

# Numerical Methods for the Conductivity Equation

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# Motivation: Magneto-Acousto-Electric-Tomography (MAET)

- ▶ Based on measurements of the electrical potential arising when an acoustic wave propagates through conductive medium placed in a magnetic field.
- ▶ MAET is an example of an inverse problem for the conductivity of biological tissue
- ▶ We focused on understanding the *direct* problem this semester.

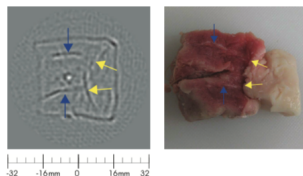


Figure: Sample Reconstruction

# Basic Physics of Current

- ▶ Let  $E(x)$  denote an electric field over  $\mathbb{R}^2$ .
- ▶ Let  $\sigma(x)$  denote the conductivity over  $\mathbb{R}^2$ .
- ▶ Current is given by  $\sigma(x)E(x)$ .
- ▶ We assume there exists a potential  $u(x)$ , for which  $E(x) = \nabla u(x)$ .
- ▶ Then current is given by  $\sigma(x)\nabla u(x)$ .
- ▶ In the absence of sinks and sources, the divergence of current is 0.

# Current Induced by Point Sources

Consider the case where conductivity is constant on  $\mathbb{R}^2$ . Suppose we inject currents  $W_j$  at a finite set of points  $y^{(j)}$ , where  $\sum_{j=1}^n W_j = 0$ . The resulting electric potential is given by

$$w(x) = -\frac{1}{\sigma_0} \sum_{j=1}^n W_j \phi(x - y^{(j)})$$

where  $\Delta\phi(x) = \delta(x)$ .  $\phi(x)$  is known as the fundamental solution to Laplace's Equation.  $\phi(x) = \frac{1}{2\pi} \ln |x|$ .

# The Problem we are Considering

- ▶ Suppose now we have a region  $\Omega$ , not containing any of the injection points  $y^{(j)}$ , where conductivity is non-constant. ie, conductivity is given by some function  $\sigma(x)$ ,  $x \in \Omega$ .
- ▶ By superposition, we can split the resulting electric potential into the potential given by the injected currents, and the potential arising when conductivity is no longer constant. Then current is given by  $\sigma(x)\nabla(u(x) + w(x))$ .
- ▶ Divergence of total current is 0 away from the injection points, and  $u(x)$  vanishes at  $\infty$ .

# The Conductivity Equation - Setup

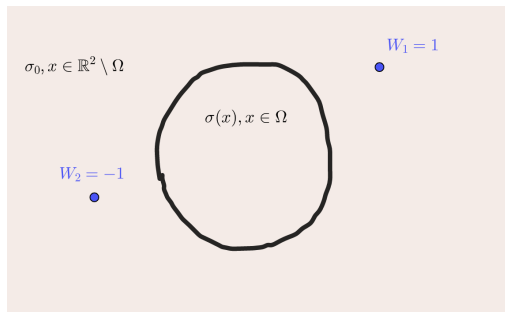


Figure: Basic Setup

$$\nabla \cdot \sigma(x) \nabla (u(x) + w(x)) = 0, \quad x \in \mathbb{R}^2 \setminus \bigcup_{j=1}^n \gamma^{(j)}.$$

# Conductivity Equation - Fredholm Reduction

- ▶ For convenience, we let  $U(x) = \Delta u(x)$ .
- ▶ It can be shown that the conductivity equation can be expressed as an integral equation of the second kind

$$U(x) + \nabla \ln(\sigma(x)) \cdot \nabla \left[ \int_{\Omega} U(y) \phi(x-y) dy \right] = -\nabla \sigma(x) \cdot \nabla w(x), x \in \Omega.$$

- ▶ Such equations are known to be well-posed.

# Primary Problem- Convolution with a Singular Kernel

- ▶ We want to make a numerical solver to reconstruct currents for arbitrary  $\sigma(x)$ .
- ▶ Since we are using the Fredholm Equation of the Second kind formulation, we need to handle convolutions with a singular kernel  $u(x) = \int_{\Omega} U(y) \frac{1}{2\pi} \ln |y - x| dy$
- ▶ Result of the convolution will satisfy the following conditions:
- ▶ (1)  $\Delta u(x) = U(x), x \in \Omega$
- ▶ (2)  $\Delta u(x) = 0, x \notin \Omega$ .
- ▶ (3)  $u(x)$  behaves like a Logarithm asymptotically.
- ▶ A function satisfying (1)-(3) will be unique.
- ▶ For simplicity, we take  $\Omega = [-a/2, a/2]^2$ , but the following approach should work for any  $\Omega$ .



