

# Foundations of Modern Optics

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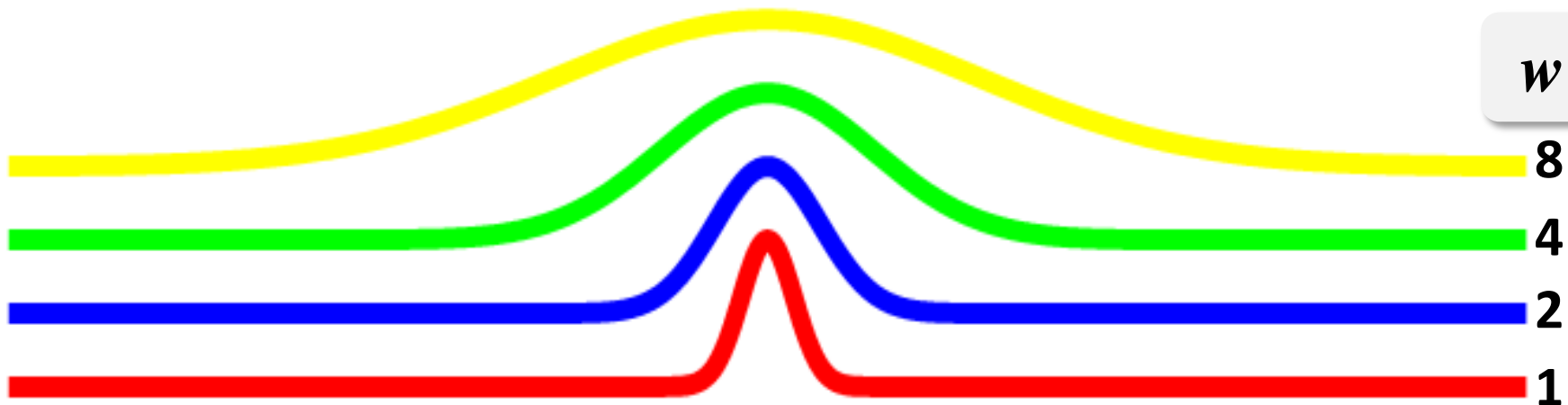
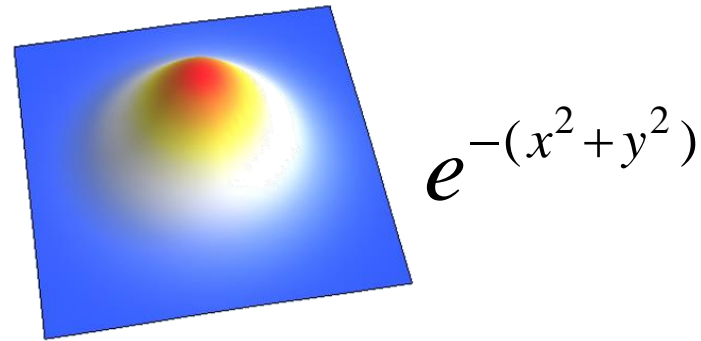
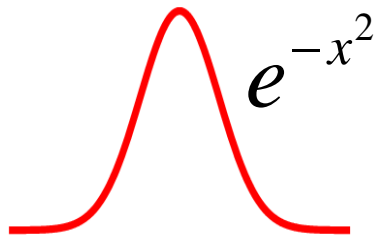
**5**

# **Propagation of Gaussian Beams**

# Description of a Gaussian Beam

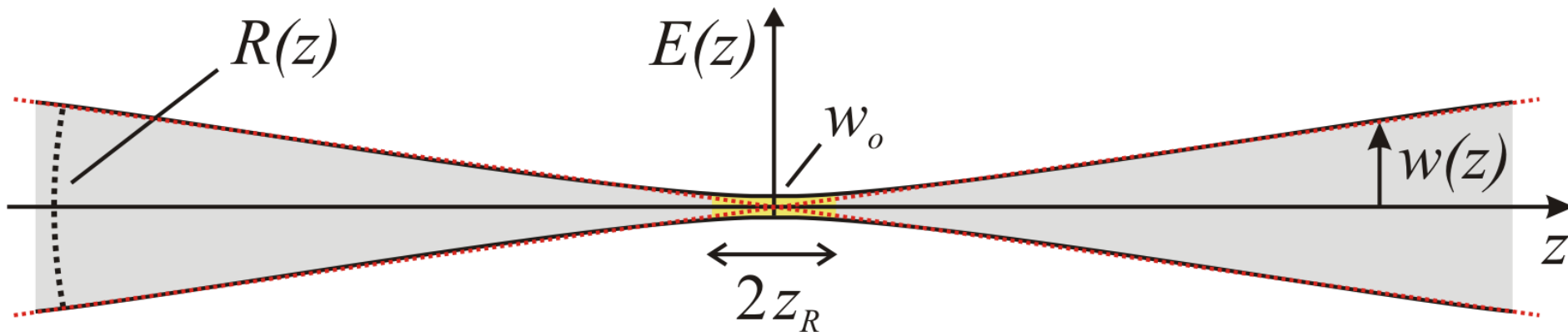
The amplitude is described by a Gaussian distribution function

$$E(x, y; z) = E_o(z) e^{-\frac{x^2 + y^2}{w^2(z)}}$$



# Propagation of a Gaussian

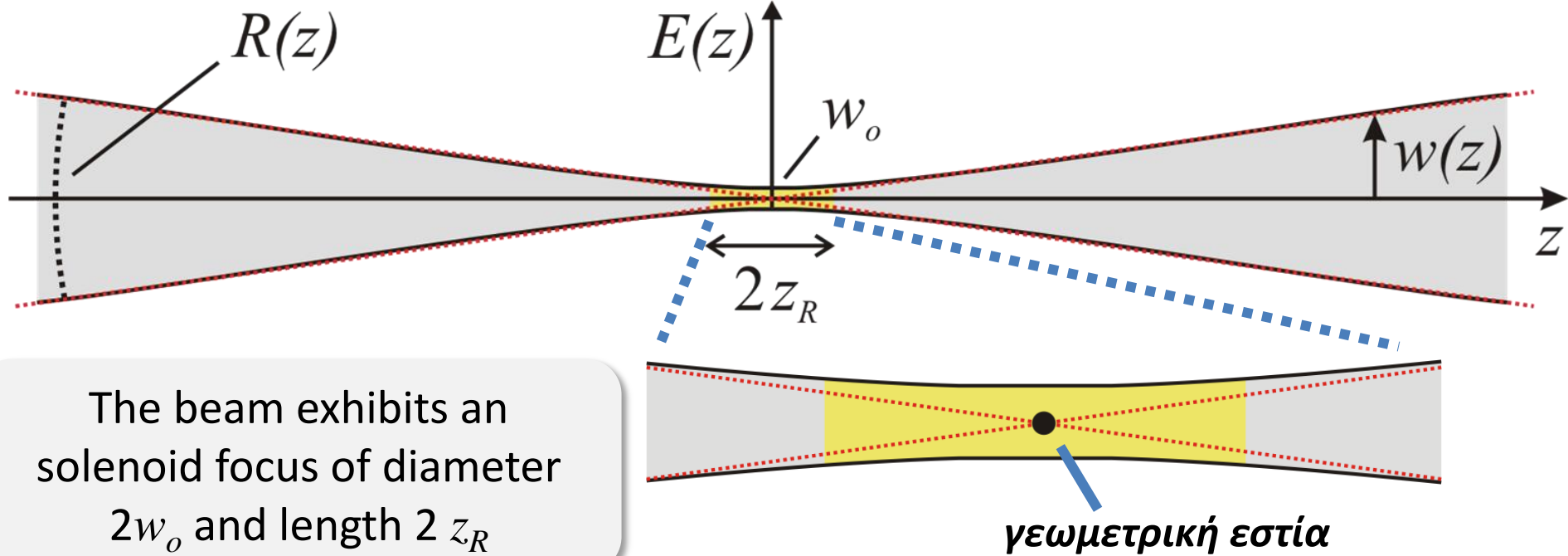
We can **analytically describe** the distribution of a Gaussian beam along its propagation



$$E(\mathbf{r}) = E_o \frac{w_o}{w(z)} \cdot e^{-\frac{x^2 + y^2}{w^2(z)}} \cdot e^{-i \frac{\pi}{\lambda} \frac{x^2 + y^2}{R(z)}} \cdot e^{-i[kz + \psi(z)]}$$

phase terms

*always Gaussian on a transverse plane*



$w(z)$

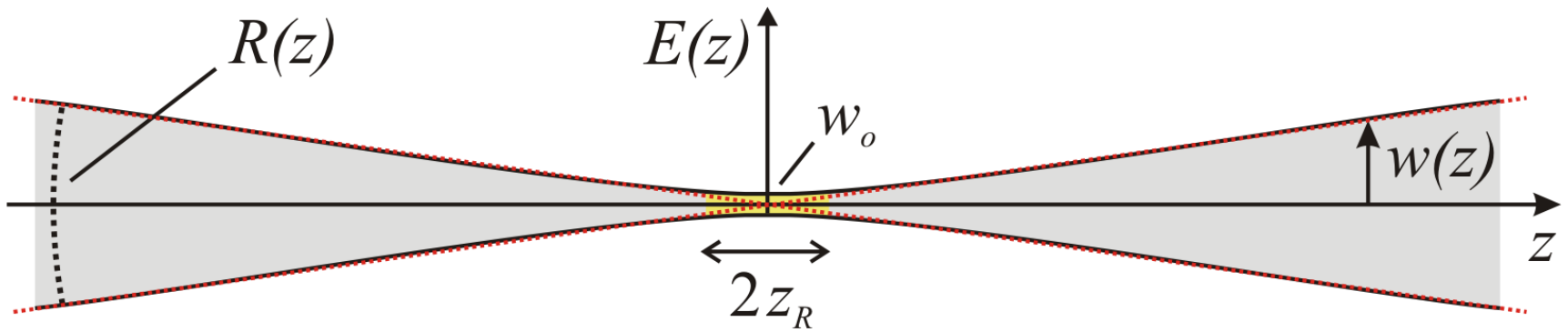
semi-diameter of the beam (*beam waist*)

