FYS4150 - Computational Physics - Project 5

Candidate number: 68
Email: eimundsm@fys.uio.no
14th November, 2014

Abstract

1 Diffusion of neurotransmitters

I will study diffusion as a transport process for neurotransmitters across synaptic cleft separating the cell membrane of two neurons, for more detail see [1]. The diffusion equation is the partial differential equation

$$\frac{\partial u(\mathbf{x}, t)}{\partial t} = \nabla \cdot \left(D(\mathbf{x}, t) \, \boldsymbol{\nabla} u(\mathbf{x}, t) \right),$$

where u is the concentration of particular neurotransmitters at location $\mathbf{x}$ and time t with the diffusion coefficient D. In this study I consider the diffusion coefficient as constant, which simplify the diffusion equation to the heat equation

$$\frac{\partial u(\mathbf{x}, t)}{\partial t} = D \nabla^2 u(\mathbf{x}, t).$$

I will look at the concentration of neurotransmitter u in two dimensions with x_1 parallel with the direction between the presynaptic to the postsynaptic across the synaptic cleft, and x_2 is parallel with both presynaptic to the postsynaptic. Hence we have the differential equation

$$\frac{\partial u\left(\{x_i\}_{i=1}^2, t\right)}{\partial t} = D \sum_{j=1}^2 \frac{\partial^2 u\left(\{x_i\}_{i=1}^2, t\right)}{\partial x_j{}^2}, \tag{1}$$

where $\{x_i\}_{i=1}^2 = (x_1, x_2) = \mathbf{x}$. The boundary and initial condition that I'm going to study is

$$\exists \{d, w\} \subseteq \mathbb{R}_{0+} \exists \{w_i\}_{i=1}^2 \subseteq \mathbb{R}_{0+}^{w-} \Big(\forall t \in \mathbb{R}_0 : \forall x_2 \in \mathbb{R}_{w_1}^{w_2} : u(0, x_2, t) = u_0$$

$$\wedge \, \forall t \in \mathbb{R} \Big(\forall x_2 \in \mathbb{R}_{0+}^{w-} : u(d, x_2, t) = 0 \wedge \forall x_1 \in \mathbb{R}_0^d : \big(u(x_1, 0, t) = 0 \wedge u(x_1, w, t) = 0 \big) \Big)$$

$$\wedge \forall x_1 \in \mathbb{R}_{0+}^{d-} \forall x_2 \in \mathbb{R}_{0+}^{w-} : u\left(\{x_i\}_{i=1}^2, 0\right) = 0 \wedge \forall x_2 \in \mathbb{R}_0^w \setminus \mathbb{R}_{w_1}^{w_2} : u(0, x_2, 0) = 0 \Big) \tag{2}$$

where d is the distance between the presynaptic and the postsynaptic, and w is the width of the presynaptic and postsynaptic. Note that the notation $\forall x \in \mathbb{R}_{a+}^{b-} \Leftrightarrow a < x < b$, where as $\forall x \in \mathbb{R}_a^b \Leftrightarrow a \le x \le b$. Note also that these boundary conditions implies that the neurotransmitters are transmitted from presynaptic at $x_1 = 0$ and $w_1 \le x_2 \le w_2$ with constant concentration u_0; the neurotransmitters are immediately absorbed at the postsynaptic $x_1 = d$; there are no neurotransmitters at boundary width $x_2 = 0$ and $x_2 = w$ of the synaptic cleft; and we have the initial condition at $t = 0$ where there are no neurotransmitters between the pre- and postsynaptic as well on the side of the
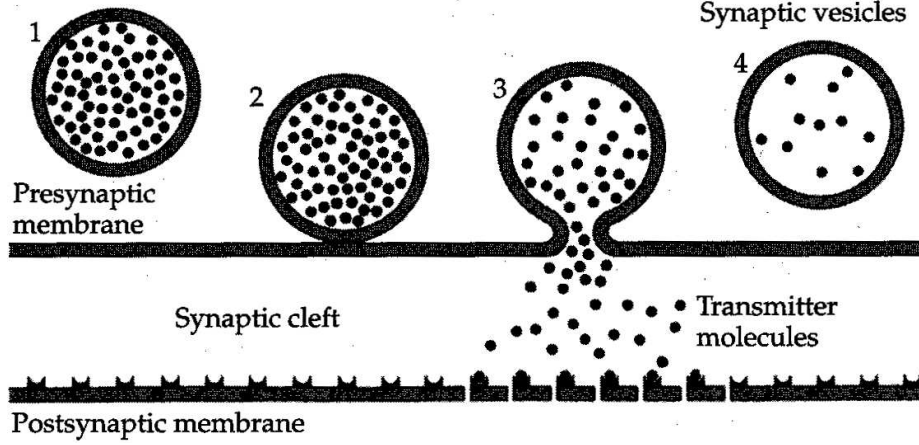
Figure 1.1: *Left: Schematic drawing of the process of vesicle release from the axon terminal and release of transmitter molecules into the synaptic cleft. (From Thompson: "The Brain", Worth Publ., 2000). Right: Molecular structure of the two important neurotransmitters glutamate and GABA.*

