear the electronic states near the Fermi level with a width of 0.05 to aid electronic convergence (ISMEAR=0, SIGMA=0.05). Lastly, van der Waals corrections were not included.

### 2.3. Structure optimization

A conjugate-gradient algorithm was used for optimization and the self-consistent-field (SCF) threshold was set as 1 meV. I also set the atomic-force component convergence thresholds to 0.01 eV/Å. The bottom two layers of the slab were frozen and the structures were relaxed. Hence, for all adsorbate simulations, I only allowed the top two layers to interact with the adsorbate to simulate a semi-infinite crystal. The calculation factor for the default values for ENCUT (PREC) was set to accurate. Lastly, the maximum number of geometry optimization steps is 100, the maximum number of electronic SC (self-consistency) steps is 100, and the minimum number of electronic SC steps is 10.

### 2.4. Thermal corrections to the free energy

Since our DFT calculations are done under 0K, I used additional thermal corrections to study the effects of temperature on the  $\text{WO}_3$  surface on one  $\text{H}_2\text{O}$  molecule's desorption energy. Typically, the free energy of a gas-phase species is its DFT total internal energy and the sum of translational, vibrational, and rotational vibrations' free energies.<sup>3</sup> Since I am working with slabs, I ignored the translational and rotational vibrations in calculations involving them.

The oxygen terminations, first layer of tungsten, and adsorbate were displaced by 0.015 Å in VASP simulations that do not apply symmetry and calculate the Hessian (IBRION = 5). All of these displacements were performed on optimized structures.

Once the simulation concluded, I used VASPKIT<sup>4</sup> to analyze the vibrational frequencies by deriving the normal mode frequencies from output Hessians in the OUTCAR and outputting free energy corrections for an inputted temperature.

## 3. DFT Results and Discussion

### 3.1. The influence of $\text{H}_2$ plasma in oxygen vacancy formation

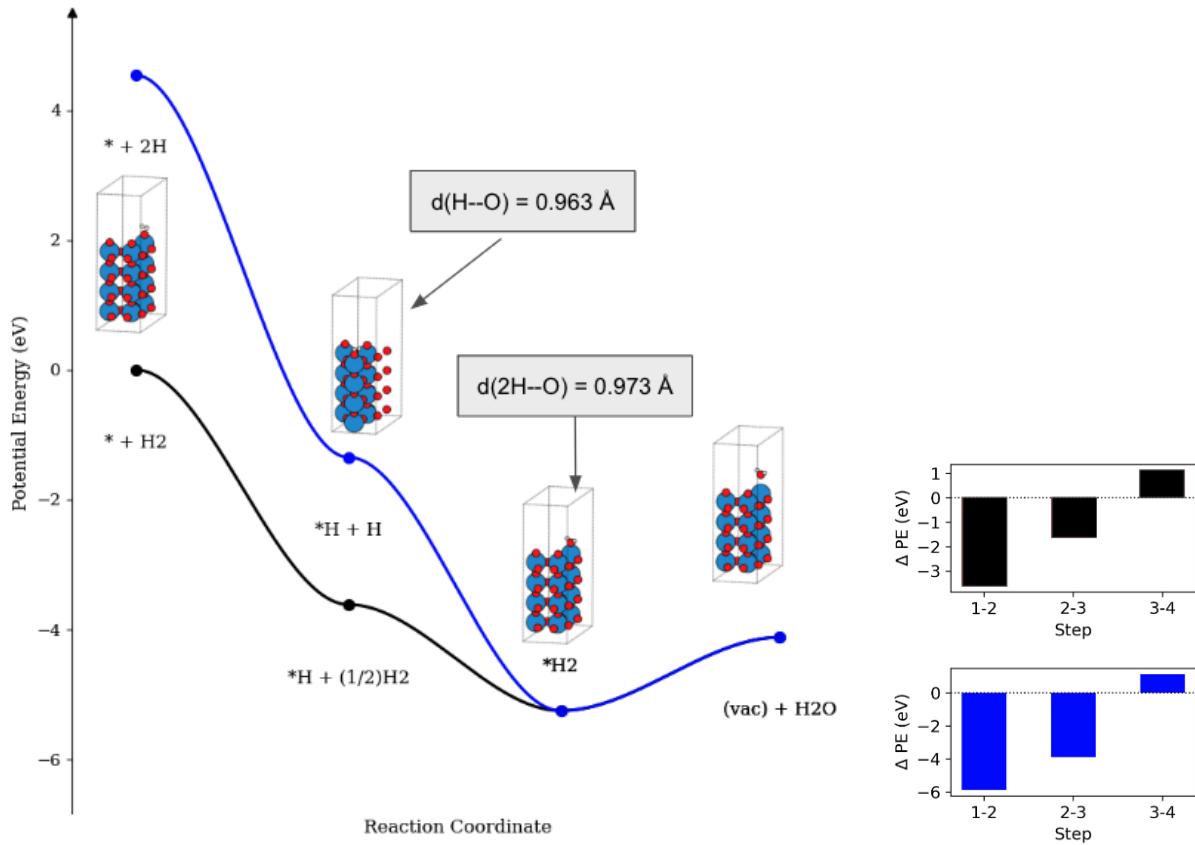
To investigate the first question regarding how  $\text{H}_2$  plasma formed oxygen vacancies, I wanted to calculate the energies of each reaction step to evaluate the chemical favorability of two