$R(z)$

*radius of curvature*

$\psi(z)$

Longitudinal phase distribution (*Gouy phase*)



**Values of the Gaussian beam parameters along propagation**

$$w(z) = w_o \sqrt{1 + \frac{z^2}{z_R^2}}$$

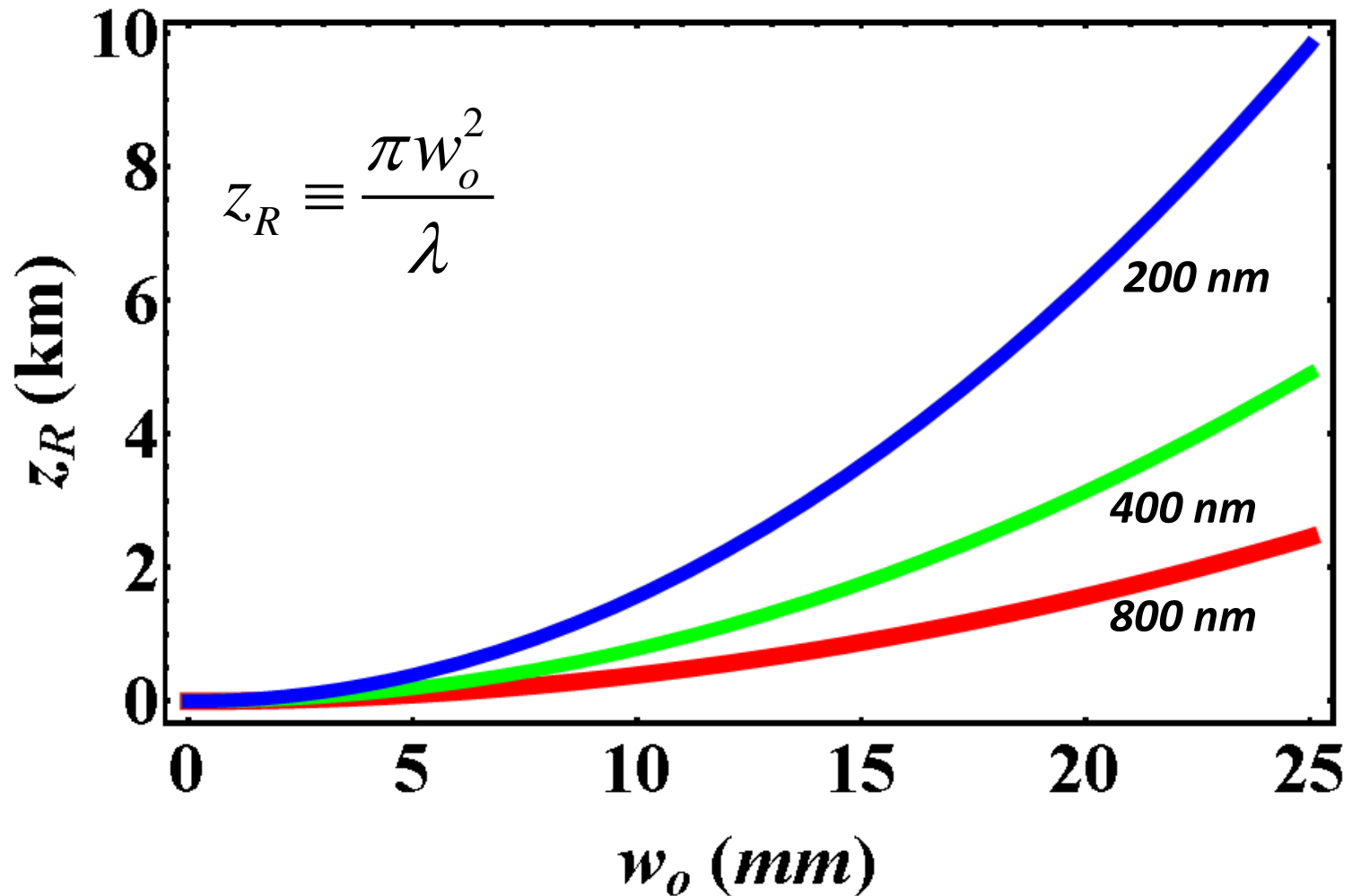
$$R(z) = z + \frac{z_R^2}{z}$$

$$\psi(z) = \arctan\left(\frac{z}{z_R}\right)$$

**Rayleigh length**

$$z_R \equiv \frac{\pi w_o^2}{\lambda}$$

Distance  $z_R$   
(often called Rayleigh length)  
equals the 1/2 of the  
solenoid focus length

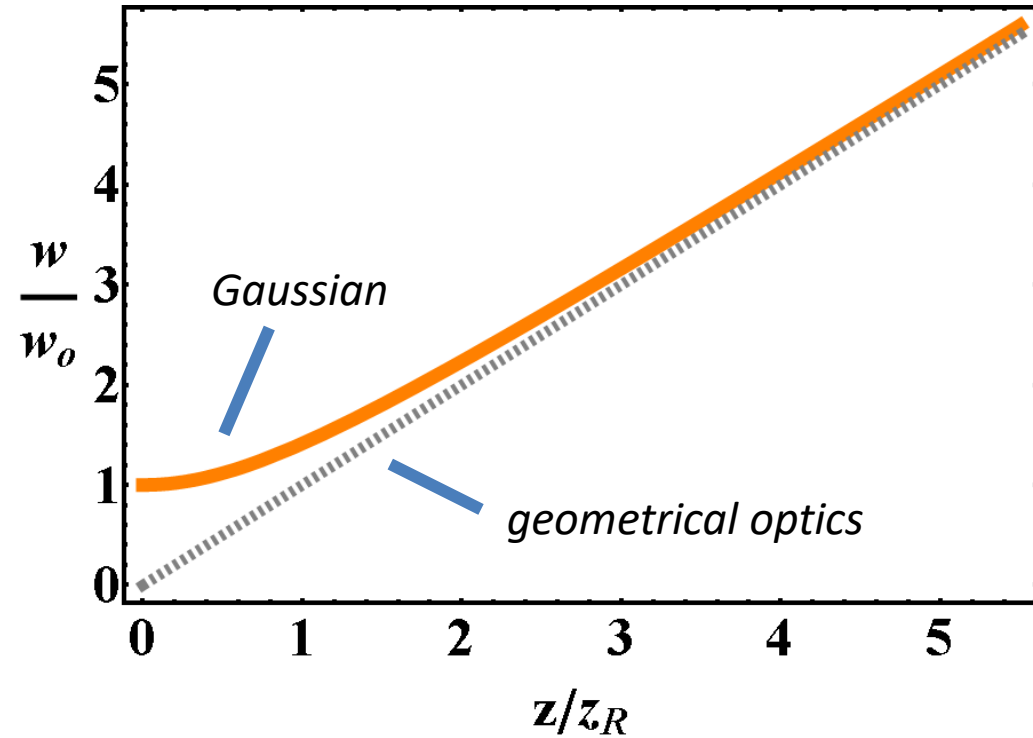


- Beams that are tightly focused (small  $w_o$ ) expand more rapidly!
- By reducing the wavelength we “slow down” the beam spatial expansion