synaptic vesicles $x_1 = 0$, $0 \leq x_2 < w_1$ and $w_2 < x_2 \leq w$.

To solve the differential equation (1) with the boundary and initial condition (2) we make an ansatz that the solution is unique, which is the case for a deterministic system. We recognize the heat equation as part of the class of partial differential equation spanned by the Poisson's equation for each time instance. The Uniqueness theorem for the Poisson's equation $\nabla^2 u = f$ [7] says that the Poisson's equation has a unique solution with the Dirichlet boundary condition, where Dirichlet boundary condition is here defined as a boundary that specifies the values the solution must have at the boundary.

Unfortunately the boundary condition in (2) is not a Dirichlet boundary, since the boundary is not specified on the side of the synaptic vesicles $x_1 = 0$, $0 \leq x_2 < w_1$ and $w_2 < x_2 \leq w$. However the Uniqueness theorem does not exclude other boundary conditions to result in a unique solution, so we make another ansatz to separate the concentration $u(\mathbf{x}, t)$ into two functions $u_1(\mathbf{x})$ and $u_2(\mathbf{x}, t)$;

$$\forall x_1 \in \mathbb{R}_0^d \forall x_2 \in \mathbb{R}_0^w \forall t \in \mathbb{R}_0 : u\left(\{x_i\}_{i=1}^2, t\right) = u_1\left(\{x_i\}_{i=1}^2\right) + u_2\left(\{x_i\}_{i=1}^2, t\right), \tag{3}$$

such that u_1 is unique and assume that u_2 satisfies

$$\frac{\partial u_2\left(\{x_i\}_{i=1}^2, t\right)}{\partial t} = D \sum_{j=1}^2 \frac{\partial^2 u_2\left(\{x_i\}_{i=1}^2, t\right)}{\partial x_j^2} \tag{4}$$

with the Dirichlet boundary condition. For u_2 to be determined by the heat equation as well, u_1 must be on the form

$$u_1\left(\{x_i\}_{i=1}^2\right) = a_0 + \sum_{i=1}^2 \sum_{j=i}^2 \prod_{k=i}^j a_{jk} x_k. \tag{5}$$

Now lets try to impose the 0th Dirichlet boundary condition on u_2

$$\forall t \in \mathbb{R} \forall x_1 \in \mathbb{R}_0^d x_2 \in \mathbb{R}_0^w : u_2(0, x_2, t) = u_2(d, x_2, t) = u_2(x_1, 0, t) = u_2(x_1, w, t) = 0. \tag{6}$$

