

# From Brachistochrone Curve to Calculus of Variations

## How Understanding of Physics Evolve with Mathematics

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# History of Mathematics and Physics

Mathematics and Physics are highly inter-correlated (but at the same time fundamentally different).

Their relationship can be split into three stages:

- 0 Aristotle, Archimedes, ...: No particular relation. Physics was highly experimental and experience-based.
- 1 Newton: Vector Analysis (and Real Analysis/Calculus).
- 2 Euler, Lagrange: Functional Analysis (Calculus of Variations).
- 3 Noether: Symmetry (Group Theory and Modern Algebra). e.g., QFT, CPT, QED, QCD.

# Physics at Newton's Time

Newton's work:

- *Philosophiæ Naturalis Principia Mathematica*, 1687, gives rise to Classical Mechanics.
- One of the first Mathematical Formulations of Physics.
- Relies on the geometry of our space (being flat).
- Kepler's Laws, elementary electromagnetism, and even elementary thermal physics can be described using forces and motion.

Restrictions:

- Highly relies on a straight-line motion, which restricts the degree of freedom.
- Uses local linear behaviour to analyse complicated paths.
- This makes thermodynamics highly inaccurate and unable to solve certain questions.

# The Brachistochrone Curve Problem

*This question is related to the **Tautochrone Curve Problem** (which has the same solution).*

## Brachistochrone Curve Problem (Galilei, 1638 and Bernoulli, 1696)

What is the **Brachistochrone Curve**, or the curve of fastest descent, which is the one lying on the plane between a point  $A$  and a lower point  $B$  (where  $B$  is not directly below  $A$ ) on which a bead slides frictionlessly under the influence of a uniform gravitational field along the curve from  $A$  to  $B$  in the shortest **time**?

Please guess what the curve could be before we continue.

# History of the Problem

- Galilei, 1638: Concluded that the arc of a circle is faster than any number of its chords and concluded incorrectly that the arc is the quickest - but warned of possible fallacies and the need for a higher science.
- Johann Bernoulli and Jakob Bernoulli, 1696: Derived the same conclusion (but Johann was incorrect). They also realised this has the same solution as the Tautochrone Problem (Huygens).
- Newton, 1697: Found the solution because he stayed up for a whole night and posted to Bernoulli.

## Abstraction and Formulation

To simplify the problem, we build a coordinate system where the origin is the starting point, the  $y$  axis points down, and the  $x$  axis points to the right. Also, assume the ending point is  $(a, b)$ .

If we denote the curve as  $y = y(x)$ , we must have boundary conditions  $y(0) = 0$  and  $y(a) = b$ .

By conservation of energy, we will have the speed  $v$  at a general point  $(x, y)$  satisfies that

$$\frac{1}{2}mv^2 = mgy \implies v = \sqrt{2gy}.$$

Also, since  $v = \frac{ds}{dt}$ , and by Pythagorean Theorem we have

$$ds = \sqrt{dx^2 + dy^2} = dx\sqrt{1 + \left(\frac{dy}{dx}\right)^2} = \sqrt{1 + (y')^2}dx.$$

This means that

$$v = \frac{ds}{dt} = \sqrt{1 + (y')^2} \frac{dx}{dt},$$

and hence

$$dt = \sqrt{\frac{1 + (y')^2}{2gy}} dx.$$

Therefore, the total time taken,  $T$ , will satisfy that

$$T = \int dt = \int_0^a \sqrt{\frac{1 + (y')^2}{2gy}} dx = \frac{1}{\sqrt{2g}} \int_0^a \sqrt{\frac{1 + (y')^2}{y}} dx.$$

Our task is to find  $y$  that it minimises such  $T$ .



# Functionals

## Definition (Functional)

A function  $\mathcal{F}$  is called a functional if it takes in some function as its parameter and outputs a scalar value as its result.

## Examples

- $\mathcal{I}[f] = f(0)$  is a functional since it takes in a function  $f$  and outputs its value at 0.
- $\mathcal{J}[f] = \int_0^1 f(x)dx$  is a functional. It outputs the area under the curve  $f(x)$  from  $x = 0$  to  $x = 1$ .
- $\mathcal{K}[f] = \frac{df}{dx}$  is not a functional. It does not output a scalar value.

# Euler-Lagrange Equation

## Theorem (Euler-Lagrange Equation)

*A functional  $\mathcal{J}$  takes in a path  $y$  with the boundary conditions  $y(a) = A$  and  $y(b) = B$ :*

$$\mathcal{J}[y] = \int_a^b L(x, y, y') dx,$$

*where  $L$  is some expression.*

*$\mathcal{J}$  takes some extremum if and only if the equation*

$$\frac{\partial L}{\partial y} - \frac{d}{dx} \frac{\partial L}{\partial y'} = 0$$

*is satisfied.*

### Example (Shortest Path)

The length of a path  $y = y(x)$  between two points  $(a, A)$  and  $(b, B)$  is

$$s = \int_a^b \sqrt{dx^2 + dy^2} = \int_a^b \sqrt{1 + (y')^2} dx.$$

The  $L$  in the equation is  $L(x, y, y') = \sqrt{1 + (y')^2}$ .