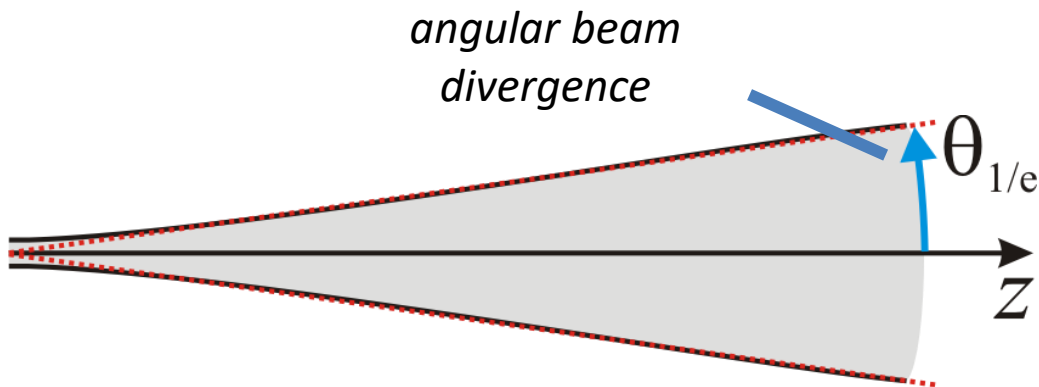


$$w(z) = w_o \sqrt{1 + \frac{z^2}{z_R^2}}$$

$$z \gg z_R \Rightarrow w(z) \cong \frac{w_o}{z_R} z$$



at large enough distances we can simply use geometrical optics

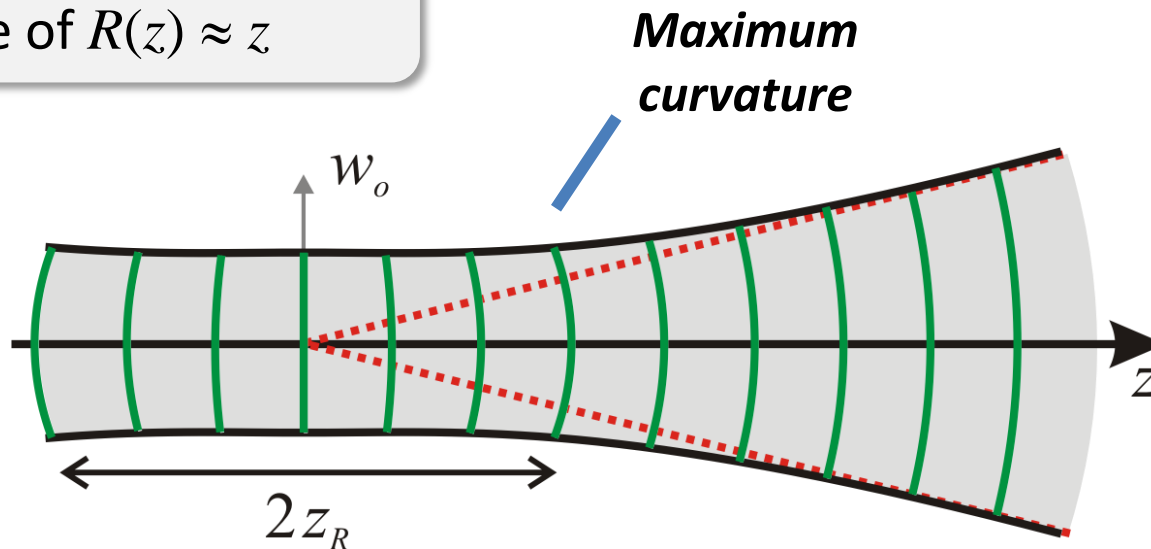
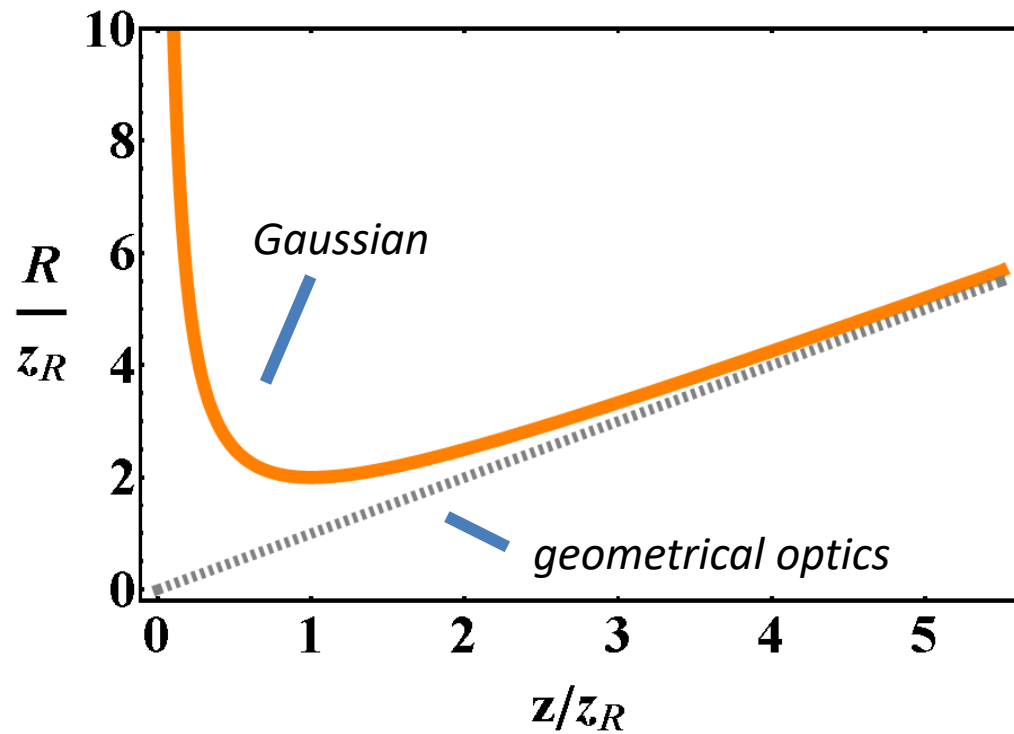


$$\theta_{1/e} \cong \frac{w_o}{z_R} = \frac{\lambda}{\pi w_o}$$

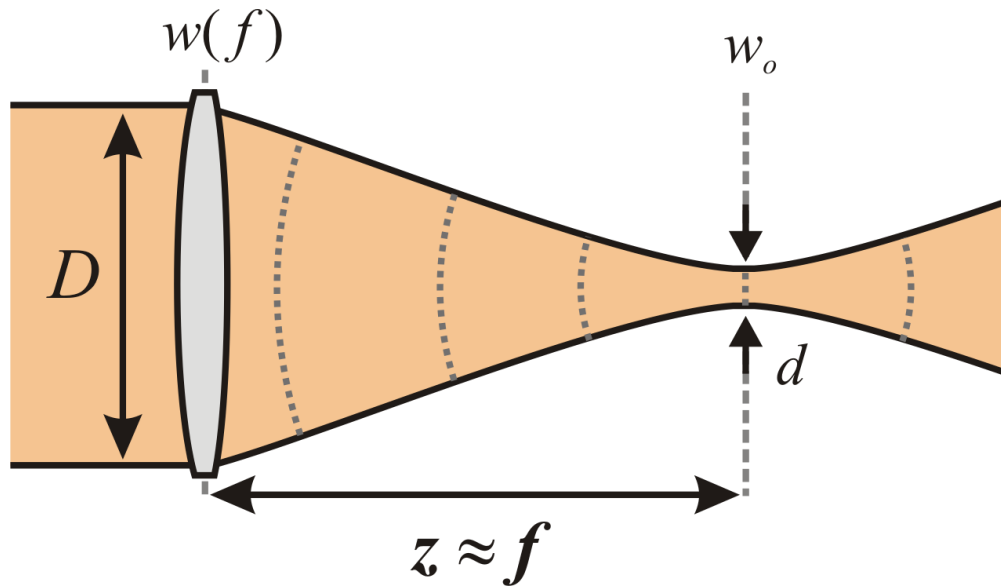
$$R(z) = z + \frac{z_R^2}{z}$$

When  $z \approx 0$ , the wavefront is flat!

Away from the focus ( $z \gg z_R$ ) the wavefront is spherical with radius of curvature of  $R(z) \approx z$



# Typical focusing of a Gaussian beam\*

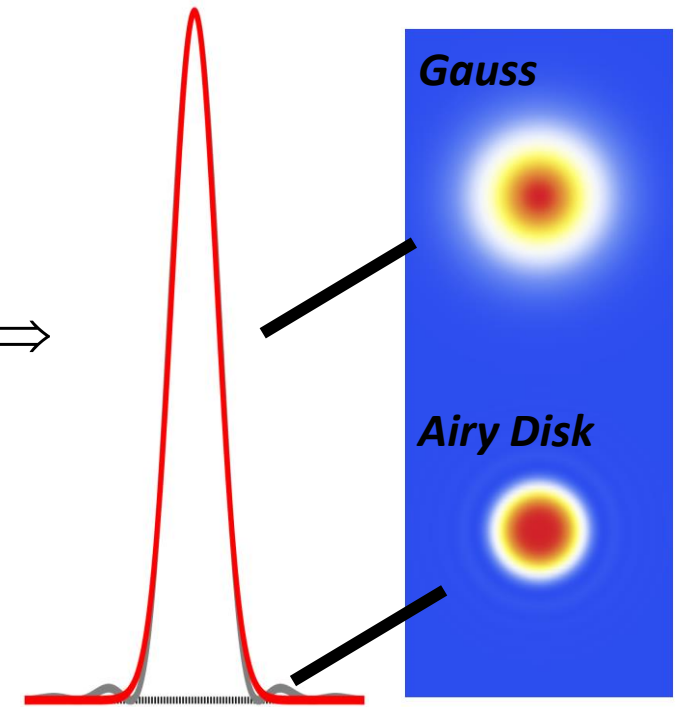


$$z_R \gg f$$

$$f \gg z'_R$$

$$D = 2w(f) = 2w_0 \sqrt{1 + \frac{f^2}{z_R'^2}} \cong \frac{2w_0}{z'_R} f = \frac{2\lambda}{\pi w_0} f \Rightarrow$$