2

When put into (3) and using the boundary condition (2) we have the following equations that u_1 must satisfied

$$u(0, x_2, t) = u_1(0, x_2) + u_2(0, x_2, t) = u_1(0, x_2) = u_0 \quad \text{when } x_2 \in \mathbb{R}_{w_1}^{w_2}$$
$$u(d, x_2, t) = u_1(d, x_2) + u_2(d, x_2, t) = u_1(d, x_2) = 0$$
$$u(x_1, 0, t) = u_1(x_1, 0) + u_2(x_1, 0, t) = u_1(x_1, 0) = 0$$
$$u(x_1, w, t) = u_1(x_1, w) + u_2(x_1, w, t) = u_1(x_1, w) = 0$$

and we can a very close fitting to (5)

$$u_1\left(\{x_i\}_{i=1}^2\right) = u_0\left(1 - \frac{x_1}{d}\right)\left(H(x_2 - w_1) - H(x_2 - w_2) + \left(1 - H(x_2 - w_1)\right)\frac{x_2}{w_1} + H(x_2 - w_2)\frac{x_2 - w}{w_2 - w}\right), \tag{7}$$

where $H(x)$ is the Heaviside step function

$$H(x) = \int_{-\infty}^{x} \delta(s) \, ds = \begin{cases} 1 & x \geq 0 \\ 0 & x < 0 \end{cases}$$

and $\delta(x)$ is Dirac delta function. The deviation that the solution of u_1 in (7) gives to the assumption in (4) is given by

$$\sum_{j=1}^{2} \frac{\partial^2 u_1\left(\{x_i\}_{i=1}^2\right)}{\partial x_j^2} = u_0\left(1 - \frac{x_1}{d}\right)\frac{\partial}{\partial x_2}\left(\delta(x_2 - w_1) - \delta(x_2 - w_2) + \frac{1 - H(x_2 - w_1) - x_2\delta(x_2 - w_1)}{w_1}\right.$$
$$\left. + \frac{H(x_2 - w_2) + (x_2 - w)\delta(x_2 - w_2)}{w_2 - w}\right)$$
$$= u_0\left(1 - \frac{x_1}{d}\right)\frac{\partial}{\partial x_2}\left(\frac{1 - H(x_2 - w_1)}{w_1} + \frac{H(x_2 - w_2)}{w_2 - w}\right)$$
$$= u_0\left(1 - \frac{x_1}{d}\right)\left(\frac{\delta(x_2 - w_2)}{w_2 - w} - \frac{\delta(x_2 - w_1)}{w_1}\right). \tag{8}$$

So we see that assumption in (4) is valid with 0th Dirichlet boundary condition for all points expect when $x_2 = w_1$ and $x_2 = w_2$;

$$\forall t \in \mathbb{R}_0 \exists \{w_i\}_{i=1}^2 \subseteq \mathbb{R}_0^w \forall x_1 \in \mathbb{R}_0^d \forall x_2 \in \mathbb{R}_0^w \setminus \{w_i\}_{i=1}^2 : \frac{\partial u_2\left(\{x_i\}_{i=1}^2, t\right)}{\partial t} = D\sum_{j=1}^{2} \frac{\partial^2 u_2\left(\{x_i\}_{i=1}^2, t\right)}{\partial x_j^2}, \tag{9}$$

which means that the solution of u_2 is unique for the condition in (9), hence if we find a solution to u_2 this is the only solution. So we make another ansatz that $u_2(\mathbf{x}, t)$ is separable into $u_3(\mathbf{x})$ and $u_4(t)$ as follows

$$\forall t \in \mathbb{R}_0 \exists \{w_i\}_{i=1}^2 \subseteq \mathbb{R}_0^w \forall x_1 \in \mathbb{R}_0^d \forall x_2 \in \mathbb{R}_0^w \setminus \{w_i\}_{i=1}^2 : u_2\left(\{x_i\}_{i=1}^2, t\right) = u_3\left(\{x_i\}_{i=1}^2\right)u_4(t), \tag{10}$$

