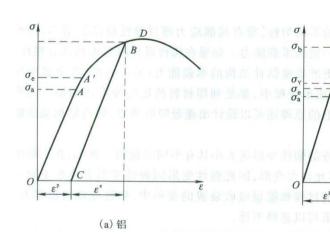
前置知识



工程应变

$$\varepsilon = \frac{l - l_0}{l}$$

名义应力

(b) 低碳钢

$$\sigma = \frac{F}{A_0}$$

弹性变形

 $\varepsilon^{\,\mathrm{e}}$

塑性变形

 $\varepsilon^{\,\mathrm{p}}$

比例极限

 $\sigma_{
m a}$

弹性极限

 $\sigma_{
m e}$

屈服应力

 $\sigma_{
m Y}$

强度极限

 $\sigma_{
m b}$

条件屈服应力(产生0.2%的塑性应变相对应的应力)

 $\sigma_{0.2}$

真应力

$$\overline{\sigma} = \frac{F}{A}$$

自然应变

$$\overline{\varepsilon} = \ln(1 + \varepsilon)$$

第一章

1. 加、卸载准则

在简单应力下, 可以写成

$$\sigma d\sigma > 0$$
 加载 $\sigma d\sigma < 0$ 卸载 $\sigma d\sigma < 0$

把简单拉伸试件在塑性阶段的应力-应变关系归纳为

$$d\sigma = E_t d\varepsilon \quad \text{対 } \tau \sigma d\sigma > 0$$
$$d\sigma = E d\varepsilon \quad \text{対 } \tau \sigma d\sigma < 0$$

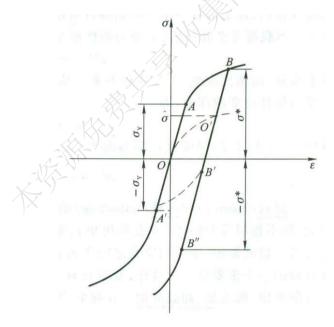
2. 后继屈服应力的函数形式

$$\sigma = H(\varepsilon^{p}) \iff \phi(\sigma, h_{a}) = 0$$

其中 h_a 是记录材料塑性加载历史的参数, ϕ 则称为加载函数.

3. Banschinger 效应

对于单晶体,在加载到强度极限后开始卸载,然后再反向加载,其反向屈服应力的绝对值比初始屈服应力 σ_Y 要大,即正向强化的反向亦强化。然而,对于一般非单晶体材料,在这种情况下反向屈服应力的绝对值比初始屈服应力 σ_Y 要小,即正向强化时反向会弱化。



4. 材料的塑性变形的共同特点

- (1)应力-应变关系呈非线性;
- (2)路劲相关性:由于材料在加载过程和卸载过程中服从不同的规律,应力与应变间不存在单值对应关系;
 - (3)耗散性:由于塑性应变不可恢复,故外力所作的塑性功具有不可逆性.

5. 静水压力试验

- (1)静水压与材料的体积改变之间近似服从线性弹性规律;
- (2)材料的塑性变形与静水压力无关.

6. 应变的对数定义

(1)可以对应变使用加法;

$$\varepsilon \neq \varepsilon_1 + \varepsilon_2 + \varepsilon_3$$

$$\overline{\varepsilon} = \overline{\varepsilon}_1 + \overline{\varepsilon}_2 + \overline{\varepsilon}_3$$

(2)体积不可压缩条件 $(1+\varepsilon_1)(1+\varepsilon_2)(1+\varepsilon_3)=1$ 可以简单地表示为

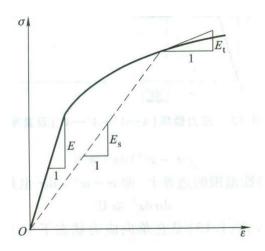
$$\overline{\varepsilon}_1 + \overline{\varepsilon}_2 + \overline{\varepsilon}_3 = 0$$

