



UNIVERSITÀ DEGLI STUDI DI MILANO
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Master degree in Physics

**Development of an open-source calibration framework for
superconducting qubits**

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Summary

Chapter 1

Notes on quantum computing

1.1 Qubits

1.2 Operation on qubits

1.2.1 Density matrix

1.3 Quantum operations

A quantum operation is a mathematical transformation that describes how a quantum state changes as a consequence of a physical process. Formally, it is a map \mathcal{E} that transforms a quantum state described by a density operator $\hat{\rho}$ into another state described by a new density operator $\hat{\rho}'$:

$$\mathcal{E}(\rho) = \rho'. \quad (1.1)$$

The simplest example of a quantum operation is the evolution of a quantum state $\hat{\rho}$ of a closed quantum system, under a unitary operator \hat{U} , which can be written as $\mathcal{E} \equiv \hat{U}\hat{\rho}\hat{U}^\dagger$.

Depolarizing channel A depolarizing channel describes a process in which the current state of the n -qubit system ρ , is replaced by $\frac{\mathbb{I}}{2^n}$, with probability d . This process can be represented with a quantum map as follows:

$$\mathcal{E}_{dc}(\rho) = d \frac{\mathbb{I}}{2^n} + (1 - d)\rho \quad (1.2)$$

1.4 Superconducting qubits

Chapter 2

Qibo

Chapter 3

Results

Tutti i risultati che sono presentati nel seguito sono stati ottenuti utilizzando il software di **Qibolab** per l'interazione con gli strumenti del laboratorio e **Qibocal** per il controllo delle operazioni sui qubit. L'hardware è un chip di QunatumWare. Durante il lavoro condotto per questo progetto di tesi entrambe le librerie, sia Qibocal che Qibolab undergo update and release, for this reason the first part of this work was realized using Qibocalv0.1 and Qibolabv0.1 while the second part of the work, dato che puntava anche allo sviluppo di routine che potessero essere utili per la calibrazione dei qubit è stato realizzato direttamente con Qibocalv0.2 e Qibolabv0.2.

3.1 RB fidelity optimization

3.1.1 Randomized Benchmarking

A strong limitation to the realization of quantum computing technologies is the loss of coherence that happens as a consequence of the application of many sequential quantum gates to the qubits. A possible approach to characterize gate error is the quantum process tomography which allows the experimenter to establish the behaviour of a quantum gates; the main drawback of this approach is that process tomography can be very time consuming since its time complexity scales exponentially with the number of qubits involved [1] and the result is affected by state preparation and measurements (SPAM) errors.

To overcome these limitations, randomized benchmarking (RB) was introduced and is currently widely used to quantify the average error rate for a set of quantum gates.

The main idea is that the error obtained from the combined action of random unitary gates drawn from a uniform distribution with respect to the Haar measure [2] and applied in sequence to the qubit will average out to behave like a depolarizing channel [3]. This last consideration simplifies the characterization of noise because it removes dependence on specific error structures and allows fidelity to be extracted through a simple exponential decay.

It was later shown that it is possible to simplify this procedure even more, by restricting the unitaries to gates in the Clifford group¹ and by not requiring that the sequence is strictly self-inverting [4].

The fundamental principle of RB is the application of sequences of randomly selected quantum gates from the Clifford group \mathcal{C} followed by an inversion gate which, in absence of noise, return the system to its initial state. For real systems, where noise is present, the observed survival probability provides an estimate of the average gate fidelity.

The standard RB protocols consist of the following steps:

1. Initialize the system in ground state $|0\rangle$
2. For each sequence-length m build a sequence of m randomly drawn Clifford gates C_1, C_2, \dots, C_m
3. Determine the inverse gate $C_{m+1} = (C_m \circ \dots \circ C_1)^{-1}$
4. Measure $C_{m+1} \circ C_m \circ \dots \circ C_1 |0\rangle$

The process must be repeated for multiple sequence of the same length and with varying length.

