Teoria dei sistemi. State Space Analysis

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Modes Analysis

Continuous Time Systems



The discrete time case

Consider the evolution

$$x(t+1) = Ax(t) + Bu(t)$$
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$$x(2) = Ax(1) + Bu(1) = A^{2}x(0) + ABu(0) + Bu(1),$$

$$\vdots$$

$$x(k) = Ax(k-1) + Bu(k-1) = A^{k}x(0) + A^{k-1}Bu(0) + \dots + ABu(k-2)$$

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...which can be written as

$$x(k) = A^{k}x(0) + \sum_{i=0}^{k-1} A^{k-1-i}Bu(i),$$

$$y(k) = CA^{k}x(0) + C\sum_{i=0}^{k-1} A^{k-1-i}Bu(i) + Du(k).$$
(1)

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Observations

Free and forces evolution

As for the continuous—time systems, we can identify a free evolution (depending on the initial condition x(0)) and a forced evolution (resulting from the application of the input function).

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Role of A^k

- Since A^k is a constant once k is fixed, it turns out that the states are combined through a time invariant linear transformation.
- ► The combination of forced and unforced response is an application of the *superposition principle* between the states and the inputs.

Connection with Z-transform

Consider the autonomous system:

$$x(k+1)=Ax(k),$$

whose Z-Transform are given by

$$\mathcal{Z}(x(k+1)) = zX(z) - zx(0)$$
, and $\mathcal{Z}(Ax(k)) = AX(z)$,

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Hence,

$$zX(z) - zx(0) = AX(z) \Rightarrow (zI - A)X(z) = zx(0)$$

\Rightarrow X(z) = (zI - A)^{-1}zx(0).

Inverse Z-transform

▶ The sequence x(k) can be recoverd via the inverse Z-trasnform.

$$x(k) = \mathcal{Z}^{-1} \left(z(zI - A)^{-1} \right) x(0),$$

Proposition

The matrix A^k can be computed through the inverse of Z-Transform:

$$A^k = \mathcal{Z}^{-1} \left(z(zI - A)^{-1} \right).$$

Example

Harmonic Oscillator
Consider the following harmonic oscillator

$$x(k+1) = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} x(k) = Ax.$$

Compute the evolution.

Since

$$zI - A = \begin{bmatrix} z & 1 \\ -1 & z \end{bmatrix},$$

• we have that the eigenvalues are $\pm i$ and hence

$$(zI - A)^{-1} = \frac{1}{z^2 + 1} \begin{bmatrix} z & -1 \\ 1 & z \end{bmatrix}$$
$$A^k = \mathcal{Z}^{-1} \left(z(zI - A)^{-1} \right)$$
$$= \mathcal{Z}^{-1} \left(\begin{bmatrix} \frac{z^2}{z^2 + 1} & -\frac{z}{z^2 + 1} \\ \frac{z}{z^2 + 1} & \frac{z^2}{z^2 + 1} \end{bmatrix} \right)$$

Example

► ...finally...

$$A^{k} = \begin{bmatrix} \cos(\frac{\pi}{2}k) & -\sin(\frac{\pi}{2}k) \\ \sin(\frac{\pi}{2}k) & \cos(\frac{\pi}{2}k) \end{bmatrix}$$

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$$\sum_{i=0}^{k-1} A^{k-1-i} Bu(i) = A^{k-1} * Bu(k),$$

As a consequence,

$$\mathcal{Z}\left(A^{k-1}*Bu(k)\right)=(zI-A)^{-1}BU(z).$$

Applying the superposition principle, we get

$$\mathcal{Z}(y(k)) = Cz(zI - A)^{-1}x(0) + C(zI - A)^{-1}BU(z) + DU(z).$$

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$$Z(y(k)) = Cz(zI - A)^{-1}x(0) + C(zI - A)^{-1}BU(z) + DU(z).$$

If we consider x(0) = 0 we find the expression for the trasfer function.

$$\frac{Y(z)}{U(z)}=C(zI-A)^{-1}B+D.$$

Example

Second order system

For the linear system

$$x(k+1) = \begin{bmatrix} 0.5 & 1 \\ 0 & -0.5 \end{bmatrix} x(k) + \begin{bmatrix} 0 \\ 1 \end{bmatrix} u(k),$$

$$y(k) = \begin{bmatrix} 1 & -1 \end{bmatrix} x(k),$$

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▶ The result is in the following function:

$$G(z) = \begin{bmatrix} 1 & -1 \end{bmatrix} (zI - A)^{-1} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \frac{-z + 1.5}{z^2 - 0.25}$$

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▶ We can restrict to $G(z) = C(zI - A)^{-1}BU(z)$.