We will realise that

$$\frac{\partial L}{\partial y} = 0, \frac{\partial L}{\partial y'} = \frac{y'}{\sqrt{1 + (y')^2}}.$$

## Example (Shortest Path)

$$\frac{\partial L}{\partial y} = 0, \frac{\partial L}{\partial y'} = \frac{y'}{\sqrt{1 + (y')^2}}.$$

If we substitute this into the equation, we will get that

$$\frac{d}{dx} \frac{y'}{\sqrt{1 + (y')^2}} = 0,$$

and therefore

$$\frac{y'}{\sqrt{1 + (y')^2}} = \text{const.}$$

This means that

$$y' = \frac{\text{const}}{\sqrt{1 - \text{const}^2}} = \text{another const},$$

which means it is a straight line!

## Solution of Problem

In the original problem, we would like to minimise

$$T = \frac{1}{\sqrt{2g}} \int_0^a \sqrt{\frac{1 + (y')^2}{y}} dx.$$

Our  $L$  here is

$$L(x, y, y') = \sqrt{\frac{1 + (y')^2}{y}}.$$

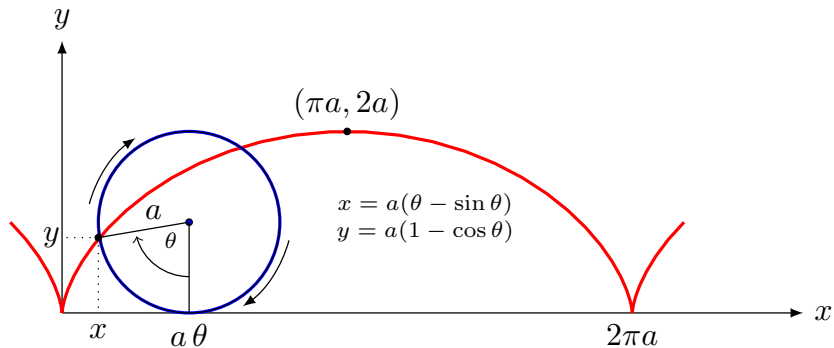
(If you trust my differentiation skills) We will have

$$\frac{\partial L}{\partial y} = -\frac{1}{2} \sqrt{\frac{1 + y'^2}{y^3}}, \quad \frac{\partial L}{\partial y'} = \frac{y'}{\sqrt{y(1 + y')^2}},$$

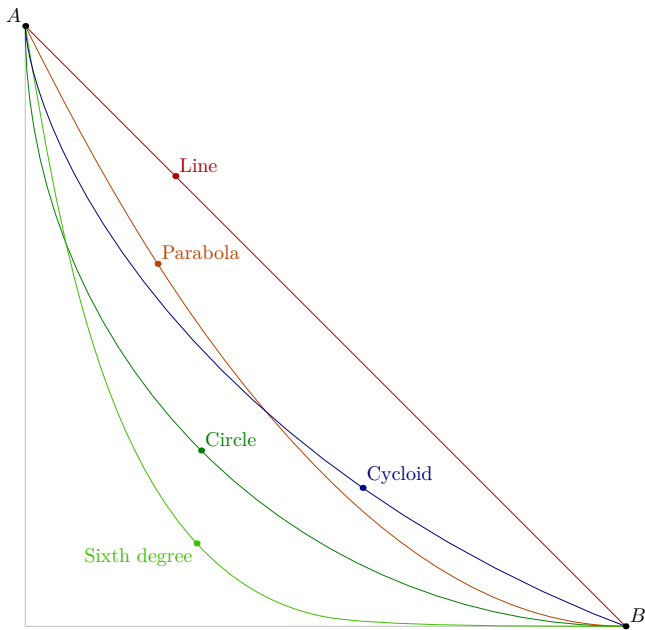
and plugging it back into the equation, mathematicians will tell us the solution is a cycloid.

# Visuallisation

A cycloid:



## Different curves:



# Calculus of Variations

Calculus of Variations is the study of how we can minimise functionals by varying functions.

The idea is similar:

- To minimise  $y = f(x)$ , we solve for  $f'(x) = 0$ .
- To minimise  $t = f(x, y, z)$ , we solve for  $\frac{\partial f}{\partial x} = \frac{\partial f}{\partial y} = \frac{\partial f}{\partial z} = 0$ .
- To minimise  $\mathcal{L}[f]$ , we solve for the Euler-Lagrange equation, which generalises to Euler-Poisson Equation and others.

