

Experiments in Physics

PHYSICS 1292
GENERAL PHYSICS II LAB

COLUMBIA UNIVERSITY



DEPARTMENT OF PHYSICS

SPRING 2015

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Introduction 2-0

General Instructions

1 Purpose of the Laboratory

The laboratory experiments described in this manual are an important part of your physics course. Most of the experiments are designed to illustrate important concepts described in the lectures. Whenever possible, the material will have been discussed in lecture before you come to the laboratory.

The sections headed Applications and Lab Preparation Examples, which are included in some of the manual sections, are *not* required reading unless your laboratory instructor specifically assigns some part. The Applications are intended to be motivational and so should indicate the importance of the laboratory material in medical and other applications. The Lab Preparation Examples are designed to help you prepare for the lab; you will not be required to answer all these questions (though you should be able to answer any of them by the end of the lab). The individual laboratory instructors may require you to prepare answers to a subset of these problems.

2 Preparation for the Laboratory

In order to keep the total time spent on laboratory work within reasonable bounds, the write-up for each experiment will be completed at the end of the lab and handed in *before the end of each laboratory period*. Therefore, it is imperative that you spend sufficient time preparing for the experiment *before* coming to laboratory. You should take advantage of the opportunity that the experiments are set up in the Lab Library (Room 506) and that TAs there are willing to discuss the procedure with you.

At each laboratory session, the instructor will take a few minutes at the beginning to go over the experiment to describe the equipment to be used and to outline the important issues. This does not substitute for careful preparation beforehand! You are expected to have studied the manual and appropriate references at home so that you are prepared when you arrive to perform the experiment. The instructor will be available primarily to answer questions, aid you in the use of the equipment, discuss the physics behind the experiment, and guide you in completing your analysis and write-up. Your instructor will describe his/her policy regarding expectations during the first lab meeting.

Some experiments and write-ups may be completed in less than the three-hour laboratory period, but under no circumstances will you be permitted to stay in the lab

after the end of the period or to take your report home to complete it. If it appears that you will be unable to complete all parts of the experiment, the instructor will arrange with you to limit the experimental work so that you have enough time to write the report during the lab period.

Note: Laboratory equipment must be handled with care and each laboratory bench must be returned to a neat and orderly state before you leave the laboratory. In particular, you must turn off all sources of electricity, water, and gas.

3 Bring to Each Laboratory Session

- A pocket calculator (with basic arithmetic and trigonometric operations).
- A pad of 8.5×11 inch graph paper and a sharp pencil. (You will write your reports on this paper, including your graphs. Covers and staplers will be provided in the laboratory.)
- A ruler (at least 10 cm long).

4 Graph Plotting

Frequently, a graph is the clearest way to represent the relationship between the quantities of interest. There are a number of conventions, which we include below.

- A graph indicates a relation between two quantities, x and y , when other variables or parameters have fixed values. Before plotting points on a graph, it may be useful to arrange the corresponding values of x and y in a table.
- Choose a convenient scale for each axis so that the plotted points will occupy a substantial part of the graph paper, but do not choose a scale which is difficult to plot and read, such as 3 or 3/4 units to a square. Graphs should usually be at least half a page in size.
- Label each axis to identify the variable being plotted and the units being used. Mark prominent divisions on each axis with appropriate numbers.
- Identify plotted *points* with appropriate symbols, such as crosses, and when necessary draw vertical or horizontal *bars* through the points to indicate the range of uncertainty involved in these points.
- Often there will be a theory concerning the relationship of the two plotted variables. A linear relationship can be demonstrated if the data points fall along a

single straight line. There are mathematical techniques for determining which straight line best fits the data, but for the purposes of this lab it will be sufficient if you simply make a rough estimate visually. *The straight line should be drawn as near the mean of the all various points as is optimal.* That is, the line need not precisely pass through the first and last points. Instead, each point should be considered as accurate as any other point (unless there are experimental reasons why some points are less accurate than others). The line should be drawn with about as many points above it as below it, and with the ‘aboves’ and ‘below’ distributed at random along the line. (For example, not all points should be above the line at one end and below at the other end).

5 Questions or Complaints

If you have a difficulty, you should attempt to work it through with your laboratory instructor. If you cannot resolve it, you may discuss such issues with:

- one of the laboratory Preceptors in Pupin Room 729;
- the Undergraduate Assistant in the Departmental Office – Pupin Room 704;
- the instructor in the lecture course, or the Director of Undergraduate Studies;
- your undergraduate advisor.

As a general rule, it is a good idea to work downward through this list, though some issues may be more appropriate for one person than another.

Experiment 2-1

Electric Fields

1 Introduction

One of the fundamental concepts used to describe electric phenomena is that of the electric field. (Later we will also deal with magnetic fields.) We explore the field concept experimentally by measuring lines of equal electric potential on carbon paper and then constructing the electric field lines that connect them.

2 Theory

2.1 Fields

A field is a function in which a numerical value associated with a physical quantity is assigned to every point in space. For example, each point in the United States may be assigned a temperature (that we measure with a thermometer at that position), as shown in figure 1. Thus, temperature as a function of location is called the “temperature field”. The temperature field is an example of a scalar field, since the value assigned to each location is a scalar.

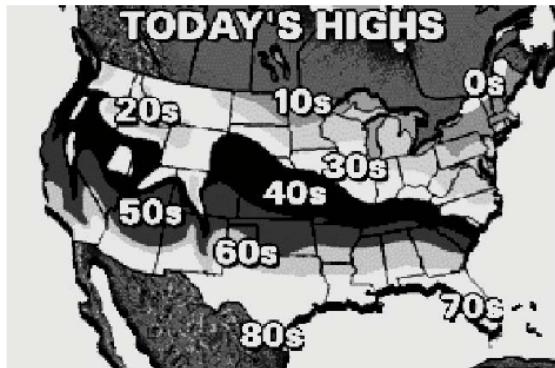


Figure 1: Temperature of Different Places in the United States

The electric field (as well as the magnetic field) is a vector field, because a vector (rather than a scalar) is associated with each position. At each point in space, the magnitude and direction of the electric field are equal to the magnitude and direction of the electric force that would act on a unit test charge at this position.

A line that is formed by connecting electric field vectors and follows the direction

of the field is called a field line. A field is considered homogeneous if the field lines are parallel. The potential associated with such a field increases linearly with the distance traveled along the field lines.

2.2 Equipotential Contours for Gravity

A topographical map is a map with lines that indicate contours of the same height. On such a map, a mountain may look like figure 2

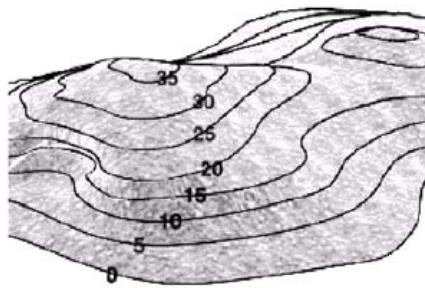


Figure 2: A Topographical Map

These contours of the same height (h) are equi-height lines or, if you think in terms of potential energy ($U = mgh$), they are equipotential lines. That is, every point on the contour has the same value of potential energy. So as you move along a contour, you do not change height, and therefore you neither gain nor lose potential energy. Only as you step up or down the hill do you change the energy.

There are a few general properties of equipotential contours:

- They never intersect. (A single point cannot be at two different heights at the same time, and therefore it cannot be on two different contours.)
- They close on themselves. (A line corresponding to constant height cannot just end.)
- They are smooth curves, as long as the topography has no sharp discontinuities (like a cliff).

2.3 Analogy between Electric and Gravitational Potentials

Electric potential is analogous to gravitational potential energy. Of course, gravitational potential energy arises from any object (with mass), whereas the electric potential arises only from charged objects.

As in mechanics, the absolute value of potential is not important. The differences in potential are the quantities that have physical meaning. In a real problem, it is usually best to choose a reference point, define it as having zero potential, and refer to potentials at all other points relative to that point. Similarly, topographic height is usually reported relative to the reference at sea level.

2.4 Field Lines

Let's exploit our analogy of mountains and valleys a bit further for electric fields. Electric field lines indicate the electric force on a test charge (magnitude and direction). Gravitational field lines tell us about the gravitational force on a test charge (magnitude and direction) – the direction it tends to roll down the hill and the acceleration it experiences during its journey. So if we carefully track a marble as we roll it down the hill in small increments, we can obtain a “field-line”.

What is the characteristic of a field line? Think about placing a marble on an inclined hillside. Which way will it roll if you release it? It will always roll in the direction of steepest descent, and this direction is always perpendicular to the contours that indicate equal height (which we have been referring to as equipotential lines).

So it is for the electric case: if we know the equipotential lines, we can draw the field lines such that each one is perpendicular to each equipotential line. Figure 3 and 4 both show equipotential lines (dashed lines) and the corresponding field lines (solid lines).

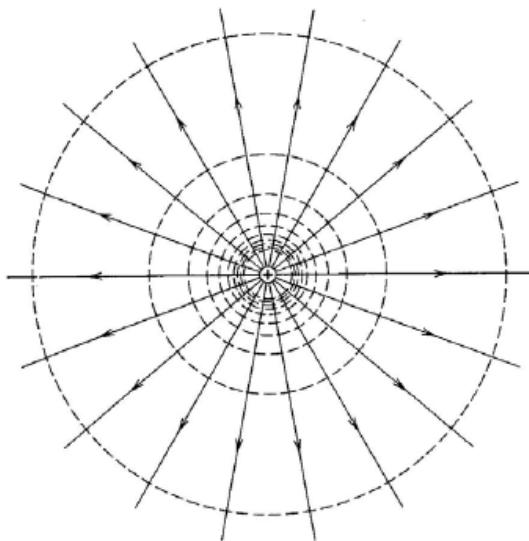


Figure 3: Equipotential Lines and Field Lines of a Single Charge

Keep in mind that the field lines (perpendicular to equipotential lines) give the biggest change in height over the shortest distance!

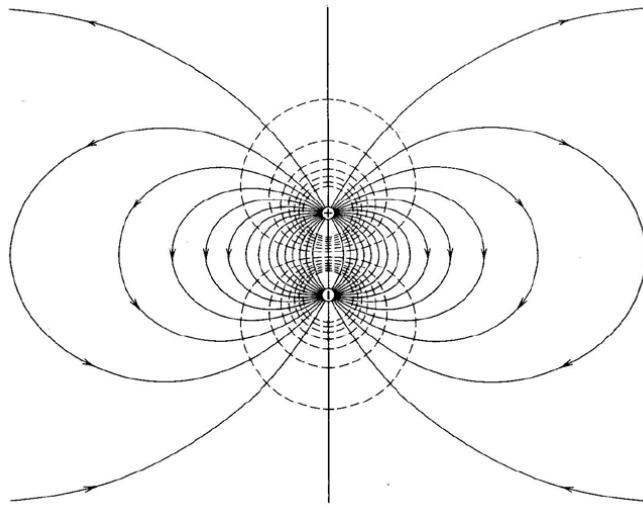


Figure 4: Equipotential Lines and Field Lines of a Pair of Opposite Charges

There are also a few rules for field lines:

- Field lines always begin and end on charges. (They often terminate on material surfaces, but that is because there are charges on those surfaces.)
- They never intersect each other (except at electric charges, where they also terminate).
- They always intersect equipotential lines perpendicularly!
- They are usually smoothly continuous (except if they terminate).

2.5 Metal Surfaces

It turns out that for electrical phenomena, metals are equipotentials. So, by using metals, we can impose some unusually shaped equipotential lines and determine the corresponding field lines.

Metals are always equipotentials because:

1. they are excellent conductors of electric charge, and
2. they have an abundant supply of freely moving negative charges (electrons).

Based on these two properties, we can understand that, if a field line were to penetrate a metal surface, these electrons would feel electric forces to move them through the material in the direction opposite to that of the electric field (since they are negative) until they would be forced to stop upon encountering the metal surface on the other side. Therefore, the free electrons ultimately distribute themselves along the surface of the metal so that there is a net negative charge on one side of the conductor and a net positive charge on the other (from where the electrons have fled). This realignment of the free electrons creates a field that precisely cancels the imposed field resulting in no net field, and the remaining free charges inside the conductor then feel no force and don't move! Hence, field lines imposed from outside always end on metal surfaces (which they intersect perpendicularly), and the surface (with the entire interior) is an equipotential.

2.6 Electric Shielding Theorem

As just discussed, whenever a piece of metal is placed in an electric field, the entire metal will remain an equipotential. That is, every point in the metal will be at the same potential as every other point. No field lines penetrate through metals.

What happens if we take an enclosed container of metal of arbitrary shape, say a tin can, and put it into an electric field? Since no field gets through the metal, there must be no electric field inside the can. This means that every point inside the container, as well as all points on the container surface must be at the same potential. (If they were not, then there would be differences in the potential and therefore an electric field.)

Lets return to our analogy of gravity equipotentials in hills and valleys, with field lines in the direction a marble would roll down the hill. Consider how a frozen lake would look on such a contour map. The surface of the lake has the same gravitational potential energy at all points, therefore if you place a marble on the frozen surface of the lake, the marble will not roll anywhere. In other words, there is no component of the gravitational field along the surface of the lake to push the marble. Similarly, there is no electric field on the inside of a closed metal container to push the charges around, and the entire interior is therefore a constant equipotential surface.

2.7 Remarks for Experts

1. The gravitational analogy operates in only two dimensions: the horizontal coordinates describing the surface of the lake. When discussing electrical fields we should, in principle, take into account all three dimensions. The analogy of the lake surface for the electrical case is the entire volume (three dimensions) of the interior of the can. For this experiment, we “cheat” a little by looking at electric field within a two dimensional world of slightly conducting paper. But the field

and potential arrangements are as one expects for electrostatic phenomena in a two dimensional system.

2. To shield electric field, we actually don't need a metallic surface that is completely closed; a closed cage of wire mesh (Faraday cage) is sufficient. So a sedan will work like a Faraday cage, but a convertible will not since it is not closed at the top.

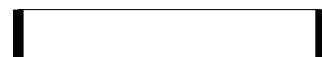
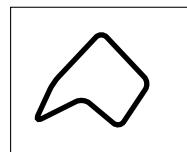
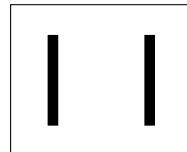
3 Experiments

In this experiment, we use pieces of slightly conductive paper, electrodes, and metallic paint to determine the equipotential and field configurations. Before we actually start our measurements we will have to draw the shapes on the conductive paper using the metal paint. The metal paint will take a few minutes to dry.

3.1 Preparation of the Conductive Paper

Take four pieces of conductive paper – three full-sized and one thin strip and paint on them with the metallic paint according to the following instructions. Each shape drawn by the metal paint is an equipotential curve. By connecting the pins to areas covered with paint, we can determine the field configurations between different equipotential shapes.

1. Leave the first piece of paper blank.
2. On the second piece of paper paint two parallel plates. Use a piece of cardboard provided to get sharp edges with the paint. Don't use the rulers!
3. On the third piece of paper paint a closed loop of any shape that does not enclose an electrode.
4. Take the thin strip and paint a line on each short end of the strip.



Hints:

1. You get the best results if you stay at least 1-2 cm away from the edges of the paper. Also make sure that the pins are always in contact with the paper. (Otherwise the experiment does not work!) Be careful not to make the holes too big or the pins will no longer be in contact with the paper.

2. Give the metal paint enough time to dry. First paint all the configurations you will use on the conductive paper. While they dry, get started on the two point-charge arrangement!

