

MAXWELL'S EQUATIONS

①

I \rightarrow Gauss's Law in Electrostatics:-

$$\oint \vec{D} \cdot d\vec{s} = \int_V \rho_v dv \Rightarrow \text{Integral Form.}$$

Applying Divergence Theorem, $\rightarrow \iint_S$ to \iiint_V

$$\oint \vec{D} \cdot d\vec{s} = \int_V (\nabla \cdot \vec{D}) dv.$$

$$\Rightarrow \boxed{\nabla \cdot \vec{D} = \rho_v} \Rightarrow \text{Point Form (or) Differential Form.}$$

The Divergence of Electric Flux density is equal to Volume charge density, ρ_v .

II \rightarrow Gauss's Law in Magnetostatics:-

Total Magnetic Flux through any closed surface = 0.

$$\oint_S \vec{B} \cdot d\vec{s} = 0 \Rightarrow \text{Integral Form.}$$

Applying Divergence Theorem,

$$\oint \vec{B} \cdot d\vec{s} = \int_V (\nabla \cdot \vec{B}) \cdot dv = 0.$$

$$\boxed{\nabla \cdot \vec{B} = 0} \Rightarrow \text{Point Form (or) Differential Form.}$$

In a closed surface, the Magnetic flux entering is equal to the Magnetic flux leaving, therefore the net divergence is zero.

III \rightarrow Maxwell's III - Eqn \Rightarrow Faraday's Law:- ⁽²⁾

$$\oint \vec{E} \cdot d\vec{l} = - \iint_S \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s} \Rightarrow \text{Integral Form.}$$

Applying Stokes' Theorem $\left\{ \begin{array}{l} \text{Relating Line Integral to} \\ \text{Surface Integral \& Vice Versa} \end{array} \right.$

$$\oint \vec{E} \cdot d\vec{l} = \iint_S (\nabla \times \vec{E}) \cdot d\vec{s} = - \iint_S \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}$$

$$\Rightarrow \boxed{\nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t}} \Rightarrow \text{Differential Form.}$$

IV \rightarrow Maxwell's IV Eqn \Rightarrow Modified Ampere's Law:-

$$\oint \vec{H} \cdot d\vec{l} = \iint_S \left(\vec{J}_c + \frac{\partial \vec{D}}{\partial t} \right) \cdot d\vec{s} \Rightarrow \text{Integral Form.}$$

Applying Stokes' Theorem,

Conduction Current \uparrow

Displacement Current \uparrow

$$\oint \vec{H} \cdot d\vec{l} = \iint_S (\nabla \times \vec{H}) \cdot d\vec{s} = \iint_S \left(\vec{J}_c + \frac{\partial \vec{D}}{\partial t} \right) \cdot d\vec{s}$$

$$\boxed{\nabla \times \vec{H} = \vec{J}_c + \frac{\partial \vec{D}}{\partial t}} \Rightarrow \text{Differential Form.}$$

The Magnetomotive force around a closed path is equal to the sum of the Conduction Current & displacement Current by the path.

Three Constitutive Relations in Field Theory Analysis \Rightarrow

$$\text{E-Field} \leftarrow \vec{D} = \epsilon_e \vec{E}$$

$$\text{H-Field} \leftarrow \vec{B} = \mu \vec{H}$$

$$\vec{J} = \sigma \vec{E}$$

Electromagnetic Wave Equation

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Wave Eqn in E-Field:-

Maxwell's Eqn from Faraday's Law, $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

$$\text{w.k.t. } \vec{B} = \mu \vec{H}$$

$$\Rightarrow \nabla \times \vec{E} = -\mu \frac{\partial \vec{H}}{\partial t} \rightarrow (1)$$

Taking Curl on both sides,

$$\nabla \times \nabla \times \vec{E} = -\mu \nabla \times \frac{\partial \vec{H}}{\partial t} \rightarrow (2)$$

According to the Identity,

$$\nabla \times \nabla \times \vec{E} = (\nabla \cdot \vec{E}) \cdot \nabla - (\nabla \cdot \nabla) \vec{E} \Rightarrow (\nabla \cdot \vec{E}) \nabla - \nabla^2 \vec{E}$$

Since, there is no electric field within the conductor,
 $\nabla \cdot \vec{E} = 0$.

$$\Rightarrow \nabla \times \nabla \times \vec{E} = -\nabla^2 \vec{E} \rightarrow (3)$$

Maxwell's Eqn from Ampere's Law, $\nabla \times \vec{H} = \vec{J}_c + \frac{\partial \vec{D}}{\partial t}$

$$\nabla \times \vec{H} = \sigma \vec{E} + \epsilon \frac{\partial \vec{E}}{\partial t}$$

Differentiating w.r.t t .

$$\nabla \times \frac{\partial \vec{H}}{\partial t} = \sigma \frac{\partial \vec{E}}{\partial t} + \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} \rightarrow (4)$$

Substituting (3) & (4) in (2) we get,

$$-\nabla^2 \vec{E} = -\mu \left[\sigma \frac{\partial \vec{E}}{\partial t} + \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} \right]$$

$$\boxed{\nabla^2 \vec{E} - \mu \sigma \frac{\partial \vec{E}}{\partial t} - \mu \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} = 0}$$

Wave Equation
in terms of E-Field.

