Swift-Hohenberg Numerics - Testing with Chebfun

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1 Operator Splitting for SH

• We have the PDE

$$w_t = -(1 + \nabla^2)^2 w + Rw - w^3.$$

• We can break this into a linear part and a non-linear part,

$$W_t = L(w) + NL(w),$$

with

$$L(w) = -(2\nabla^2 + \nabla^4)w$$
$$NL(w) = (R - 1)w - w^3.$$

- Basic procedure is as follows:
 - (i) Let $A = (-\nabla^4 2\nabla^2)$. Consider the pair of PDEs

$$w_t = Aw$$
$$w_t = (R - 1)w - w^3.$$

(ii) Handling the linear and nonlinear PDEs separately, we get the relations

$$w(t + \Delta t) = \Delta t A w(t) + w(t)$$

$$\implies w(t + \Delta t) \approx e^{A\Delta t} w(t).$$

And for the nonlinear part,

$$w(t + \Delta t) \approx \Delta t \left[(R - 1)w(t) - w(t)^3 \right] + w(t).$$

(iii) Apply operator splitting (strang splitting) now

The code is implemented in python as follows:

```
def non_lin_rhs(w,R):
    return (R-1)*w - w**3
def integrateSH(w0,R,dt,nSteps,L,o4=False):
     :param w0: initial temperature surface
     :param R: bifurcation parameter- can be a constant, or of same shape as w0
     :param dt: time step length
     :param nSteps: number of time steps to take
     :param L: Length of square over which w0 is defined
     :return w0: time evolution of w0 at time 0+dt*nSteps
     Ideally, the size of w0 is fft friendly, ie 2^n x 2^n
     print("Starting time integration of Swift Hohenberg")
     ny, nx = np.shape(w0)
     print("Dimensions of w0:", nx, ny)
     kx = (2.*np.pi/L)*sp.fft.fftfreq(nx,1./nx)
     ky = (2.*np.pi/L)*sp.fft.fftfreq(ny,1./ny)
     Kx, Ky = np.meshgrid(kx,ky)
     fourierLaplacian = -(Kx**2+Ky**2)
     A = -(fourierLaplacian*fourierLaplacian)-2*fourierLaplacian
     for i in range(0,nSteps):
         # if i%100 == 0:
             # print("step number:",i)
         w1 = np.real(sp.fft.ifft2(np.exp(A*.5*dt)*sp.fft.fft2(w0)))
         #rk4 version
         if o4:
            k1 = dt*w1
            k2 = dt*non_lin_rhs(w1+.5*k1, R)
            k3 = dt*non_lin_rhs(w1+.5*k2, R)
            k4 = dt*non_lin_rhs(w1+k3, R)
            w2 = (k1+2*k2+2*k3+k4)/6 + w1
         #fwd euler version
         else:
            w2 = dt*((R-1)*w1-w1**3)+w1
         w0 = np.real(sp.fft.ifft2(np.exp(A*.5*dt)*sp.fft.fft2(w2)))
     return w0
```

Note: It seems the Forward Euler scheme is working better than RK4 Scripts

2 Comparison with Chebfun

2.1 Chebfun vs Python 1

Chebfun 1 code

```
dom = [-16 \ 16 \ -16 \ 16]
tspan = [0 \ 10];
S = spinop2(dom, tspan)
S.lin = Q(u) -2*lap(u)-biharm(u)
R = .5
S.nonlin = Q(u) (R-1)*u - u.^3;
u0 = 1/20 \cdot \text{chebfun2}(@(x,y) \cos(x) + \sin(2*x) + \sin(y) + \cos(2*y), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-5*pi).^2 + (y-5*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-5*pi).^2 + (y-15*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-15*pi).^2 + (y-15*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-15*pi).^2 + (y-5*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-10*pi).^2 + (y-10*pi).^2)), dom, 'trig');
S.init = u0;
u = spin2(S, 256, .1, 'plot', 'off');
plot(u), view(0,90), axis equal, axis off
saveas(gcf,'/Users/edwardmcdugald/Research/convection_patterns_matlab/figs/sh_tsts_1018/
close all
```

python 1 code

```
x = np.linspace(-16,16,256)
y = np.linspace(-16,16,256)
X,Y = np.meshgrid(x,y)
w0 = (1./20.)*(np.cos(X)+np.sin(2*X)+np.sin(Y)+np.cos(2*Y))
w0 = w0 + (np.exp(-((X-5*np.pi)**2 +(Y-5*np.pi)**2)))
w0 = w0 + (np.exp(-((X-5*np.pi)**2 +(Y-15*np.pi)**2)))
w0 = w0 + (np.exp(-((X-15*np.pi)**2 +(Y-15*np.pi)**2)))
w0 = w0 + (np.exp(-((X-15*np.pi)**2 +(Y-5*np.pi)**2)))
w0 = w0 + (np.exp(-((X-10*np.pi)**2 +(Y-10*np.pi)**2)))
dt = .1
R=.5
L=x[len(x)-1]-x[0]
nSteps = 100
```