7. 材料塑性行为的基本假设

- (1)材料的塑性行为与时间、温度无关;
- (2)材料具有无限的韧性,即认为材料可以无限地变形而不出现断裂;
- (3)变形前材料是初始各向同性的,且拉伸和压缩的 $\bar{\sigma} \bar{\epsilon}$ 曲线一致;
- (4)关于卸载和后继屈服的假设:在产生塑性变形后卸载载荷,材料服从弹性规律;重新加载后的后继屈服应力等于卸载前的应力,这就是说重新加载达到屈服后的 $\sigma \varepsilon$ 曲线是卸载前 $\sigma \varepsilon$ 曲线的延伸线;
- (5)关于弹性和塑性的假设:在任何条件下的应变总能分解为弹性和塑性两部分,即 $\varepsilon=\varepsilon^{\rm e}+\varepsilon^{\rm p}$;材料的弹性性质不因塑性变形而改变,即 $\varepsilon^{\rm e}=\frac{\sigma}{E}$,其中弹性模量 E 是与塑性变形无关的常数;
- (6)塑性变形是在体积不变(不可压缩)条件下发生的. 静水压力只产生体积的弹性变化, 不产生塑性变形;
 - (7)关于材料稳定性的假设: 当应力单调变化时, 假设 $\sigma-\varepsilon$ 曲线具有以下不等式

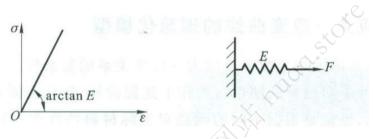
$$E > E_{\rm s} \geqslant E_{\rm t} \geqslant 0$$

其中 $E_{\rm s} = \frac{\sigma}{E}$ 和 $E_{\rm t} = \frac{{
m d}\sigma}{{
m d}\varepsilon}$ 分别是 $\sigma - \varepsilon$ 曲线的割线模量和切线模量.

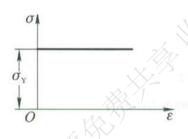


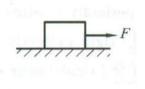
8. 应力-应变曲线的理想化模型

(1)理想弹性模型

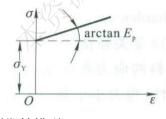


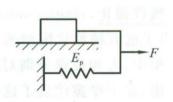
(2)理想刚塑性模型



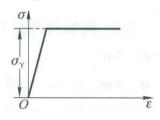


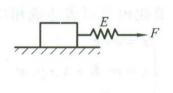
(3)刚-线性强化模型





(4)理想弹塑性模型





(5)弹-线性强化模型



9. 强化模型

(1)等向强化模型:不论是拉伸或压缩,应变强化总是相等地产生和发展.其表达式为

$$|\sigma| = \sigma^* = \psi \left(\int |d\varepsilon^p| \right) = F \left(\int dW^p \right)$$

其中 σ^* 是 $|\sigma|$ 在此前的塑性变形历史中曾经达到过的最大值,且 $\sigma^*>\sigma_{\rm Y}$; $\int |{
m d} \epsilon^{
m p}|$ 表示塑性 应变按照绝对值进行积累,而函数 ψ 可以根据材料的拉伸试验得出; $\int {
m d} W^{
m p}$ 表示塑性比功,函数F可以根据材料的拉伸试验得出. σ^* 、 $\int |{
m d} \epsilon^{
m p}|$ 和 $\int {
m d} W^{
m p}$ 都是单调变化的正数,都可以选作记录材料塑性变形历史的参数.