$$d = 2w_0 \cong \frac{4\lambda}{\pi D} f$$



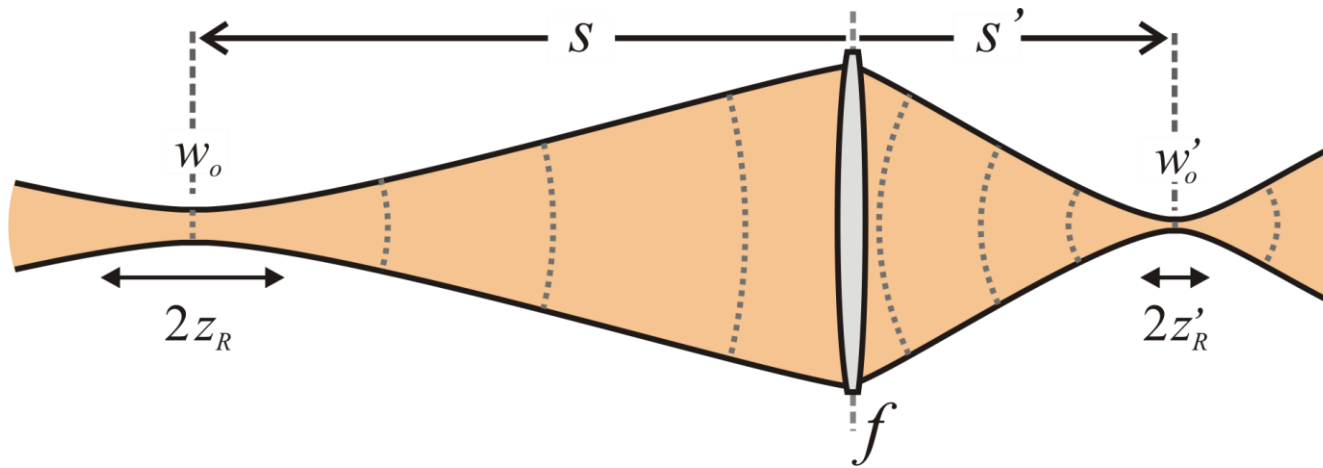
Σύγκριση εστίας Γκάους με κατανομή Airy

\* Idealized calculation. We assume that the input beam waist is on the lens entrance

# Imaging Gaussian beams



# Imaging relations



We assume that the input Gaussian beam focus is the **object**, while the output Gaussian beam is the **image**

*Geometrical optics*

$$\frac{1}{s} + \frac{1}{s'} = \frac{1}{f}$$

$$\frac{1}{s + \frac{z_R^2}{s - f}} + \frac{1}{s'} = \frac{1}{f}$$

*input beam parameters*



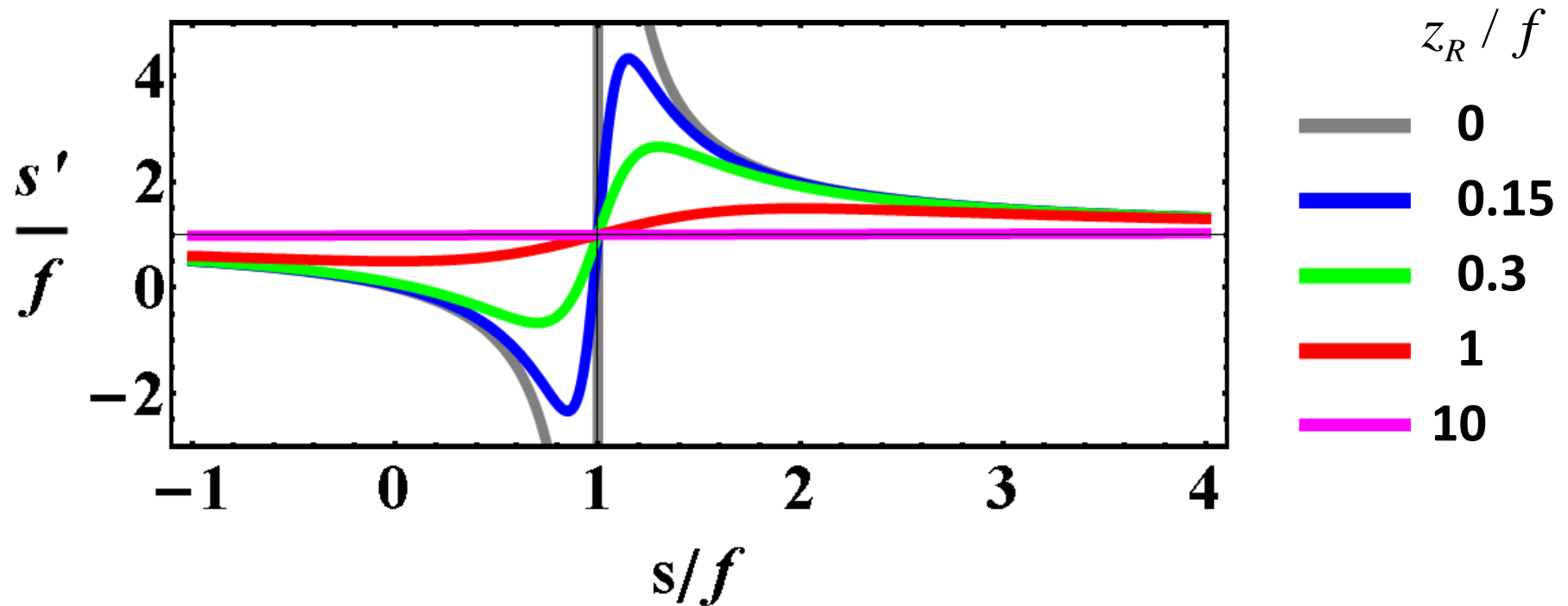
*Gaussian beams*

$$\frac{1}{s} + \frac{1}{s' + \frac{z_R'^2}{s' - f}} = \frac{1}{f}$$

*output beam parameters*

In contrast to the classical geometrical optics the imaging formulas are **not symmetric** to the input and output beam parameters

# Position of the image



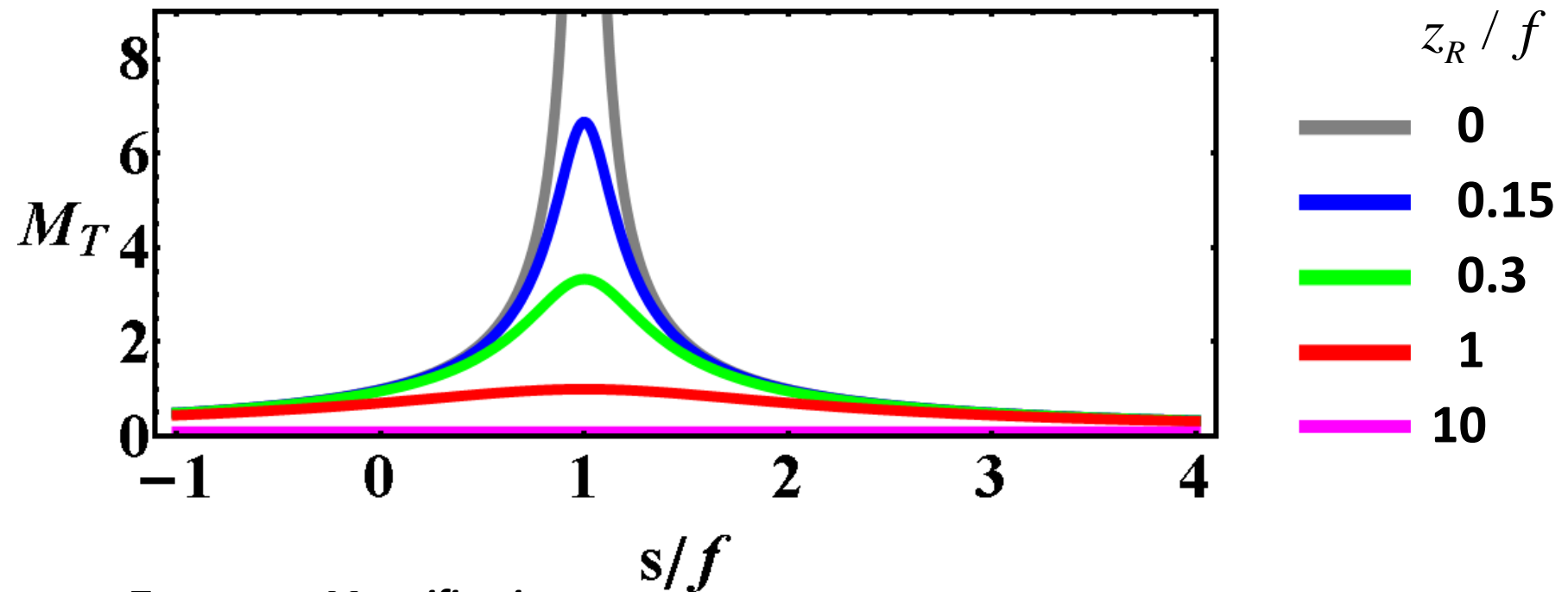
There is a **minimum** and a **maximum** distance for the image position in Gaussian beams