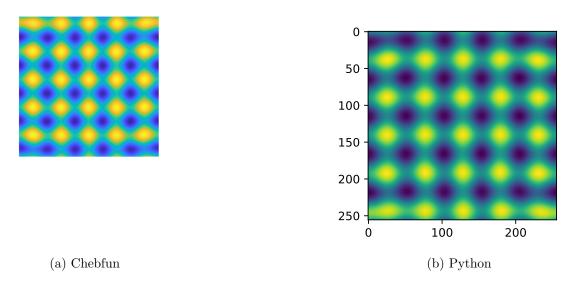


Figure 1: Chebfun v Python 1

```
W1 = integrateSH(w0,R,dt,nSteps,L)
fig1, ax1 = plt.subplots(nrows=1, ncols=1, figsize=(3,3))
ax1.imshow(W1)
plt.savefig("/Users/edwardmcdugald/Research/
convection_patterns/code/figs/sh_num_tsts_1018/mySH1.pdf")

#Running the Methods on the Ellipse
R=.5*(-1./(1.+np.exp(-(X**2+2*Y**2-256.)))+1.)

W1_Ell = integrateSH(w0,R,dt,nSteps,L)
fig1_Ell, ax1_Ell = plt.subplots(nrows=1, ncols=1, figsize=(3,3))
ax1_Ell.imshow(W1_Ell)
plt.savefig("/Users/edwardmcdugald/Research/
convection_patterns/code/figs/sh_num_tsts_1018/mySH1_Ell.pdf")
```

Test 1: Chebfun v Python Results

Python Result on Ellipse

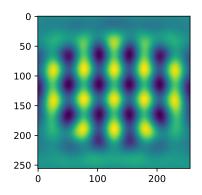


Figure 2: Python Result on Ellipse

2.2 Chebfun vs Python 2

y = np.linspace(0,20*np.pi,128)

X,Y = np.meshgrid(x,y)

Chebfun 2 code

```
dom = [0 \ 20*pi \ 0 \ 20*pi];
tspan = [0 \ 200];
S = spinop2(dom, tspan);
S.lin = Q(u) -2*lap(u) - biharm(u);
r = 1e-2;
S.nonlin = @(u) (-1 + r)*u - u.^3;
u0 = 1/20 \cdot \text{chebfun2}(0(x,y) \cos(x) + \sin(2*x) + \sin(y) + \cos(2*y), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-5*pi).^2 + (y-5*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-5*pi).^2 + (y-15*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-15*pi).^2 + (y-15*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-15*pi).^2 + (y-5*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-10*pi).^2 + (y-10*pi).^2)), dom, 'trig');
S.init = u0;
plot(S.init), view(0,90), axis equal, axis off
close all
u = spin2(S, 96, 2e-1, 'plot', 'off');
plot(u), view(0,90), axis equal, axis off
saveas(gcf,'/Users/edwardmcdugald/Research/
convection_patterns_matlab/figs/sh_tsts_1018/cf2.pdf');
close all
python 2 code
x = np.linspace(0,20*np.pi,128)
```

```
w0 = (1./20.)*(np.cos(X)+np.sin(2*X)+np.sin(Y)+np.cos(2*Y))
w0 = w0 + (np.exp(-((X-5*np.pi)**2 + (Y-5*np.pi)**2)))
w0 = w0 + (np.exp(-((X-5*np.pi)**2 + (Y-15*np.pi)**2)))
w0 = w0 + (np.exp(-((X-15*np.pi)**2 + (Y-15*np.pi)**2)))
w0 = w0 + (np.exp(-((X-15*np.pi)**2 + (Y-5*np.pi)**2)))
w0 = w0 + (np.exp(-((X-10*np.pi)**2 + (Y-10*np.pi)**2)))
dt = 2e-1
R= 1e-2
L=x[len(x)-1]-x[0]
nSteps = 1000
W2 = integrateSH(w0,R,dt,nSteps,L)
fig2, ax2 = plt.subplots(nrows=1, ncols=1, figsize=(3,3))
ax2.imshow(W2)
plt.savefig("/Users/edwardmcdugald/Research/
convection_patterns/code/figs/sh_num_tsts_1018/mySH2.pdf")
#Running the Methods on the Ellipse
R=.5*(-1./(1.+np.exp(-((X-10*np.pi)**2+2*(Y-10*np.pi)**2-256.)))+1.)
W2_Ell = integrateSH(w0,R,dt,nSteps,L)
fig2_Ell, ax2_Ell = plt.subplots(nrows=1, ncols=1, figsize=(3,3))
ax2_Ell.imshow(W2_Ell)
plt.savefig("/Users/edwardmcdugald/Research/
convection_patterns/code/figs/sh_num_tsts_1018/mySH2_Ell.pdf")
```