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<sup>3</sup> Martirez, J. M. P.; Carter, E. A. First-Principles Insights into the Thermocatalytic Cracking of Ammonia-Hydrogen Blends on Fe(110): 1. Thermodynamics. *J. Phys. Chem. C* 2022, 126 (46), 19733– 19744, doi: 10.1021/acs.jpcc.2c06003

<sup>4</sup> <https://vaspkit.com/>

hydrogen atoms, or one H<sub>2</sub> molecule to absorb on a surface oxygen atom along the oxygen terminations sites and then desorb from the surface of S to form a vacancy and water molecule.



**Figure 1.** *Left:* Potential surface of surface and H<sub>2</sub> and H reaction at 0 K, *Right:* Potential energy differences between each step.

After verifying the convergence of all structures involved in the specific reaction steps, I extracted the energy of the configurations and plotted them along a potential surface graph in Figure 1 and scaled everything relative to the energy of the surface + H<sub>2</sub>. For steps 2 and 3, the bond distance between \*H and O<sub>t</sub> in the optimized structure is included to indicate a typical hydroxide bond length and water bond length.

Note that the H<sub>2</sub> surface does not include the reaction steps involved in the dissociation of H<sub>2</sub> into two H atoms, which is where we would expect an energy barrier that would require external work to be applied to the system to advance the reaction.

Without including that reaction pathway, we can still infer, from the difference in magnitude of change in potential energy between the reaction steps, that H atoms adsorbing onto O-terminations are more chemically favorable than sourcing \*H from H<sub>2</sub>. This demonstrates that hydrogen plasma, which gives the system access to individual H atoms, would have more vacancy formation than only H<sub>2</sub> gas.

### 3.2. H<sub>2</sub>O desorption energy barrier

For both pathways in Figure 1, there is a slight uphill (~1 eV) desorption energy of H<sub>2</sub>O from the surface. Since room temperature, 300K, can only account for 0.03 eV in contribution, other factors could lower this desorption energy barrier. Therefore, I investigated the effects of temperature on this energy barrier and the effects of H<sub>2</sub>O coverage concerning this uphill magnitude as well.

#### 3.2.1. Effects of Temperature

I calculated the desorption energies by finding the negative of the adsorption energy. Namely,

$$\Delta E_{desorption} = -\Delta E_{ads} = -[E_{WO3/X} - (E_{WO3} + E_X)].$$

Where X is the adsorbate chemical symbol.

The free energy corrections were then calculated with VASPKIT at a specific temperature and added to the calculated internal potential energy found in VASP outputs to create the table below.