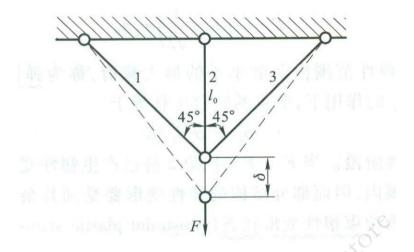
(2)随动强化模型:假定由于 Bauschinger 效应减少了反方向加载时的屈服应力,而总的 弹性范围的大小保持不变. 其表达式为

$$|\hat{\sigma} - arepsilon| = \sigma_{
m Y}$$

(3)组合强化模型:将上述两种强化模型组合起来,写成

第二章

1. 只施加竖直力的理想弹塑性材料的三杆桁架



设各杆的初始截面积均为 A_0 ,第 2 杆的初始长度为 I_0 ,相邻杆间夹角为 45° ,则第 1 杆和第 3 杆的长度均为 $\sqrt{2}\,I_0$,用 N_1 、 N_2 和 N_3 分别表示杆的内力,则由节点平衡方程得出

$$N_1 = N_3$$
, $\frac{N_1 + N_3}{\sqrt{2}} + N_2 = F$

用应力表示可得

$$\sqrt{2}\,\sigma_1 + \sigma_2 = \frac{F}{A_0}$$

变形协调条件

$$\delta = \varepsilon_2 l_0 = 2\varepsilon_1 l_0$$

于是有

$$\varepsilon_2 = 2\varepsilon_1$$

(1)弹性阶段:应力关系有

$$\sigma_2 = 2\sigma_1$$

解得

$$\left\{egin{aligned} \sigma_1 &= rac{1}{2+\sqrt{2}}\cdotrac{F}{A_0}\ \sigma_2 &= rac{2}{2+\sqrt{2}}\cdotrac{F}{A_0} \end{aligned}
ight.$$

所以第2杆先屈服,此时

$$\sigma_2 = \sigma_{
m Y}$$
 , $F_{
m e} = \left(1 + rac{1}{\sqrt{2}}\right) \sigma_{
m Y} A_0$, $\delta_{
m e} = arepsilon_2 l_0 = rac{\sigma_{
m Y} l_0}{E}$

(2)约束塑性阶段:此时有

$$\left\{egin{aligned} \sigma_2 &= \sigma_{
m Y} \ \sigma_1 &= rac{1}{\sqrt{2}} \left(rac{F}{A_0} - \sigma_{
m Y}
ight) \end{aligned}
ight.$$

(3)塑性流动阶段:第1杆屈服,此时

$$\sigma_{1} = \sigma_{2} = \sigma_{Y}$$
, $F_{Y} = (1 + \sqrt{2})\sigma_{Y}A_{0}$, $\delta_{y} = 2\varepsilon_{1}l_{0} = \frac{2\sigma_{Y}l_{0}}{E} = 2\delta_{e}$

(4)卸载阶段:由增量法

$$riangle \sigma_1 = rac{1}{2+\sqrt{2}} \cdot rac{ riangle F}{A_0} = E riangle arepsilon_1, \ \ riangle \sigma_2 = rac{2}{2+\sqrt{2}} \cdot rac{ riangle F}{A_0} = E riangle arepsilon_2, \ \ riangle F = \left(1+\sqrt{2}
ight) \sigma_{
m Y} A_0$$

因此有

$$riangle \sigma_{
m l} = rac{1}{\sqrt{2}}\,\sigma_{
m Y}\,,\;\; riangle \sigma_{
m 2} = \sqrt{2}\,\sigma_{
m Y}\,,\;\; riangle arepsilon_{
m l} = rac{1}{\sqrt{2}}\,rac{\sigma_{
m Y}}{E}\,,\;\; riangle arepsilon_{
m 2} = \sqrt{2}\,rac{\sigma_{
m Y}}{E}$$