$$s = f + z_R \Rightarrow s'_{\max} = f \left( 1 + \frac{f}{2z_R} \right)$$

The image is located at the maximum distance from the lens when  $s = f + z_R$  and not at  $s = f$

$$s = f \Rightarrow s' = f$$

# Magnification



*Transverse Magnification*

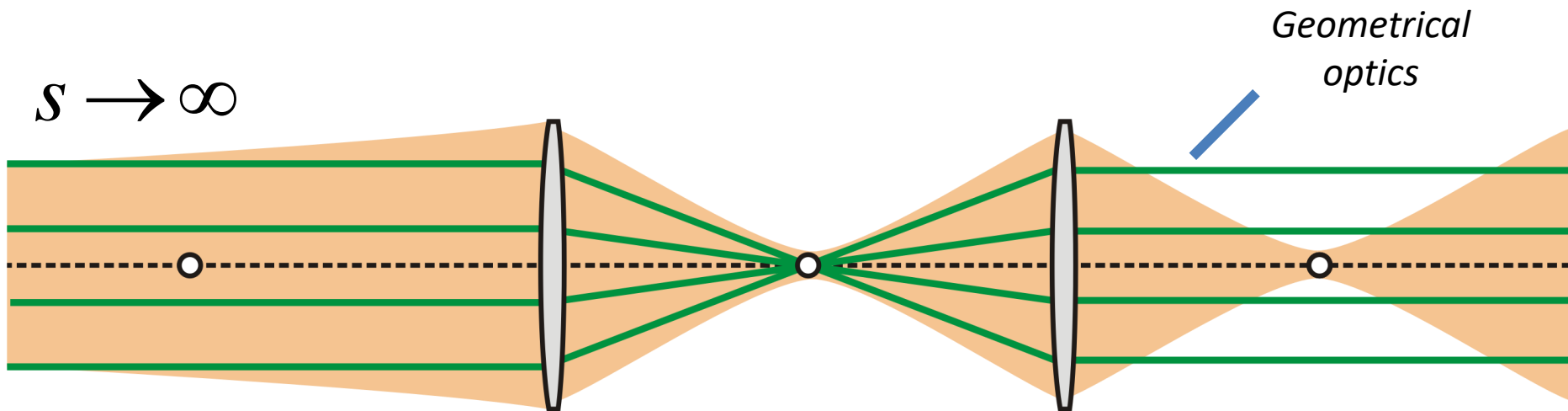
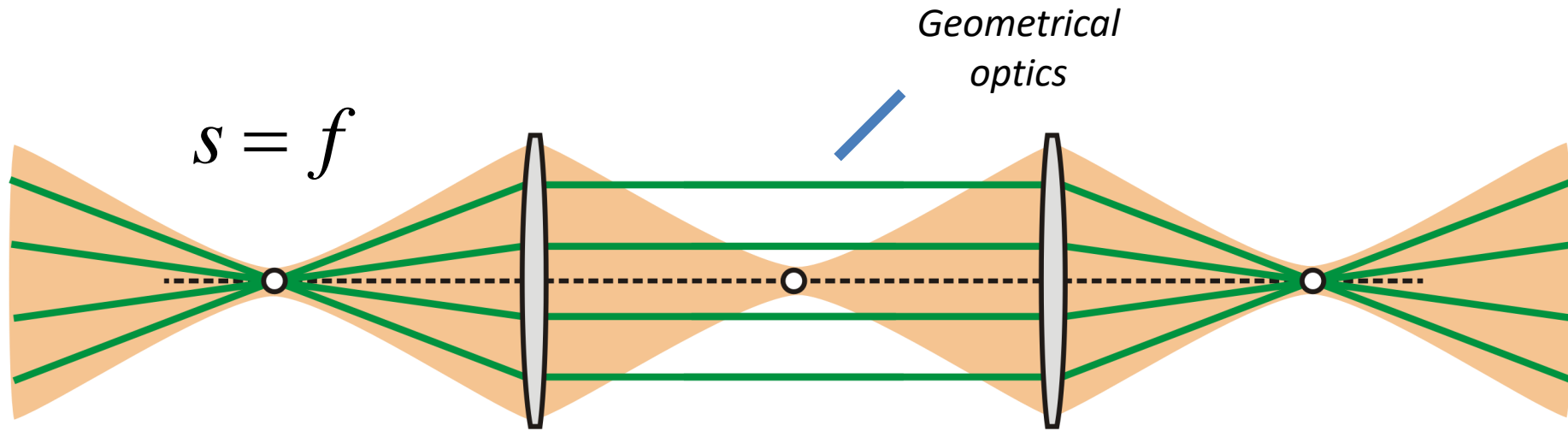
$$M_T \equiv \frac{w'_o}{w_o} = \frac{1}{\sqrt{(1 - s/f)^2 + z_R^2 / f^2}}$$

$$s = f \Rightarrow M_T^{\max} = \frac{f}{z_R}$$

*Longitudinal magnification*

$$M_L \equiv \frac{z'_R}{z_R} = M_T^2$$

# Examples of peculiar behavior





# **Generic analysis**

# Normalization

it is advantageous to normalize all spatial values on the lens focal length

*input Gaussian beam*

*output Gaussian beam*

*distance from  
the lens*

$$\xi = \frac{s}{f}$$

$$\xi' = \frac{s'}{f}$$

*Rayleigh  
Length*

$$\zeta_R = \frac{z_R}{f} = \frac{\pi w_o^2}{f \cdot \lambda}$$

$$\zeta'_R = \frac{z'_R}{f} = \frac{\pi w_o'^2}{f \cdot \lambda}$$

*normalized distance  
of a plane lying at  
distance z*

$$z_n = \frac{z}{f}$$

# Normalized imaging equations

*position of  
the output  
Gaussian  
Beam*

$$\xi' = 1 + \frac{\xi - 1}{\zeta_R^2 + (\xi - 1)^2}$$

$$\zeta_R' = \frac{\zeta_R}{\zeta_R^2 + (\xi - 1)^2}$$

*Rayleigh Length  
of the output  
Gaussian Beam*

*waist of the  
output Gaussian  
Beam*

$$w_o' = \frac{w_o}{\sqrt{(\xi - 1)^2 + \zeta_R^2}}$$

**Transverse  
Magnification**

$$M_T \equiv \frac{w'_o}{w_o} = \frac{1}{\sqrt{(\xi - 1)^2 + \zeta_R^2}}$$

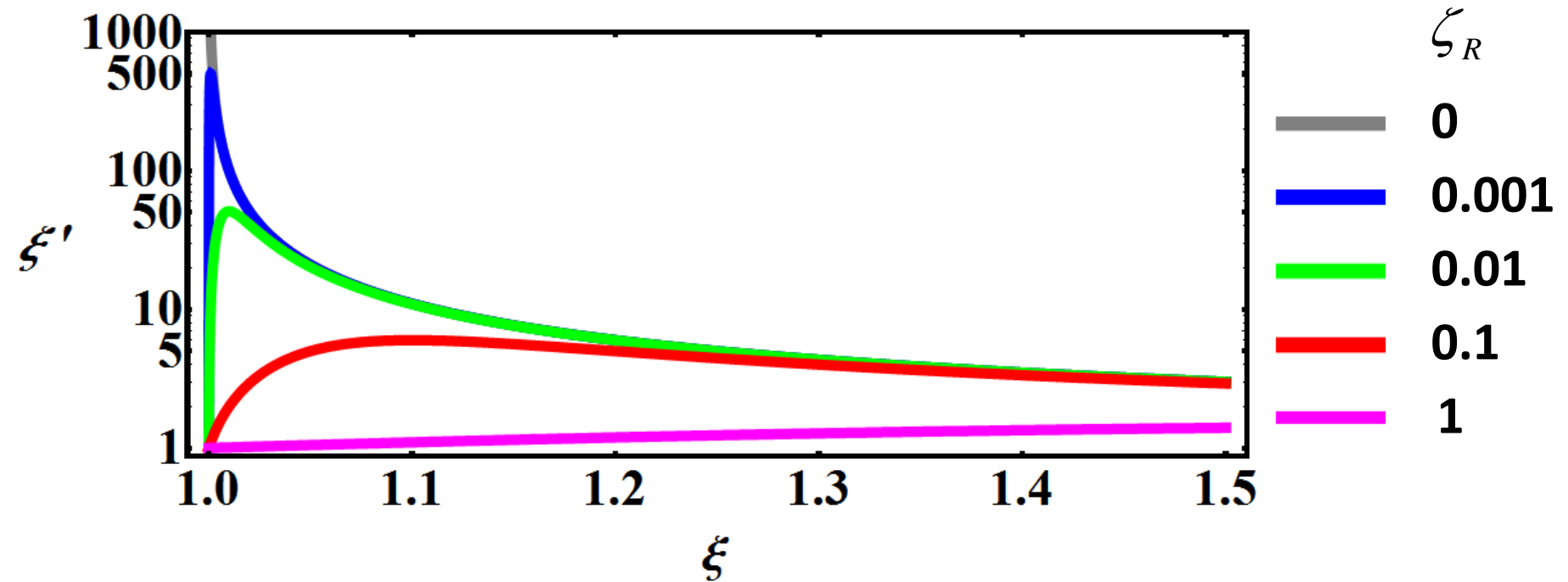
**Longitudinal  
Magnification**

$$M_L \equiv \frac{z'_R}{z_R} = \frac{\zeta'_R}{\zeta_R} = M_T^2 = \frac{1}{(\xi - 1)^2 + \zeta_R^2}$$

output beam semi-diameter  $w_{out}$  size at a distance  $z_n$  from the lens:

$$\frac{w_{out}^2(\xi, z_n)}{w_o^2} = 1 + \frac{\xi^2}{\zeta_R^2} + \frac{z_n^2}{\zeta_R^2} [\zeta_R^2 + (\xi - 1)^2] - 2 \frac{z_n}{\zeta_R^2} [\zeta_R^2 + (\xi - 1)\xi]$$