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**Table 1. Desorption energy calculated with free energy corrected values at corresponding temperatures.**

Temperature	H <sub>2</sub> O Desorption Energy (eV)
0K	1.1285
400K	1.1528
3000K	2.9848

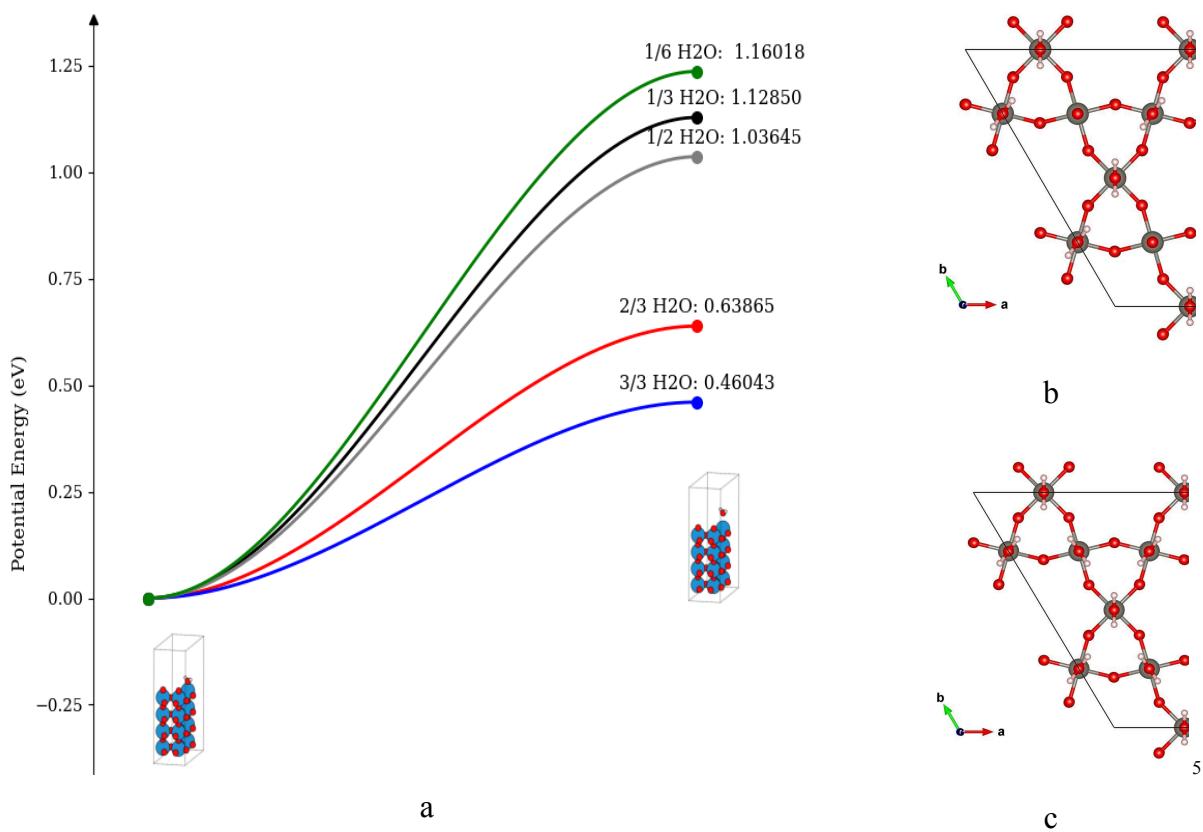
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The increasing desorption energies with an increase in temperature do indicate that the reaction step \* H<sub>2</sub> → H<sub>2</sub>O + vac is slightly exothermic. Additionally, since the final goal of the study is to develop a system that would operate at room temperature, the 3000K serves only as a reference for the desorption energy at higher temperatures. Lastly, the 0.0243eV difference between 0K and 400K indicates that temperature does not significantly affect the desorption temperatures, and hence the amount of vacancies formed.

The vacancy formation energy of S is 1.60426 eV and calculated with the equation

$$\Delta E_{vac} = E_{WO3} - (E_{WO3-v} + \frac{1}{2}E_{O2}).$$

This vacancy formation energy is 0.45146 eV greater than the desorption energy of H<sub>2</sub>O at 400 K, which suggests that even though H<sub>2</sub>O has positive adsorption energy, it is still more energetically favorable to desorb an H<sub>2</sub>O molecule than remove an oxygen atom along the surface to form a vacancy. Thus, supporting that H<sub>2</sub> plasma aids in forming O<sub>t</sub> vacancies.



**Figure 2.** a) Comparison of H<sub>2</sub>O desorption energies at different coverage of the O<sub>t</sub> sites. b) 3/3 coverage on the 001 surface. c) 2/3 coverage on the 001 surface.