叠加上初始应力可得

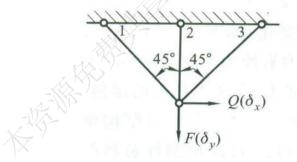
上例始近月時

$$\sigma_{1}^{\circ} = \sigma_{Y} - \frac{1}{\sqrt{2}}\sigma_{Y} = \left(1 - \frac{1}{\sqrt{2}}\right)\sigma_{Y} > 0, \quad \sigma_{2}^{\circ} = \sigma_{Y} - \sqrt{2}\sigma_{Y} = \left(1 - \sqrt{2}\right)\sigma_{Y} < 0$$

$$\varepsilon_{1}^{\circ} = \frac{\sigma_{Y}}{E} - \frac{1}{\sqrt{2}}\frac{\sigma_{Y}}{E} = \left(1 - \frac{1}{\sqrt{2}}\right)\frac{\sigma_{Y}}{E} > 0, \quad \varepsilon_{2}^{\circ} = 2\frac{\sigma_{Y}}{E} - \sqrt{2}\frac{\sigma_{Y}}{E} = \left(2 - \sqrt{2}\right)\frac{\sigma_{Y}}{E} > 0$$

$$\delta_{y}^{\circ} = \frac{2\sigma_{Y}l_{0}}{E} - 2 \cdot \frac{1}{\sqrt{2}}\frac{\sigma_{Y}}{E} \cdot l_{0} = \left(2 - \sqrt{2}\right)\frac{\sigma_{Y}l_{0}}{E} = \left(2 - \sqrt{2}\right)\delta_{e} > 0$$

2. 先施加竖直力再施加水平力的理想弹塑性材料的三杆桁架



对桁架先仅仅施加F直到极限载荷 $F_Y = (1 + \sqrt{2})\sigma_Y A_0$,同时保持Q = 0.此时

$$\sigma_1 = \sigma_2 = \sigma_3 = \sigma_Y$$
, $\delta_y = 2\varepsilon_1 l_0 = \frac{2\sigma_Y l_0}{E} = 2\delta_e$, $\delta_x = 0$

然后在保持节点竖直位移 δ_v 不变的情况下增加Q,这时F将有相应的改变. 用增量法有

$$\triangle \sigma_2 + rac{\triangle \sigma_1 + \triangle \sigma_3}{\sqrt{2}} = rac{\triangle F}{A_0}, \ rac{\triangle \sigma_1 - \triangle \sigma_3}{\sqrt{2}} = rac{\triangle Q}{A_0}$$

变形协调条件

$$\triangle \varepsilon_1 = \frac{\triangle \delta_y + \triangle \delta_x}{2l_0}, \ \triangle \varepsilon_2 = \frac{\triangle \delta_y}{l_0}, \ \triangle \varepsilon_3 = \frac{\triangle \delta_y - \triangle \delta_x}{2l_0}$$

可得

$$\triangle \varepsilon_2 = \triangle \varepsilon_1 + \triangle \varepsilon_3$$

现在保持 $\triangle \delta_v = 0$ 施加Q,故

$$\triangle \delta_x = \delta_x > 0$$
, $\triangle \varepsilon_1 = \frac{\triangle \delta_x}{2l_0} > 0$, $\triangle \varepsilon_2 = 0$, $\triangle \varepsilon_3 = -\frac{\triangle \delta_x}{2l_0} < 0$

即第1杆继续伸长,第2杆长度不变,第3杆发生卸载.于是

$$\triangle \sigma_1 = \triangle \sigma_2 = 0$$
, $\triangle \sigma_3 = E \triangle \varepsilon_3 = -\frac{E \delta_x}{2 I_0} < 0$

进而求得

$$\frac{\triangle F}{A_0} = \frac{\triangle \sigma_3}{\sqrt{2}} < 0$$
, $\frac{\triangle Q}{A_0} = -\frac{\triangle \sigma_3}{\sqrt{2}} > 0$, $\triangle F = -\triangle Q$