Test 2: Chebfun v Python Results

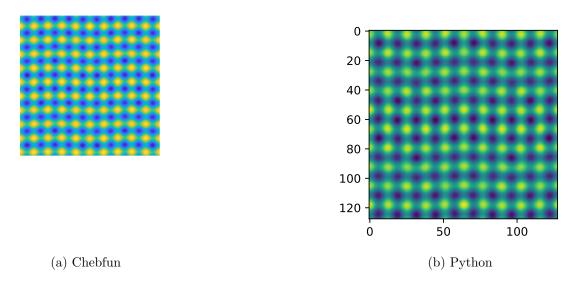


Figure 3: Chebfun v Python 2

Python Result on Ellipse

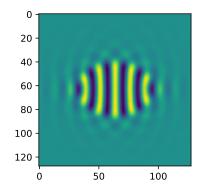


Figure 4: Python Result on Ellipse

2.3 Chebfun vs Python 3

Chebfun 3 code

```
dom = [0 20*pi 0 20*pi];
tspan = [0 200];
```

```
S = spinop2(dom, tspan);
S.lin = Q(u) -2*lap(u) - biharm(u);
u0 = 1/20 \cdot \text{chebfun2}(@(x,y) \cos(x) + \sin(2*x) + \sin(y) + \cos(2*y), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-5*pi).^2 + (y-5*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-5*pi).^2 + (y-15*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-15*pi).^2 + (y-15*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-15*pi).^2 + (y-5*pi).^2)), dom, 'trig');
u0 = u0 + chebfun2(@(x,y) exp(-((x-10*pi).^2 + (y-10*pi).^2)), dom, 'trig');
S.init = u0;
r = 7e-1;
S.nonlin = @(u) (-1 + r)*u - u.^3;
u = spin2(S, 256, 2e-1, 'plot', 'off');
plot(u), view(0,90), axis equal, axis off
saveas(gcf,'/Users/edwardmcdugald/Research/
convection_patterns_matlab/figs/sh_tsts_1018/cf3.pdf');
close all
python 3 code
x = np.linspace(0,20*np.pi,256)
y = np.linspace(0,20*np.pi,256)
X,Y = np.meshgrid(x,y)
w0 = (1./20.)*(np.cos(X)+np.sin(2*X)+np.sin(Y)+np.cos(2*Y))
w0 = w0 + (np.exp(-((X-5*np.pi)**2 + (Y-5*np.pi)**2)))
w0 = w0 + (np.exp(-((X-5*np.pi)**2 + (Y-15*np.pi)**2)))
w0 = w0 + (np.exp(-((X-15*np.pi)**2 + (Y-15*np.pi)**2)))
w0 = w0 + (np.exp(-((X-15*np.pi)**2 + (Y-5*np.pi)**2)))
w0 = w0 + (np.exp(-((X-10*np.pi)**2 + (Y-10*np.pi)**2)))
dt = 2e-1
R= 7e-1
L=x[len(x)-1]-x[0]
nSteps = 1000
W3 = integrateSH(w0,R,dt,nSteps,L)
fig3, ax3 = plt.subplots(nrows=1, ncols=1, figsize=(3,3))
ax3.imshow(W3)
plt.savefig("/Users/edwardmcdugald/Research/
convection_patterns/code/figs/sh_num_tsts_1018/mySH3.pdf")
#Running the Methods on the Ellipse
R=.5*(-1./(1.+np.exp(-((X-10*np.pi)**2+2*(Y-10*np.pi)**2-256.)))+1.)
W3_Ell = integrateSH(w0,R,dt,nSteps,L)
fig3_Ell, ax3_Ell = plt.subplots(nrows=1, ncols=1, figsize=(3,3))
```

