

## 18.369 Problem Set 5

### Problem 1: Group Velocity and Material Dispersion

In class, we showed (following the book) that the group velocity  $d\omega/dk$ , computed via  $d(\omega^2)/dk = \langle H_k, \frac{\partial \hat{O}_k}{\partial k} H_k \rangle / \langle H_k, H_k \rangle$ , was equal to Poynting flux divided by energy density (both averaged over the unit cell).

Now, go through a similar Hellman–Feynman derivation, but in this case assume that we have a lossless *dispersive* material with a real  $\epsilon(\mathbf{x}, \omega)$ . In this case, when you take the  $k$  derivative, apply the chain rule to obtain a term with  $\frac{\partial \epsilon}{\partial \omega} \frac{d\omega}{dk}$  on the right-hand side. Solve for  $d\omega/dk$  and show that it is Poynting flux divided by energy density, but the energy density is now the “Brillouin” energy density of a lossless dispersive medium, which we gave in the notes on time evolution and energy conservation:

$$\frac{1}{4} \left[ \frac{\partial(\omega \epsilon)}{\partial \omega} |\mathbf{E}|^2 + |\mathbf{H}|^2 \right]$$

(for  $\mu = 1$ , where we have an additional  $1/2$  factor from the time average).

### Problem 2: Transmission spectra

Based on the sample code in the Jupyter notebook, compute the TM transmission spectrum for planewaves incident upon  $N_x$  layers of the square lattice of rods

- Compute the transmission spectrum for frequencies from  $0.2c/a$  to  $0.8c/a$  as a function of  $N_x$ , for  $N_x = 1, 2, 3, 5, 6$ , and plot them (on a single plot). The transmission spectrum should be normalized by dividing by the transmission for  $N_x = 0$  (no holes). Relate the features of this transmission spectrum to the band diagram in the sample code.
- Predict analytically at what frequency  $\omega_0$  you should start to see additional diffracted orders in the reflected wave (i.e. reflected waves at angles in *addition* to the normal  $0^\circ$  reflection). Now, modify the simulation to use a TM *continuous-wave* (CW) source and output  $E_z$  at the end and show that there is a qualitative change in the reflected field pattern if you put in a frequency *just* below  $\omega_0$  versus a frequency *just* above  $\omega_0$ . If you look *just* below  $\omega_0$ , then you will have to increase the “pad” parameter in order to see an undisturbed  $0^\circ$  reflection pattern far from the crystal—why?