3

which satisfies the 0th Dirichlet boundary condition of u_2 when u_3 also have the 0th Dirichlet boundary condition;

$$u_3(0, x_2) = u_3(d, x_2) = u_3(x_1, 0) = u_3(x_1, w) = 0.$$

Now putting (10) into (9) yields

$$u_3\left(\{x_i\}_{i=1}^2\right) \frac{\partial u_4(t)}{\partial t} = Du_4(t) \sum_{j=1}^2 \frac{\partial^2 u_3\left(\{x_i\}_{i=1}^2\right)}{\partial x_j^2}.$$

and separating the variables

$$\frac{1}{Du_4(t)} \frac{\partial u_4(t)}{\partial t} = \frac{1}{u_3\left(\{x_i\}_{i=1}^2\right)} \sum_{j=1}^2 \frac{\partial^2 u_3\left(\{x_i\}_{i=1}^2\right)}{\partial x_j^2} = -\lambda^2,$$

where λ is a constant, because t and x can vary independently. These two equations have the following solution

$$u_3\left(\{x_i\}_{i=1}^2\right) = \sum_{i=1}^2 \sum_{j=i}^2 \prod_{k=i}^j A_{ijk} \sin\left(\lambda_{ijk} x_k + \varphi_{ijk}\right) \qquad \text{and}$$

$$u_4(t) = Ce^{-D\lambda^2 t}.$$

The boundary condition $u_3(0, x_2) = u_3(x_1, 0) = 0$ gives $A_{111} = A_{222} = \varphi_{121} = \varphi_{122} = 0$, and $u_3(d, x_2) = u_3(x_1, w) = 0$ gives

$$\lambda_{121} = \frac{n_1 \pi}{d} \qquad \text{and} \qquad \lambda_{122} = \frac{n_2 \pi}{w} \qquad \text{for } \{n_i\}_{i=1}^2 \in \mathbb{N}/\{0\},$$

where I let $\mathbb{N} = \mathbb{N}_{-\infty}^\infty$ represent all positive and negative integers, which much satisfies

$$\lambda^2 = \lambda_{121} \lambda_{122}.$$

Hence

$$\forall t \in \mathbb{R}_0 \exists \{w_i\}_{i=1}^2 \subseteq \mathbb{R}_0^w \forall x_1 \in \mathbb{R}_0^d \forall x_2 \in \mathbb{R}_0^w \setminus \{w_i\}_{i=1}^2:$$

$$u_2\left(\{x_i\}_{i=1}^2, t\right) = \sum_{n_1, n_2 = 1} A_{n_1 n_2} \sin\left(n_1 \pi \frac{x_1}{d}\right) \sin\left(n_2 \pi \frac{x_2}{w}\right) \exp\left(-D \frac{n_1 n_2 \pi^2}{dw} t\right), \tag{11}$$

where the negative values of n_1 and n_2 are absorbed into the coefficient $A_{n_1 n_2}$. Applying the initial condition from (2) on the concentration in (4) with (11) we must satisfy the following equation

$$u\left(\{x_i\}_{i=1}^2, 0\right) = u_1\left(\{x_i\}_{i=1}^2\right) + \sum_{n_1, n_2 = 1} A_{n_1 n_2} \sin\left(n_1 \pi \frac{x_1}{d}\right) \sin\left(n_2 \pi \frac{x_2}{w}\right) = 0. \tag{12}$$

4

We need to determine the coefficients $A_{n_1 n_2}$, and the trick is to do something with the equation above such that we isolate the $A_{n_1 n_2}$ coefficients. To achieve this we use the fact that $\sin\left(n\pi\frac{x}{d}\right)$ is orthogonal with $\sin\left(m\pi\frac{x}{d}\right)$ under integration

$$\int \sin\left(m\pi\frac{x}{d}\right)\sin\left(n\pi\frac{x}{d}\right) dx = -\frac{d}{m\pi}\cos\left(m\pi\frac{x}{d}\right)\sin\left(n\pi\frac{x}{d}\right) + \frac{n}{m}\int \cos\left(m\pi\frac{x}{d}\right)\cos\left(n\pi\frac{x}{d}\right) dx$$

$$= -\frac{d}{m\pi}\cos\left(m\pi\frac{x}{d}\right)\sin\left(n\pi\frac{x}{d}\right) + \frac{n}{m}\left(\frac{d}{\pi m}\sin\left(m\pi\frac{x}{d}\right)\cos\left(n\pi\frac{x}{d}\right) + \frac{n}{m}\int \sin\left(m\pi\frac{x}{d}\right)\sin\left(n\pi\frac{x}{d}\right) dx\right),$$