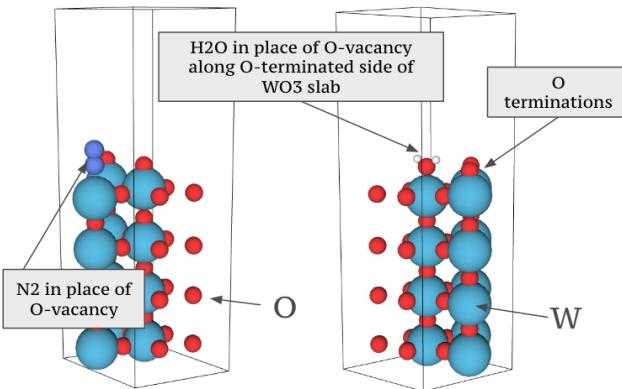
### 3.2.2. Effects of H<sub>2</sub>O Coverage

This study compared the desorption energies of one H<sub>2</sub>O molecule with its oxygen facing down at O<sub>t</sub> at different adsorbed H<sub>2</sub>O coverages to study the effects of H<sub>2</sub>O coverage on these sites. Specifically, we conducted the simulations on *S*.

From *Figure 2*, the desorption energy of H<sub>2</sub>O at minimal coverage (1/6) is 1.16 eV, which is more than two times the desorption energy of full coverage (3/3) 0.46 eV. There is also decreasing desorption energy as H<sub>2</sub>O coverage on the oxygen termination sites increases. Hence, the coverage effects on combating the H<sub>2</sub>O desorption energy barrier are non-negligible and significant.

Qualitatively examining these results also yields strengthened confidence, since more water absorbed on the surface will mean many more hydrogen-bond interactions and electron repulsion between the partial-negatives and partial-positives. It is hypothesized that these electron repulsions contribute to the smaller desorption energy.

<sup>5</sup> K. Momma and F. Izumi, "VESTA 3 for three-dimensional visualization of crystal, volumetric and morphology data," *J. Appl. Crystallogr.*, 44, 1272-1276 (2011).



**Figure 3.** Initial configurations. Left:  $\text{N}_2$  on top of  $\text{O}_t$  vacancy. Right:  $\text{H}_2\text{O}$  on top of  $\text{O}_t$  vacancy.

**Table 2. Calculated adsorption energies of  $\text{H}_2\text{O}$  and  $\text{N}_2$  into  $\text{O}_t$  vacancies.**

Adsorbate	Adsorption Energy (eV)
$\text{H}_2\text{O}$	-1.1285 <sup>6</sup>
$\text{N}_2$	-0.3556

Thus, I hypothesize that as the reaction continues with more atomic H adsorbed onto  ${}^*\text{H}$  and  ${}^*$  to form  ${}^*\text{H}_2$  ( ${}^*\text{H}_2\text{O} + \text{vac}$ ) and  ${}^*\text{H}$  respectively, there would be an increased rate at which vacancies form since the reaction barrier for  $\text{H}_2\text{O}$  desorption decreases with more adsorbed water. Overall, it is clear that higher  $\text{H}_2\text{O}$  coverage will contribute to more vacancies.

### 3.3. $\text{H}_2\text{O}$ versus $\text{N}_2$ adsorption onto $\text{O}_t$ vacancies

The second question is looked into by simply calculating the adsorption energy of water versus nitrogen at 1/3 coverage on S at an  $\text{O}_t$  vacancy site. The nitrogen molecule in its initial configuration is oriented upright while the water molecule is oriented so that its oxygen atom points downwards into the vacancy.

In *Table 2*, the calculated adsorption energy of  $\text{H}_2\text{O}$  is less than  $\text{N}_2$ , which means that under conditions listed in the methods,  $\text{O}_t$  vacancies will prefer water more than nitrogen since the adsorption is more chemically favorable. However, it is necessary to note that accounting for the electrochemical origins of water in the simulation environment may significantly affect these values since the hydrogen bonding may lead water to be less attracted to surface vacancies. Additionally, the water adsorption energy may decrease due to coverage effects, which have not been tested yet.

These results reveal that future experimental setups may find it more difficult to treat  $\text{WO}_3$  with  $\text{N}_2$  plasma to obtain  $\text{WO}_x\text{N}_y$  ( $x+y < 3$ ) or  $\text{W}_x\text{N}_y$ .