# Position of the image



There is a **minimum** and a **maximum** distance for the image position in Gaussian beams

$$\xi = 1 \Rightarrow \xi' = 1$$

$$\xi = 1 + \zeta_R \Rightarrow \begin{cases} \xi'_{max} = 1 + \frac{1}{2\zeta_R} \\ \xi'_{min} = 1 - \frac{1}{2\zeta_R} \end{cases}$$

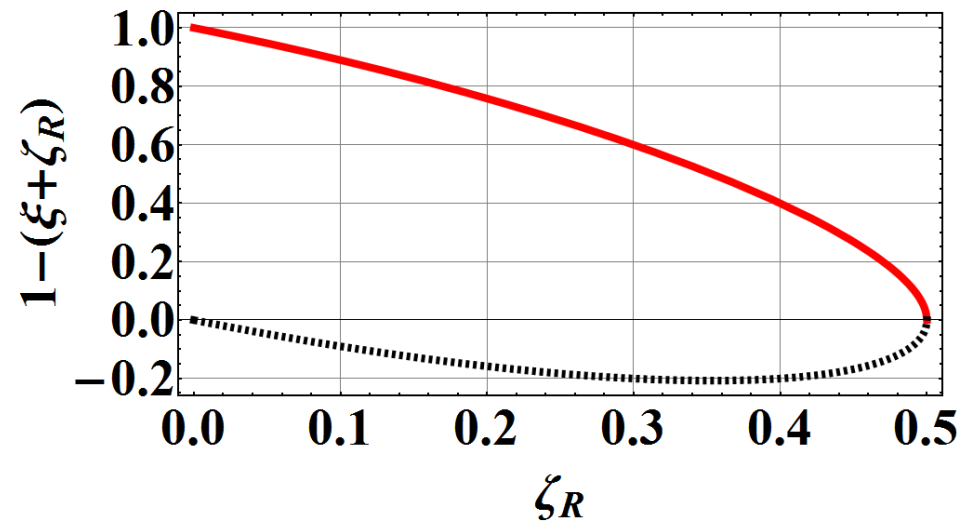
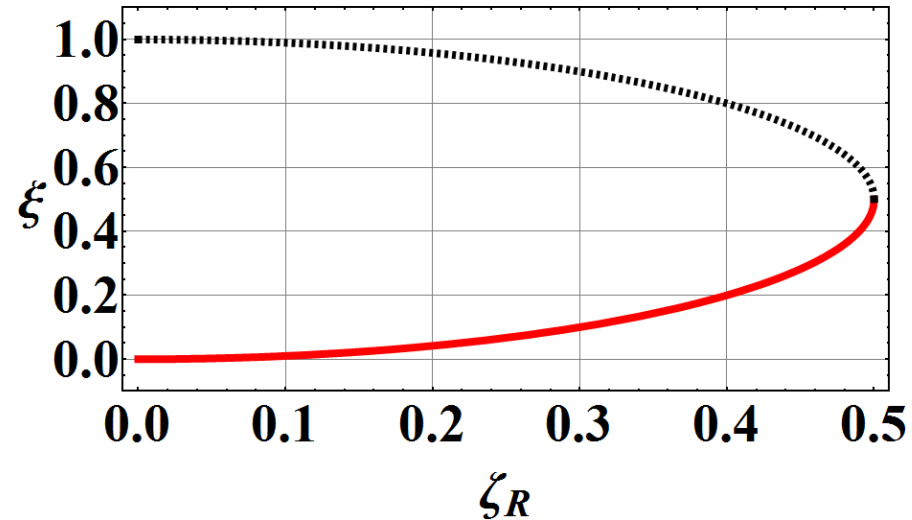
# Image on the lens $\xi' = 0$

$$\xi' = 0 \Rightarrow \xi = \frac{1 \pm \sqrt{1 - 4\zeta_R^2}}{2}$$

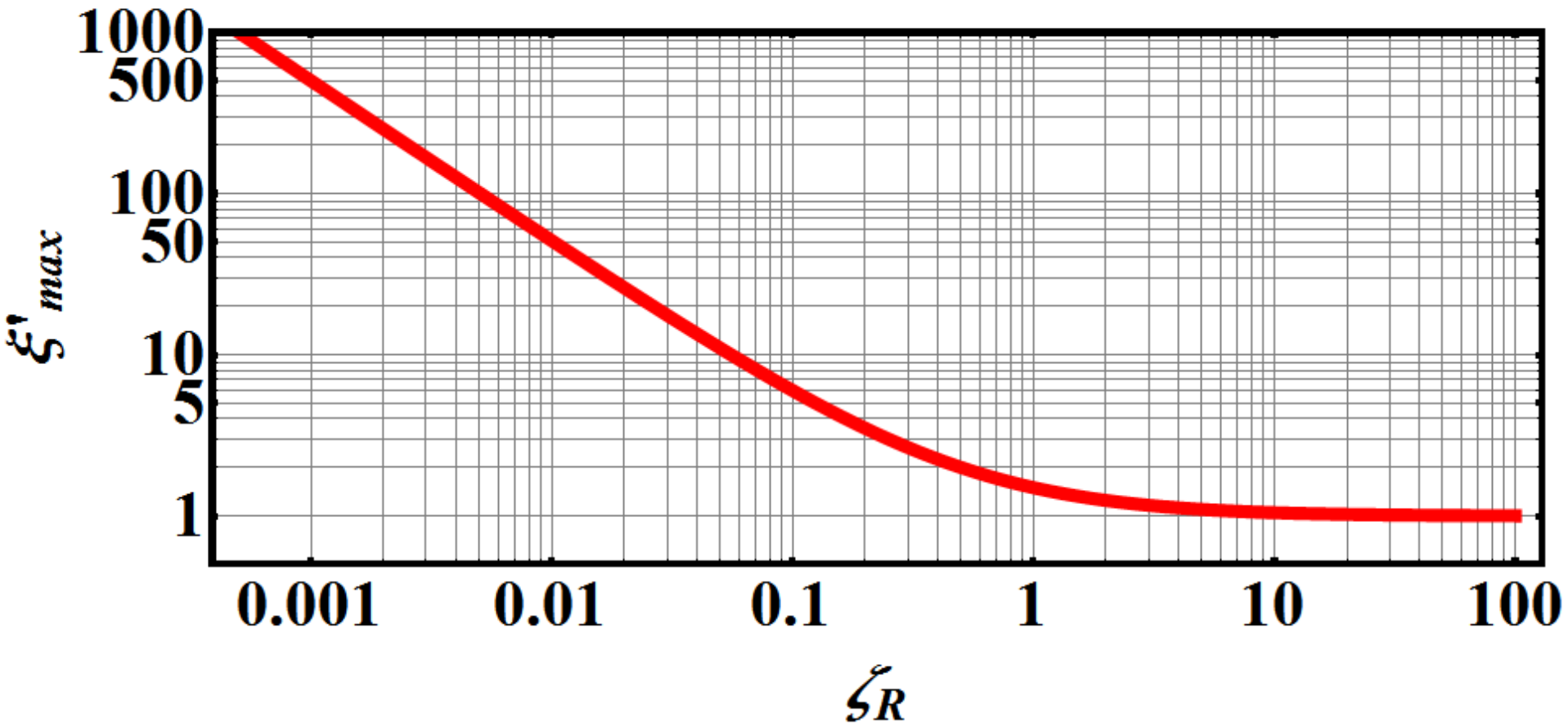
There are two solutions in this problem that merge at  $\zeta_R = 0.5$

“Geometrical regime”

$$\zeta_R \rightarrow 0 \Rightarrow \begin{cases} \xi \rightarrow 0 \\ \xi \rightarrow 1 \end{cases}$$