where I have used integration by parts $\int u\dot{v} = uv - \int \dot{u}v$. Solving this equation with regard to the integral we get

$$\int \sin\left(m\pi\frac{x}{d}\right)\sin\left(n\pi\frac{x}{d}\right) dx = \frac{d}{\pi\left(n^2 - m^2\right)}\left(n\sin\left(m\pi\frac{x}{d}\right)\cos\left(n\pi\frac{x}{d}\right) - m\cos\left(m\pi\frac{x}{d}\right)\sin\left(n\pi\frac{x}{d}\right)\right).$$

We want to make the result of this integral zero for $n \neq m$, which is the result if $x = \frac{kd}{2}$ and $x = \frac{\ell d}{2}$ where $k, \ell \in \mathbb{N}$ (for zero, negative and positive integers), because then $\sin\cos$ parts above becomes zero. This means that we need to integrate from $x = \frac{kd}{2}$ and $x = \frac{\ell d}{2}$. However the result above is not defined for $n = m$ because we get $\frac{0}{0}$. So we redo the integration for $n = m$;

$$\int_{\frac{kd}{2}}^{\frac{\ell d}{2}} \sin^2\left(n\pi\frac{x}{d}\right) dx = -\left[\frac{d}{\pi n}\cos\left(n\pi\frac{x}{d}\right)\sin\left(n\pi\frac{x}{d}\right)\right]_{\frac{kd}{2}}^{\frac{\ell d}{2}} + \int_{\frac{kd}{2}}^{\frac{\ell d}{2}} \cos^2\left(n\pi\frac{x}{d}\right) dx = \int_{\frac{kd}{2}}^{\frac{\ell d}{2}} \cos^2\left(n\pi\frac{x}{d}\right) dx$$

$$= \int_{\frac{kd}{2}}^{\frac{\ell d}{2}} \left(1 - \sin^2\left(n\pi\frac{x}{d}\right)\right) dx = [x]_{\frac{kd}{2}}^{\frac{\ell d}{2}} - \int_{\frac{kd}{2}}^{\frac{\ell d}{2}} \sin^2\left(n\pi\frac{x}{d}\right) dx = \frac{d}{2}\left(\ell - k\right) - \int_{\frac{kd}{2}}^{\frac{\ell d}{2}} \sin^2\left(n\pi\frac{x}{d}\right) dx = \frac{d}{4}\left(\ell - k\right),$$

where I solve the equation with regard to in $\int \frac{d}{2}\left(\ell - k\right)\sin^2\left(n\pi\frac{x}{d}\right) dx$ in the last step. I have also used integration by parts $\int u\dot{v} = uv - \int \dot{u}v$ and the Pythagoras trigonometric relation $\sin^2 x + \cos^2 x = 1$. So the solution of the following integral is

$$\int_{\frac{kd}{2}}^{\frac{\ell d}{2}} \sin\left(m\pi\frac{x}{d}\right)\sin\left(n\pi\frac{x}{d}\right) dx = \frac{d}{4}\left(\ell - k\right)\delta_{mn},$$

where δ_{mn} is the Kronecker delta, and hence I have showed the orthogonality of the above integral. Applying this to (12);

$$\int_{\frac{k_1 d}{2}}^{\frac{\ell_1 d}{2}} \int_{\frac{k_2 w}{2}}^{\frac{\ell_2 w}{2}} \sum_{n_1, n_2 = 1} A_{n_1 n_2} \sin\left(n_1\pi\frac{x_1}{d}\right)\sin\left(n_2\pi\frac{x_2}{w}\right)\sin\left(m_1\pi\frac{x_1}{d}\right)\sin\left(m_2\pi\frac{x_2}{w}\right) dx_2 dx_1$$