▶ the (i,j) element of $(zI - A)^{-1}$ can be computed using the Cramer's rule

$$(-1)^{i+j} \frac{\det \operatorname{Adj}_{i,j}}{\mathcal{P}(A)},$$

where $\mathcal{P}(A)$ is the *characteristic polynomial* of A, defined as

$$\mathcal{P}(A) = \det(zI - A).$$

▶ det $Adj_{i,j}$ has degree less than n and there is still the possibility of having a *cancellation* between the roots of $\mathcal{P}(A)$ and the roots of det $Adj_{i,j}$.

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- ▶ If $v \in \mathbb{R}^n$ is an eigenvector of a matrix A associated with the eigenvalue λ , i.e., $Av = \lambda v$, than it immediately follows that v is also an eigenvector of the matrix A^k associated with the eigenvalue λ^k , i.e., $A^kv = \lambda^kv$;

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- ▶ $\det(A^k) = [\det(A)]^k$. From this property, from $\det(AB) = \det(A) \det(B)$ it follows that $\det(AA^{-1}) = \det(A) \det(A^{-1}) = 1$, which is of course $\det(I) = 1$. Hence, the *eigenvalues does not change for similarity transformation*.

Matrix Powers for Diagonal Matrices

If a matrix is diagonalisable, the matrix power can be easily obtained. Indeed, let

$$T^{-1}AT = \Lambda = \begin{bmatrix} \lambda_1 & 0 & \dots & 0 \\ 0 & \lambda_2 & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & \lambda_n \end{bmatrix},$$

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We have that

$$A^{k} = T\Lambda^{k}T^{-1} = T \begin{vmatrix} \lambda_{1}^{k} & 0 & \dots & 0 \\ 0 & \lambda_{2}^{k} & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & \lambda_{n}^{k} \end{vmatrix} T^{-1}.$$

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- The diagonal blocks are associated to each complex and conjugated pair.
- As for the continuous time case, since $AV = V\Lambda$ we have $AVH = V\Lambda H = VHH^{-1}\Lambda H$, we have

$$A[v_r + jv_c | v_r - jv_c]H = A[v_r | v_c] = \begin{bmatrix} v_r | v_c \end{bmatrix} \begin{bmatrix} \sigma & \omega \\ -\omega & \sigma \end{bmatrix},$$

The matrix can be further transformed using the modulus $\rho = \sqrt{\sigma^2 + \omega^2}$ and phase $\theta = \arctan\left(\frac{\omega}{\sigma}\right)$, i.e.,

$$A[v_r|v_c] = [v_r|v_c]\rho \begin{bmatrix} \cos(\theta) & \sin(\theta) \\ -\sin(\theta) & \cos(\theta) \end{bmatrix}.$$

...which leads us to...

$$A^{k}[v_{r}|v_{c}] = [v_{r}|v_{c}]\rho^{k}\begin{bmatrix}\cos(k\theta) & \sin(k\theta)\\ -\sin(k\theta) & \cos(k\theta)\end{bmatrix}.$$

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A 4 X 4 example

let A be a matrix defined in $\mathbb{R}^{4\times 4}$, with λ_1 and λ_4 real and $\lambda_{2,3}=\sigma\pm j\omega$, we have

$$A^{k} = T\Lambda^{k}T^{-1} = T \begin{bmatrix} \lambda_{1}^{k} & 0 & 0 & 0 \\ 0 & \rho^{k}\cos(k\theta) & \rho^{k}\sin(k\theta) & 0 \\ 0 & -\rho^{k}\sin(k\theta) & \rho^{k}\cos(k\theta) & 0 \\ 0 & 0 & 0 & \lambda_{4}^{k} \end{bmatrix} T^{-1}.$$