<sup>6</sup> This result is similar to hexagonal calculations done by [other literature](#), where the adsorption energy of  $\text{H}_2\text{O}$  is -1.22 eV.

### 3.4. Additional results

#### 3.4.1. Spin-Polarized Calculations

I have attached below the spin-polarized calculations. They yield significantly different results, which may be due to the electrons within the WO<sub>2</sub> structure itself being fully accounted for in the spin-polarized calculations.

Adsorption/desorption (for H<sub>2</sub>O) energies

	SP	no SP
n-o	-6.300231	-7.352501
no-vac	-0.839013	-1.777893
n2-o	-1.521460	-1.624960
n2o-vac	-0.209161	-0.199271
n2-bridge0	-0.030220	-0.032640
h-00	-4.743186	-5.883536
n-vac	-0.264831	-1.171646
n2-vac	-0.364560	-0.355680
h2o 1/3	1.134578	1.128501
h2o 2/3	0.405528	0.638658
h2o 3/3	0.055608	0.460138
vac_form	1.490868	1.604258

Energy of systems, and comparison to whether SP decreased it.

	SP	no SP	SP < no SP?
wo3 energy	-349.51698	-348.464710	True
wo3 v energy	-346.06652	-345.127640	True
2 h2o 1 vac energy	-355.52345	-354.450970	True
3 h2o 1 vac energy	-362.21825	-360.323120	True
1 h2o energy	-361.43218	-360.487223	True
2 h2o energy	-370.16006	-369.320710	True
3 h2o energy	-376.50494	-375.014340	True
n o0 energy	-358.94094	-358.940940	False
n2 o0 energy	-367.68835	-366.739580	True
n2 bridge0 energy	-366.19711	-365.147260	True
h energy	-355.37725	-355.465330	False
n vac energy	-349.45508	-349.423015	True
n2 vac energy	-363.08099	-362.133230	True

#### 3.4.2. Effect of N<sub>2</sub> coverage on N<sub>2</sub> adsorption energies

Since the nitrogen adsorption energy was more than two times less than the water adsorption energy into an O<sub>t</sub> vacancy, a few additional simulations were run to compare coverage and adsorption energy.

The adsorption energy for N<sub>2</sub> is calculated with

$$\Delta E_{ads\ N2} = E_{WO3(vac)/N2} - (E_{WO3(vac)} + E_{N2}).$$

**Table 3. N<sub>2</sub> coverage on 001 surface above O<sub>t</sub> vacancies versus adsorption energy.**

Coverage	Adsorption Energy (eV)
1/3	-0.3556
1/12	-0.3964

It is found that with fewer adsorbates, it is more chemically favorable for N<sub>2</sub> to adsorb onto a vacancy since 1/12 adsorption energy is less than 1/3. However, the difference of 0.0408 eV indicates that the coverage effects for N<sub>2</sub> are less significant than H<sub>2</sub>O.

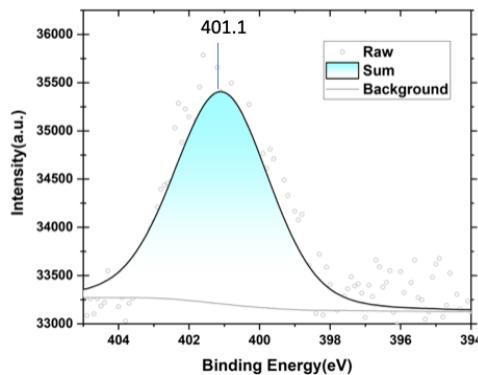
#### 3.4.3. Nitrogen species adsorption investigation

In addition to N<sub>2</sub> versus H<sub>2</sub>O adsorption into O<sub>t</sub> vacancies, I investigated N and N<sub>2</sub> adsorption onto O<sub>t</sub> and a bridge oxygen atom (O<sub>b</sub>), NO, and N<sub>2</sub>O adsorption into O<sub>t</sub> vacancies.

All these species were placed in their upright orientations for initial structures on  $S$ . The motivation is to investigate pathways nitrogen species may take during  $\text{N}_2$  plasma treatment.