$$= -\int_{\frac{k_1 d}{2}}^{\frac{\ell_1 d}{2}} \int_{\frac{k_2 w}{2}}^{\frac{\ell_2 w}{2}} u_1\left(\{x_i\}_{i=1}^2\right)\sin\left(m_1\pi\frac{x_1}{d}\right)\sin\left(m_2\pi\frac{x_2}{w}\right) dx_2 dx_1,$$

we can isolate $A_{n_1 n_2}$ at $n_i = m_i$ because of the Kronecker deltas $\delta_{n_i m_i}$;

$$A_{n_1 n_2} = -\frac{16}{dw \prod_{i=1}^2 \left(\ell_i - k_i\right)} \int_{\frac{k_1 d}{2}}^{\frac{\ell_1 d}{2}} \int_{\frac{k_2 w}{2}}^{\frac{\ell_2 w}{2}} u_1\left(\{x_i\}_{i=1}^2\right)\sin\left(n_1\pi\frac{x_1}{d}\right)\sin\left(n_2\pi\frac{x_2}{w}\right) dx_2 dx_1.$$

Since $x_1 \in \mathbb{R}_0^d$ and $x_2 \in \mathbb{R}_0^w \setminus \{w_i\}_{i=1}^2$ leads to $k_i, \ell_i \in \{0, 2\}$ and $k_i \neq \ell_i$ for $A_{n_1 n_2}$ to apply to the whole interval of x_1 and x_2, which is independent of each other.

$$
\begin{aligned}
A_{n_1 n_2} &= -\frac{4}{dw} \int_0^d \sin\left(n_1 \pi \frac{x_1}{d}\right) \left(\int_0^{w_1} u_1\left(\{x_i\}_{i=1}^2\right) \sin\left(n_2 \pi \frac{x_2}{w}\right) dx_2 \right. \\
&\quad + \left. \int_{w_1}^{w_2} u_1\left(\{x_i\}_{i=1}^2\right) \sin\left(n_2 \pi \frac{x_2}{w}\right) dx_2 + \int_{w_2}^{w} u_1\left(\{x_i\}_{i=1}^2\right) \sin\left(n_2 \pi \frac{x_2}{w}\right) dx_2 \right) dx_1 \\
&= -\frac{4u_0}{dw} \int_0^d \left(1 - \frac{x_1}{d}\right) \sin\left(n_1 \pi \frac{x_1}{d}\right) \left(\int_0^{w_1} \frac{x_2}{w_1} \sin\left(n_2 \pi \frac{x_2}{w}\right) dx_2 + \int_{w_1}^{w_2} \sin\left(n_2 \pi \frac{x_2}{w}\right) dx_2 \right. \\
&\quad + \left. \int_{w_2}^{w} \frac{x_2 - w}{w_2 - w} \sin\left(n_2 \pi \frac{x_2}{w}\right) dx_2 \right) dx_1 \, ,
\end{aligned}
$$

where I have used (7) for u_1. To solve the integral above we make use of integration by parts $\int u \dot{v} = uv - \int \dot{u} v$, and we start by

$$
\begin{aligned}
\int_0^{w_1} \frac{x_2}{w_1} \sin\left(n_2 \pi \frac{x_2}{w}\right) dx_2 &= \frac{w}{n_2 \pi w_1} \left(-\left[x_2 \cos\left(n_2 \pi \frac{x_2}{w}\right) \right]_0^{w_1} + \int_0^{w_1} \cos\left(n_2 \pi \frac{x_2}{w}\right) dx_2 \right) \\
&= \frac{w}{n_2 \pi w_1} \left(-w_1 \cos\left(n_2 \pi \frac{w_1}{w}\right) + \frac{w}{n_2 \pi} \left[\sin\left(n_2 \pi \frac{x_2}{w}\right) \right]_0^{w_1} \right) = \frac{w}{n_2 \pi w_1} \left(\frac{w}{n_2 \pi} \sin\left(n_2 \pi \frac{w_1}{w}\right) - w_1 \cos\left(n_2 \pi \frac{w_1}{w}\right) \right) ,
\end{aligned}
$$