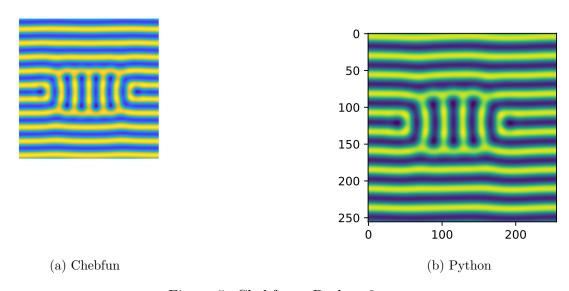


Figure 5: Chebfun v Python 3

ax3_Ell.imshow(W3_Ell)
plt.savefig("/Users/edwardmcdugald/Research/
convection_patterns/code/figs/sh_num_tsts_1018/mySH3_Ell.pdf")

Test 3: Chebfun v Python Results

Python Result on Ellipse

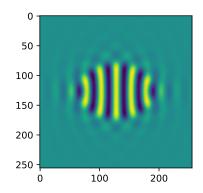


Figure 6: Python Result on Ellipse

3 Shankars Code

Implementation based on the following ideas

- Global description of patterns far from onset: a case study
- fourth order time stepping for stiff pdes

The original code

```
%% Swift-Hohenberg equation
% We use a pseudospectral discretization along with Exponential
% time-differencing/RK4 to solve the Swift-Hohenberg equation on
% "arbitrary" domains.
% The equation we are solving is
   u_t = -(\Delta + 1)^2 u + R u - u^3
% To use a spectral discretization, we fit the domain inside a rectangle
% [0,L_x] X [0, L_y] and extend periodically. Since the patterns only form
% for R > 0, we choose R = -0.5 + R \cdot \{0 \}, where \chi is a smoothed
% characteristic function of \Omega.
% To prevent aliasing due to the nonlinearity u^3, we need to keep twice as
% many sample points in each direction as the number of useful (i.e.
% low-frequency) modes. We can keep R "sharp" on the scale of the grid
% spacing, although we will typically use a low pass filtered version of R
% since it shouldn't play a big role in determining the pattern, and a
% "softer" boundary seems more robust.
% This idea borrows from "Global description of patterns far from onset: a
% case study" by Ercolani, Indik, Newell and Passot for the setup of
% Swift-Hohenberg, and from "Fourth-order time-stepping for stiff PDEs" by
% Kassam and Trefethen for the implementation of the ETD.
%% Spatial
             grid
                     and
                           initial
                                     condition:
               % Size of enclosing periodic rectangle
Lx = 120*pi;
Ly = 60*pi;
% we want N ~ L_x because max wavenumber is like k_x ~ pi*N/Lx which we
% want to be about 3. Likewise for Ly.
```

```
Nx = 512;
Ny = 256;
% Number of grid points in x and y. We will keep them as powers of 2.
% Useful number of modes = Nx Ny/4
beta = 0.45; % Less than 0.5, sets relative size of the ellipse in the box
amplitude = 0.1; % for initial pattern. want this "small" compared to sqrt(R)
xx = Lx*(-Nx/2+1:Nx/2)'/Nx;
yy = Ly*(-Ny/2+1:Ny/2)'/Ny;
[X,Y] = ndgrid(xx,yy);
R = 0.5*tanh(sqrt(Lx^2+Ly^2)*(beta-sqrt((X/Lx).^2+(Y/Ly).^2))/2);
if (init_flag == 1)
    u0 = randn(N);
    u0 = amplitude*u0/norm(u0, inf);
elseif (init_flag == 2)
    u0 = amplitude*sin(Y);
elseif (init_flag == 3)
    u0 = ellipse_init(X,Y,beta*Lx,beta*Ly,amplitude);
end
figure(21)
surf(X,Y,u0.*(R+0.5),'mesh','none');
axis equal;
view(2);
%%
     Precompute
                            ETDRK4
                                               quantities:
                  various
                                      scalar
h = 3;
         %
             time
                    step
kx = 2*pi*[0:Nx/2 -Nx/2+1:-1]'/Lx;
                                     %
                                          wave
                                                 numbers