3.2 Setup of the Voltage Source

- “OUTPUT A” of the power supply will be used in this experiment to provide a constant potential difference of 10 volts.
- Turn the “A VOLTAGE”, the “A CURRENT”, the “B VOLTAGE”, and the “B CURRENT” control knobs located on the right side of the front panel counter-clockwise to the “MIN” position.
- Set the “A/B OUTPUT” switch located on the upper right of the front panel to the “INDEPENDENT” setting, and set the “A/B METER” switch located between the two meters on the front panel to the “A” setting.
- Turn the power on by using the button located on the left of the front panel.
- Slowly turn the “A VOLTAGE” output knob until the voltage reads 10 volts. After doing this, do not make any further changes on the power supply for the remainder of the experiment. (The current meter will stay at a zero reading.)
- Connect two cables to the “A OUTPUT (0 ~ 20 V 0.5 A)” of the power supply (red and black poles) and then connect these cables to two of the yellow-metallic pins.

NOTE: To use the digital multimeter as an appropriate voltmeter for this experiment: turn the multimeter power on, select “DC” (button in out position), select “VOLTS” (push button in), set range to “2 VOLTS”, and connect your voltage-measuring cables to the “V- Ω ” and the “COM” jacks.

3.3 Equipotential and Field Lines

- Connect two cables to the output of the power supply (black and red poles) and two of the yellow-metallic pins as described in the last step of the previous part. These are your output pins.
- Take one additional cable and connect one end to the black pole of the multimeter and the other to a different metallic pin. This is your reference or fixed-point probe. Take the pen-like pin and plug it into the red pole of the multimeter. This is your measuring probe.

- Take an empty sheet and put the two output pins on it to act as point-like sources. Push the pins in only a little bit and not all the way through.
- Insert the reference probe at an arbitrary position and move the measuring probe until the multimeter shows a value of zero. This means that the two points are at the same potential. Mark this position with one of the red markers provided, so as not to confuse it with the probe. Search for more equipotential points until you have enough to draw an equipotential contour.
- Move the reference pin to a new position and construct another equipotential line.
- Repeat for a total of about 5 equipotential lines.
- Given these equipotential lines, draw about 5 field lines using the yellow pen. Remember that the field lines and equipotential lines always intersect perpendicularly.
- Do your equipotential and field lines show the symmetries you would expect for this system? Just what symmetries does/should this system have?
- Obtain the equipotential and field lines for the parallel-plate configuration as well. (Stick the two pins from the function generator into the metal paint-covered part of the paper, after the paint has dried!).
- Are the equipotential lines as you expect them?
- What the major problems and uncertainties in this part?

3.4 Electric Shielding Theorem

Take the sheet with the conducting loop on it, and set up an electric field. You need the reference pin as in the previous part. Put it somewhere either within the loop or on the rim of it.

Put the reference pin anywhere within the closed surface or on its rim, and check that any point within the loop has the same potential as the reference pin. This demonstrates that all points within the closed loop have the same potential.

- What do you conclude about the shielding theorem? Does it hold or not?
- What would happen with your readings if you had drawn the loop badly such that the loop was not perfectly closed?

3.5 Linear Increase in Potential

The purpose is to verify that the potential increases linearly with distance as you move the probe from one end to the other.

Take the small strip and connect the two output pins to the metal covered ends. This produces a potential difference between the ends. Take the reference probe and remove the pin from the end. (You only need the cable and not the pin.) Connect it to one of the output pins. Take the measurement probe and measure how the voltage increases as you move further away from the reference point.

Plot a voltage vs. distance graph. Draw a best-fit line.

- Do the points fall on a reasonably straight line?
- How do you interpret the slope and intercept of this line?
- Is the field between the two poles homogeneous?

4 Applications

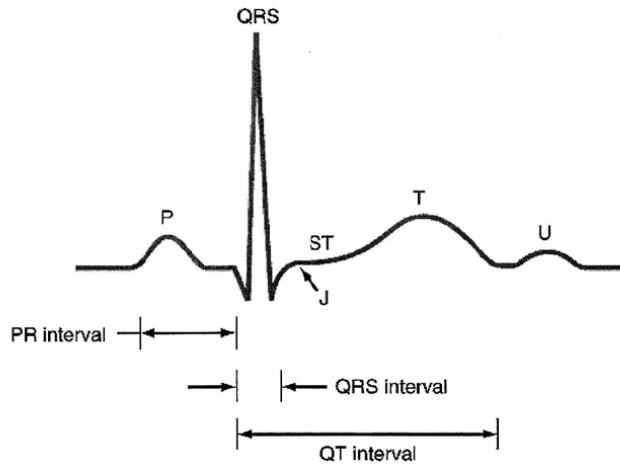


Figure 5: Sample Voltage Pattern for ECG

The canonical example in medicine for measuring voltage potentials and displaying them (with an oscilloscope) are ECG (Electrocardiography) and EEG (Electroencephalography). The basic idea of these two standard devices is fairly simple: by measuring the potential difference in time between two points (or usually several pairs of points) one gets a nice direct insight into the activity of the heart (brain) and can therefore detect dysfunctions easily. Let us, for example, take a closer look at how an ECG

works: If you record the electrical potential between two points along your chest you can record the voltage pattern shown in figure 5 (in time).¹

As you can see, the pattern observed can be split into different parts: At first the atrial cells are depolarized, giving the first signal (P wave). (This signal is relatively weak due to the small mass of the atrium.) After a delay one gets the QRS complex, which indicates the depolarization (wave) of the ventricles. After another delay one gets the T wave, which comes from the repolarization of the heart.² What do all these results, obtained from the heart as a total, mean in terms of processes going on in the single cells? The next figure 6 shows the measurement by your ECG in comparison to the potential in a single cell in a ventricle (obtained using a different method). First we recall that the interior and the exterior of the cell in their standard state have different ion concentrations and are therefore at different electrical potentials. (That is the zero line in the graph.) In the depolarization phase the ion channels in the cell membrane open and a flow of K^+ ions changes the potential inside the cell very rapidly. After this rapid change in potential one gets a plateau, which is mainly due to the inflow of Ca^{++} ions into the cell. (The Ca^{++} ions are much bigger than the K^+ ions, because they have a much larger hydrogen cover surrounding them and they diffuse much slower through the ion channels.) Finally in the repolarization phase the ion channels close again and the ion pumps in the cell membrane reestablish the initial ion concentrations.

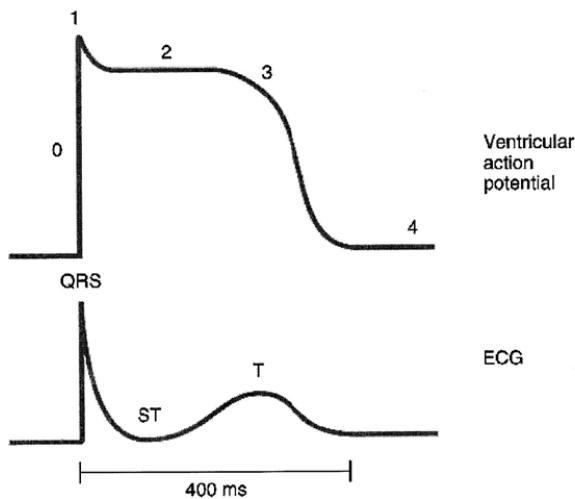


Figure 6: Comparison with Ventricular Potential

How can you now take advantage of this method for a diagnosis? First of all you

¹Remember: The ECG only shows the electrical polarization of the heart muscle. It does not show the contraction of the heart muscle!

²The interpretation of the U wave is still not yet 100% understood

don't want to take the measurements only along one single axis or plane, since e.g. if an infarction occurs on the front or back wall of your heart you are probably going to miss it. These days the usual way to get a 3-D picture of the position of the heart axis³ is obtained by measuring with multiple channels simultaneously between the points indicated in figure 7. (The contacts with an R are placed on the patients back.)

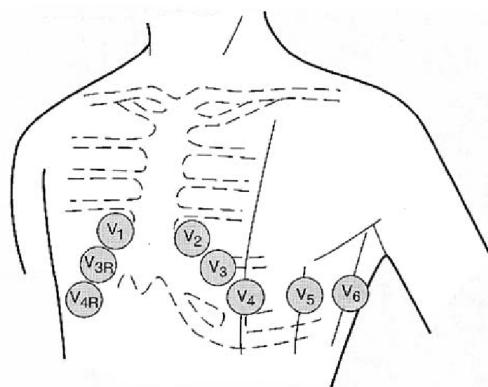


Figure 7: Channels of Measurement

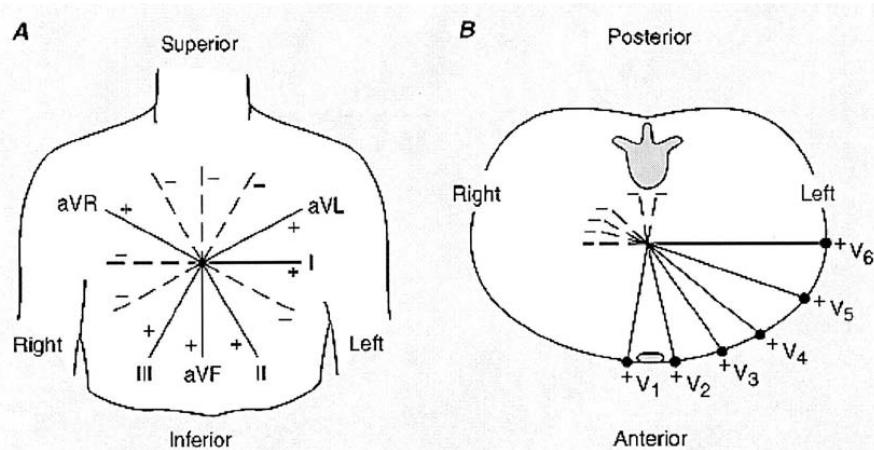


Figure 8: Electric Potentials and Information

From the location and amplitude of the main vector (axis) one can see for example if the muscle mass is increased on one side of the heart (hypertrophy). In that case the

³The heart axis is a simplified concept of the locations of the electrical potentials in the heart. One can think of the heart axis as a vector symbolizing the (physical) axis of the heart.

main vector is tilted.⁴ Another thing to look for is if the depolarization and repolarization was performed properly. For example if the ion channels in a certain region of the heart are destroyed by an infarction, then the electric potential between the depolarization and repolarization phase does not reach the zero level. By looking at your data you can not only locate the infarction, but also read off additional information, e.g. if the infarction cilled the tissue through all of the heart wall or only parts of it. (That determines your treatment of the patient!)

You can see that you can get a lot of useful information if you look at electrical potentials (and their change in time), as shown in figure 8.

Textbook references:

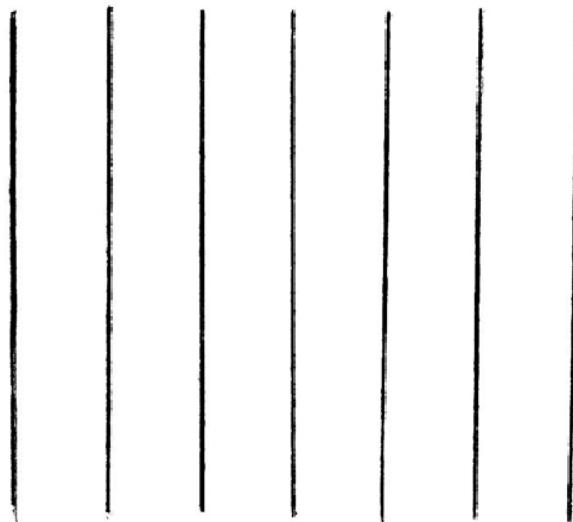
- Stein: *Internal Medicine*
- Harrison's *Principles of Internal Medicine*

⁴Also in pregnant women the heart as a total is slightly repositioned and therefore the electrical axis is a little bit off.

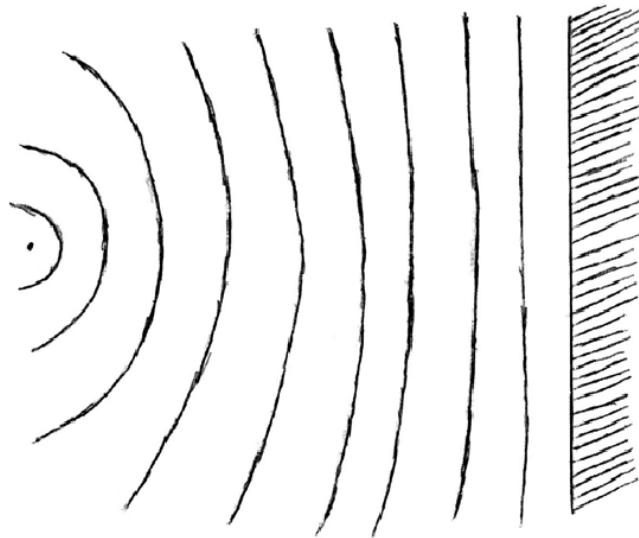
5 Lab Preparation Examples

Field Lines: Given the following equipotential lines, draw 5-10 field lines on each diagram.

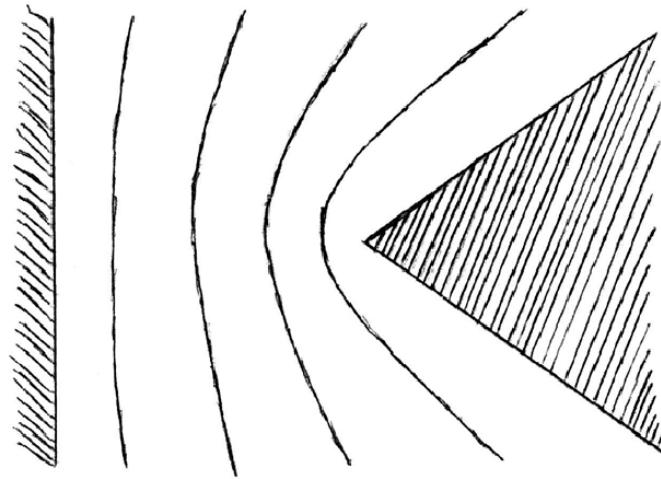
1.



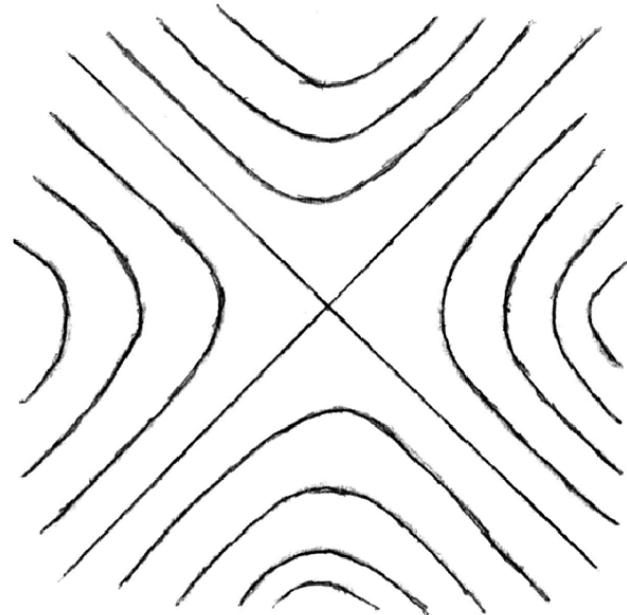
2.



3.



4.

Shield Theorem

5. Explain in a few sentences why you might be safe inside your new Volkswagen Beetle even when struck by a bolt of lightning.

Experiment 2-2

DC Circuits

1 Introduction

The purposes of this lab will be:

1. to investigate the three variables involved in a mathematical relationship known as Ohms Law;
2. to experiment with the variables that contribute to the operation of an electrical circuits; and
3. to experimentally demonstrate Kirchhoffs Rules for electrical circuits.