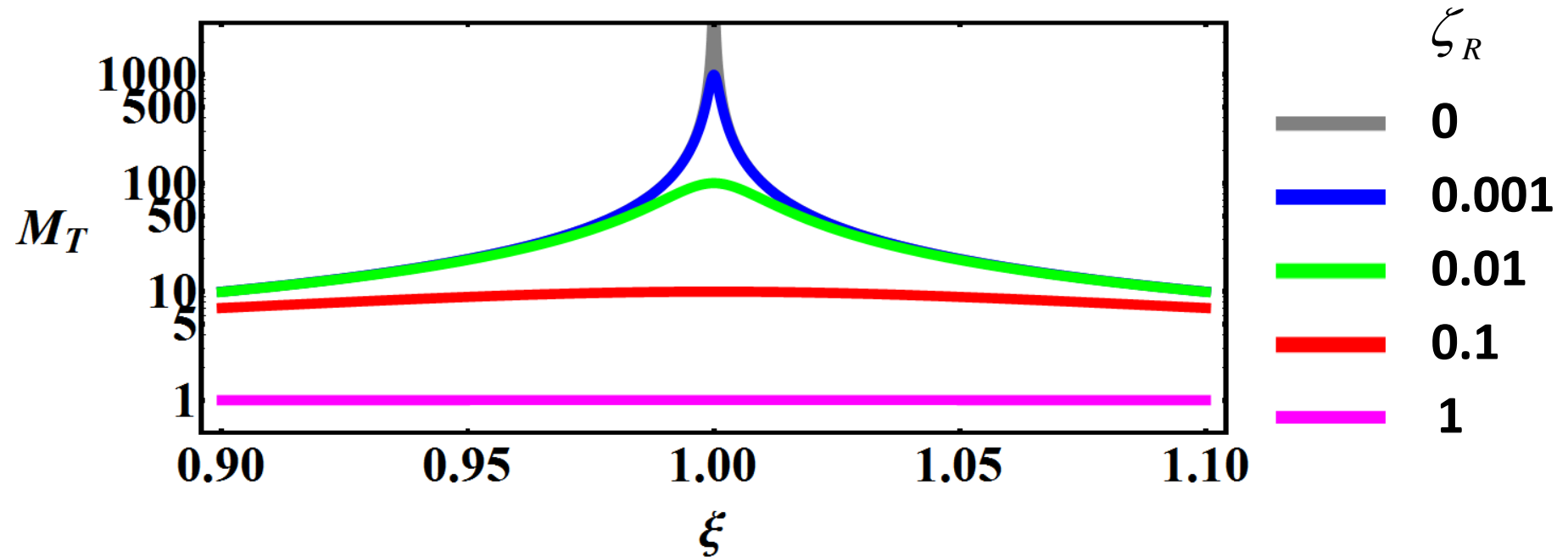


$$\xi'_{max} = 1 + \frac{1}{2\zeta_R}$$



maximal position of the image as a function of the normalized Rayleigh length  
of the input Gaussian Beam

# Magnification



*Transverse  
Magnification is  
maximized at*

$$\xi = 1 \Rightarrow M_T^{\max} = \frac{1}{\zeta_R}$$



# At what distance $\xi$ is the output beam collimated?

*One safe criterion is that the output Gaussian beam is **collimated** when it's Rayleigh length is **maximized**.*

*Since the only varying parameter is the distance of the input Gaussian beam to the lens, it is sufficient to estimate the maximum of the function  $\zeta_R'$  in respect to  $\xi$ :*

$$\frac{\partial \zeta_R'}{\partial \xi} = 0 \Rightarrow \xi_{col} = 1$$

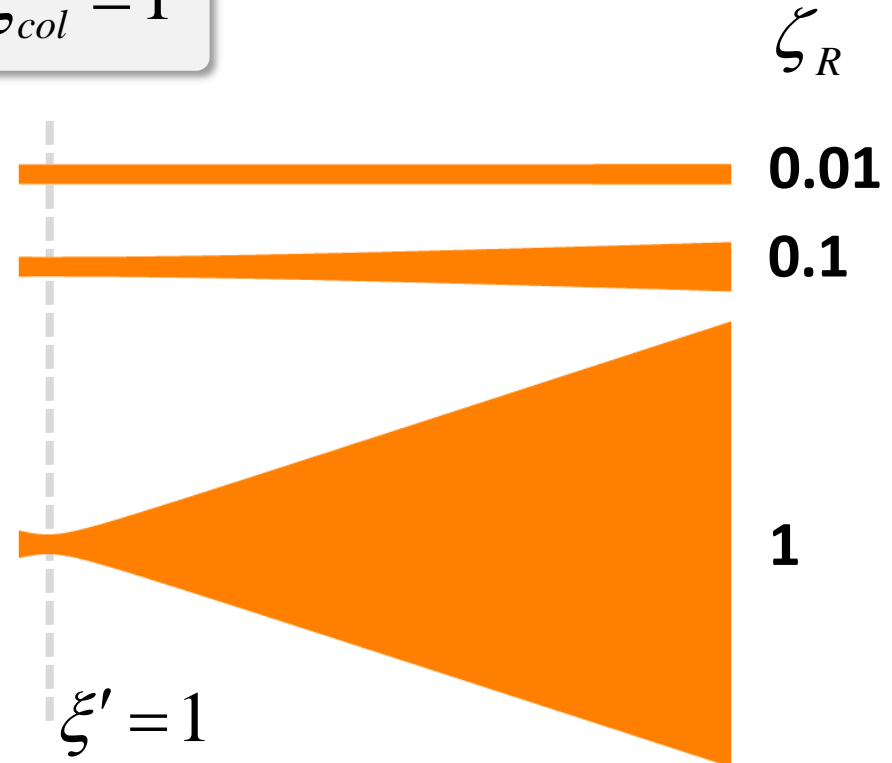
*In this case:*

$$\xi' = 1$$

$$\zeta_R' = \frac{1}{\zeta_R}$$

$$M_T = \frac{1}{\zeta_R}$$

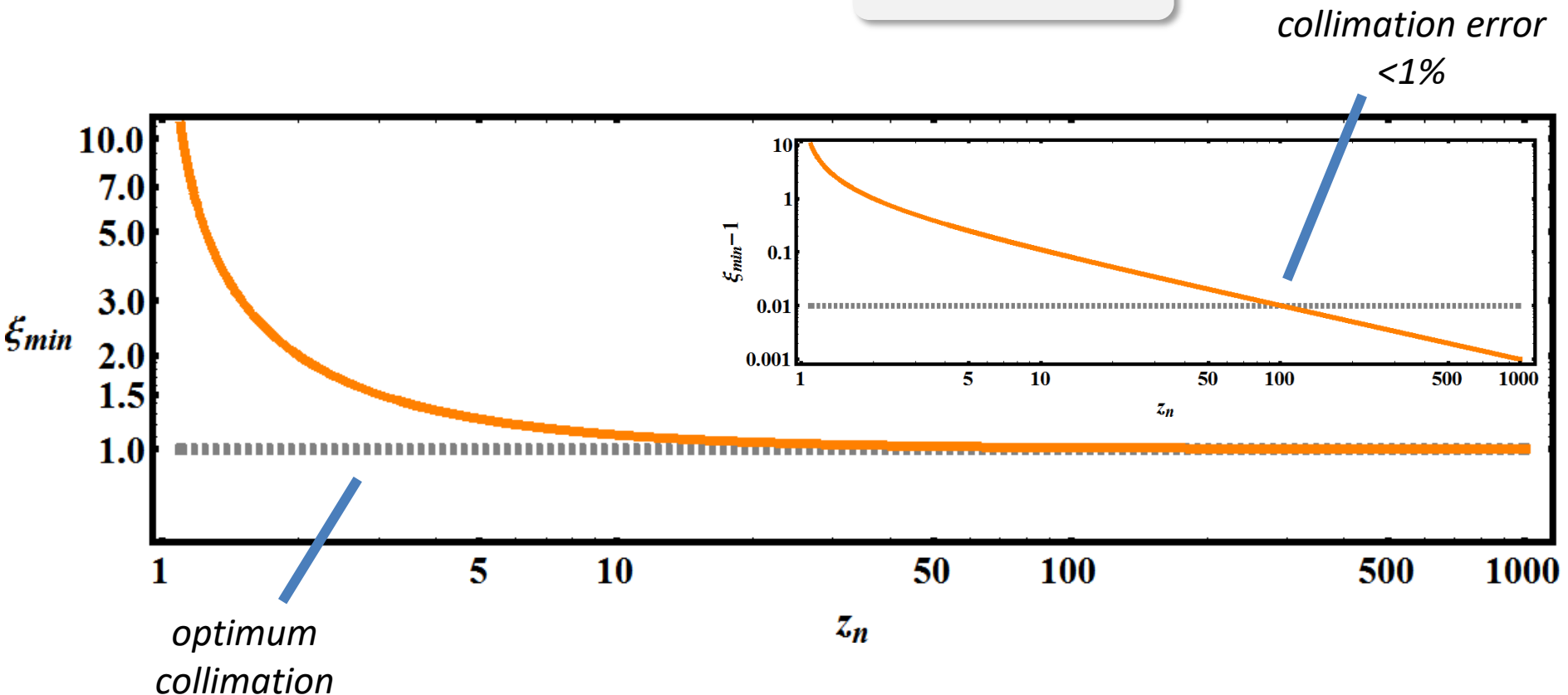
$$M_L = \frac{1}{\zeta_R^2}$$



# Minimum beam diameter at a distance $z_n$

Since the only varying parameter is the distance of the input Gaussian beam to the lens, it is sufficient to estimate the minimum of the function  $w_{out}$  in respect to  $\xi$ :

$$\frac{\partial w_{out}(\xi, \zeta_R; z_n)}{\partial \xi} = 0 \Rightarrow \xi_{min} = \frac{z_n}{z_n - 1}$$