**Table 4. Adsorption energies of different nitrogen species**

Location	Adsorbate	Adsorption Energy (eV)
$O_t$	N	-7.353
$O_t$	$\text{N}_2$	-1.777
$O_{t-v}$	NO	-1.625
$O_{t-v}$	$\text{N}_2\text{O}$	-0.199
$O_b$	$\text{N}_2$	-0.356



**Figure 4.** N 1s XPS plot on  $\text{N}_2$  treated  $\text{WO}_3$  surface.

The fairly negative adsorption energies of atomic N and  $\text{N}_2$  on  $O_t$  indicate that these reaction pathways are chemically favorable. Additionally, these adsorption energies are less than the adsorption energy of  $\text{N}_2$  on  $O_t$  vacancies, as seen in *Table 2*, which suggests that experimental observations may find more O-N bonds rather than W-N bonds.

Indeed, in Figure 4 the peak at 401.1eV from Prof. He's experiments indicates that there are O-N bonds observed on the only  $\text{N}_2$  plasma-treated  $\text{WO}_3$  001 surface.

## 4. Potential Energy Surface with DeepMD

### 4.1. Methodology

I used the open-source DeepMD-kit<sup>7</sup> framework that uses deep neural networks (DNNs) to train and create our model. The DNN takes both angular and radial atomic configurations into account while fitting the data. Not only can it use the distance between atoms as input, but also

<sup>7</sup> Han Wang, Linfeng Zhang, Jiequn Han, and Weinan E. “DeePMD-kit: A deep learning package for many-body potential energy representation and molecular dynamics.” Computer Physics Communications 228 (2018): 178-184.

atomic local frame inputs. The main purpose of these descriptors is the aim to model and then predict the physics and interactions between atoms of different types in given training set systems (or other systems if the model is built to extrapolate).

To train the model, additional DFT data was used to train the model. The model was trained with data from ab initio molecular dynamics (AIMD) simulations and optimizations involving initial configurations generated by Python scripts. These DFT calculations were performed by VASP with the Perdew-Burke-Ernzerhof (PBE) exchange-correlation (XC) functional, as mentioned in DFT Methods.

For the model, it is currently trained on 3394 geometries of a  $\text{WO}_3(\text{vac}) + \text{N}_2$  system, 53 geometries of a  $\text{WO}_3 + \text{N}$  system, and 61 geometries of a  $\text{WO}_3 + \text{N}_2$  system to experiment with DeepMD-kit. Ninety percent of the geometries were assigned to be training data, with the rest as validation data. The model was then trained for 1,000,000 steps, taking approximately 6 hours on an NVIDIA GPU partition. The learning rate of the process is an exponential decay, with the decay steps at 10,000 and the initial learning rate at 0.001, which yields the equation