# Convolution Method - Description

Our goal is to compute  $u(x) = \int_{\Omega} U(y) \frac{1}{2\pi} \ln |y - x| dy$ .

- ▶ Step 1- Compute the solution having 0 boundary conditions on a larger square via Sine series.
- ▶ Step 2- Smoothly cutoff the corners of the resulting solution.
- ▶ Step 3- Compute the laplacian of the resulting function.
- ▶ Step 4- Set the laplacian to 0 over the small square.
- ▶ Step 5- Compute the convolution of the resulting with the fundamental solution, taking the convolution over the small square.
- ▶ Step 6- Subtract the result of the convolution from the smoothly cutoff solution.

# Convolution Method - In Pictures

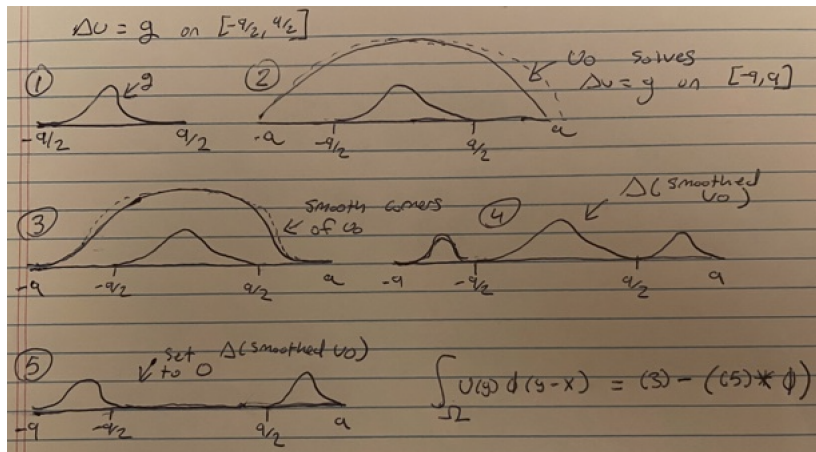


Figure: Convolution Method

# Why it Works + Benefits

- ▶  $\Delta u_{\text{final}} = U(x)$  inside of  $\Omega$ .
- ▶  $\Delta u_{\text{final}} = 0$  outside of  $\Omega$ .
- ▶  $u_{\text{final}}$  behaves like a logarithm asymptotically.
- ▶ Therefore,  $u_{\text{final}} = \int_{\omega} U(y) \phi|y - x|$ .
- ▶ We completely avoid the singularity in the kernel  $\phi(x)$ .
- ▶ All non-trivial computations - computing the laplacian, inverting the laplacian, taking the convolution, computing partial derivatives- can be done via FFT methods.

# Convolution Method- Numerical Example

We test a numerical scheme using the function  $(1 - \sin(\hat{x})) \ln(|x|)$ . This function has a laplacian which we can exactly compute, and behaves like a logarithm asymptotically.

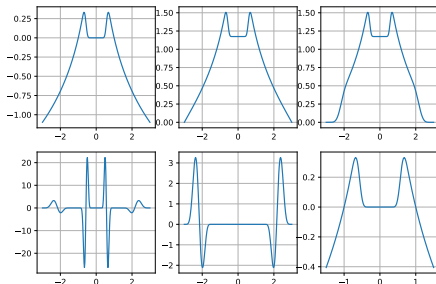


Figure: Numerical Example

# Computing the Convolution- Error

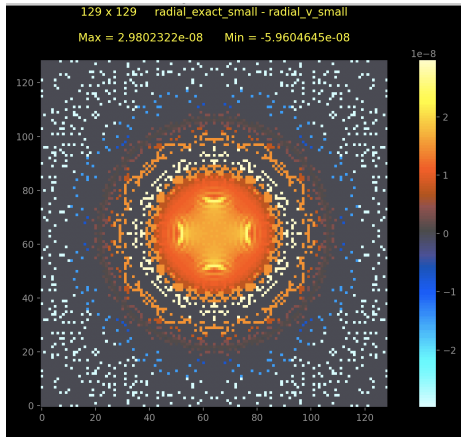


Figure: Difference between  $u_{\text{final}}$  and exact solution