▶ We can still use the Jordan form

- We can still use the Jordan form
- ▶ Indeed for a block diagonal matrix we have:

$$\begin{bmatrix} A_1 & 0 & \dots & 0 \\ 0 & A_2 & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & A_n \end{bmatrix}^k = \begin{bmatrix} A_1^k & 0 & \dots & 0 \\ 0 & A_2^k & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & A_n^k \end{bmatrix}.$$

► Hence, the name of the game becomes to understand how to compute the power of *Jordan miniblocks*.

We can see that

$$J_{i,l}^{k} = \begin{bmatrix} \lambda_{i} & 1 & 0 & \dots & 0 & 0 \\ 0 & \lambda_{i} & 1 & \dots & 0 & 0 \\ \vdots & \vdots & \ddots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & \lambda_{i} & 1 \\ 0 & 0 & 0 & \dots & 0 & \lambda_{i} \end{bmatrix}^{l} = (\Lambda_{i} + \overline{J}_{i,l})^{k} = \sum_{q=0}^{k} C_{q}^{k} \lambda_{i}^{k-q} \overline{J}_{i,l}^{q},$$

where

$$\begin{cases} C_q^k = \binom{k}{q} = \frac{k!}{q!(k-q)!} & \text{if } k \ge q, \\ C_q^k = 0 & \text{if } k < q, \end{cases}$$

that is a polynomial function of k of degree q, i.e., the binomial coefficient.

▶ Therefore, for a miniblock of dimension $p = m_{i,k}$, we have

$$J_{i,l}^k = \begin{bmatrix} \lambda_i^k & C_1^k \lambda_i^{k-1} & C_2^k \lambda_i^{k-2} & \dots & C_{q-2}^k \lambda_i^{k-q+2} & C_{q-1}^k \lambda_i^{k-q+1} \\ 0 & \lambda_i^k & C_1^k \lambda_i^{k-1} & \dots & C_{q-3}^k \lambda_i^{k-q+3} & C_{q-2}^k \lambda_i^{k-q+2} \\ \vdots & \vdots & \ddots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & \lambda_i^k & C_1^k \lambda_i^{k-1} \\ 0 & 0 & 0 & \dots & 0 & \lambda_i^k \end{bmatrix}.$$

An example

A 3 X 3 Matrix Let us consider the matrix

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$$A = \begin{bmatrix} \lambda & 1 & 0 \\ 0 & \lambda & 1 \\ 0 & 0 & \lambda \end{bmatrix}.$$

By applying the formulas, it follows that

$$A^{k} = \begin{bmatrix} \lambda^{k} & k\lambda^{k-1} & \frac{k(k-1)}{2}\lambda^{k-2} \\ 0 & \lambda^{k} & k\lambda^{k-1} \\ 0 & 0 & \lambda^{k} \end{bmatrix}.$$

▶ If we has generalised eigenvectors associated with complex eigenvalues, we can repeat the same reasoning of the continuous time case.

- ▶ If we has generalised eigenvectors associated with complex eigenvalues, we can repeat the same reasoning of the continuous time case.
- ▶ For the case p = 2 we have

$$J_r = H^{-1}JH = \begin{bmatrix} \sigma & \omega & 1 & 0 \\ -\omega & \sigma & 0 & 1 \\ 0 & 0 & \sigma & \omega \\ 0 & 0 & -\omega & \sigma \end{bmatrix} = \begin{bmatrix} W & I \\ 0 & W \end{bmatrix}.$$

For the matrix power, we recall that

$$J_r^k = \begin{bmatrix} W^k & kW^{k-1} \\ 0 & W^k \end{bmatrix} = \begin{bmatrix} W^k & C_1^k W^{k-1} \\ 0 & W^k \end{bmatrix},$$

Which corresponds to

$$J_r^k = \begin{bmatrix} \rho^k \begin{bmatrix} \cos(k\theta) & \sin(k\theta) \\ -\sin(k\theta) & \cos(k\theta) \end{bmatrix} & k\rho^{k-1} \begin{bmatrix} \cos((k-1)\theta) & \sin((k-1)\theta) \\ -\sin((k-1)\theta) & \cos((k-1)\theta) \end{bmatrix} \\ 0 & \rho^k \begin{bmatrix} \cos(k\theta) & \sin(k\theta) \\ -\sin(k\theta) & \cos(k\theta) \end{bmatrix} \end{bmatrix}.$$

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We can generalise this to p > 2...