2 Procedure

2.1 Ohm's Law

1. Choose one of the resistors that you have been given. Using the chart below, decode the resistance value and record that value in the first column of Table 2.1.
2. **MEASURING CURRENT:** Construct the circuit shown in PUT FIGURE HERE by pressing the leads of the resistor into two of the springs in the Experimental Section on the Circuits Experiment Board.
3. Set the multimeter to the 200 mA range, noting any special connections needed for measuring current. Connect the circuit and read the current that is flowing through the resistor. Record this value in the second column of Table 2.1.
4. Remove the resistor and choose another. Record its resistance value in Table 2.1 then measure and record the current as in Steps 2 and . Continue this process until you have completed all of the resistors you have been given. As you have more than one resistor with the same value, keep them in order as you will use them again in the next steps. (Note: Always disconnect the circuit whenever you are done with your measurements.)
5. **MEASURING VOLTAGE:** Disconnect the multimeter and connect a wire from the positive lead (spring) of the battery directly to the first resistor you used as shown in FIGURE HERE. Change the multimeter to the 2V DC scale and connect the leads as shown also in FIGURE HERE. Measure the voltage across the resistor and record it in Table 2.1.

Resistance (Ω)	Current (A)	Voltage (V)	Voltage/Resistance ($V/\Omega = A$)

Table 2.1: Table for data obtained in Section 2.1.

- Remove the resistor and choose the next one you used. Record its voltage in Table 2.1 as in Step 5. Continue this process until you have completed all of the five resistors.

2.2 Resistances in Circuits

- Choose three resistors of the same value (330Ω or 1000Ω). Enter those sets of colors in Table 2.2 on the next page. We will refer to one as #1, another as #2 and the third as #3.
- Determine the coded value of your resistors. Enter the value in the column labeled Coded Resistance in Table 2.2. Enter the Tolerance value as indicated by the color of the fourth band under Tolerance.
- Use the multimeter to measure the resistance of each of your three resistors. Enter these values in Table 2.2.
- Determine the percentage experimental error of each resistance value and enter it in the appropriate column.

$$\text{ExperimentalError} = \frac{|\text{Measured} - \text{Coded}|}{\text{Coded}} \times 100\% \quad (1)$$

- Now connect the three resistors into the SERIES CIRCUIT, FIGURE HERE using the spring clips on the Circuits Experiment Board to hold the leads of the resistors together without bending them. Measure the resistances of the combinations as indicated on the diagram by connecting the leads of the multimeter between the points at the ends of the arrows. Use the $20K$ range if using 1000Ω and $2K$ if using 330Ω .
- Construct a PARALLEL CIRCUIT below, first using combinations of two of the resistors, and then using all three. Measure and record your values for these circuits.

	Colors				Coded	Measured	% Error	Tolerance
	1st	2nd	3rd	4th	Resistance (Ω)	Resistance (Ω)		
#1								
#2								
#3								

Table 2.2: Table for data obtained in Section 2.2.

Experiment 2-3

Capacitance

1 Introduction

A capacitor is a device for storing electric charge and energy. For simplicity, an ideal capacitor can be considered as a pair of parallel conducting metal plates, as shown in figure 1. When a charge $+Q$ is placed on the upper plate and $-Q$ on the lower plate, a potential difference V is established between the plates, and the quantities Q and V are related by the expression:

$$Q = CV \quad (1)$$

where the capacitance C is determined by the size and separation of the plates.

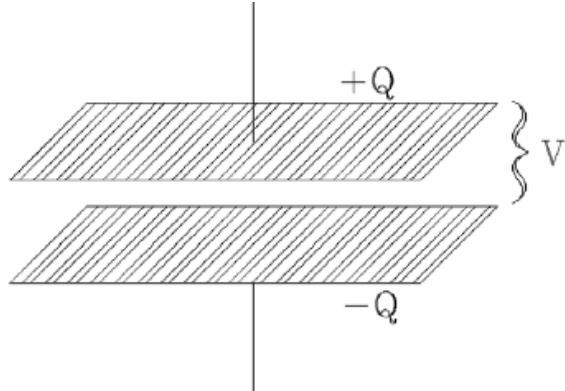


Figure 1: An Ideal Capacitor

1.1 Discharging a Capacitor

If a wire is attached to the upper plate and then touched to the lower plate, charge will flow through the wire until the charge q and potential difference V are zero. (We will use Q to indicate the initial charge and q for the time dependent charge.) This process, known as discharging the capacitor, does not occur instantaneously. Instead, as the charge flows from one plate to the other, the potential difference decreases, and the current (or rate of change of charge) gradually diminishes. Since the rate of change is proportional to the amount of charge remaining, the current is initially large and then *decays exponentially* with increasing time. Exponential growth or decay occurs whenever the rate of change of a quantity is proportional to the quantity present at that time. Exponentials are also characteristic of the growth rate of cells or organisms and the decay of radioactive isotopes.

The changing current which discharges a capacitor can be calculated as a function of time by considering the circuit shown in figure 2. The capacitor C is initially charged by a battery of emf ε , placing a charge $Q = C\varepsilon$ on the plates. When the switch is closed, a current I flows in the circuit, and according to Ohm's Law the voltage drop across the resistance is IR .

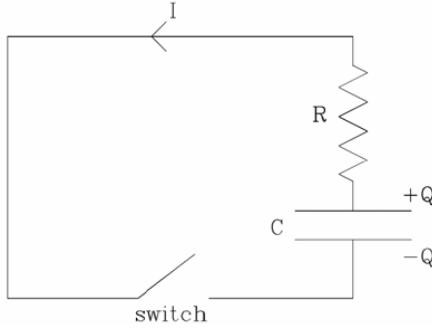


Figure 2: Discharging Circuit of a Capacitor

The resistance may be a separate circuit element or merely the resistance of the wire itself. Since the sum of the voltage drops around the circuit must be zero, we obtain the equation:

$$IR + \frac{Q}{C} = 0 \quad (2)$$

The equation can be rewritten in terms of charge alone from the definition of current, $I = dQ/dt$ ($= \lim_{\Delta t \rightarrow 0} \frac{\Delta Q}{\Delta t}$),

$$R \frac{dQ}{dt} + \frac{Q}{C} = 0 \quad (3)$$

Equation (3) is a differential equation for Q and has a solution which gives the time dependence of the charge

$$Q(t) = C\varepsilon e^{-t/RC} \quad (4)$$

Students familiar with calculus can verify this result by differentiating equation (4) and combining $Q(t)$ and dQ/dt to show that equation (3) is satisfied. From the definition of current, $I = dQ/dt$, we find that

$$I(t) = -\frac{\varepsilon}{R} e^{-t/RC} \quad (5)$$

Note that at $t = 0$, when the switch is just closed, $I = -\varepsilon/R$. As time increases, the current $I(t)$ gets smaller and reaches zero as time goes to infinity.

1.2 Charging a Capacitor

The process of charging an uncharged capacitor has many similarities with the process of discharging described above. In this case, a battery with an emf of ε volts is connected in series with a resistance R and the capacitance C , as indicated in figure 3.

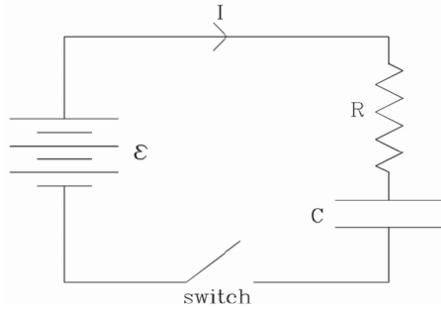


Figure 3: Charging Circuit of a Capacitor

When the switch is first closed, there is no charge on the capacitor and thus no voltage across it, and the full voltage ε appears across R , causing maximum current I to flow. As the capacitor becomes charged, the voltage-drop IR (and thus the current I) gradually decreases. I approaches zero as the capacitor is charged toward its maximum potential difference of $\varepsilon = Q/C$.

The equation which indicates the sum of voltage drops around the circuit loop at any time during the charging is:

$$\varepsilon = IR + \frac{Q}{C} \quad (6)$$

Again, this can be converted to a *differential equation*, the solution of which is:

$$I(t) = \frac{\varepsilon}{R} e^{-t/RC} \quad (7)$$

Comparing equation (5) with (7), we note that the current for discharging a capacitor from a given potential ε decreases in time identically to the current for charging the capacitor through the same resistance to the same final voltage ε . The difference in the sign of I indicates, of course, that the two currents flow in opposite directions.

1.3 Graphical Presentation of Exponential Decay

In equations (5) and (7), the symbol e stands for the constant $2.71828\cdots$, the *base of natural logarithms*. A graph of the function $y = ae^{-x/b}$ is given in figure 4, where $y = a$ at $x = 0$ and then decays exponentially as x increases, reaching a value of $y = ae^{-1} = 0.37a$ at $x = b$.

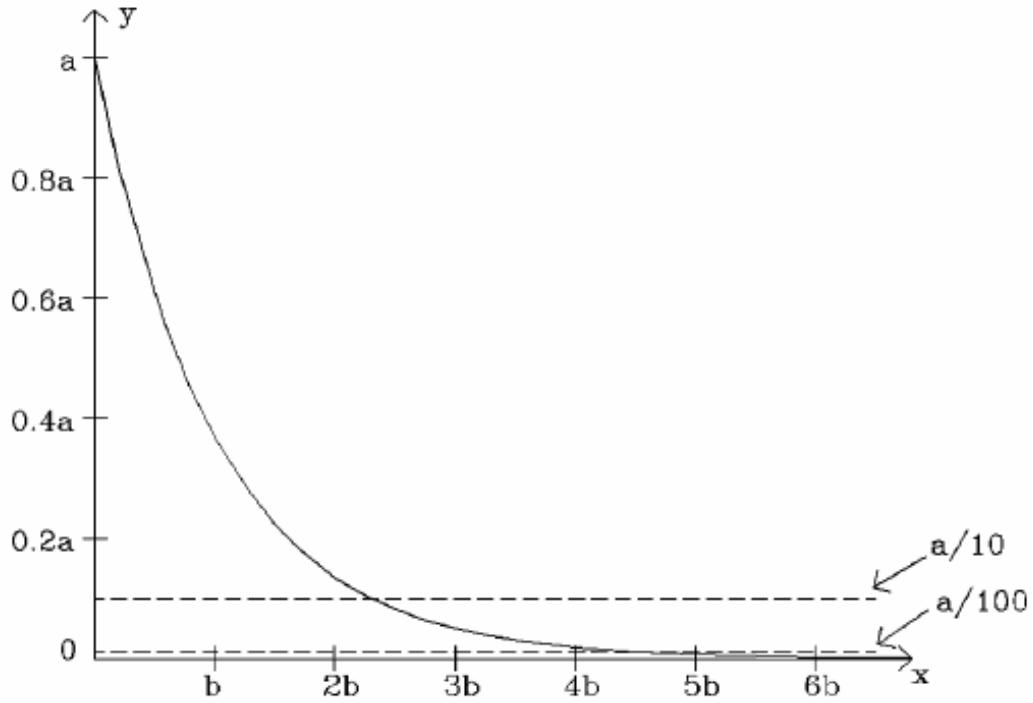


Figure 4: Graph of $y = ae^{-x/b}$

Although it is possible to determine the values of a and b from experimental results plotted on such a curve, it is more informative to plot $\log y$ versus x , or to use *semi-log* graph paper, where the horizontal lines and scales of the y -axis are spaced in equal increments of $\log y$. Then, taking the logarithm (base 10) of $y = ae^{-x/b}$ yields:

$$\log y = \log a - \frac{x}{b} \log e = \log a - \left(\frac{\log e}{b} \right) x = \log a - \left(\frac{0.434}{b} \right) x \quad (8)$$

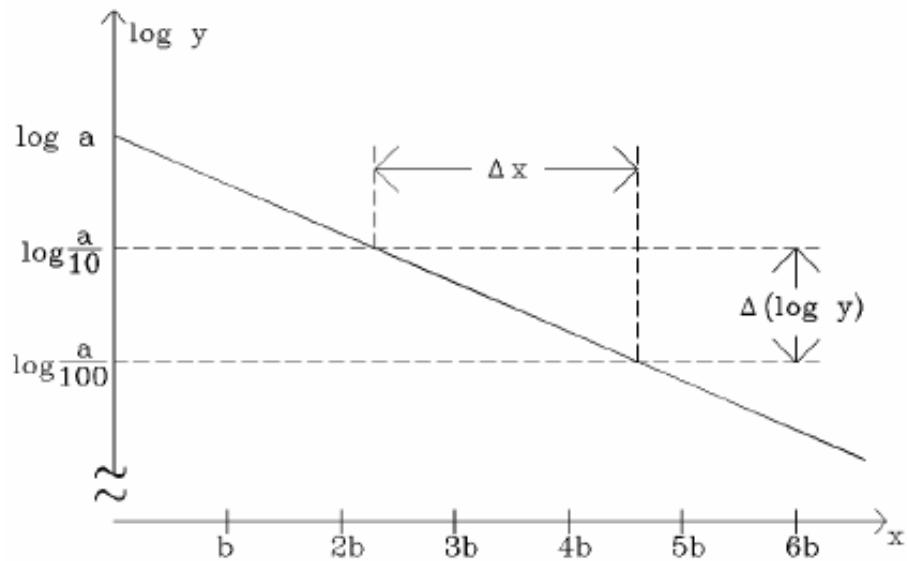
which is a linear equation in x representing a straight line with slope $-0.434/b$, if $\log y$ is plotted against x as in figure 5.

A simple way to determine b from a straight line drawn through the experimental points on such a plot is to note the values of x at which the line crosses values of y differing by a factor of 10. Then, using the values shown in figure 5,

$$\Delta(\log y) = \log \frac{a}{100} - \log \frac{a}{10} = (\log a - \log 100) - (\log a - \log 10) = -2 + 1 = -1 \quad (9)$$

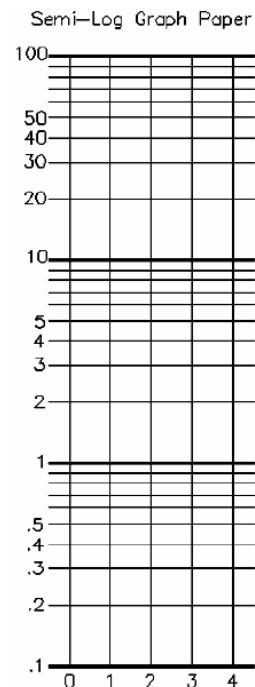
Since from equation (8) the slope $\Delta(\log y)/\Delta x$ equals $-0.434/b$, therefore

$$\frac{-1}{\Delta x} = \frac{-0.434}{b}, \quad \text{or} \quad b = 0.434 \Delta x \quad (10)$$

Figure 5: Plot of $\log y$ versus x

In the present experiment, y represents the current I ; x is the time t ; a is the initial value of I ; and b is the time constant RC .

Note: You may produce a graph like the one in figure 5 by either: 1) Plotting $\log y$ on the y axis of normal graph paper; or 2) Plotting y on semi-log graph paper (shown in the right). Semi-log graph paper (provided in the lab) has the y axis already logarithmically compressed, so you do not need to calculate $\log y$ for each point – you can simply plot the y values themselves. With semi-log graphs you can verify quickly whether the data behaves exponentially or not; $y = 1/x$, for example, will not produce a straight line on a semi-log plot. Note also that on semi-log graphs y approaches 0 only exponentially – there is no $y = 0$.



2 Procedure

2.1 Large RC-charging

Figure 6 shows the circuit to be wired for measuring the current I as a function of time for charging a capacitor C through a resistance R to a voltage ε . For this part of the lab you will use the bank of three large capacitors which are attached to a piece of wood with three switches and the label “30 MFd”. (In this case the prefix “M” in MFd indicates micro, μ , or 10^{-6} .) Be sure to use a low-voltage (15 volts max.) power supply and to connect the positive terminal on the power supply to the positive terminal on the microammeter. Before closing the switch to the power supply and starting each new measurement, momentarily connect a low resistance across the capacitor in order to start with no charge on the capacitor plates.