② wave Eqn in Magnetic field:-

④

The Maxwell's equation from Ampere's law is,

$$\nabla \times \vec{H} = \vec{J}_c + \frac{\partial \vec{D}}{\partial t}.$$

$$\Rightarrow \nabla \times \vec{H} = \sigma E + \epsilon \frac{\partial E}{\partial t}.$$

Taking Curl on both sides;

$$\nabla \times \nabla \times H = \sigma (\nabla \times E) + \epsilon \left[\nabla \times \frac{\partial E}{\partial t} \right] \rightarrow \textcircled{1}$$

According to the Identity,

$$\nabla \times \nabla \times H = (\nabla \cdot H) \cdot \nabla - (\nabla \cdot \nabla) H \Rightarrow \nabla \cdot (\nabla \cdot H) - \nabla^2 H$$

From Gauss Law, w.k.t $\nabla \cdot B = 0$.

$$\mu [\nabla \cdot H] = 0.$$

$$\Rightarrow \nabla \times \nabla \times H = -\nabla^2 H. \xrightarrow{\nabla \cdot H = 0} \textcircled{2}$$

w.k.t. Maxwell's Equation from Faraday's law is

$$\nabla \times \vec{E} = -\mu \frac{\partial H}{\partial t} \rightarrow \textcircled{3}$$

Differentiating w.r.t "t".

$$\nabla \times \frac{\partial E}{\partial t} = -\mu \frac{\partial^2 H}{\partial t^2} \rightarrow \textcircled{4}$$

Substituting $\textcircled{2}, \textcircled{3}, \textcircled{4}$ in $\textcircled{1}$ we get,

$$-\nabla^2 H = -\sigma \mu \frac{\partial H}{\partial t} - \mu \epsilon \frac{\partial^2 H}{\partial t^2}$$

$$\boxed{\nabla^2 H - \sigma \mu \frac{\partial H}{\partial t} - \mu \epsilon \frac{\partial^2 H}{\partial t^2} = 0}$$

Wave Equation
in terms of H-Field.

Wave Equation in Free Space :-

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$$\Rightarrow \sigma = 0 ; \rho = 0.$$

Conductivity $= 0$

charge density $= 0$.

The wave Equations can be written as,

$$\nabla^2 E - \mu \epsilon \frac{\partial^2 E}{\partial t^2} = 0 ; \nabla^2 H - \mu \epsilon \frac{\partial^2 H}{\partial t^2} = 0.$$

For Free space

$$\mu = \mu_0 \mu_r \Rightarrow \mu = \mu_0 \quad \boxed{\mu_r = 1}$$

$$\epsilon = \epsilon_0 \epsilon_r \Rightarrow \epsilon = \epsilon_0 \quad \boxed{\epsilon_r = 1}$$

$$\boxed{\nabla^2 E - \mu_0 \epsilon_0 \frac{\partial^2 E}{\partial t^2} = 0 ; \nabla^2 H - \mu_0 \epsilon_0 \frac{\partial^2 H}{\partial t^2} = 0}$$

w.k.t. $\frac{1}{\sqrt{\mu_0 \epsilon_0}} = v_0 = 3 \times 10^8 \text{ m/sec} \therefore \mu_0 \epsilon_0 = \frac{1}{v_0^2}$

$$\boxed{\begin{aligned} \nabla^2 E - \frac{1}{v_0^2} \frac{\partial^2 E}{\partial t^2} &= 0 \\ \nabla^2 H - \frac{1}{v_0^2} \frac{\partial^2 H}{\partial t^2} &= 0. \end{aligned}}$$

Wave Equations
in Free Space.

Electromagnetic wave Equation in Phasor Form: ^(b)

In phasor Form, $\vec{E} = \vec{E}_0 e^{j\omega t}$

$$\frac{\partial \vec{E}}{\partial t} = j\omega \vec{E}_0 e^{j\omega t} \quad \left\| \quad \frac{\partial^2 \vec{E}}{\partial t^2} = j\omega E_0 e^{j\omega t} (j\omega) \right.$$

$$\boxed{\frac{\partial \vec{E}}{\partial t} = j\omega \vec{E}} \quad \& \quad \boxed{\frac{\partial^2 \vec{E}}{\partial t^2} = (j\omega)^2 \vec{E}} \quad = (j\omega)^2 \cdot E_0 \cdot e^{j\omega t}$$

$$\boxed{\begin{aligned} \therefore \nabla^2 \vec{E} &= \left[j\omega \mu (\sigma + j\omega \epsilon) \right] \vec{E} \\ \nabla^2 \vec{H} &= \left[j\omega \mu (\sigma + j\omega \epsilon) \right] \vec{H} \end{aligned}}$$

