# Calculating Hermite-Gauss Modes in Python

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PauLisa.py was used to calculate Hermite-Gauss(HG) modes. This document presents intensity profiles for "pure" modes and phase and intensity slices for higher-order modes. Beam tilts and misalignments relative to the optical axis can be used to model tilt-to-length(TTL) coupling in LISA's long arm.

### 1. EQUATIONS FOR HG MODES

The Living Review Article [1] gives the general expression for HG modes as

$$u_{nm}(x,y,z) = (2^{n+m-1}n!m!\pi)^{-1/2} \frac{1}{w(z)} H_n\left(\frac{\sqrt{2}x}{w(z)}\right) H_m\left(\frac{\sqrt{2}y}{w(z)}\right) \exp\left(\frac{-ik(x^2+y^2)}{2R_c(z)} - \frac{x^2+y^2}{w(z)^2}\right). \tag{1}$$

In terms of the Gaussian beam parameter,

$$q(z) = iz_R + z - z_0 = q_0 + z - z_0 , (2)$$

this may be expressed as

$$u_{nm}(x,y,z) = u_n(x,z)u_m(y,z)$$

$$= \left(\frac{2}{\pi}\right)^{1/4} \left(\frac{1}{2^n n! w_0}\right)^{1/2} \left(\frac{q_0}{q(z)}\right)^{1/2} \left(\frac{q_0 q * (z)}{q_0 * q(z)}\right)^{n/2} H_n\left(\frac{\sqrt{2}x}{w(z)}\right) \exp\left(\frac{-ik(x^2)}{2q_z}\right) \times u_m(y,z) , \qquad (3)$$

where the first three Hermite polynomials are given by

$$H_n(x) = \begin{cases} 1 & (n=0) \\ 2x & (n=1) \\ 4x^2 - 2 & (n=2) \end{cases}$$

For n = 0, 1, 2 and m = 0 at the beam waist(letting  $z_0 = 0$ ):

$$u_{00}(x,y,0) = \left(\frac{2}{\pi}\right)^{1/2} \left(\frac{1}{w_0}\right) \exp\left[-\left(\frac{x^2 + y^2}{w_0^2}\right)\right] , \qquad (4)$$

$$u_{10}(x,y,0) = \left(\frac{2}{\pi}\right)^{1/2} \left(\frac{2x}{w_0^2}\right) \exp\left[-\left(\frac{x^2 + y^2}{w_0^2}\right)\right] , \qquad (5)$$

$$u_{20}(x,y,0) = \left(\frac{2}{\pi}\right)^{1/2} \left[ \left(\frac{2\sqrt{2}x^2}{w_0^3}\right) - \frac{\sqrt{2}}{2} \right] \exp\left[ -\left(\frac{x^2 + y^2}{w_0^2}\right) \right] , \tag{6}$$

Expressing  $u_{10}$  and  $u_{20}$  in terms of  $u_{00}$  at the beam waist,

$$u_{10}(x,y,0) = \frac{2x}{w_0} u_{00}(x,y,0) \tag{7}$$

$$u_{20}(x,y,0) = \left[2\sqrt{2}\left(\frac{x^2}{w_0^2}\right) - \frac{\sqrt{2}}{2}\right]u_{00}(x,y,0) \tag{8}$$

In PauLisa.py, Eq. 1 is represented by the *calculate* function and Eq. 3 as *calculate\_q*. The outputs of both functions agree exactly.

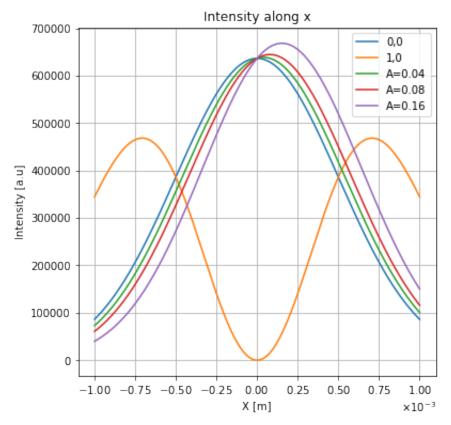


FIG. (1) Intensity at the beam waist and x=0 for HG(0,0) and HG(1,0) modes alongside combined modes. The variable A represents the scale of the HG(1,0) mode, where 1 is the HG(0,0) mode coefficient.

### 2. INTENSITY PROFILES

In the appendix, intensity profiles produced by PauLisa.py for  $u_{00}$  to  $u_{33}$  are shown in Fig. A.1.

## 3. SHIFTED BEAM

For a small shift of the input axis,  $a << w_0$ , in the +x-direction relative to the cavity axis, the shifted  $u_{00}$  mode can be solved in terms of an added  $u_{10}$  mode up to a constant factor:

$$u_{00}(x - a, y, 0) = \left(\frac{2}{\pi}\right)^{-1/2} \left(\frac{1}{w_0}\right) \exp\left(-\frac{(x - a)^2 + y^2}{w_0^2}\right)$$

$$= \left(\frac{2}{\pi}\right)^{-1/2} \left(\frac{1}{w_0}\right) \exp\left(-\frac{y^2}{w_0^2}\right) \exp\left(-\frac{(x - a)^2}{w_0^2}\right)$$

$$= u_{00}(x, y, 0) \times \exp\left(\frac{2ax + a^2}{w_0^2}\right)$$

$$= u_{00}(x, y, 0) \left[1 + \frac{2ax}{w_0^2} + \mathcal{O}\left(\frac{a}{w}\right)^2\right]$$

$$\approx u_{00} + \left(\frac{2ax}{w_0^2}\right) u_{00}$$

$$= u_{00}(x, y, 0) + \frac{a}{w_0} u_{10}(x, y, 0) . \tag{9}$$

Results for HG(0,0) and HG(1,0) addition are shown in Fig. 1.