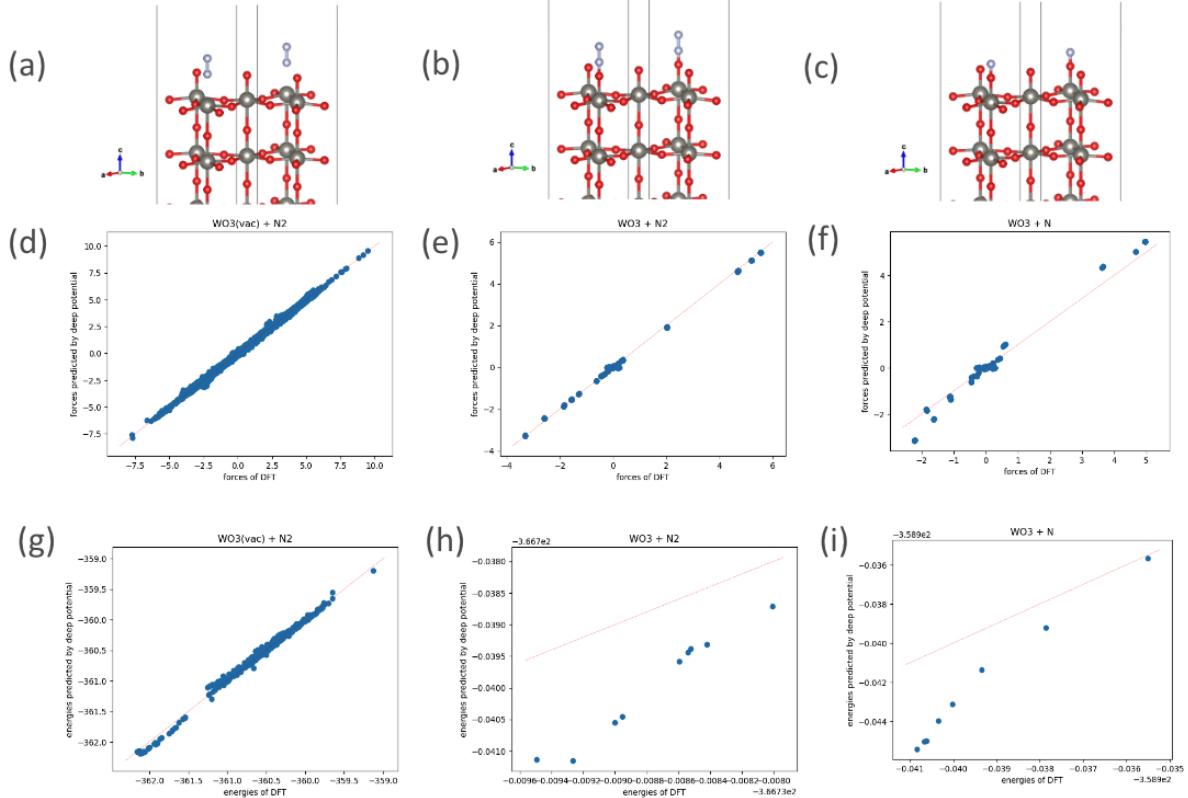
$$lr = 0.001 e^{-\frac{\text{step}}{10,000}}$$
. Additionally, the descriptor for this initial testing model is assigned to be *se\_e2\_a*, which is constructed of both angular and radial atomic confirmations and takes the distance between atoms as input for the neural network. The cut-off radius used is 7 Å. The model hidden layer sizes, in order, are 25, 50, 100, 100, and 100 with an axis neuron of 16. This model is referred to as *Model1*

*Model1* was assessed with the validation data of systems  $\text{WO}_3 + \text{N}$ ,  $\text{WO}_3 + \text{N}_2$ , and  $\text{WO}_3(\text{vac}) + \text{N}_2$  to evaluate its performance in capturing the physics of the systems and extrapolation capabilities. This model was frozen and then used to run optimization simulations with LAAMPS<sup>8</sup> for two initial configurations containing  $\text{WO}_3 + \text{N}_2$  and  $\text{WO}_3(\text{vac}) + \text{N}_2$  for additional comparisons with DFT calculations.

To verify that the *se\_e2\_a* descriptor can be used for the model to learn the physics in systems involving N, W, and O species, I trained a model with the same configuration as above but only used data for the  $\text{WO}_3(\text{vac}) + \text{N}_2$  system training data. Freezing this model and then using LAAMPS to minimize and test this model in a  $\text{WO}_3 + \text{N}_2$  system, we found that the optimized structure resembled the final configuration from DFT calculations. Therefore, we proceeded with this descriptor for this project.

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<sup>8</sup> LAMMPS - a flexible simulation tool for particle-based materials modeling at the atomic, meso, and continuum scales, A. P. Thompson, H. M. Aktulga, R. Berger, D. S. Bolintineanu, W. M. Brown, P. S. Crozier, P. J. in 't Veld, A. Kohlmeyer, S. G. Moore, T. D. Nguyen, R. Shan, M. J. Stevens, J. Tranchida, C. Trott, S. J. Plimpton, Comp Phys Comm, 271 (2022) 10817.



**Figure 5.** Comparison between energies and forces predicted by *Model1* and values from training data. The red line represents exact matches. (a)  $\text{WO}_3(\text{vac}) + \text{N}_2$  system. (b)  $\text{WO}_3 + \text{N}_2$  system. (c)  $\text{WO}_3 + \text{N}$  system. (d,e,f) Forces predicted by *Model1*. (g,h, i) Energies predicted by *Model1*. All the graphs' corresponding systems are the ones depicted in the same column it is found in.