$$
\int_{w_1}^{w_2} \sin\left(n_2 \pi \frac{x_2}{w}\right) dx_2 = -\frac{w}{n_2 \pi} \left[\cos\left(n_2 \pi \frac{x_2}{w}\right) \right]_{w_1}^{w_2} = \frac{w}{n_2 \pi} \left(\cos\left(n_2 \pi \frac{w_1}{w}\right) - \cos\left(n_2 \pi \frac{w_2}{w}\right) \right) ,
$$

$$
\begin{aligned}
\int_{w_2}^{w} \frac{x_2 - w}{w_2 - w} \sin\left(n_2 \pi \frac{x_2}{w}\right) dx_2 &= \frac{w}{n_2 \pi (w_2 - w)} \left(-\left[(x_2 - w) \cos\left(n_2 \pi \frac{x_2}{w}\right) \right]_{w_2}^{w} + \int_{w_2}^{w} \cos\left(n_2 \pi \frac{x_2}{w}\right) dx_2 \right) \\
&= \frac{w}{n_2 \pi (w_2 - w)} \left((w_2 - w) \cos\left(n_2 \pi \frac{w_2}{w}\right) + \frac{w}{n_2 \pi} \left[\sin\left(n_2 \pi \frac{x_2}{w}\right) \right]_{w_2}^{w} \right) \\
&= \frac{w}{n_2 \pi (w_2 - w)} \left((w_2 - w) \cos\left(n_2 \pi \frac{w_2}{w}\right) - \frac{w}{n_2 \pi} \sin\left(n_2 \pi \frac{w_2}{w}\right) \right)
\end{aligned}
$$

and

$$
\begin{aligned}
\int_0^d \left(1 - \frac{x_1}{d}\right) \sin\left(n_1 \pi \frac{x_1}{d}\right) dx_1 &= \frac{d}{n_1 \pi} \left(-\left[\left(1 - \frac{x_1}{d}\right) \cos\left(n_1 \pi \frac{x_1}{d}\right) \right]_0^d - \frac{1}{d} \int_0^d \cos\left(n_1 \pi \frac{x_1}{d}\right) dx_1 \right) \\
&= \frac{d}{n_1 \pi} \left(1 - \frac{1}{n_1 \pi} \left[\sin\left(n_1 \pi \frac{x_1}{d}\right) \right]_0^d \right) = \frac{d}{n_1 \pi} \, ,
\end{aligned}
$$

which gives

$$
A_{n_1 n_2} = -\frac{4u_0 w}{n_1 n_2^2 \pi^3} \left(\frac{\sin\left(n_2 \pi \frac{w_1}{w}\right)}{w_1} - \frac{\sin\left(n_2 \pi \frac{w_2}{w}\right)}{w_2 - w} \right) \qquad \text{for } x_1 \in \mathbb{R}_0^d \text{ and } x_2 \in \mathbb{R}_0^w \setminus \{w_i\}_{i=1}^2. \qquad (13)
$$

So what about the points $x_2 \in \{w_i\}_{i=1}^2$? We make another ansatz that (7) and (11) with (13) in the ansatz (3) is a solution to heat equation (1) with the boundary condition (2) for $x_2 \in \{w_i\}_{i=1}^2$ as well. Then we have the concentration

$$u(x_1, w_i, t) \tag{14}$$

$$= \lim_{x_2 \to w_i} u_0 \left(1 - \frac{x_1}{d} - \sum_{n_1,n_2=1} \frac{4w}{n_1 n_2^2 \pi^3} \left(\frac{\sin\left(n_2 \pi \frac{w_1}{w}\right)}{w_1} - \frac{\sin\left(n_2 \pi \frac{w_2}{w}\right)}{w_2 - w} \right) \sin\left(n_1 \pi \frac{x_1}{d}\right) \sin\left(n_2 \pi \frac{x_2}{w}\right) \exp\left(-D \frac{n_1 n_2 \pi^2}{dw} t\right) \right),$$