即Q增加时F必须减小. 当 $\triangle \sigma_3 = -2\sigma_Y$ 使 $\sigma_3 = -\sigma_Y$ 时,第 3 杆进入反向屈服,整个桁架再次进入塑性流动状态(保持 δ_y 不变即增加了约束,所以需要 3 根杆都屈服结构才进入塑性流动状态),此时

$$Q = Q_{\rm Y} = \sqrt{2} \, \sigma_{
m Y} \, A_0 \,, \; F = F_{
m Y} + \triangle F = \left(1 + \sqrt{2} \,\right) \sigma_{
m Y} \, A_0 - \sqrt{2} \, \sigma_{
m Y} \, A_0 = \sigma_{
m Y} \, A_0$$

叠加上初始状态可得

$$\begin{split} \sigma_{1} &= \sigma_{2} = \sigma_{\mathrm{Y}} \;,\;\; \sigma_{3} = -\;\sigma_{\mathrm{Y}} \;,\;\; \delta_{y} = 2\varepsilon_{1}l_{0} = \frac{2\sigma_{\mathrm{Y}}l_{0}}{E} \;,\;\; \delta_{x} = -\;2\triangle\varepsilon_{3}l_{0} = \frac{4\sigma_{\mathrm{Y}}l_{0}}{E} \\ \varepsilon_{3} &= \varepsilon_{\mathrm{Y}} - \frac{2\sigma_{\mathrm{Y}}l_{0}}{E} = -\;\varepsilon_{\mathrm{Y}} \;,\;\; \varepsilon_{1} = \varepsilon_{\mathrm{Y}} + \frac{2\sigma_{\mathrm{Y}}l_{0}}{E} = 3\varepsilon_{\mathrm{Y}} \;,\;\; \varepsilon_{2} = 2\varepsilon_{\mathrm{Y}} \end{split}$$

3. 按照比例加载的理想弹塑性材料的三杆桁架

按照 $F: Q=1: \sqrt{2}$ 单调比例增加,直到结构达到塑性极限状态。有协调关系

$$\varepsilon_2 = \varepsilon_1 + \varepsilon_3$$

于是有

$$\sigma_2 = \sigma_1 + \sigma_3$$

有应力方程的平衡方程

$$\sigma_2 + \frac{\sigma_1 + \sigma_3}{\sqrt{2}} = \frac{F}{A_0}, \ \frac{\sigma_1 - \sigma_3}{\sqrt{2}} = \frac{Q}{A_0}, \ Q = \sqrt{2} F$$

解得

$$\sigma_1 = \left(rac{\sqrt{2}}{1+\sqrt{2}} + 2
ight) rac{F}{2A_0} > 0 \; , \; \; \sigma_2 = rac{\sqrt{2}}{1+\sqrt{2}} rac{F}{A_0} > 0 \; , \; \; \sigma_3 = \left(rac{\sqrt{2}}{1+\sqrt{2}} - 2
ight) rac{F}{2A_0} < 0 \; .$$

必背公式

1. 三根桁架的应变方程(变形协调方程)

$$\triangle \varepsilon_1 = \frac{\triangle \delta_y + \triangle \delta_x}{2l_0}, \ \triangle \varepsilon_2 = \frac{\triangle \delta_y}{l_0}, \ \triangle \varepsilon_3 = \frac{\triangle \delta_y - \triangle \delta_x}{2l_0}$$

2. 应力边界条件

$$S_{Ni} = \sigma_{ij} l_j$$

3. 平均正应力(静水压力)

$$\sigma_{
m m}=rac{\sigma_{
m 11}+\sigma_{
m 22}+\sigma_{
m 33}}{3}=rac{1}{3}\,\sigma_{kk}$$

4. 应力张量分解

$$\sigma_{ij} = \sigma_{
m m} \delta_{ij} + s_{ij}$$

5. 应力张量的应力不变量

$$J_1 = \sigma_1 + \sigma_2 + \sigma_3$$
, $J_2 = -(\sigma_1 \sigma_2 + \sigma_2 \sigma_3 + \sigma_3 \sigma_1)$, $J_3 = \sigma_1 \sigma_2 \sigma_3$