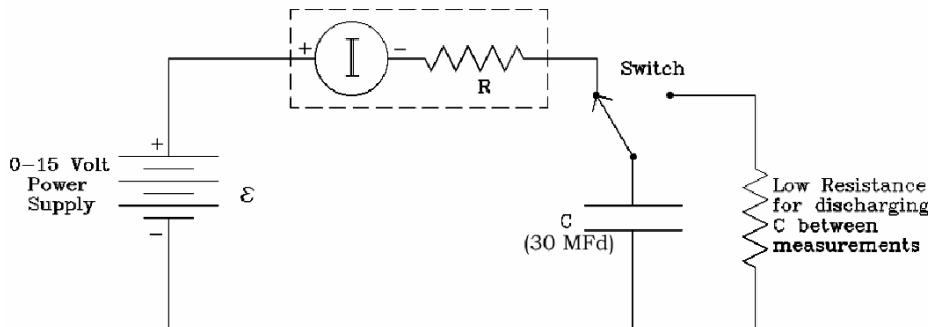


Figure 6: Large RC-charging Circuit

Measure I at a series of time intervals after closing the switch. Plot the results on semi-log paper and determine the RC time constant.

Repeat the above for a different value of C . (Note that the “capacitor” supplied is in fact a bank of capacitors with the provision for varying C by switching in different numbers of the capacitors in parallel).

2.2 Large RC-discharging

Use the circuit shown in figure 7 to measure the discharge current through the resistance R of the same capacitors C used in Part 1. You will use the same equipment as you used in Part 1 of the lab.

Before closing the switch to begin the measurement, momentarily connect the power supply directly across the capacitor in order to give it an initial charge. Then disconnect

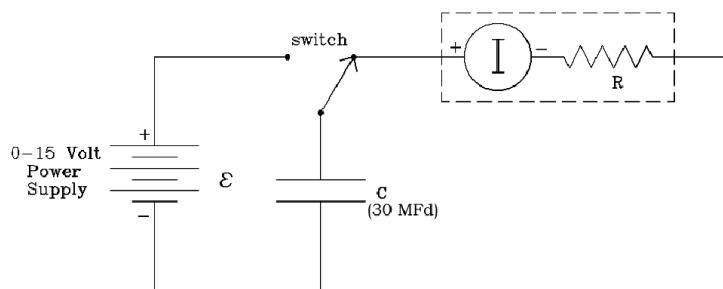


Figure 7: Large RC-discharging Circuit

the power supply, close the switch, and measure I as a function of time. Again, plot the results on semi-log paper and determine RC . Compare with the results found in Part 1 for the same values of C .

2.3 A Relaxation Oscillator

A neon bulb has the property of having a very high resistance (almost infinite) until the voltage applied to it is high enough to “break down” the gas, at which point the bulb lights and its resistance becomes very low. In this part of the lab you will use a high voltage power supply. Thus, *it is essential that you are sure to use the small capacitor with a value of $0.082 \mu\text{F}$* rather than the bank of capacitors you use in the first two parts of the lab. Wire a neon bulb in parallel with the capacitor in the circuit of figure 8.

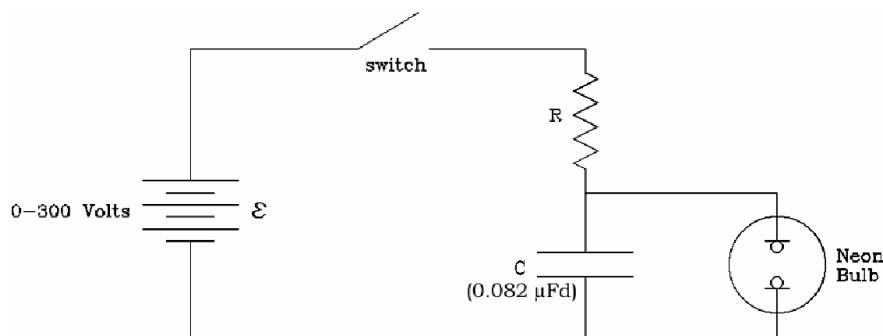


Figure 8: Relaxation Oscillator Circuit

CAUTION: Since a relatively high-voltage power supply is necessary for this part of the experiment, it is important not to touch any exposed metal parts of the circuit while the power supply is connected to the circuit and turned on, whether or not the

switch is closed. Since the capacitor stores charge, *it may be charged even if the voltage supply has been removed*. Be sure to discharge the capacitor fully by simultaneously touching an insulated wire to each end of the capacitor before touching any of the metal parts of the circuit.

When the switch is closed, the capacitor will start to be charged as in Part 1, with the time constant RC , and the high resistance of the neon bulb will have negligible effect on the circuit. However, when the capacitor is charged to sufficiently high voltage, the neon bulb will light. The capacitor will then be quickly discharged through the bulb. If R is sufficiently large, there will not be sufficient current to keep the bulb lit after the capacitor is discharged. The bulb will then be extinguished; it will return to a state of high resistance, and the charging process will start again. For a given applied voltage ε , and a neon bulb with a given breakdown voltage, the period of this repetitive “oscillation” will thus be determined by the value of RC .

Measure the period for different combinations of R and C , and verify its dependence on the product RC .

2.4 Demonstration of Short RC Time (to be set up by Lab Instructor)

If the RC time constant of a circuit is too short to be detected using an ammeter, the effect can be displayed on an oscilloscope as illustrated in figure 9. In this case, it is convenient to use a square-wave generator, which switches a voltage on and off at a variable repetition rate.

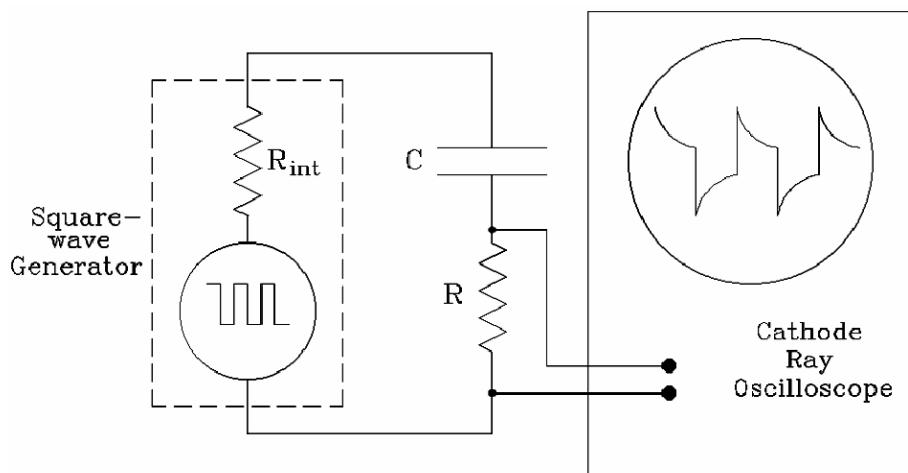


Figure 9: Using an Oscilloscope to Study the Time-dependence of Voltage

Use the smallest values of R and C available. From observation of the oscilloscope trace, plot I versus time on semi-log graph paper and determine the value of RC . Compare the result with the value of RC calculated from the labels on the circuit elements. (Remember that the total R of the circuit includes any internal resistance in the square-wave generator).

Experiment 2-4

The Magnetic Field

OBJECT: To determine the strength of the magnetic field in the gap of an electromagnet by measuring: (I) the force on a current-carrying rod, (II) the effect of the change of magnetic flux through a coil when the coil is inserted into or removed from the region containing the magnetic field.

CAUTION: (1) Always reduce the current through the electromagnet to zero before opening the circuit of the magnet coils. (2) Remove wrist watches before placing hands near the magnet gaps.

1 Part I - Force on a Current-carrying Wire

1.1 Physical Principles

Consider a rigid length of wire rod L , held horizontally and normal to the direction of a uniform, horizontal magnetic field B . If a current i is passed through the wire, as indicated in Figure 1, then there will be a vertical force $F = iLB$ on the wire rod.

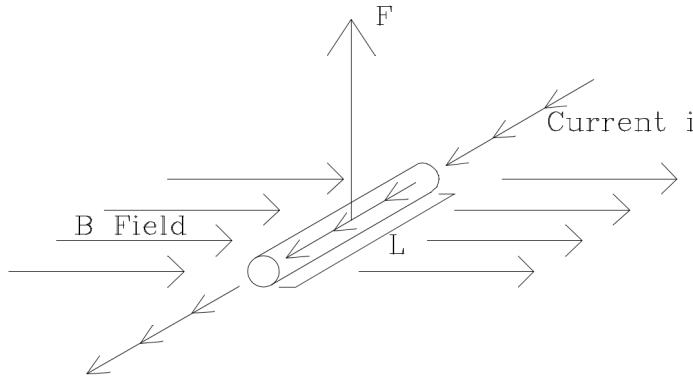


Figure 1: Force on a Current-carrying Wire

1.2 Experimental Apparatus

The experimental set-up is shown in Figure 2. The horizontal magnetic field B is produced in the air gap of a “C-shaped” iron electromagnet. The strength of B is determined by I , the current in the magnet coils, which is supplied by an adjustable low-voltage power supply.

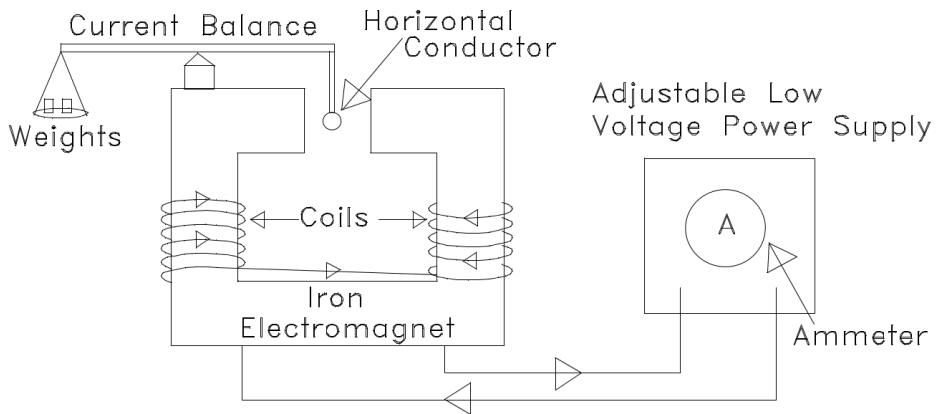


Figure 2: Experimental Set-up

A more detailed drawing of the balance and electro-magnet arrangement is shown in Figure 3.

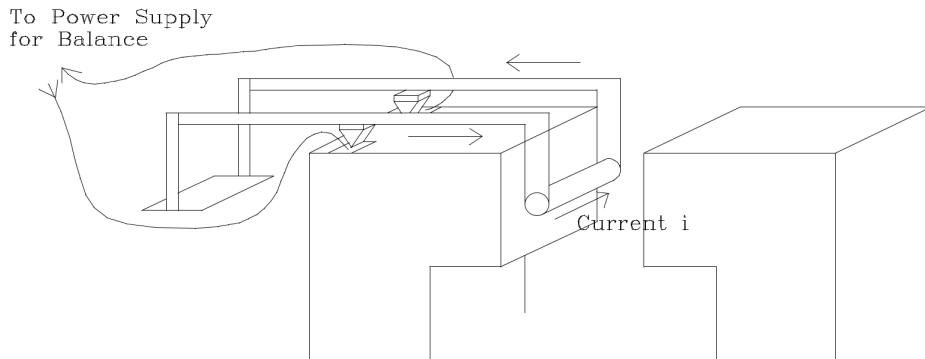


Figure 3: Setup of the Wire and Balance for Force Measurement

Note that the current i through the horizontal conductor is supplied by a separate power supply. The current balance is constructed out of conducting and insulating materials such that current can enter through one side of the knife edge, flow through one side of the balance arm to the horizontal conductor, and then flow back through the other side of the balance arm and out through the other knife edge.

The current i in the balance is provided by the HP E3610A power supply, which can operate either in constant voltage or constant current mode. The voltage dial sets the *maximum voltage* the device will supply to the circuit; the current dial likewise sets the *maximum current*, and if the circuit tries to draw more current, the power supply will reduce the voltage until it reaches whatever value is needed to maintain the maximum current (by $V = IR$).

To use the power supply in constant current mode, begin with the current dial

turned all the way down (counter-clockwise) and the voltage dial turned all the way up (clockwise). Set the range to 3 Amps, and connect the leads to the + and – terminals. You can now set the current to the desired level—note that the digital meters on the power supply show the *actual* voltage and current being supplied, so you will not normally see any current unless the leads are connected to a complete circuit. (If you want to set the current level without closing the circuit, you can hold in the CC Set button while you turn the current dial.) Once you close the circuit, the voltage adjusts automatically to maintain the constant current level, and the CC (Constant Current) indicator light should be on.

1.3 Procedure

Set the current I through the electromagnet at 5 amperes. Then determine the weights needed to bring the balance to equilibrium for several different values of the current i through the balance. (Note that it is easier to make the final adjustment on i once weights have been selected for the approximate current value.) Present the measurements of force versus balance current graphically, and from the graph find a value for B . Repeat the above for magnet currents of 4 amperes, 3 amperes, and 2 amperes. Determine B in each case. All your results may be plotted on the same graph. Draw a diagram similar to Figure 3, and indicate the direction of i , F , and B for your set-up.

2 Part II - EMF Induced in a Moving Coil

2.1 Physical Principles

A more accurate method of measuring B is to use a **charge integrating circuit** to determine the total charge which flows through the circuit of a small search coil when it is suddenly inserted into (or removed from) the region of magnetic field. The EMF ε induced across the coil by a change of magnetic flux $\Delta\Phi$ through the coil is given by Faraday's Law:

$$\varepsilon = N \frac{\Delta\Phi}{\Delta t} \quad (1)$$

where N is the number of turns in the search coil.

If the total resistance of the coil and its circuit is R , then the resulting instantaneous current $i = \Delta Q / \Delta t$ flowing through the circuit will be

$$\frac{\Delta Q}{\Delta t} = i = \frac{\varepsilon}{R} = \frac{N}{R} \frac{\Delta\Phi}{\Delta t}. \quad (2)$$

The instantaneous induced EMF and current will thus depend on the time Δt , i.e. the speed with which the search coil is moved into the field. However, the total charge

which flows during the motion is independent of Δt :

$$\Delta Q = \frac{N}{R} \Delta \Phi. \quad (3)$$

If the search coil is started in a field-free region and is plunged into the magnet gap so that B is normal to the plane of the coil, then $\Delta \Phi = A_{\text{coil}}B$, where A_{coil} is the area of the coil, and therefore:

$$\Delta Q = \frac{N}{R} A_{\text{coil}} B. \quad (4)$$

Note that both the inner and outer diameter of the search coil should be measured, since the innermost windings have a different area than the outermost ones.

2.2 Experimental Apparatus

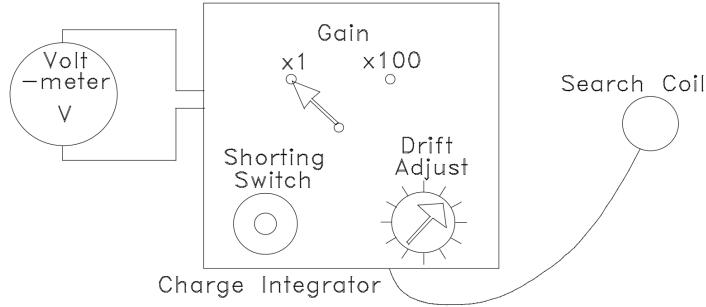


Figure 4: The Magnetic Field Module

A charge integrator (the Magnetic Field Module shown in Figure 4) is used to measure the ΔQ produced by the EMF induced in the search coil. A capacitor in the module stores the charge ΔQ , and the voltage across this capacitor (read on the external voltmeter shown) is proportional to ΔQ . Therefore

$$V = K \Delta Q = K \frac{N}{R} \Delta \Phi \quad (5)$$

where K is a constant that depends on the capacitance and gain of the integrator circuit. Instead of trying to calculate a value of KN/R in terms of the components, it is more direct to calibrate the combination of the search coil and the integrator circuit by measuring V for a known $\Delta \Phi$. This known magnetic flux can be created by passing a measured current I_{sol} through a long air-core solenoid of n turns per meter and with a cross-sectional area of A_{sol} so that

$$\Delta \Phi_{\text{sol}} = B_{\text{sol}} A_{\text{sol}} = \mu_0 n I_{\text{sol}} A_{\text{sol}} \quad (6)$$

where μ_0 is the permeability of free space ($\mu_0 = 4\pi \times 10^{-7} \text{ T} \cdot \text{m/A}$).