The idea of variation of a function is by doing  $y(x) = y_0(x) + \epsilon\eta(x)$  where  $\eta(x)$  is an arbitrary function and  $\epsilon \rightarrow 0$ .



# Lagrangian and Hamiltonian Mechanics

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \\ & + i\bar{\psi}\not{D}\psi + h.c. \\ & + \bar{\psi}_i\gamma_{ij}\psi_j\phi + h.c. \\ & + |D_\mu\phi|^2 - V(\phi)\end{aligned}$$

is the Lagrangian of the standard model.

Doing something similar to (but much more complicated than) the Euler-Lagrange equation and solving it will enable you to explain everything apart from gravity.

Hamilton, 1833 provides an alternative formulation using the Hamiltonian  $\mathcal{H}$ .

These will be covered next week by Dara Daneshvar - come if you are interested!

# Benefits

If we look back at our derivation of the solution:

- we used energy rather than forces,
- we used path rather than points,

and we are considering an action along the path.

- It does not assume any geometry of the space, and does not ignore interior structures of objects, does not use simple linear approximations.
- It converts all problems to finding the extremum of some quantity - and applying variations to such quantity will enable us to form a sophisticated but self-consistent system.

It provides a unified way to solve all kinds of different physical problems.

# Applications

- Direct applications in mechanics include the Principle of Least Action and Hamilton's Principle.
- We can also apply variation to maximise entropy - or minimise free energy - in thermodynamics, and they lead to the same result.
- We can also apply variation on the Lagrangian density of an electromagnetic field to deduce Maxwell's Equations (and this even works for QED).
- We can use calculus of variations to deduce Einstein's field equations based on Einstein-Hilbert action in general relativity.

# Restrictions of this Mathematical Formulation

Mathematics should not rely on intuition, but Physics should, or at least an intuitive explanation should exist alongside an abstract formulation.

So what is

$$L(x, y, y') = \sqrt{\frac{1 + (y')^2}{y}},$$

and what is

$$\frac{\partial L}{\partial y} - \frac{d}{dx} \frac{\partial L}{\partial y'} = 0?$$

What is action, what is entropy, what is free energy?

This is very much less intuitive than forces and speed in Newtonian Mechanics.

# Symmetry, Group Theory, and Noether's Theorem

## Theorem (Noether's Theorem)

*Every continuous symmetry of the action of a physical system with conservative forces has a corresponding conservation law.*

## Examples

- Energy conserves since physical laws have time symmetry (i.e., physical laws do not change over time).
- Momentum conserves since physical laws have space symmetry (i.e., physical laws do not change under translation).

Under this formulation, the Lagrangian we see can simply take the mathematical structure of a group, specifically the group

$$SU(3) \times SU(2) \times U(1)$$

This is how Physics in the recent century developed.

*The chief forms of beauty are order and symmetry and definiteness,  
which the mathematical sciences demonstrate in a special degree.*  
– Aristotle

## Solving the Differential Equation

If we recall

$$\frac{\partial L}{\partial y} = -\frac{1}{2} \sqrt{\frac{1 + y'^2}{y^3}}, \quad \frac{\partial L}{\partial y'} = \frac{y'}{\sqrt{y(1 + y'^2)}},$$

we can simplify the equation to

$$\frac{1}{2}(1 + y'^2) + yy'' = 0.$$

If we multiply  $2y'$  on both sides,

$$y'(1 + y'^2) + 2yy'y'' = 0 \implies \frac{d}{dx} [y(1 + y'^2)] = 0.$$

This means that for some constant  $k$ , we have

$$y(1 + y'^2) = k \implies y' = \frac{dy}{dx} = \sqrt{\frac{k - y}{y}}.$$

If we let  $y = k \sin^2 \theta = \frac{1}{2}k(1 - \cos 2\theta)$ , we notice that

$$\begin{aligned}x &= \int \sqrt{\frac{k-y}{y}} dy \\&= \int \sqrt{\frac{k \sin^2 \theta}{k - k \sin^2 \theta}} d(k \sin^2 \theta) \\&= \int \frac{\sin \theta}{\cos \theta} 2k \sin \theta \cos \theta d\theta \\&= \int 2k \sin^2 \theta d\theta \\&= k \int (1 - \cos 2\theta) d\theta \\&= k\theta - \frac{1}{2}k \sin 2\theta.\end{aligned}$$

Now we let  $a = \frac{1}{2}k$ ,  $t = 2\theta$ , we will get the parametric equation of  
 $(x, y) = (a(t - \sin t), a(1 - \cos t)).$