6. 应力偏张量主值(主偏应力)

$$s_i = \sigma_i - \sigma_{\rm m}$$

7. 应力偏张量的应力不变量

$$\begin{split} J_1' &= s_1 + s_2 + s_3 = 0 \,, \ J_3' = s_1 s_2 s_3 = (\sigma_1 - \sigma_m) \, (\sigma_2 - \sigma_m) \, (\sigma_3 - \sigma_m) \\ J_2' &= - \left(s_1 s_2 + s_2 s_3 + s_3 s_1 \right) = \frac{1}{2} \left(s_1^2 + s_2^2 + s_3^2 \right) = \frac{1}{2} \left(s_{11}^2 + s_{22}^2 + s_{33}^2 + 2 s_{12}^2 + 2 s_{23}^2 + 2 s_{31}^2 \right) \\ &= \frac{1}{2} s_{ij} s_{ij} \\ &= \frac{1}{6} \left[\left(\sigma_1 - \sigma_2 \right)^2 + \left(\sigma_2 - \sigma_3 \right)^2 + \left(\sigma_3 - \sigma_1 \right)^2 \right] \\ &= \frac{1}{6} \left[\left(\sigma_x - \sigma_y \right)^2 + \left(\sigma_y - \sigma_z \right)^2 + \left(\sigma_z - \sigma_x \right)^2 + 6 \left(\tau_{xy}^2 + \tau_{yz}^2 + \tau_{zx}^2 \right) \right] \\ &= \frac{1}{6} \left[\left(\sigma_r - \sigma_\theta \right)^2 + \left(\sigma_\theta - \sigma_z \right)^2 + \left(\sigma_z - \sigma_r \right)^2 + 6 \left(\tau_{r\theta}^2 + \tau_{\theta z}^2 + \tau_{zr}^2 \right) \right] \end{split}$$

8. 八面体应力分量,正应力和剪应力

$$|F_8|^2 = \frac{1}{3} \left(\sigma_1^2 + \sigma_2^2 + \sigma_3^2\right), \ \ \sigma_8 = \frac{\sigma_{11} + \sigma_{22} + \sigma_{33}}{3} = \sigma_{\mathrm{m}}, \ \ au_8 = \sqrt{|F_8|^2 - \sigma_8^2} = \sqrt{\frac{2}{3}J_2'}$$

9. 等效应力

$$\overline{\sigma} = \sqrt{3J_2'} = \frac{1}{\sqrt{2}}\sqrt{(\sigma_1 - \sigma_2)^2 + (\sigma_2 - \sigma_3)^2 + (\sigma_3 - \sigma_1)^2}$$

$$\overline{\tau} = \sqrt{J_2'} = \frac{1}{\sqrt{6}}\sqrt{(\sigma_1 - \sigma_2)^2 + (\sigma_2 - \sigma_3)^2 + (\sigma_3 - \sigma_1)^2}$$

10. 最大剪应力

$$au_{ ext{max}} = rac{\sigma_1 - \sigma_3}{2}$$

11. 等效应变

$$ar{arepsilon} = \sqrt{rac{4}{3}I_2'} = rac{\sqrt{2}}{3}\sqrt{\left(arepsilon_1 - arepsilon_2
ight)^2 + \left(arepsilon_2 - arepsilon_3
ight)^2 + \left(arepsilon_3 - arepsilon_3 - arepsilon_3
ight)^2 + \left(arepsilon_3 - arepsilon_3
ight)^2 + \left(arepsilon_3 - arepsilon_3 - arepsilon_3
ight)^2 + \left(arepsilon_3 - arepsilon_3 - arepsilon_3
ight)^2 + \left(arepsilon_3 - arepsilon_3 - arepsilon_3
ight)^2 + \left(arepsilo$$