2.3 Procedure

Connect the apparatus as shown in Figure 4. Turn on the power supply. Before any measurements are made, depress the shorting switch, then release the shorting switch and turn the drift adjust control to minimize the drift in the output voltage as observed on your meter. The shorting switch must be used to discharge the integrating capacitor prior to each measurement. Also, the drift adjust setting should be checked occasionally. If the gain setting on the Magnetic Field Module is changed, the drift adjust control must be reset.

Magnetic field measurements are made by inserting or removing the search coil from the region containing the field to a field-free region or by leaving the search coil stationary and turning the field on or off.

Since the magnetic field in the large iron-core electromagnet is much greater than the field in an air-core solenoid, the Magnetic Field Module was designed with two gain settings. In the gain=100 position, the module is 100 times as sensitive as in the gain=1 position. For measuring fields generated by the air-core solenoid, set the gain at 100. For fields generated by the large electromagnet, set the gain at 1.

Connect the solenoid to the + and – terminals of the HP Power Supply. A three position (on-off-on) reversing switch is part of the solenoid circuit. Flip the switch to one of the on positions and adjust the power supply current such that two amperes is flowing through the solenoid.

Slide the search coil over the solenoid and, while holding the search coil at the center of the solenoid, discharge the Magnetic Field Module and then turn the current through the solenoid off using the three position switch (or turn the current from off to on). Take several readings and record the voltage on the integrator and the current through the solenoid. Then the result is

$$V_{\text{sol}} = 100 \cdot K \cdot \frac{N}{R} \cdot \Delta\Phi_{\text{sol}} = 100 \cdot K \cdot \frac{N}{R} (\mu_0 n I_{\text{sol}}) A_{\text{sol}} \quad (7)$$

Now set the gain to 1 on the Magnetic Field Module, and readjust the drift controls. Set the current for the large electromagnet to 5 amps, and use the search coil to measure the resulting B . Move the search coil *gently and smoothly* into the region of the magnetic field (do not move the coil hastily as you may damage it by striking against the magnet itself). Record V_{mag} . The corresponding equation is:

$$V_{\text{mag}} = K \cdot \frac{N}{R} A_{\text{coil}} B_{\text{mag}} \quad (8)$$

Combine the results of Eq(7) and Eq(8) to determine the value of B for the large electromagnet when the current is 5 amps, and then do the same for the other magnet currents of 4, 3, and 2 amps. Compare these values for B with those obtained in Part I.

Experiment 2-5

e/m of The Electron

1 General Discussion

The “discovery” of the electron by J. J. Thomson in 1897 refers to the experiment in which it was shown that “cathode rays” behave as beams of particles, all of which have the same ratio of charge to mass, e/m . Since that time, a number of methods have been devised for using electric and magnetic fields to make a precise measurement of e/m for the electron. When combined with the value of the electron’s charge, which is measured in the Millikan Oil Drop Experiment, the determination of e/m leads to an accurate value of the mass of the electron. In the present experiment, electrons are emitted at a very low velocity from a heated filament, then accelerated through an electrical potential V to a final velocity v , and finally bent in a circular path of radius r in a magnetic field B . The entire process takes place in a sealed glass tube in which the path of the electrons can be directly observed. During its manufacture, the tube was evacuated, and a small amount of mercury was introduced before the tube was sealed off. As a result, there is mercury vapor in the tube. When electrons in the beam have sufficiently high kinetic energy (10.4 eV or more), a small fraction of them will collide with and ionize mercury atoms in the vapor. Recombination of the mercury ions, accompanied by the emission of characteristic blue light, then occurs very near the point where the ionization took place. As a result, the path of the electron beam is visible to the naked eye.

The tube is set up so that the beam of electrons travels perpendicular to a uniform magnetic field B . B is proportional to the current I through a pair of large diameter coils (so-called “Helmholtz Coils”) in which the coil separation is selected to produce optimum field uniformity near the center.

2 Experimental Apparatus

2.1 The Sealed Glass Tube

Figure 1 shows the filament surrounded by a small cylindrical plate. The filament is heated by passing a current directly through it. A variable positive potential difference of up to 100 volts is applied between the plate and the filament in order to accelerate the electrons emitted from the filament. Some of the accelerated electrons come out as a narrow beam through a slit in the side of the cylinder. The entire tube is located inside a set of coils, which produce a uniform magnetic field B perpendicular to the electron beam. The magnitude of the field can be adjusted until the resultant circular

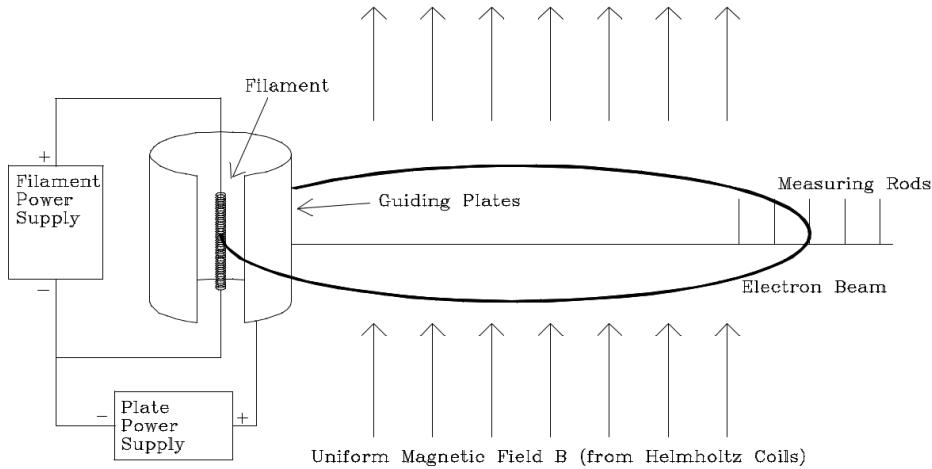


Figure 1: Diagram of the Interior of the Sealed Glass Tube

path of the electron beam just reaches one of the measuring rods. These rods are located along a cross bar, which extends from the cylinder in a direction perpendicular to that in which the electron beam was emitted—i.e., along a diameter of the circular orbits.

2.2 The Helmholtz Coils and Uniform Magnetic Field

The magnetic field produced at the position of the electron beam by a current I flowing through the coils must be computed. For a single turn of wire of radius R , the field on the axis at a distance x from the plane of the loop is given by:

$$B' = \frac{\mu_0 R^2 I}{2(R^2 + x^2)^{3/2}}. \quad (1)$$

For the arrangement in Figure 2, there are two loops with N turns each, separated by a distance equal to the coil radius R . The coils contribute equally to the field at the center:

$$B_I = \frac{\mu_0 R^2 N I}{[R^2 + (R/2)^2]^{3/2}} = \frac{4\pi \times 10^{-7} N I}{R(1 + \frac{1}{4})^{3/2}} = \text{constant} \times I \text{ Tesla} \quad (2)$$

where $N = 72$ is the number of turns of each coil and $R = 33 \text{ cm}$ is the radius of the coils used.

This arrangement, called a pair of **Helmholtz coils**, yields a remarkably uniform field in the region at the center.

The net field B in which the electrons move is not B_I alone, but the resultant of the earth's field B_e and B_I . If the equipment is oriented so that the field of the Helmholtz

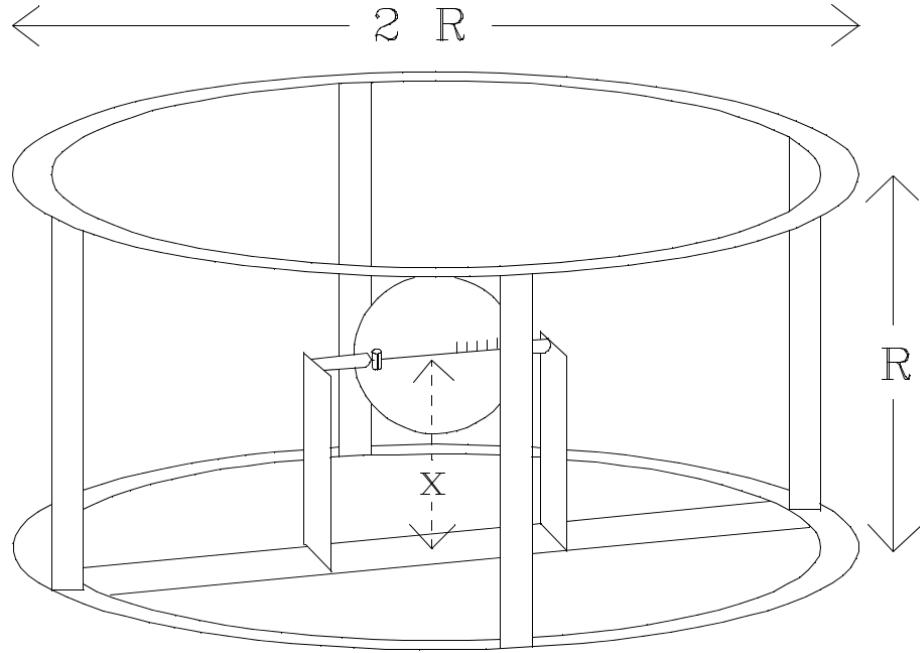


Figure 2: Helmholtz coils used to produce a uniform magnetic field.

coils is parallel to that of the earth, and if the current through the coils causes B_I to be directed opposite to B_e , then

$$B = B_I - B_e \quad (3)$$

2.3 The Trajectories of the Electrons in the Glass Tube

If an electron of charge e and mass m starts nearly from rest and is accelerated through a potential difference V to a final velocity v , then

$$\frac{1}{2}mv^2 = eV \quad \text{or} \quad \frac{e}{m} = \frac{v^2}{2V} \quad (4)$$

If the electron then enters a uniform magnetic field B which is perpendicular to its velocity, it will move in a circular orbit of radius r , where

$$\frac{mv^2}{r} = evB \quad \text{or} \quad \frac{e}{m} = \frac{v}{Br} \quad (5)$$

If it were possible to measure the velocity directly, then e/m could be determined by measurements of either the electric or magnetic field alone. Since a direct measurement of v is not feasible in this experiment, e/m can be determined from the combination of electric and magnetic fields used. Specifically, by eliminating v from equations (4) and (5), e/m can be expressed directly in terms of V , B , and r .

Instead of determining e/m from a single measurement of r for given values of V and B , however, it is preferable to measure the variation of r with B (or I) at fixed values of V . In particular, the data can be presented in convenient form by plotting the **curvature** $1/r$ as a linear function of I .

$$\frac{1}{r} = \sqrt{\frac{e}{m}} \frac{1}{2V} B_I - \sqrt{\frac{e}{m}} \frac{1}{2V} B_e \quad (6)$$

Derive (6) from (3), (4), and (5) for your report before coming to laboratory.

Note that equations (4), (5), and (6) apply strictly only to electrons with trajectories on the *outside edge* of the beam – i.e., the most energetic electrons. There are two reasons why some electrons in the beam will have less energy:

1. There is a small potential difference across the filament caused by the heating current. Only electrons leaving the negative end of the filament are accelerated through the whole potential difference V .
2. Some of the electrons in the beam will lose energy through collisions with mercury atoms.

3 Procedure

3.1 Orientation of the Coil and Tube Setup

For reasons already explained, we would like to orient the Helmholtz coils such that their axes are parallel to the ambient magnetic field.

Please exercise extreme care in the following section as you align the Helmholtz coils: the cathode ray tube is very delicate and may break if the coil support is not very firmly secured and falls. It is recommended that two people handle the coil frame at all times while the coils are being aligned. Also take care not to touch the dip needle, which is easily bent.

In order to align the coils so that their axis is aligned with the ambient magnetic field proceed as follows:

- With the coils in the horizontal position, rotate the horizontal arm of the dip needle support so that the needle itself and the plastic protractor are horizontal. Avoid touching the needle or plastic protractor, and instead turn the arm using the attached lever.
- Allow the compass needle to come to a rest. It is now pointing in the direction of the horizontal component of the ambient field.

- Rotate the entire frame by turning the base, until the compass needle is aligned along the 90° - 270° line on the protractor. The horizontal axis of the cathode ray tube (coaxial with the metal rod) is now aligned with the horizontal component of the ambient field.
- Rotate the horizontal arm of the dip needle support so that the needle itself and the plastic protractor are now in a vertical plane. Avoid touching the needle or plastic protractor, and instead turn the arm using the attached lever.
- Allow the needle to come to a rest. It is now pointing in the direction of the ambient field.
- Loosen the wingnut and **gently** raise one side of the coils. You want to increase the angle until the dip needle is aligned along the 0° - 180° line on the protractor. **Hold the wooden frame rather than the coils as you raise the setup.**
- Securely **tighten the wingnut so that the coils remain in position**. One person should be supporting the frame while a second person tightens the wingnut.
- The coil axis should now be aligned (coaxial) with the ambient magnetic field.

3.2 Preliminary Adjustments

The supplies and controls for the Helmholtz coils and the filament are permanently wired on a board and are designed to minimize the possibility of damage to the tube or coils. Locate each control, and note the qualitative effects observed when the control is varied. In particular:

- (a) Figure 1 shows that the filament and its associated lead wires form a small loop. Since a 4 amp current is required to heat the filament, this loop creates a measurable field. The filament coil reversing switch permits you to study the effect of this field. The effect can be minimized in the experiment by rotating the tube slightly in its mounting so that the electrons come out parallel to the plane of the coils.
- (b) Note the direction of the coil current for each position of its reversing switch by using the dip needle to check the direction of the resultant field. Knowing the field direction, check the sign of the charge of the particles in the deflected beam. Also, determine whether the earth's field adds to or subtracts from the coil field.
- (c) The beam will have a slight curvature in the earth's field when the coil current is zero. Make a preliminary measurement of the earth's field by adjusting the coil current to remove this curvature. The special Meter Switch and low current meter (200 mA) will enable you to measure the relatively small current needed, and the straight line trajectory can be checked by comparison with the light emitted from the filament.

3.3 Measurement of the Circular Orbits

With the accelerating voltage at an intermediate value, the current in the Helmholtz coils can be adjusted so that the outside edge of the beam strikes the outside edge of each bar in turn. Measure field current as a function of radius for the highest voltage V which allows you to adjust the beam with respect to all five bars.

For one measurement, test the reproducibility of the current setting as an aid to error analysis. The tube manufacturer supplies the following values for the *diameters* from the filament to the *outside* of each bar in succession:

$$6.48 \text{ cm}, \quad 7.75 \text{ cm}, \quad 9.02 \text{ cm}, \quad 10.30 \text{ cm}, \quad 11.54 \text{ cm}$$

3.4 Calculations

Plot a graph of $1/r$ versus I , and draw the straight line that gives a best fit to the five measured points. Use equation (6) to calculate e/m from the slope of this line. Write your report up to this point and then proceed to the following:

3.5 Further Considerations

- (a) Calculate the percentage difference between your value of e/m and the accepted value which is 1.758×10^{11} coulombs/kg. What do you think is your largest source of error? Can you account for this much error by a numerical estimate?
- (b) Calculate the actual maximum velocity of the electrons in your beam.
- (c) Make another run with a lower accelerating voltage. Do you find a consistent value of e/m ?
- (d) Compare the intercept of your graph with the value of coil current you obtained by balancing the earth's field. If these numbers are not roughly the same, you may have made an error. Note that this current is not an important number, but it makes a good check on your technique. Physicists frequently check the consistency of their data by computing numbers which they do not "need".
- (e) Calculate the actual value of B_e .
- (f) Test the maximum error in your readings which could be caused by a change in the field of the nearest neighboring coil.