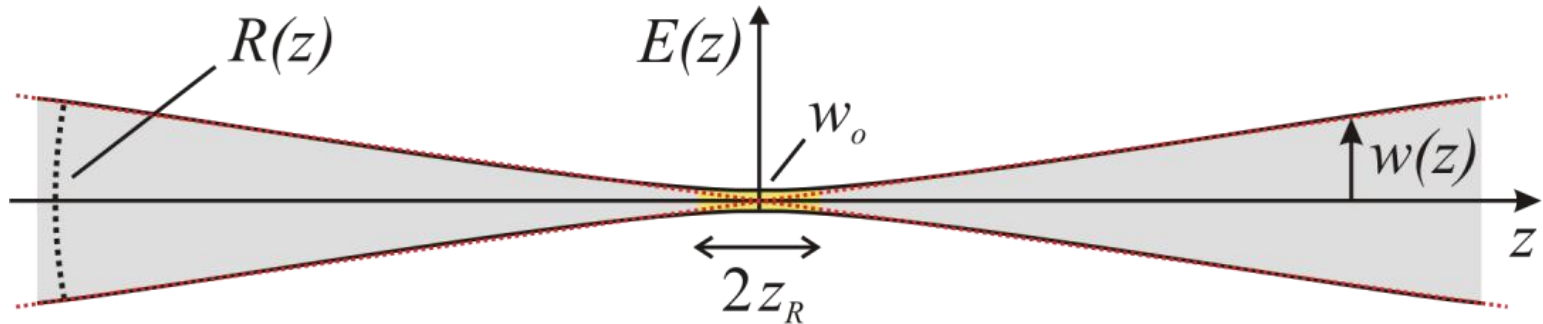


# ABCD matrix theory and Gaussian beams



$\mathcal{M}$

# Complex parameter $q$



$$E(\mathbf{r}) = E_o \frac{e^{-i[kz + \psi(z)]}}{w(z)} \cdot e^{-\frac{x^2 + y^2}{w^2(z)} - i \frac{\pi}{\lambda} \frac{x^2 + y^2}{R(z)}} \equiv E_o \frac{e^{-i[kz + \psi(z)]}}{w(z)} e^{-i \frac{\pi}{\lambda} \frac{x^2 + y^2}{q(z)}}$$

$$\frac{1}{q(z)} \equiv \frac{1}{R(z)} - i \frac{\lambda}{\pi} \frac{1}{w^2(z)}$$

The complex parameter  $q(z)$  describes the Gaussian beam completely

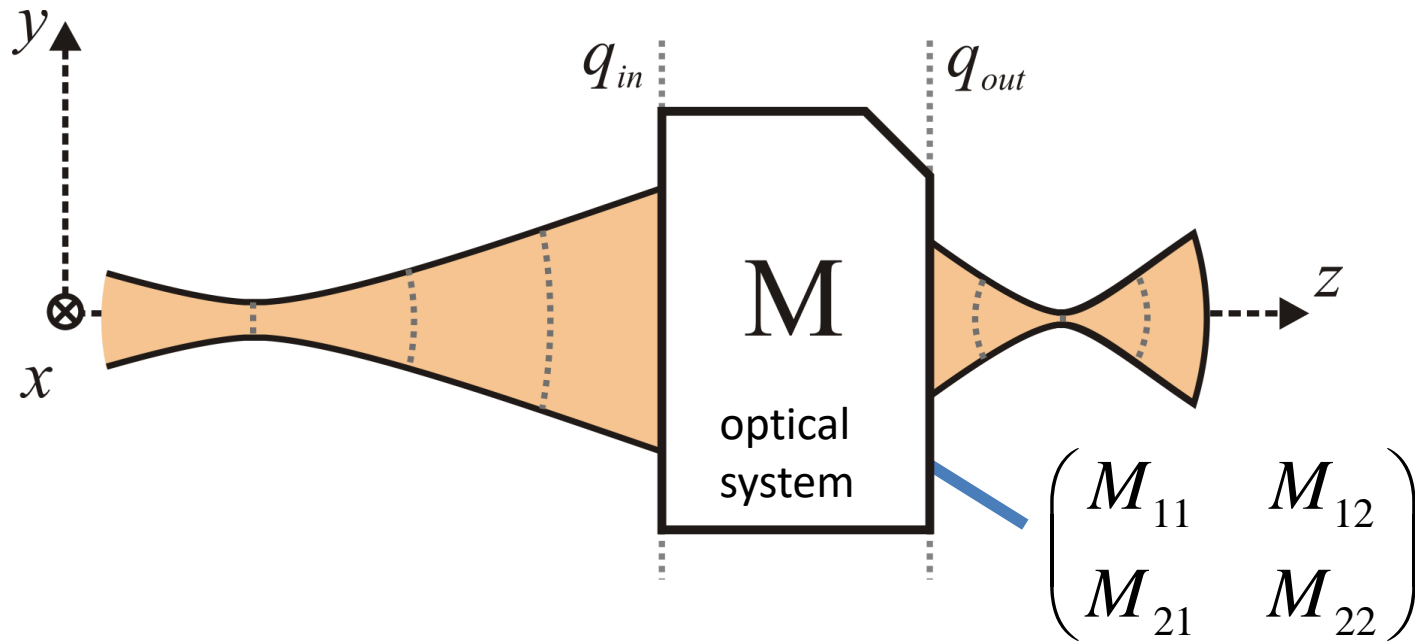
By replacing the known values of the  $R(z)$ ,  $w(z)$  parameters for the free space propagation we get:

$$\left. \begin{aligned} w(z) &= w_o \sqrt{1 + \frac{z^2}{z_R^2}} \\ R(z) &= z + \frac{z_R^2}{z} \end{aligned} \right\} \Rightarrow \frac{1}{q(z)} = \frac{1}{z + z_R^2 / z} - i \frac{1}{z_R + z^2 / z_R}$$

$$\frac{1}{q(z)} = \frac{z_R + z^2 / z_R - i(z + z_R^2 / z)}{(z_R + z^2 / z_R)(z_R + z^2 / z_R)} \Rightarrow$$

$$q(z) = z + i z_R$$

# Propagation through an optical system



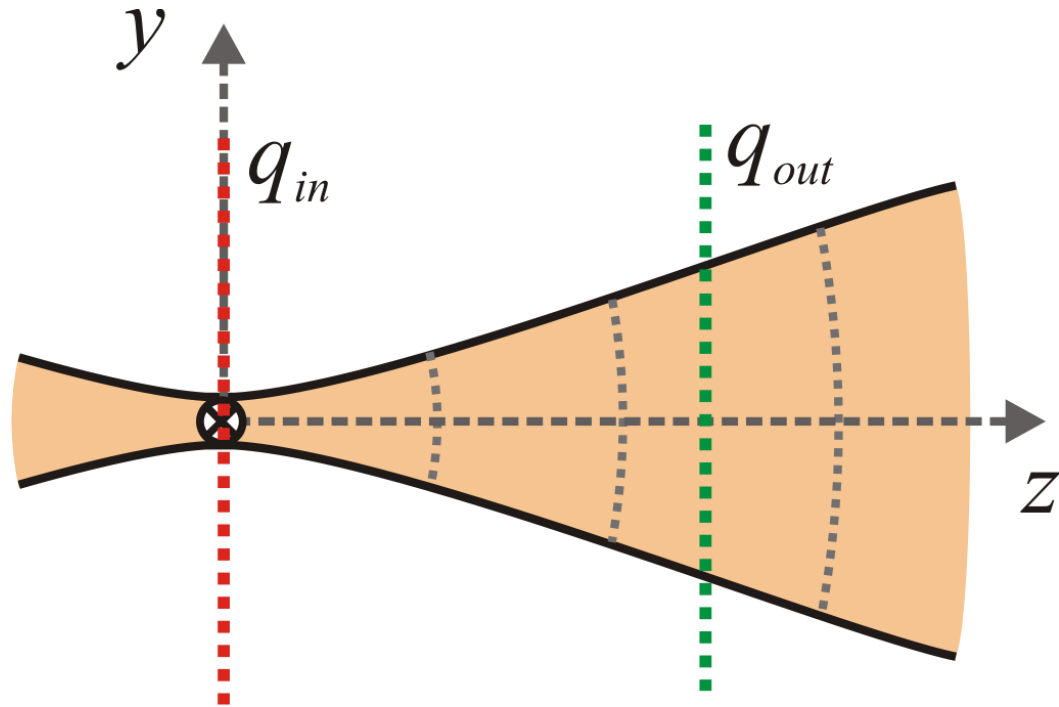
$$q_{out} \equiv \frac{M_{11} q_{in} + M_{12}}{M_{21} q_{in} + M_{22}}$$

The parameter  $q$  in the exit of an optical system is defined by the optical transfer matrix of the system  $M$  and its value in the input of the system.

## Application in free space propagation

*From the ray matrix theory we know that free space propagation is described as a simple transport over  $z$*

$$M = \begin{pmatrix} 1 & z \\ 0 & 1 \end{pmatrix}$$



$$\left. \begin{array}{l} q(z) \equiv q_{out} \\ q(0) = q_{in} = i z_R \end{array} \right\} \Rightarrow q(z) = \frac{1 \cdot q(0) + z}{0 \cdot q(0) + 1} = z + i z_R$$