ky = 2*pi*[0:Ny/2 -Ny/2+1:-1]'/Ly;
[xi,eta] = ndgrid(kx,ky);
L = -(1-xi.^2-eta.^2).^2; %
                               Fourier
                                          multipliers
E = \exp(h*L);
E2= \exp(h*L/2);
M = 16;
          %
              no.
                    of points
                                  for
                                         complex
r = \exp(1i*pi*((1:M)-.5)/M);
                              %
                                   roots
                                            of
                                                 unity.
% This circle should suffice since max eigenvalue of Lh is 0, at k=0, Lh =
% -1/8
```

```
L2 = L(:); % convert to a single column
LR = h*L2(:,ones(M,1)) + r(ones(Nx*Ny,1),:); % adding r(j) to the jth column
Q = h*real(mean((exp(LR/2)-1)./LR,2)); % means in the 2 direction.
f1 = h*real(mean((-4-LR+exp(LR).*(4-3*LR+LR.^2))./LR.^3,2));
f2 = h*real(mean((2+LR+exp(LR).*(-2+LR))./LR.^3,2));
f3 = h*real(mean((-4-3*LR-LR.^2+exp(LR).*(4-LR))./LR.^3,2));
                       % convert from columns to N by N matrices
f1=reshape(f1,Nx,Ny);
f2=reshape(f2,Nx,Ny);
f3=reshape(f3,Nx,Ny);
Q=reshape(Q,Nx,Ny);
%clear LR;
%% Dealiasing
Fx = false(Nx,1); \% Fx = 1 for high frequencies which will be set to zero
Fy = false(Ny, 1);
Fx((Nx/2+1-round(Nx/4):1+Nx/2+round(Nx/4)))=true; % For dealiasing with u^3
Fy((Ny/2+1-round(Ny/4):1+ Ny/2+round(Ny/4))) = true;
[alxi,aleta] = ndgrid(Fx,Fy);
ind = alxi | aleta; %de-aliasing index.
%filter R and u0
Rhat = fftn(R);
Rhat(ind) = 0;
R = real(ifftn(Rhat));
vv = fftn(u0);
vv(ind) = 0;
u0 = real(ifftn(vv));
Q(ind) = 0; % Q is the only term that multiplies the nonlinear factors
%%
           time-stepping
     Main
                            loop:
           1000;
tmax
       =
       = round(tmax/h);
nmax
nplt
          floor((tmax/25)/h);
```

```
tt = zeros(1,round(nmax/nplt)+1);
uu = zeros(Nx,Ny,round(nmax/nplt)+1);
% vv = fftn(u0);
ii = 1;
tic
for n
           1:nmax
   t
           n*h;
       = fftn(R.*u0 - u0.^3);
   Νv
       = E2.*vv +
                      Q.*Nv;
%
   a(ind) = 0;
   ua = real(ifftn(a));
      = fftn(R.*ua - ua.^3);
           E2.*vv + Q.*Na;
   b(ind) = 0;
   ub = real(ifftn(b));
   Nb = fftn(R.*ub - ub.^3);
       = E2.*a + Q.*(2*Nb-Nv);
%
   c(ind) = 0;
   uc = real(ifftn(c));
   Nc = fftn(R.*uc - uc.^3);
   vv = E.*vv + Nv.*f1 + 2*(Na+Nb).*f2 +
                                                     Nc.*f3;
   vv(ind) = 0;
%
          real(ifftn(vv));
       mod(n,nplt) == 0
   if
       uu(:,:,ii) =
                      u0;
       tt(ii) = t;
       ii = ii+1;
   end
end
toc
figure(1)
clf;
```

```
surf(X,Y,u0,'mesh','none');
axis equal;
view(2);
%% Processing to find the locations of the convex-concave pairs
xrng = xx(round((0.5-beta)*Nx):round((0.5+beta)*Nx));
onaxis = u0(round((0.5-beta)*Nx):round((0.5+beta)*Nx),Ny/2);
figure(2)
plot(xrng,onaxis)
zci = find(onaxis(2:end).*onaxis(1:end-1) <= 0);</pre>
                                                       % Returns Zero-Crossing Indices Of
upval = onaxis(zci+1);
downval = onaxis(zci);
zx = (upval.*xrng(zci)-downval.*xrng(zci+1))./(upval-downval);
figure(2)
clf