which we see satisfies the required boundary condition $u(0, w_i, t) = u_0$ and $u(d, w_i, t) = 0$. Then $u_2(0, w_i, t) = u_2(d, w_i, t) = 0$ which means that we have satisfied the 0th Dirichlet boundary condition of u_2 for all $x_1 \in \mathbb{R}_0^d$ and $x_2 \in \mathbb{R}_0^w$ and we uniqueness of the solution of u_2, and hence a uniqueness of to u as well if indeed this u is a solution to (1). What we have exploited in the expression above that u_1 is continues and therefore

$$\forall i \in \mathbb{N}_1^2 \forall x_2 \in \mathbb{R}_{w_i-}^{w_i+} : u_1\left(\{x_i\}_{i=1}^2\right) = u_0 \left(1 - \frac{x_1}{d}\right) + \epsilon,$$

and we can therefore make ϵ as small as we want by making the w_i- and w_i+ go closer and closer together. Therefore we can use (14) as approximation to the actual solution in (1) with (2) for $x_2 \in \{w_i\}_{i=1}^2$, where we can approximate as close to the actual solution as we want. If we put this approximation into the heat equation (1), then we see that

$$\forall i \in \mathbb{N}_1^2 \forall x_2 \in \mathbb{R}_{w_i-}^{w_i+} : \sum_{j=1}^2 \frac{\partial^2 u_1\left(\{x_i\}_{i=1}^2\right)}{\partial x_j^2} = 0,$$

and using this approximation into (8) gives us

$$\sum_{j=1}^2 \frac{\partial^2 u_1\left(\{x_i\}_{i=1}^2\right)}{\partial x_j^2} = 0$$

for our entire domain of x_1 and x_2. I have now verified all the ansatzes that I made, and therefore the unique analytical solution to the heat equation (1) with the boundary condition (2) is given by

$$\forall t \in \mathbb{R}_0 \forall x_1 \in \mathbb{R}_0^d \forall x_2 \in \mathbb{R}_0^w : u\left(\{x_i\}_{i=1}^2, t\right) \tag{15}$$

$$= u_0 \left(1 - \frac{x_1}{d}\right) \left(H(x_2 - w_1) - H(x_2 - w_2) + \left(1 - H(x_2 - w_1)\right) \frac{x_2}{w_1} + H(x_2 - w_2) \frac{x_2 - w}{w_2 - w} \right)$$

$$- \sum_{n_1,n_2=1} \frac{4wu_0}{n_1 n_2^2 \pi^3} \left(\frac{\sin\left(n_2 \pi \frac{w_1}{w}\right)}{w_1} - \frac{\sin\left(n_2 \pi \frac{w_2}{w}\right)}{w_2 - w} \right) \sin\left(n_1 \pi \frac{x_1}{d}\right) \sin\left(n_2 \pi \frac{x_2}{w}\right) \exp\left(-D \frac{n_1 n_2 \pi^2}{dw} t\right).$$

2 Attachments

The source files developed are

3 Resources

1. QT Creator 5.3.1 with C11

2. Eclipse Standard/SDK - Version: Luna Release (4.4.0) with PyDev for Python

3. Ubuntu 14.04.1 LTS

4. ThinkPad W540 P/N: 20BG0042MN with 32 GB RAM

References

[1] Morten Hjorth-Jensen, *FYS4150 - Project 5* - *Diffusion in two dimensions*, University of Oslo, 2014

[2] Morten Hjorth-Jensen, *Computational Physics - Lecture Notes Fall 2014*, University of Oslo, 2014

[3] http://en.wikipedia.org/wiki/Diffusion_equation

[4] http://en.wikipedia.org/wiki/Heat_equation

[5] http://en.wikipedia.org/wiki/Dirichlet_boundary_condition

[6] http://en.wikipedia.org/wiki/Poisson%27s_equation

[7] http://en.wikipedia.org/wiki/Uniqueness_theorem_for_Poisson%27s_equation

[8] http://en.wikipedia.org/wiki/Ansatz

[9] http://en.wikipedia.org/wiki/Heaviside_step_function

[10] http://en.wikipedia.org/wiki/Dirac_delta_function