Experiment 2-6

Optics I: Reflection, Refraction, and Lens Equation

1 Introduction

This experiment is the first of two experiments introducing the basic ideas of geometrical optics. These experiments are unrelated to the material you will cover in class this semester.¹ They are designed to present the ideas of geometrical optics from an empirical point of view.

In this first experiment, we utilize a simple model describing light as a bundle of rays. We then explore the behavior of these rays as they reflect from smooth surfaces and as they are transmitted, or refracted, through transparent media. We experimentally verify Snell's law of refraction and observe the phenomenon of total internal reflection as a consequence of Snell's law.

We will subsequently use the concept of light rays to describe ray diagrams and derive the lens equation². These are powerful tools to draw conclusions based on geometrical optics. In this experiment, we will verify the lens equation for real and virtual images. In the next experiment, we will use these tools to build a magnifying glass, microscope, and telescope.

Why is optics important (besides the fact that the human eye works by the rules we describe here)? Optical instruments are used in many real-life situations in which an image of a system is needed. Such cases apply whether the system is large or small, and whether easily accessible or not. For example, whenever you want insight into how a biological system works, you will likely need to use imaging methods based on geometrical optics. Even if you are only interested in the final data from some fancy imaging system, it will be essential for the quality of your work that you know how to evaluate what you are seeing and what the limitations of the method are.

Remark: It is strongly recommended that you look over the setup in the Lab Library before coming to the lab. This lab has many subparts and getting stuck at one point may make it difficult to finish the lab. It is difficult to understand optical phenomena in the abstract so simply reading the lab manual is insufficient preparation for this lab.

¹The material in these labs is discussed in the assigned lecture course textbook (Fundamentals of Physics, 10th Ed., Halliday, Resnick & Walker, Chapters 34-38) In particular, we adopt the sign conventions used in the course textbook.

²Many approximations are made with "thin lenses". However, even the most sophisticated treatments of optics usually begin with these "thin lens approximations".

Be sure to stop by the Lab Library and let the TA show you the apparatus!

2 Theory

2.1 Geometrical Optics

Geometrical optics is a simplified way of describing light phenomena. It is valid as long as we do not consider cases in which light passes through small pinholes or slits, or examine the edges of shadows. In the second semester lecture and lab, you will encounter effects such as diffraction and interference, which cannot be explained by the theory of geometrical optics. The assumption for geometrical optics is that rays of light propagate along straight lines until they are reflected, refracted, or absorbed at a surface.

A very simple, but important, application of geometrical optics is X-rays. (X-rays are a type of ‘light’, with wavelengths that are much shorter than those characteristic of visible light.) The shadow image of the skeleton made in an X-ray may be understood by simply propagating straight lines from the source to the detector, except for those rays absorbed by some tissue (e.g. bone) in the path.

2.2 Light as Rays

The concept of light rays may become more plausible to you if you imagine being in a dark, dusty room with light from the outside entering through a small hole. You can “see” a ray of light traveling in a straight line from the hole to wherever it hits the wall. (Your eye actually sees the light scattered by dust particles along the path of the ray.)

We describe light using the following simple model. A light source emits rays of light in all possible directions. Each ray propagates along a *straight line* until one of the following happens. It is:

1. reflected if it encounters a reflective material.
2. refracted if it encounters a transparent material.
3. absorbed if it encounters a non-reflecting, non-transparent material.

For most materials in the real world, combinations of these effects occur.

In geometrical optics, we assume that each ray of light travels along a straight line indefinitely unless it strikes a boundary, like a mirror or an interface between air and another material. A source of light, such as a flashlight, produces many rays that travel

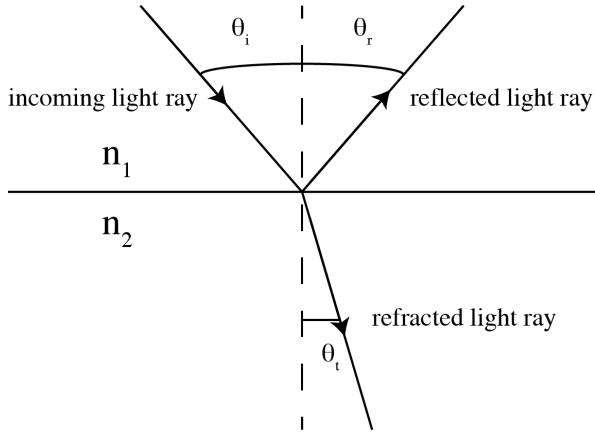


Figure 1: Reflection and Refraction of Light Ray Across a Surface

nearly parallel to each other on their way to the illuminated spot on the boundary. If a ray is not absorbed when it strikes the boundary, it splits into parts as shown in the figure above. Both the reflected and refracted rays again travel in straight lines until they encounter another boundary. The angles θ_i , θ_r , and θ_t are measured with respect to the normal (perpendicular) to the boundary surface.

Although in actuality a light ray that encounters a surface experiences all three of these effects, there is usually one that dominates. If the boundary is a metallic mirror, only the reflected ray is relevant, and if the boundary is a transparent medium such as water, only the refracted ray is relevant. In sections 2.4 and 2.6, we describe the quantitative relationships between the incident ray and the reflected and refracted rays. In particular, we write the relations between the angles θ_i and θ_r , and between the angles θ_i and θ_t shown in the figure.

2.3 Absorption/Color (a brief aside)

The colors we perceive in light are directly correlated with the wavelength of the light ray (for example, blue light has a shorter wavelength than red light). White light is the response your brain perceives if light of all colors strikes your eye's receptors.

Materials that absorb light can either decrease the reflected intensity for all colors equally (which means that the light simply appears less bright than before) or they can selectively decrease the intensity of some colors. An example of the second case is white light from the sun falling onto the grass where the grass absorbs all colors except green (which it reflects). This is why grass appears to you as green. If an object absorbs all the light, it appears to be black.

All materials absorb to some extent, even when the light appears to pass through

or reflects. (The best commercial mirrors reflect about 99.99% of the incoming light.)

2.4 Reflection

The ray reflected from a surface emerges at an angle equal to the angle of incidence of the original light ray. Quantitatively, this law of reflection is expressed in terms of the angles shown in the figure as:

$$\theta_r = \theta_i \quad (1)$$

There are two general types of reflection – diffuse and specular. Diffuse reflection is exemplified by the reflection from the surface of this page. On a microscopic scale, the surface of the paper is quite rough; consequently, parallel rays striking even nearby parts of the paper's surface are characterized by different angles of incidence. Each of the initially parallel rays is reflected in a different direction. Therefore, when we shine a flashlight on a piece of paper, irrespective of our position, we see a bright spot on the paper.³ In this case, it is not feasible to determine the relationship between θ_i and θ_r for any particular ray.

If parallel rays reflect from a surface which is very flat and smooth (such as that found on a mirror or on window glass), there is a unique angle of incidence and therefore a unique angle of reflection. This type of reflection is referred to as specular or regular. A flashlight beam reflected specularly will only be observed if it is viewed along the direction of reflection. We restrict our attention here to specular reflection, since only in that case the relationship between the incident and reflected rays can be understood well enough to be used in optical instruments.

You may have seen fresh snow glitter in the sunlight, and wonder which type of reflection that is. The explanation is that some of the small flat surfaces of the snowflake are smooth and act like mirrors that reflect several rays in the same direction, which just happens to be where your eye is. But the totality of the surface of white snow is irregular and reflection from the snow more often looks much like the diffuse reflection from this page.

Remark for Experts: There is of course no perfectly smooth surface in nature. Every surface is rough on the atomic level. To get specular reflection, it is sufficient that surface irregularities are small compared to the wavelength of light. So when you polish anything, from optical instruments to your furniture, the desired successful appearance is accomplished when you have made all surface irregularities small compared to the wavelength of light (which is about 5×10^{-5} cm, or 500 nm). Since atoms are much, much smaller than that, there is no contradiction.

³Only a few of the light rays are reflected into our eyes!

2.5 Real and Virtual Images

When you look in a mirror, you see your own image. It looks as if a copy of you is standing behind the mirror. But if you put a screen behind the image at the position where your copy seems to stand, you would obviously not capture the image of yourself on the paper. Your copy in the mirror or any image you cannot capture on a screen at its apparent location is called a virtual image. If you can record the image on a screen at its apparent location, it is called a real image. Real images are caused by the actual convergence of light rays at a given point; virtual images are caused by light rays that diverge in such a way that they seem to, but do not actually, converge at a certain point.

Mirrors produce virtual images. In the next two labs, we will deal with many different examples of real and virtual images.

2.6 Refraction and Snell's Law

Light waves (rays) propagate through vacuum with a fixed velocity c equal to about 300,000 km/s or 3×10^8 m/s. One of the consequences of our understanding of electromagnetism, detailed in Einstein's theory of special relativity, is that nothing can travel faster than this speed.

As the waves travel through any material, interactions of the light with the atoms result in a velocity v that is smaller than c . The ratio of these speeds is called the refractive index n , and is a specific constant associated with the medium:

$$n = \frac{c}{v} \quad (2)$$

For visible light traversing through most transparent media, the index varies roughly between 1 and 2.5, depending on the material. Although in actuality, light travels through air slightly slower than it does through a vacuum, for our purposes we will consider the refractive index to be 1. Typical glass, for example, has a refractive index of about 1.5.

When a ray of light in air encounters a medium, the change in wave velocity requires it to change direction.⁴ The new angle relative to the normal, shown in the figure, is the angle of refraction θ_t given by Snell's Law

$$n_i \sin \theta_i = n_t \sin \theta_t, \quad (3)$$

which reduces to the following equation when the incident ray travels through vacuum/air

$$\sin \theta_t = \frac{\sin \theta_i}{n_t} \quad (4)$$

⁴The directional change follows from a simple physical argument.

For example, if an incoming ray at an angle $\theta_i = 15^\circ$ encounters glass (with an refractive index of $n = 1.5$), then the refracted ray will change direction and travel through the glass at an angle of $\theta_t = 9.9^\circ$. Since n is always larger than unity, the refracted ray entering a material is closer to the normal than the incident ray.

Remark for Experts - A Physical Argument for Snell's Law: Fermat's principle states that a traveling light wave will follow the path of least time. Imagine two points, where light is moving from Point A to Point B. Now if the space between them is empty the light will travel at c for the entire duration, and thus the path of least time is equivalent to that of shortest distance, which we know to be a straight line. Thus in the case of empty space Fermat's principle reduces to the trivial answer that light will move in a straight line. However, when the speed at which the light travels is not constant for the entire journey, the picture can change dramatically. If we now let Points A and B be inside materials with indices of refraction $n_1 = \frac{c}{v_1}$ and $n_2 = \frac{c}{v_2}$, respectively, we can imagine a theoretically infinite number of paths the light can take between the two. If we now calculate the time it takes the light to traverse each path, we would find the least amount of time occurs on the path where $n_1 \sin \theta_1 = n_2 \sin \theta_2$ or Snell's Law! Note that if we divide each side by c we now get

$$\frac{1}{v_1} \sin \theta_1 = \frac{1}{v_2} \sin \theta_2, \quad (5)$$

which relates the speeds of light in each medium to the path the light follows. You may wonder, does this have applications outside of light? The answer: yes! Imagine you're playing frisbee at the beach and your friend throws it over your head and into the water. You're on the sand, where you know you run at speed v_1 ; likewise you know you swim at speed v_2 . If you want to minimize the time it takes you to retrieve your frisbee, just follow the path defined by Equation (5) (of course, it may take you longer once you carry out the necessary calculations and measurements). The only difference with light is it *must* "choose" this path, whereas you have the free will to consider other options.

2.7 Total Internal Reflection and Critical Angle

Snell's Law describes the paths light rays follow when traversing interfaces between two media. Though in the previous section we considered light passing from a lower index of refraction to a higher, it is equally valid to consider the reverse. We would therefore see that when light passes from glass to air, the ray exits with a larger θ_t than θ_i . In this example Snell's Law can be reduced to: $n \sin \theta_i = \sin \theta_t$. The refracted ray must have a larger angle to the normal than the incident ray, but this angle must always be less than 90° in order for the refracted ray to emerge on the opposite side of the interface from the incident ray. Therefore, at some maximum incident angle the refracted ray travels along the surface; for angles beyond this, no refracted ray emerges

and the ray is reflected entirely back into the glass. This angle is called the critical angle, θ_c , and from Snell's Law is defined by

$$\sin \theta_c = \frac{1}{n}. \quad (6)$$

An easy way to determine the index of refraction of a medium is by measuring the critical angle and using this equation. As you can see, total internal reflection can only happen when a light beam passes from a medium with a higher refractive index to a medium with a lower refractive index.

2.8 Lenses

Lenses are made of glass, or a similar transparent material which refracts light at its surfaces. They are shaped so that they cause a bundle of parallel rays to converge to or (seem to) diverge from, a single point, called the focal point. The focal length, f , is defined as the distance between the lens and the focal point. The focal length is a fixed characteristic of a given lens, and depends only on the lens material and shape. A lens that converges light is called convex, and its value of f is positive. A lens that diverges light is called concave, and its value of f is negative. We deal only with converging lenses in this lab.

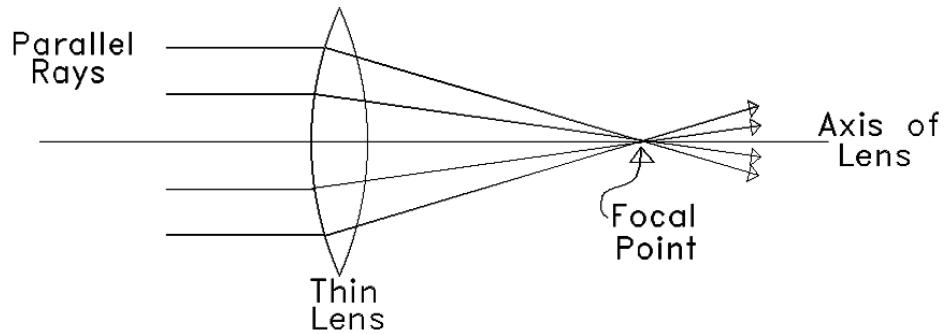


Figure 2: Converging Lens

Remark: Sometimes lenses have a number printed on them which indicates the focal length, in millimeters. A positive number indicates a convex/converging lens while a negative number indicates a concave/diverging lens.

Remark for Experts: If a bundle of parallel rays falls on the lens but not along the axis of the lens, then the focal point will be shifted upward or downward along a line perpendicular to the axis (in the “focal plane”). Note that the focal distance along the axis stays the same.

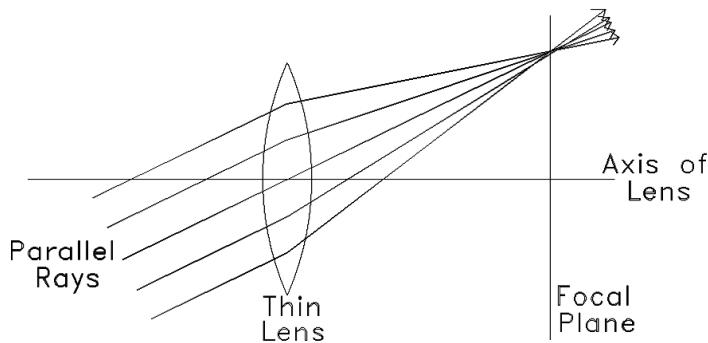


Figure 3: Focal Plane

2.9 Ray Diagrams

Ray diagrams allow us to trace the paths of rays. They make it easier to understand how images are formed and what lenses do.