plot(xrng,onaxis,'-r')
hold on
plot(zx,zeros(size(zci)),'bp')
hold off
grid
legend('Signal', 'Approximate Zero-Crossings');
semi_mjr = beta*Lx;
semi_mnr = beta*Ly;
xc = semi_mjr*min(1,max(-1,semi_mjr*zx/(semi_mjr^2-semi_mnr^2)));
rho = sqrt((zx-xc).^2 + semi_mnr^2*(1-xc.^2/semi_mjr^2));
figure(3)
if (init_flag < 4)</pre>
    clf
end
plot(rho/pi,'-o');
hold on;
figure(4)
if (init_flag < 4)</pre>
    clf
end
```

```
plot(0.5+(1:length(zx)-1),diff(rho)/pi,'-*');
hold on
%% Compute energy density
eloc = (1-xi.^2-eta.^2).*fftn(u0);
eloc(ind) = 0;
eloc = ifftn(eloc).^2;
u0sq = fftn(u0.^2);
u0sq(ind) = 0;
u0sq = ifftn(u0sq);
u04th = fftn(u0sq.^2);
u04th(ind) = 0;
u04th = ifftn(u04th);
edens = 0.5*(eloc-R.*u0sq + 0.5*u04th);
figure(5)
clf;
surf(X,Y,edens,'mesh','none');
axis equal;
view(2);
figure(6)
clf;
surf(X,Y,u0sq,'mesh','none');
axis equal;
view([0 0 30]);
%% Processing the phase
%phase = unwrap(angle(hilbert(0.1+u0(round((0.5-beta)*Nx):round((0.5+beta)*Nx),1:Ny/2)')
phase = zeros(Nx/2-round((0.5-beta)*Nx)+1,Ny);
for ii = round((0.5-beta)*Nx):Nx/2
    ind = find(R(ii,:) > 0);
    mindx = ind(1);
    mxndx = ind(end);
   u1 = (0.5+R(ii,mindx:mxndx)).*u0(ii,mindx:mxndx);
    uext = [u1 \ 0 \ -u1 \ 0];
```

```
phase_loc = unwrap(angle(hilbert(uext)))/(2*pi);
    phase_loc = phase_loc(1:mxndx-mindx+1);
    reverse = phase_loc(end:-1:1);
    phase_loc = min(phase_loc,reverse);
    phase(ii+1-round((0.5-beta)*Nx),:) = ...
         [phase_loc(1)*ones(1,mindx-1) phase_loc phase_loc(end)*ones(1,Ny-mxndx)];
end
nmx = 256;
q = 2*pi*(1:nmx)/nmx;
bdry = beta*[Lx*cos(q);Ly*sin(q)];
imx = length(xrng);
phase0 = zeros(imx,1);
for ii = 1:imx
    phaseO(ii) = sqrt(min((xrng(ii)-bdry(1,:)).^2+(bdry(2,:)).^2))/(2*pi);
end
figure(7);
clf
plot(xrng(1:imx),phase0);
hold on;
plot(X(round((0.5-beta)*Nx):Nx/2,Ny/2),phase(:,Ny/2));
% filter to compute amplitude
rho = 2*real(ifftn(exp(-(xi.^2+eta.^2)).*fftn(u0.^2)));
figure(8)
surf(X,Y,rho,'mesh','none');
%axis equal;
view(2);
% figure(9)
% \operatorname{surf}(X(\operatorname{round}((0.5-\operatorname{beta})*Nx):\operatorname{round}((0.5+\operatorname{beta})*Nx),1:Ny/2),...
      Y(round((0.5-beta)*Nx):round((0.5+beta)*Nx),1:Ny/2),...
      phase,'mesh','none');
% %axis equal;
% view([0 0 30]);
phase = [phase; phase(end-1:-1:1,:)];
figure(9)
mesh(X(round((0.5-beta)*Nx):round((0.5+beta)*Nx),:),...
```

```
Y(round((0.5-beta)*Nx):round((0.5+beta)*Nx),:),...
    2*pi*phase,-phase)
colormap colorcube
%, 'mesh', 'none', 'colormap', 'lines');
%axis equal;
view(2);
figure(1)
clf;
surf(X(round((0.5-beta)*Nx):round((0.5+beta)*Nx),:),...
    Y(round((0.5-beta)*Nx):round((0.5+beta)*Nx),:),...
    u0(round((0.5-beta)*Nx):round((0.5+beta)*Nx),:),'mesh','none');
axis equal;
view([0 0 30]);
figure(10)
surf(X(round((0.5-beta)*Nx):round((0.5+beta)*Nx),:),...
    Y(round((0.5-beta)*Nx):round((0.5+beta)*Nx),:),...
    (0.5+R(round((0.5-beta)*Nx):round((0.5+beta)*Nx),:))...
    .*cos(2*pi*phase), 'mesh', 'none');
axis equal;
view([0 0 30]);
%% Postscript
% I'm going to split the initialization code from the time stepping. This
% makes it easier to continue an existing solution for a longer time, or
% restarting as appropriate.
```