(1,0) Scale	Pred. Shift[ $\times 10^{-5}m$ ]	Act. Shift[ $\times 10^{-5}m$ ]	%Error
0.04	4.0	3.9849	0.37
0.08	8.0	7.8998	1.25
0.16	16.0	15.2546	4.65

#### 4. TILTED BEAM

A relative tilt between the input and cavity axis can be expressed as an addition of a first-order mode in quadrature phase,

$$u_{tilt}(x, y, 0) = u_{00}(x, y, 0) + \frac{ia\pi w_0}{\lambda} u_{10}(x, y, 0) . \tag{10}$$

This addition varies the phase of the wave in the x-direction

$$u_{tilt}(x, y, 0) = u_{00}(x, y, 0) \exp(i\phi)$$
 (11)

Expressing this in terms of a quadrature phase addition of  $u_{10}$ , with  $\alpha << \frac{\lambda}{w_0 \pi}$ :

$$u_{tilt(0,0)} = u_{00} \exp(i\phi)$$

$$= u_{00} \exp[ikx \sin(\alpha)]$$

$$\approx u_{00} \exp[ikx\alpha]$$

$$= u_{00} \exp\left[i\left(\frac{2\pi x\alpha}{\lambda}\right)\right]$$

$$\approx u_{00} \left[1 + i\left(\frac{2\pi x\alpha}{\lambda}\right)\right]$$

$$= u_{00}(x, y, 0) + i\left(\frac{\pi w_0 \alpha}{\lambda}\right)u_{10}.$$
(12)

#### 5. TILTING A MISALIGNED BEAM

Building on the two previous sections, a tilt to an already shifted beam is

$$u(x, y, 0) = u_{tilt(0,0)} + \left(\frac{a}{w_0}\right) u_{tilt(1,0)} . \tag{13}$$

Following the same approximations for a general tilted beam  $(\alpha << \frac{\lambda}{w_0 \pi})$ , the first term is the same as in Eq. 9, while the second term is:

$$\left(\frac{a}{w_0}\right) u_{tilt(1,0)}(x,y,0) = \left(\frac{a}{w_0}\right) u_{10} \exp(i\phi) 
= \left(\frac{a}{w_0}\right) u_{10} \exp\left[ikx \sin(\alpha)\right] 
\approx \left(\frac{a}{w_0}\right) u_{10} \exp\left[ikx\alpha\right] 
= \left(\frac{a}{w_0}\right) u_{10} \exp\left[i\left(\frac{2\pi x\alpha}{\lambda}\right)\right] 
\approx \left(\frac{a}{w_0}\right) u_{10} \left[1 + i\left(\frac{2\pi x\alpha}{\lambda}\right)\right]$$
(14)

In terms of  $u_{00}$ , the second term is:

$$i\left(\frac{2\pi x\alpha}{\lambda}\right)u_{10} = i\left(\frac{4\pi x^2\alpha a}{w_0^2\lambda}\right)u_{00} \tag{15}$$

Rewriting in terms of  $u_{20}$ :

$$i\left(\frac{4\pi x^2 a\alpha}{w_0^2 \lambda}\right) u_{00} = i\left(\frac{2\alpha a}{\sqrt{2}\lambda}\right) \left[u_{20} + \frac{\sqrt{2}}{2}u_{00}\right]$$
(16)

Therefore, Eq. 13 is

$$u(x,y,0) = \left[1 + i\left(\frac{a\alpha}{\lambda}\right)\right]u_{00} + \left(\frac{a}{w_0}\right)u_{10} + i\left(\frac{2a\alpha}{\sqrt{2}\lambda}\right)u_{20} \tag{17}$$

# 6. MODELING TTL IN LISA'S LONG ARM

Jitter of a tilted beam, as outlined in the previous section, can be used to model TTL coupling in LISA's long arm

[1] C. Bond, D. Brown, A. Freise, and K. A. Strain, Living Reviews in Relativity 19, 3 (2017), ISSN 1433-8351, URL https://doi.org/10.1007/s41114-016-0002-8.

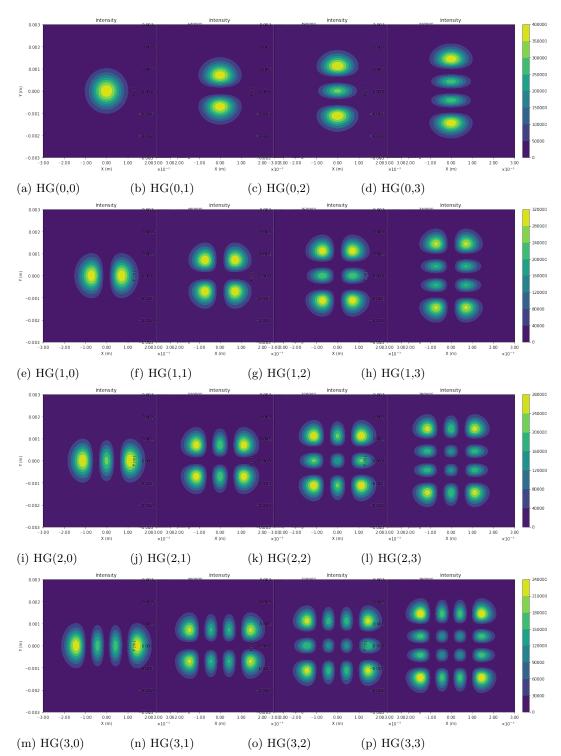


FIG. (A.1) Intensity profiles for HG modes  $u_{00}$  to  $u_{33}$  at the beam waist.