To draw ray diagrams, follow a few simple rules illustrated in the figure below:

1. Mark the focal points of the lens on each side of the lens.
2. If two rays from a source point intersect at another point, then all rays from that source point will intersect at the second point.
3. When drawing ray diagrams for convex lenses, it is most convenient to use three specific kinds of rays from the tip of the object (P) labeled a , b , and c in the figure and described here:
 - (a) Rays that enter the lens parallel to the axis pass through the focal point behind the lens.
 - (b) Rays that pass through the focal point in front of the lens leave the lens parallel to the axis of the lens.
 - (c) Rays going through the center of a lens will go through the lens in a straight line and do not bend. (These are the only rays that have no net refraction from the lens)

2.10 Lens Equation

The ray diagram described in the previous section provides a graphical method for locating images. Using the simple geometrical argument given here, with the notation indicated in the figure below, we derive the lens equation. This equation relates the image distance S' , the object distance S , and the focal length f .

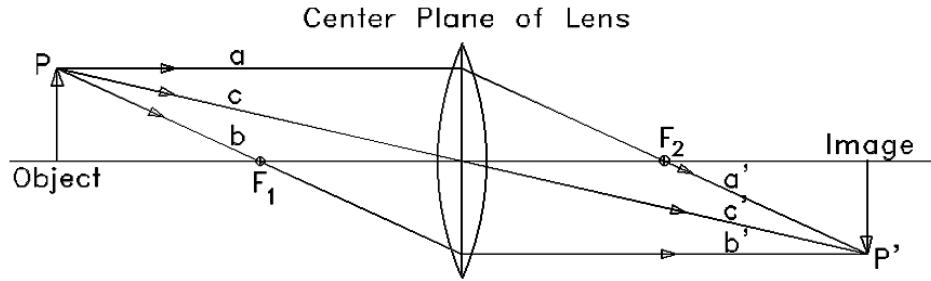
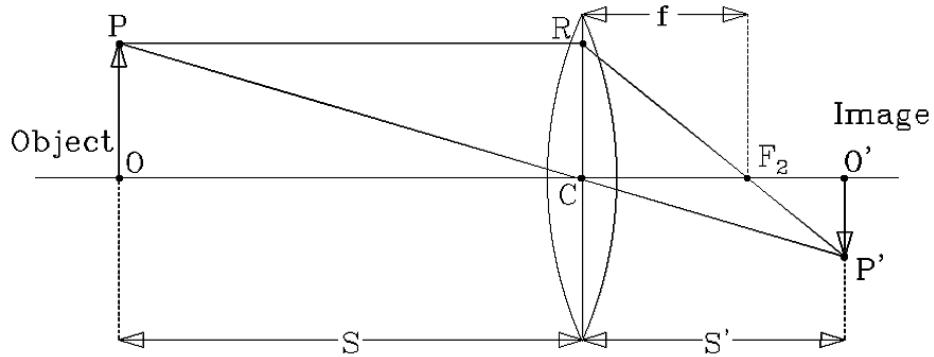


Figure 4: Ray Diagram

Figure 5: An image P' of an object P is formed by a lens

Since triangle COP is similar to triangle $CO'P'$:

$$\frac{S'}{S} = \frac{CO'}{CO} = \frac{O'P'}{OP} \quad (7)$$

Since triangle $F_2O'P'$ is similar to triangle F_2CR :

$$\frac{S' - f}{f} = \frac{O'F_2}{CF_2} = \frac{O'P'}{RC} = \frac{O'P'}{OP} \quad (8)$$

Combining equations (7) and (8):

$$\frac{S'}{S} = \frac{S' - f}{f} \quad (9)$$

which reduces to:

$$S'f = SS' - Sf \quad (10)$$

and finally, dividing by fSS' and rearranging, we obtain the thin lens equation:

$$\frac{1}{f} = \frac{1}{S} + \frac{1}{S'} \quad (11)$$

The sign convention that we have used in deriving this equation, sets f is positive for a converging lens, S is positive on the side of the lens where the object is placed,

and S' is positive on the other side of the lens. For $S > f$, S' is also positive and a real image exists on the other side of the lens, as shown in the figure above. If the object were placed inside the focal length, i.e. if $S < f$, then the solution of the lens equation would give a negative value for S' and the image would appear on the same side of the lens as the object.

The magnification of the image is the ratio of the image height to the object height. Eq. (7) shows that this magnification is numerically equal to S'/S , the ratio of image distance to object distance. The figure above also shows that the real image, formed when $S > f$, is inverted.

3 Experiments



Figure 6: Equipment for the Experiment

3.1 Law of Reflection and Snell's Law

In the first part of the lab we verify the Law of Reflection and Snell's Law with the materials seen in Figure 6. To begin, set the D-shaped acrylic lens in in the similarly-designed outline on the Ray Table. Rotate the knob on the front of the Versatile Light

Source until one light ray emerges, and arrange it so it perpendicularly crosses the flat side of the center of the acrylic lens (you can arrange it to align with the “Normal” arrow on the Ray Table). Now rotate the Ray Table such that the ray still crosses the center of the lens, but is no longer orthogonal, as in Figure 7. You should see transmitted and reflected rays. Using the angle markings on the Ray Table, measure and record these angles along with that of the incident ray. Repeat this process for two other incident angles.

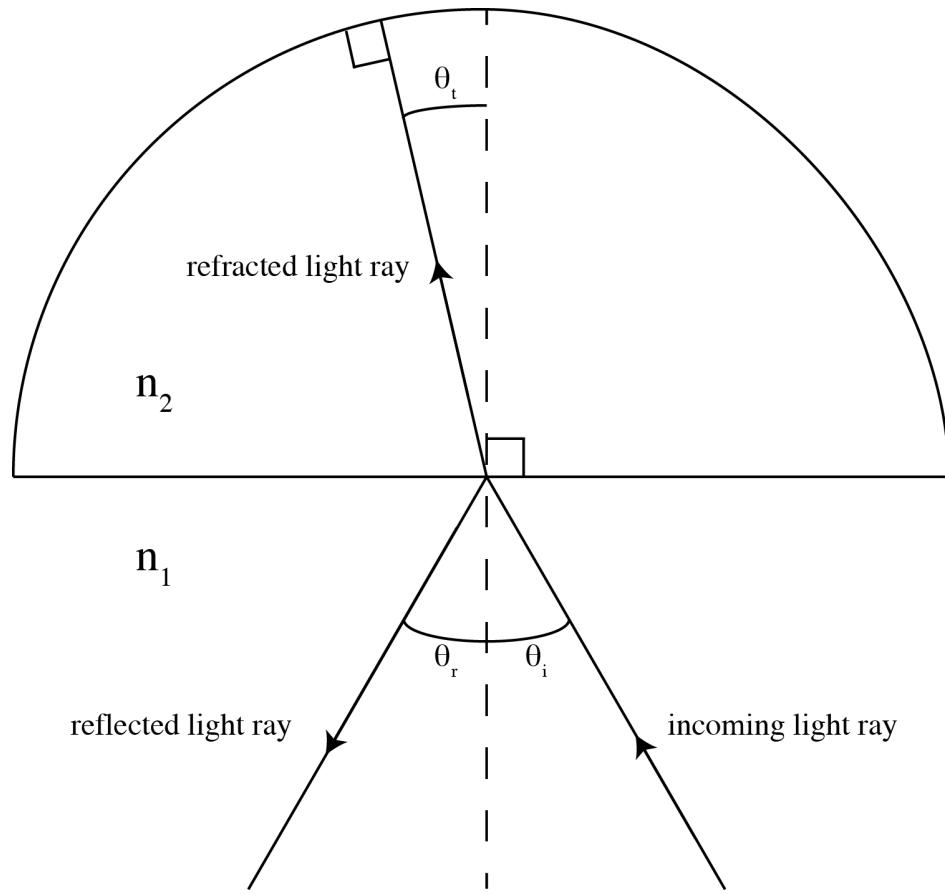


Figure 7: This diagram demonstrates what should happen as a bundle of light rays incident on the face of a D-shaped acrylic lens at angle θ_i are reflected and refracted at angles θ_r and θ_t , respectively.

Consider θ_i vs. θ_r :

- Do the incident and reflected rays intersect at the reflecting edge of the acrylic?
- Estimate a reasonable uncertainty for these angles, considering the technique used to construct and measure them.
- Are the two angles equal within uncertainty?

- What were the main errors in performing this experiment?
- Plot θ_i vs. θ_r . Draw a best-fit line. What are the slope and intercept (and associated uncertainties)? What should the slope and intercept be? Explain.

Consider $\sin \theta_i$ vs. $\sin \theta_t$:

- Calculate the refractive index explicitly from each pair of rays in your data. Calculate the mean and estimate the uncertainty of a single measurement of the refractive index.
- Plot your values for $\sin \theta_i$ vs. $\sin \theta_t$. Explain why your data points should now appear on a straight line.
- Draw a best-fit line. What are the slope and intercept (and associated uncertainties)? What should the intercept be? Explain. Determine the refractive index of acrylic from the graph.
- The accepted value for the refractive index of acrylic is approximately 1.49. Does the book value fall within one or two uncertainties of the mean?
- Note the main sources of error.

3.2 Total Internal Reflection and Critical Angle

In the previous section, we found $\theta_i > \theta_t$ as the light passed from air to acrylic (Snell's Law guaranteed this because $n_{\text{air}} < n_{\text{acrylic}}$). Total internal reflection may only occur when light passes from a medium with a higher index of refraction to one with a lower ($n_i > n_t$). For this portion of our experiment we want to setup a scenario in which light passes from acrylic to air. Figure 8 shows a schematic for how we can accomplish this.

Begin with the ray of light from the Versatile Light Source orthogonally penetrating the curved portion of the D-shaped acrylic lens along the "Normal" line that is transcribed on the Ray Table. You should see the ray of light emerge perpendicularly to the flat surface of the acrylic along the opposite side's "Normal" tracing. Next begin rotating the Ray Table. In doing so, you should notice the emerging ray of light from the acrylic's flat side begins to bend. Though the effect is initially subtle, eventually you should reach an angle where the refracted light is now at $\theta_t = 90^\circ$. This incident angle has reached the critical angle θ_c - the point beyond which we have total internal reflection, and no refracted light escapes from the acrylic.

Measure the angle θ_c at which total internal reflection occurs.

- From θ_c determine the index of refraction for the acrylic lens.

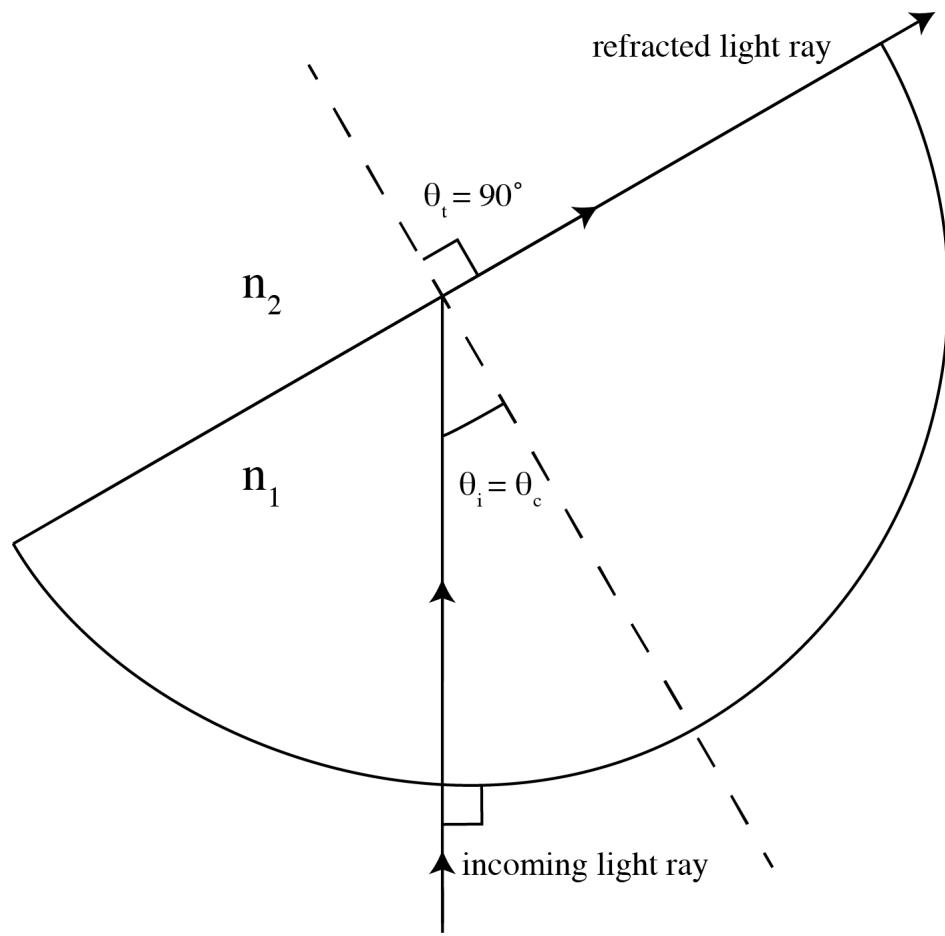


Figure 8: This demonstrates how one achieves total internal reflection in the case where $n_1 > n_2$ and thus $\theta_i < \theta_t$. Note because a line running from a circle's center to any point is necessarily orthogonal to the circumference, some simple geometry allows us to easily determine the θ_i that leads to the critical angle θ_c . It is recommended you study this figure until you are convinced of this.

- How does this value for n compare to your findings in Section 3.1?
- Note the main sources of error. How does your error compare to that from Snell's Law?

An alternate way to achieve total internal reflection: Begin with the ray of light from the Versatile Light Source orthogonally aligned with the flat face of the D-shaped acrylic lens, as we did with Snell's Law. This time instead of rotating the Ray Table, simply slide it sideways so that the light beam remains perpendicular to the flat face of the lens, but now crosses on either the left or right side of the center. You should notice that the light ray remains perpendicular to the face once in the acrylic. However, it will now hit the curved portion of the D-shaped lens at an angle less than 90° , and as a

result is refracted out the back. Eventually you will reach a distance where you notice the light is no longer refracted, but appears to bounce around the circular portion of the lens, due to recurring total internal reflections.

3.3 Measuring the Focal Length of a Lens

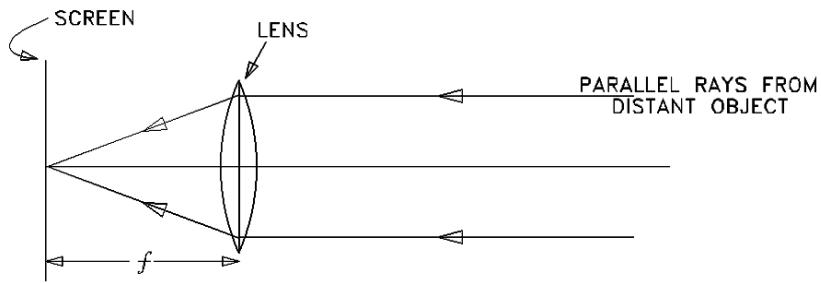


Figure 9: Measuring the Focal Length

Using a lens and a piece of paper, focus the image of an object outside the laboratory window (such as a tree or light outside the window) onto the paper, like illustrated in Figure 9. If it is dark outside or the window blinds are closed, you can use a lightbulb in the room. Based on the fact that light from a distant source may be safely assumed to contain only parallel rays, the distance between the lens and the paper should be the focal length of the lens. You may already suspect that this method is not very precise and that you will get large uncertainties.

- Measure this distance and estimate uncertainty for two lenses of your choosing.
- Take note of the orientation of the image on your paper.
- How could you improve this simple procedure?

3.4 Lens Equation

Set up your optical bench in the following way: stand the Versatile Light Source upright on one end of the optical bench such that the crossed arrow target points down the bench. Place one convex lens on the bench a distance that is greater than its focal length f from the Light Source. Finally, place the Viewing Screen on the opposite side of the lens and move it forward and backward until the crossed arrow target comes into focus. By measuring S and S' , you can verify the lens equation.

- Measure S and S' for five different pairs, including uncertainties. Plot a graph of $1/S'$ vs. $1/S$, with the uncertainty represented as error bars. Draw a best-fit line and get the value of f from the intercepts with the axes.