12. Tresca 屈服条件

$$rac{1}{2}\left(\sigma_{1}-\sigma_{3}
ight)=k$$
 $k=rac{1}{2}\sigma_{Y}$

单向拉伸

$$k = \frac{1}{2}\sigma_{\gamma}$$

纯剪切

$$k = \tau_{\scriptscriptstyle Y}$$

13. Mises 屈服条件

$$J_{2}' = \frac{1}{6} \left[\left(\sigma_{1} - \sigma_{2} \right)^{2} + \left(\sigma_{2} - \sigma_{3} \right)^{2} + \left(\sigma_{3} - \sigma_{1} \right)^{2} \right] = C$$

单向拉伸

$$J_2'=rac{1}{3}\sigma_{\mathrm{Y}}^2=C$$

纯剪切

$$J_2' = au_Y^2 = C$$

14. 受拉扭组合的薄壁圆管

Mises 屈服条件

$$\left(\frac{\sigma}{\sigma_{\rm Y}}\right)^2 + 3\left(\frac{\tau}{\sigma_{\rm Y}}\right)^2 = 1$$

Tresca 屈服条件

$$\left(\frac{\sigma}{\sigma_{\rm Y}}\right)^2 + 4\left(\frac{\tau}{\sigma_{\rm Y}}\right)^2 = 1$$

16. 平面主应力公式

$$\sigma' = \frac{\sigma_x + \sigma_y}{2} + \frac{1}{2}\sqrt{(\sigma_x - \sigma_y)^2 + 4\tau_{xy}^2}, \ \ \sigma'' = \frac{\sigma_x + \sigma_y}{2} - \frac{1}{2}\sqrt{(\sigma_x - \sigma_y)^2 + 4\tau_{xy}^2}, \ \ \sigma''' = 0$$

17. 广义 Hooke 定律

$$arepsilon_{ij} = rac{\sigma_{ij}}{2G} - rac{3v}{E} \sigma_{
m m} \delta_{ij}$$

18. 增量理论

Levy-Mises 理论(理想刚塑性,复杂加载, $\mu = \frac{1}{2}$)

$$\mathrm{d} \varepsilon_{ij} = \mathrm{d} \lambda s_{ij} \,, \; \mathrm{d} \lambda = \frac{3 \overline{\mathrm{d} \varepsilon}}{2 \sigma_{\mathrm{Y}}} \,, \; \frac{\mathrm{d} \varepsilon_{1}}{s_{1}} = \frac{\mathrm{d} \varepsilon_{2}}{s_{2}} = \frac{\mathrm{d} \varepsilon_{3}}{s_{3}}$$

Prandtl-Reuss 理论(理想弹塑性,复杂加载, $\mu \leq \frac{1}{2}$)

性,复杂加载,
$$\mu \leqslant \frac{1}{2}$$
) $\mathrm{d} \varepsilon_{ij} = \frac{\mathrm{d} s_{ij}}{2G} + \mathrm{d} \lambda s_{ij}$, $\mathrm{d} \lambda = \frac{3 \, \mathrm{d} W}{2\sigma_{\mathrm{Y}}^2}$ 载, $\mu = \frac{1}{2}$)

伊柳辛理论(幂强化,简单加载, $\mu = \frac{1}{2}$)

$$e_{ij} = \frac{1}{2G_s} s_{ij}, \ G_s = \frac{\overline{\sigma}}{3\overline{\varepsilon}}$$

19. 薄壁圆管受拉扭作用的量纲归一化

$$\sigma^2 + \tau^2 = 1$$

20. 阶梯变形路径积分

γ为常数

$$\varepsilon - \varepsilon_0 = \frac{1}{2} \ln \left(\frac{1+\sigma}{1+\sigma_0} \cdot \frac{1-\sigma_0}{1-\sigma} \right)$$

 ε 为常数

$$\gamma - \gamma_0 = \frac{1}{2} \ln \left(\frac{1+\tau}{1+\tau_0} \cdot \frac{1-\tau_0}{1-\tau} \right)$$