Mode Analysis

- ▶ The unforced response of a linear system is just a linear composition of *only* the functions that can be found in the real Jordan form matrix.
- ▶ Such functions are called the *modes* of the system.

Continuous Time Systems

- Simple exponential functions $e^{\lambda t}$ given by Jordan miniblocks of dimension one with real eigenvalue λ
 - converge to 0 if $\lambda < 0$, constant if $\lambda = 0$, diverges if $\lambda > 0$
- ▶ Composite exponential functions $t^k e^{\lambda t}$ given by Jordan miniblocks of dimension m > 1 with real eigenvalue λ .
 - converge to 0 if $\lambda <$ 0, diverge polynomially if $\lambda =$ 0, diverge exponentially if $\lambda >$ 0
- ▶ Oscillating functions of the type $e^{\sigma t}\cos(\omega t)$ and $e^{\sigma t}\sin(\omega t)$ given by two miniblocks of dimension 1 associated to complex and conjugated eigenvalues $\lambda = \sigma \pm j\omega$.
 - converge to 0 if σ < 0, persistently oscillate if λ = 0, diverge exponentially if σ > 0



Continuous Time Systems

- ▶ Oscillating functions of the type $t^k e^{\sigma t} \cos(\omega t)$ and $t^k e^{\sigma t} \sin(\omega t)$ given by Jordan miniblocks of dimension m>1 (and hence $k=0,\ldots,m-1$) with complex and conjugated eigenvalues $\lambda=\sigma\pm j\omega$.
 - ▶ converge σ < 0, diverge polynomially if σ = 0 and k > 0 or diverge exponentially if σ > 0

- ▶ Powers λ^k given by Jordan miniblocks of dimension one with real eigenvalue λ .
 - converge $|\lambda|<1$, constant if $|\lambda|=1$ or diverge exponentially if $|\lambda|>1$. We have an oscillation if $\lambda<0$

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- Series of functions $C_q^k \lambda^{k-q}$ given by Jordan miniblocks of dimension m>1 (and hence $k=0,\ldots,m-1$) with real eigenvalue λ .
 - ▶ Converge if $|\lambda| < 1$, diverge polinomially if $\lambda = 1$, diverge exponentially of $|\lambda| > 1$. We still have an oscillation if $\lambda < 0$.

- ▶ Powers λ^k given by Jordan miniblocks of dimension one with real eigenvalue λ .
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 - ▶ Converge if $|\lambda| < 1$, diverge polinomially if $\lambda = 1$, diverge exponentially of $|\lambda| > 1$. We still have an oscillation if $\lambda < 0$.
- ▶ Oscillating sequence of the type $\rho^k \cos(k\theta)$ and $\rho^k \sin(k\theta)$ given by two miniblocks of dimension 1 associated to complex and conjugated eigenvalues $\lambda = \sigma \pm j\omega = \rho \mathrm{e}^{\pm j\theta}$.
 - \blacktriangleright Converge if $\rho<1,$ oscillate persistently if $\rho=1,$ diverge exponentially id $\rho>1$



- ▶ Oscillating series of the type $C_q^k \rho^{k-q} \cos((k-q)\theta)$ and $C_q^k r h o^{k-q} \sin((k-q)\theta)$ given by Jordan miniblocks of dimension m>1 (and hence $k=0,\ldots,m-1$) with complex and conjugated eigenvalues $\lambda=\sigma\pm j\omega=\rho e^{\pm j\theta}$.
 - ▶ Converge if ρ < 1, diverge polynomially if ρ = 1 and k > 0 or diverge exponentially if ρ > 1.