**Table 5. Energy and Force RMSEs of *Model1* in  $\text{WO}_3 + \text{N}_2$  system validation data.**

	$\text{WO}_3(\text{vac}) + \text{N}_2$	$\text{WO}_3 + \text{N}_2$	$\text{WO}_3 + \text{N}$
Energy RMSE (eV/Natom)	0.006994	0.00002557349	0.00006749
Force RMSE (eV/A)	0.05660	0.04169796	0.1732

#### 4.2. Results and Discussion

Acceptable RMSE values in DNN validation systems that are similar to my project are typically less than 0.001 eV and 0.1 eV/A for energy and forces respectively. From *Table 5*, our values for the  $\text{WO}_3$  (vac) +  $\text{N}_2$  and  $\text{WO}_3 + \text{N}_2$  systems meet these requirements, while the other one only meets the requirement for energy.

*Figure 5* shows the comparison between energies and forces predicted by my model and energies from training data. For the first column (a,d,g) with the system  $\text{WO}_3(\text{vac}) + \text{N}_2$ , I find

that data points are not uniformly distributed among the force and energy spectrums, with most geometries clustered around 1.5 eV and -360.5 eV for energies and forces respectively. Despite this trend, the general adherence to the exact match line demonstrates the creation of a robust and accurate model for this specific system. For the non-vacancy systems, in *Figure 5* the graphs for forces indicate a poor fit since most data points are extremely far from the exact match line.

I hypothesize that these discrepancies can be explained by the lack of training data for those two systems in comparison to the amount of data for the  $\text{WO}_3(\text{vac}) + \text{N}_2$  system. Due to how we set up our model with fitting energy taking priority over fitting forces, the model fits extremely well on a small amount of training data for the two non-vacancy systems for energy while for forces lacking the necessary amount to train well.

## 5. Summary

### 5.1. Conclusion

I studied the effects of  $\text{H}_2$ -plasma and  $\text{N}_2$ -plasma and water versus nitrogen interactions on hexagonal  $\text{WO}_3(001)$  surfaces via periodic DFT-PBE calculations. Specifically, I investigated the reaction pathway atomic H, a particle readily available with  $\text{H}_2$ -plasma, takes with the  $\text{O}_t$  surface to conclude that  $\text{H}_2$ -plasma aids more in the formation of terminating oxygen vacancies than room atmosphere through forming and desorbing  $\text{H}_2\text{O}$ . In addition, I explored the effects of coverage and temperature on the previously noted reaction pathway.

To investigate the reaction mechanisms of vibrational  $\text{N}_2$  ( $\text{N}_2(\text{v})$ ), I also trained a machine learning model on systems involving nitrogen species and  $\text{WO}_3(001)$  surfaces with and without vacancies. For systems that have large amounts of data, the fit is acceptable.

Thus, these DFT calculations are necessary for the quantitative and qualitative understanding of the interactions that will help the internal team I work with explain their experimental results and offer additional confidence in their findings.

### 5.2. Challenges and Setbacks

In the initial stages due to a programming error, we (my advisor and I) determined that the most stable termination for hexagonal  $\text{WO}_3(001)$  is the W-terminated side. However, once simulations were initially conducted and our hypothesis was challenged, I reviewed additional literature to find that the O-terminated side was the true stable surface.

Additionally, we found that including spin-polarization had a significant impact on the slab calculations after all calculations were finished and the report was written. I hypothesize that this is because of the asymmetry nature of the stoichiometric composition of  $\text{WO}_3$ , which causes the polarity to vary in the slab. This finding affected the accuracy of most to nearly all calculations involving the  $\text{WO}_3(001)$  slab by  $\pm 1\text{-}4$  eV, which may explain the large magnitude of adsorption energy of atomic species on the surface.

### 5.3. Path Forward

I am aiming to finish the DFT simulations project by the end of August, and to do so involves re-submitting most configurations with spin-polarization included. I am in the process

of this at the moment. After the simulations are concluded, I will update this report for the internal team.

For the DNN, I aim to generate more training data and investigate the impact of different descriptors. In the process, the project will also be looked into through an automation lens, specifically using the Dp-Gen<sup>9</sup> package provided by Deepmd-kit's creators.

## 6. Acknowledgements

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<sup>9</sup> Yuzhi Zhang, Haidi Wang, Weijie Chen, Jinzhe Zeng, Linfeng Zhang, Han Wang, and Weinan E, DP-GEN: A concurrent learning platform for the generation of reliable deep learning based potential energy models, Computer Physics Communications, 2020, 107206.