In ideal systems without noise we should have

$$C_{m+1} \circ C_m \circ \dots \circ C_1 |0\rangle = (C_m \circ \dots \circ C_1)^{-1} \circ (C_m \circ \dots \circ C_1) |0\rangle = |0\rangle \quad (3.1)$$

¹unitary rotations mapping the group of Pauli operators in itself

In realsystems, where noise is present, eq. ?? does not hold; instead randomization with Clifford gates behave as a depolarizing channel 1.2 with depolarization probability d . The survival probability of the initial state $|0\rangle$ for different sequence lengths follows the exponential decay model

$$F(m) = Ap^m + B, \quad (3.2)$$

where $1 - p$ is the rate of depolarization and A and B capture the state preparation and measurement error but not the rate of decay p . Note that the exponential form arises naturally due to the assumption that each gate introduces independent noise.

The parameter p is directly related to the depolarization probability d through the average gate fidelity F which, for a depolarizing channel, is given by

$$F = 1 - \frac{d}{2^n - 1}. \quad (3.3)$$

For the details of the calculations to obtain eq. 3.3 see 4

Now we can derive the average error per Clifford gate $\epsilon_{Clifford}$

$$\epsilon_{Clifford} = 1 - F, \quad (3.4)$$

where F is the average gate fidelity. Substituting in 3.4 the formula for the average gate fidelity 3.3 we obtain

$$\epsilon_{Clifford} = \frac{d}{2^n - 1} = \frac{1 - p}{1 - 2^{-n}}, \quad (3.5)$$

which shows how the average error per Clifford gate is directly connected to the exponential decay rate.

QUA Randomized Benchmarking

For the results we present in the following the technique used slightly differs from the one described in section 3.1.1

3.1.2 Optimization methods

I primi metodi che abbiamo provato per l'ottimizzazione dei parametri sono quelli standard implementati nella libreria `Scipy` [5] evitando metodi gradient-based considerato il landscape potenzialmente complicato della funzione RB. The first gradient-free optimization method to be tested was Nelder-Mead since in letteratura era già stato riportato il suo utilizzo per obiettivi simili [6].

The Nelder-Mead or simplex method is a widely used numerical optimization technique for unconstrained problems in multidimensional spaces. It was originally introduced by Nelder and Mead in 1965 [7]; this derivative-free method is particularly effective for minimizing non-differentiable, noisy functions where gradient-based methods struggle.

The Nelder-Mead method operates using simplex, which is a polytope of $n + 1$ vertices in a n -dimensional space. The algorithm iteratively updates the simplex by replacing its worst-performing vertex with a new candidate point, thereby guiding the search towards an optimal solution. If the goal is to minimize a given function $f(\mathbf{x})$ where $\mathbf{x} \in \mathbb{R}^n$ the algorithms proceeds with the following steps:

1. If not otherwise initialized, $n + 1$ points are sampled for building the initial simplex
2. **Order** the test points according to their values at vertices: $f(\mathbf{x}_1) \leq f(\mathbf{x}_2) \leq \dots \leq f(\mathbf{x}_{n+1})$ and check whether the algorithm should terminate.
3. **Calculate** \mathbf{x}_0 , the centroid of all points except \mathbf{x}_{n+1} .
4. **Reflection**: Compute the reflected point $\mathbf{x}_r = \mathbf{x}_0 + \alpha(\mathbf{x}_0 - \mathbf{x}_{n+1})$ with $\alpha > 0$. If \mathbf{x}_r satisfies $f(\mathbf{x}_1) \leq f(\mathbf{x}_r) < f(\mathbf{x}_n)$, then a new simplex is obtained by replacing the worst-performing point \mathbf{x}_{n+1} with \mathbf{x}_r and then go to step 1.
5. **Expansion**: If \mathbf{x}_r is the current best point, meaning that $f(\mathbf{x}_r) < f(\mathbf{x}_1)$, then the expanded point is computed: $\mathbf{x}_e = \mathbf{x}_0 + \gamma(\mathbf{x}_r - \mathbf{x}_0)$ with $\gamma > 1$. If \mathbf{x}_e satisfies $f(\mathbf{x}_e) < f(\mathbf{x}_r)$, then a new simplex is obtained by replacing \mathbf{x}_{n+1} with the expanded point \mathbf{x}_e and then go to step 1. If instead $f(\mathbf{x}_e) \geq f(\mathbf{x}_r)$, the new simplex is obtained by replacing \mathbf{x}_{n+1} with \mathbf{x}_r , and then go to step 1.
6. **Contraction**: In this case is certain that $f(\mathbf{x}_r) \geq f(\mathbf{x}_n)$ then:

- If $f(\mathbf{x}_r) < f(\mathbf{x}_{n+1})$: compute the contracted point $\mathbf{x}_c = \mathbf{x}_0 + \rho(\mathbf{x}_r - \mathbf{x}_0)$ with $0 < \rho \leq 0.5$. If \mathbf{x}_c satisfies $f(\mathbf{x}_c) < f(\mathbf{x}_r)$, then a new simplex is obtained by replacing \mathbf{x}_{n+1} with \mathbf{x}_c and go to step 1. Else go to step 6.
- If $f(\mathbf{x}_r) \geq f(\mathbf{x}_{n+1})$: compute the contracted point $\mathbf{x}_c = \mathbf{x}_0 + \rho(\mathbf{x}_{n+1} - \mathbf{x}_0)$ with $0 < \rho \leq 0.5$. If \mathbf{x}_c satisfies $f(\mathbf{x}_c) < f(\mathbf{x}_{n+1})$, then a new simplex is constructed with \mathbf{x}_c and go to step 1. Else go to step 6.

7. **Shrinkage**: Replace all points except the best, \mathbf{x}_1 , with $\mathbf{x}_i = \sigma(\mathbf{x}_i - \mathbf{x}_1)$

The algorithm terminates when the standard deviation of the function values of the current simplex fall below a user-initialized tolerance. When the cycle stops the point of the simplex associated to the lower function value is returned as proposed optimum

The values of the parameters α, γ, ρ and σ were left to default of `scipy`: $\alpha = 1, \gamma = 2, \rho = 0.5, \sigma = 0.5$.

Optuna [8]

CMA-ES [9]

3.2 RX90 calibration

Another possible source of error is ...

3.3 Flux pulse correction

3.3.1 Notes on signal analysis

3.3.2 Cryoscope

The experiment that we describe in this section was first introduced in [10], the goal is to determine predistortions that needs to be applied to a flux pulse signal so that the qubit receives the flux pulse as intended by the experimenter.

3.3.3 Filter determination

IIR

FIR

for description and notes on CMA-ES see section 3.1.2

Output filters in QM

Chapter 4

Conclusions

Appendix A

Starting with the definition of average gate fidelity:

$$F = \int d\psi \langle \psi | U^\dagger \mathcal{E}(|\psi\rangle \langle \psi|) U | \psi \rangle$$

The quantum maps that represent a depolarizing channel is

$$\mathcal{E}(\rho) = (1-d)U\rho U^\dagger + d\frac{I}{2^n} \quad (4.1)$$

When we substitute it in the average gate fidelity definition we get:

$$\begin{aligned} F &= \int d\psi \langle \psi | U^\dagger \left[(1-d)U | \psi \rangle \langle \psi | U^\dagger + d\frac{I}{2^n} \right] U | \psi \rangle \\ &= \int d\psi \langle \psi | U^\dagger [(1-d)U | \psi \rangle \langle \psi | U^\dagger] U | \psi \rangle + \int d\psi \langle \psi | U^\dagger \left[d\frac{I}{2^n} \right] U | \psi \rangle \\ &= (1-d) \int d\psi \langle \psi | \psi \rangle \langle \psi | \psi \rangle + \frac{d}{2^n} \int d\psi \langle \psi | U^\dagger I U | \psi \rangle \\ &= (1-d) + \frac{d}{2^n} \int d\psi \langle \psi | \psi \rangle \\ &= (1-d) + \frac{d}{2^n} = 1 - d + \frac{d}{2^n} = 1 - d \left(1 - \frac{1}{2^n} \right) \\ &= 1 - d \frac{2^n - 1}{2^n} \end{aligned}$$

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