WAVE PARAMETERS:-

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1. Propagation Constant:-

The phasor form of EM wave equation is.,

$$\nabla^2 \vec{E} = [j\omega\epsilon(\sigma + j\omega\epsilon)] \vec{E} \rightarrow \textcircled{1}$$

$$\nabla^2 \vec{H} = [j\omega\mu(\sigma + j\omega\epsilon)] \vec{H} \rightarrow \textcircled{2}$$

propagation Constant, $\gamma^2 = j\omega\mu(\sigma + j\omega\epsilon) \rightarrow \textcircled{3}$

$$\gamma = \sqrt{j\omega\mu\sigma - \omega^2\mu\epsilon}$$

$\gamma = \alpha + j\beta$

propagation
Constant. Attenuation
Constant Phase
Constant.

$$(\alpha + j\beta)^2 = -\omega^2\mu\epsilon + j\omega\mu\sigma \rightarrow \textcircled{4}$$

$$\alpha^2 + 2j\alpha\beta - \beta^2 = -\omega^2\mu\epsilon + j\omega\mu\sigma$$

Equating Real parts, $\alpha^2 - \beta^2 = -\omega^2\mu\epsilon \rightarrow \textcircled{5}$

Taking Modulus of $(\alpha + j\beta)^2 \Rightarrow \{\text{eqn } \textcircled{4}\}$.

$$|\alpha + j\beta|^2 = |-\omega^2\mu\epsilon + j\omega\mu\sigma|$$

$$\alpha^2 + \beta^2 = \sqrt{\omega^4\mu^2\epsilon^2 + \omega^2\mu^2\sigma^2} \rightarrow \textcircled{6}$$

Adding $\textcircled{5}$ & $\textcircled{6}$

$$2\alpha^2 = -\omega^2\mu\epsilon + \sqrt{\omega^4\mu^2\epsilon^2 + \omega^2\mu^2\sigma^2}$$

$$2\alpha^2 = -\omega^2\mu\epsilon + \sqrt{\omega^4\mu^2\epsilon^2\left(1 + \frac{\sigma^2}{\omega^2\epsilon^2}\right)} \quad (8)$$

$$\alpha^2 = \frac{-\omega^2\mu\epsilon}{2} + \frac{\omega^2\mu\epsilon}{2} \sqrt{1 + \frac{\sigma^2}{\omega^2\epsilon^2}}$$

$$\alpha^2 = \frac{\omega^2\mu\epsilon}{2} \left[\sqrt{1 + \frac{\sigma^2}{\omega^2\epsilon^2}} - 1 \right]$$

$$\alpha = \omega \sqrt{\frac{\mu\epsilon}{2} \left[\sqrt{1 + \frac{\sigma^2}{\omega^2\epsilon^2}} - 1 \right]}$$

⑥-⑤ we get .,

$$\alpha^2 + \beta^2 - \alpha^2 + \beta^2 = \sqrt{\omega^4\mu^2\epsilon^2 + \omega^2\mu^2\sigma^2} - (-\omega^2\mu\epsilon)$$

$$2\beta^2 = \sqrt{\omega^4\mu^2\epsilon^2 + \omega^2\mu^2\sigma^2} + \omega^2\mu\epsilon$$

$$= \omega\mu \sqrt{\omega^2\epsilon^2 + \sigma^2} + \omega^2\mu\epsilon$$

$$\beta^2 = \frac{\omega\mu}{2} \sqrt{\omega^2\epsilon^2 + \sigma^2} + \frac{\omega^2\mu\epsilon}{2}$$

$$\beta = \sqrt{\frac{\omega\mu}{2} \sqrt{\omega^2\epsilon^2 + \sigma^2} + \frac{\omega^2\mu\epsilon}{2}}$$

$$\beta = \sqrt{\frac{\omega^2 \mu \epsilon_e}{2} \left[\frac{2}{\omega^2 \mu \epsilon_e} \frac{\omega \mu}{2} \sqrt{\omega^2 \epsilon_e^2 + \sigma^2} + 1 \right]} \quad (9)$$

$$= \sqrt{\frac{\omega^2 \mu \epsilon_e}{2} \left[\frac{1}{\omega \epsilon_e} \sqrt{\omega^2 \epsilon_e^2 + \sigma^2} + 1 \right]}$$

$$\beta = \omega \sqrt{\frac{\mu \epsilon_e}{2} \left[\sqrt{\frac{\omega^2 \epsilon_e^2}{\omega^2 \epsilon_e^2} + \frac{\sigma^2}{\omega^2 \epsilon_e^2}} + 1 \right]}$$

$$\beta = \omega \sqrt{\frac{\mu \epsilon_e}{2} \left(\left(\sqrt{1 + \frac{\sigma^2}{\omega^2 \epsilon_e^2}} \right) + 1 \right)}$$

$$\beta = \omega \sqrt{\frac{\mu \epsilon_e}{2} \left(\sqrt{1 + \frac{\sigma^2}{\omega^2 \epsilon_e^2}} + 1 \right)}$$