- Is the value of f consistent with the value from the previous measurements? Is the value of the slope of the line what it is supposed to be?
- Comment on how well the line fits your data points. Which of the two methods used so far should give you the best estimate for f ? Explain why.
- Give the main sources of error and how you could improve this experiment.

4 Applications

Swimming pool experience:

The next time you go swimming, do the following experiment (safely!):

From underwater, look at people or trees outside the water. You will see that you can see them when you view with your eyes at small angles relative to the normal to the water surface. But as you look at larger angles, you find that the water surface suddenly appears weird, somewhat like liquid mercury. This occurs because the water surface reflects like a mirror so that you are not able to observe anything outside the pool. This is a consequence of total internal reflection. (If you will not be swimming for a while, you can also check this by looking at the water-air surface of a fish tank from below.)

Medical diagnoses:

A standard method of medical imaging utilizes ultrasonic sound echocardiography: ultrasonic sound waves produce images of inner body organs (along with many other applications). A wave of ultrasonic sound sent into the body is refracted and reflected as it passes through parts of the body of different density. The wave reflected from a specific organ is observed by a detector and produces an image. The manner in which the ultrasonic sound waves get reflected and refracted is similar to the case for light that we have considered.

An important difference between light and sound waves is that sound waves propagate faster in water than in air. (With light, it is the other way around!) This means that, for sound, water has a lower refractive index than air! Therefore, total internal reflection occurs going from air to water. To enable the sound waves to get into the body requires having no air between the source and the body, usually accomplished by putting a layer of gel between the sound head and the body. This also means you cannot send ultrasonic sound through air-filled organs (like the lungs).

The picture below shows an echocardiogram of a human heart:

LA = left atrium, LV = left ventricle, RA = right atrium, RV = right ventricle.

(Picture from: Marvin Berger: Doppler Echocardiography in Heart Disease.)

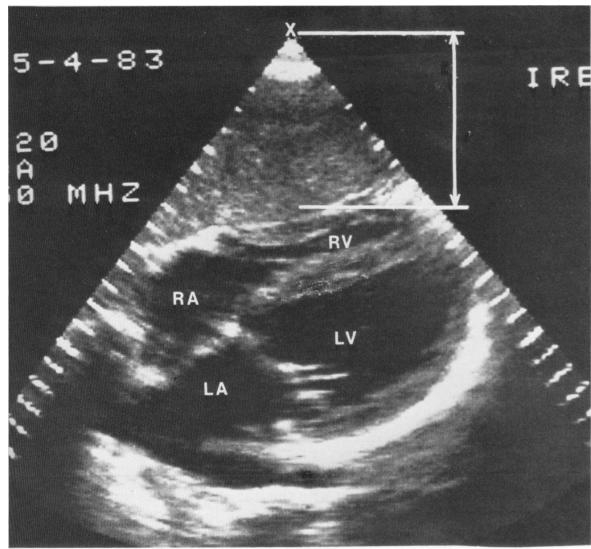


Figure 10: Echocardiogram of a Human Heart

The human eye:

The most important application of lenses (to us) is the human eye! The retina is located a fixed distance from the lens. We need the ability to focus objects located at different distances in front of the eye onto the retina. These specifications require us to have an adjustable lens (with variable focal length). Adjusting the focal length is accomplished by deformation of the lens through contraction of the ciliary muscles. If the eye views a distant object, the muscles are relaxed and the lens is rather flat, with a long focal length. If the eye must examine a nearby object, the muscles contract and the lens becomes rounder, with a shorter focal length. With advancing age, the lens loses its flexibility so that the eye loses much of its ability to adapt to objects at close distances. (A common misperception is that this can be compensated by “eye exercises”, which would be the case if the problem were muscles. But the problem is not in lost vigor of the muscles, but in decreased flexibility of the lens!)

The two most common optical defects are nearsightedness and farsightedness. In a nearsighted (myopic) eye, the focal length is too short even when the ciliary muscles are completely relaxed. Thus, parallel rays from a distant object come to focus in front of the retina and fail to form a sharp image on the retina. Vision of distant objects is blurred. Eyeglasses with diverging lenses correct this condition.

In a farsighted (hyperopic) eye, the focal length is excessively long, even when the ciliary muscles are fully contracted. Hence, rays from a nearby object converge toward an image beyond the retina and fail to form a sharp image on the retina. Eyeglasses

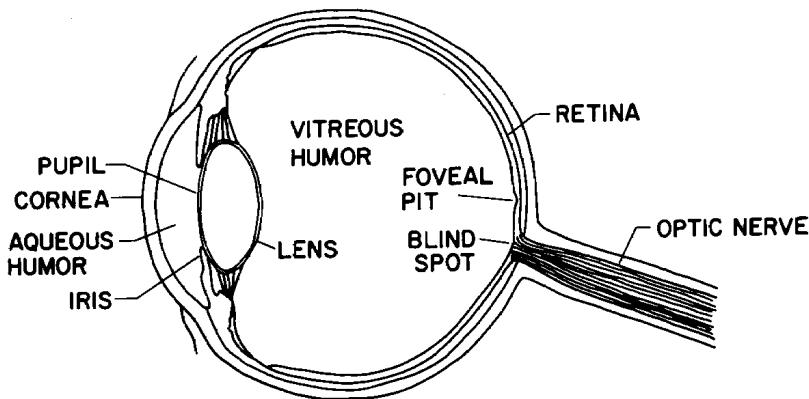


Figure 11: Human Eye

with converging lenses can correct this condition.

Reference: see Physics, by Ohanian, for further information.

Picture from: Daniel Malacara: Geometrical and Instrumental Optics.

5 Lab Preparation Problems

Absorption:

- How many times does a ray of light get reflected on a parallel pair of mirrors (which reflect 99% of the incoming light) such that the total intensity is down to 50%?

Reflection:

- Explain in your own words the difference between specular and diffuse reflection using the following system and observations:

On a day with no wind, the surface of the sea can be very smooth and shiny. On the other hand, the tops of the waves appear white in a strong storm.

Refractive index and Snell's Law:

- What is the refractive index if a medium can slow down light to $v = 2/3 c$?
- If you measure an incoming ray from the air to be at an angle of 15 degrees to the normal, what angle will it have in the medium with an refractive index of $n = 1.2$?
- If you measure $\theta_i = 25$ degrees (air) and $\theta_t = 20$ degrees, what is the refractive index of the medium in which you measured θ_t ?

6. Given the following pairs of values for θ_i and θ_t , draw a graph of $\sin \theta_t$ vs. $\sin \theta_i$. Determine the value of n using only the slope of the graph and not the individual values from the table.

i in degrees	0	15	30	45	60
r' in degrees	0	11.5	22.6	33.0	41.8

7. Given the following measured values for the refractive index, what is the mean and uncertainty (using the 2/3 rule)?

$$n = [1.24 \quad 1.21 \quad 1.29 \quad 1.23 \quad 1.27 \quad 1.26]$$

Total Internal Reflection and Critical Angle:

8. Given an refractive index of diamond $n = 2.42$, what is its critical angle?

9. You want to measure the refractive index of water. For that you take a glass of water or go to your fish tank and look upwards at the air-water boundary through the side of your glass. When look at the boundary with a small enough angle, the surface will be (almost) 100% reflective and you cannot see out through the boundary. With this setup, you measure the critical angle to be 60 degrees. What does this data give for the refractive index of water? Is it the same as the value you find in books?

Experiment 2-7

Optics II: Applications of Lenses

1 Introduction

In this experiment, we will apply our knowledge of lenses and geometrical optics to build a microscope and telescope. Microscopes are common in the medical field, and allow doctors to closely examine things that our eyes cannot resolve. Understanding the basic operation is important to knowing when they are useful, and how they can be improved. Telescopes allow a user to view a far away object as if it were closer, again allowing us to resolve details that the naked eye cannot. Both devices are useful in a wide variety of fields. Throughout this experiment, we will be using the ‘thin lens’ approximation¹.

2 Theory

Why do we need optical instruments? The human eye is incredible in its ability to capture a wide range of colors and focus over a wide range of distances. However, there are some objects that are either too small, or too far away for the eye to resolve. Doctors and scientists require optical instruments to take improve our natural ability.

2.1 The Near Point

In the human eye, light is focused on the retina by the elastic lens. Though the lens can change its focal length (through contractions of the ciliary muscles), the range of focus is limited. The closest point at which the eye can get a sharp image of an object is called the near point. The near point is different for every person, but the value of 25 cm is an average over a random segment of the population. We use the standard reference of 25 cm here. Young people with good eyesight may have a near point as close as 10 – 15 cm, and very young children have an even smaller one. You will determine the near point for your eyes in the course of this experiment.

2.2 Resolution

In art museums, you probably observed the following effect. As you look at an impressionist painting from a distance, the picture appears sharp and easily identifiable. But

¹Many approximations are made with “thin lenses”. However, even the most sophisticated treatments of optics usually begin with these “thin lens approximations”.

as you approach the picture, you observe that it is made up of a number of smaller objects, large dots or blobs, and is quite coarse-grained.

What happened? Your retina consists of an array of light receptors, somewhat like the pixels in a video camera. When you are far enough away from the picture, all light rays from an object like a big dot fall onto a single receptor. Your visual impression from a long distance therefore consists of various sharp points. As you approach the picture, rays from this object begin to strike several receptors and you recognize that the blob is an extended object. If you look at the same picture, but shrunk by a factor of 10, you observe that you can get much closer before seeing the individual blobs. Such effects are related to the resolution limits of the eye.

The figure on the next page illustrates how an image appears on the retina. Note that the image is inverted on the retina. We actually see objects upside down. (Our brains do the work of inverting it back!) When we view an object, its size is determined by how big the objects image is on our retina.

Figure 1 below, we see that the size of the retinal image is directly proportional to $\tan \alpha = h/s$. Clearly, there is a minimum value of α that can be resolved on the retina.

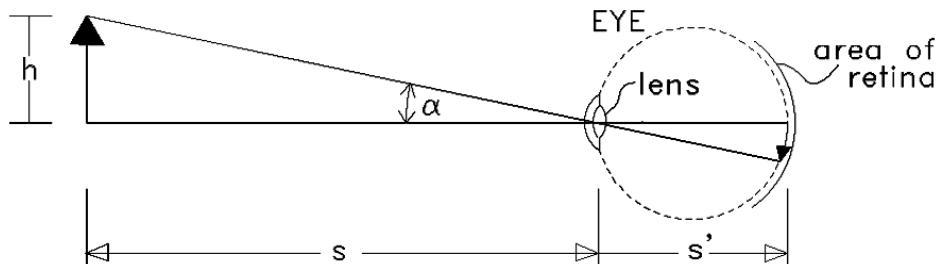


Figure 1: The Formation of Images on the Retina

2.3 Need for Magnification

Suppose a person with a near point of 25 cm examines the object shown in the figure above. To see the most detail, the person should view the object so that it appears to be as large as possible while remaining in focus. As discussed above, the size of the retinal image varies as $\tan \alpha = h/s$. This means that the retinal image gets larger as s gets smaller (i.e., as the object is brought closer to the eye). The largest focused image this person can see is obtained when $s = 25$ cm. The size of the object when viewed at 25 cm is defined to have a magnification of unity. Is there any way to see more detail on the object, or, equivalently, can the object be made to appear larger? The answer is yes, and the simplest means of making the object appear larger is to use a magnifying glass.

2.4 Virtual Images vs. Real Images

In the lens experiments we have been covering so far, it has been possible to place a screen at the location of the image and since the light rays focus at that screen, a clear image is formed on the screen. Such images are real. However, consider a ray diagram with a lens, where the object is placed between the focal point and the lens, as in Figure 2.

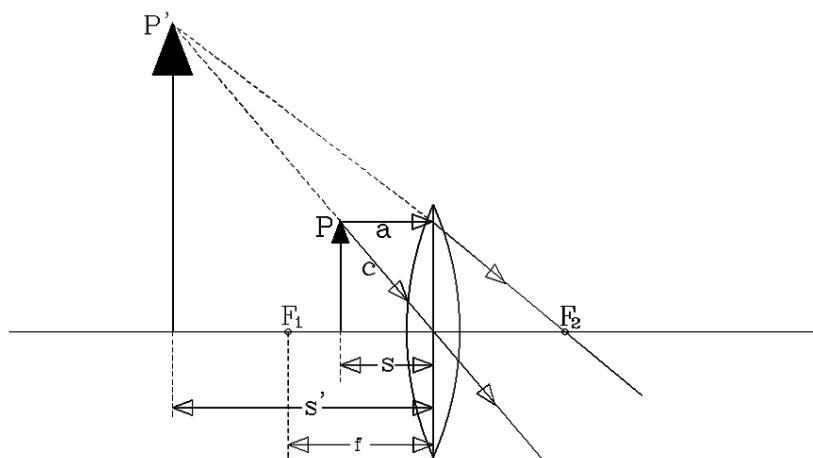


Figure 2: Convergent Lens: Virtual Image

When the object is placed exactly at the focal point, the rays exit parallel (and the image is infinitely far away). (Remember that this is how we define the focal point, and how we determined the focal length in Experiment 4) As the object gets even closer to the lens, as shown in Figure 2, the rays must diverge as they exit. If you extend the outgoing rays backward (behind the lens), you find that they intersect behind the original object. A virtual image is formed. It cannot be captured on a screen, but you can see it as you look through the lens - the object appears at the position of the virtual image and enlarged.

This same phenomena happens when you look into a mirror and see a clear reflection of yourself. The rays originating from your body never really converge, so it is not possible to see the image on a screen.

2.5 Parallax

In everyday life, we are able to tell distance by a combination of seeing how large things are relative to other objects, and how they move relative to each other. But what happens when we are not sure what the size of an object is supposed to be? In optics experiments, both virtual and real images can be magnified in non-intuitive ways, so using the size of the image we see compared to the original object is not a

good way to tell which object is closer or further. We must use the method of parallax. Parallax is the apparent change in distance between two objects that depends on where the observer is.

Imagine walking along a road and seeing two trees in the distance, as in Figure 3. At point A in the left image, you see an apparent distance between the two trees given by L_A . But as you walk to point B, you would notice that that distance changes dramatically, and the trees only look lined up when you are standing directly in front of them. However, if the trees are in nearly the same location, the apparent change between L_A and L_B is very small as you walk along the road.

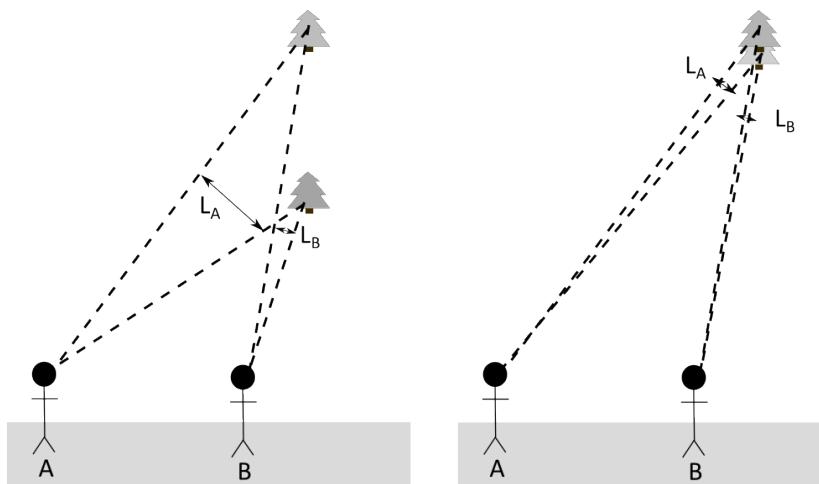


Figure 3: An example of parallax. In the image on the left, the difference in observed distance between the trees changes dramatically as you move along the road, hence the large difference between L_A and L_B . In the image on the right, where the trees are nearly in the same location, L_A and L_B are basically unchanged.

2.6 The Magnifying Glass

A magnifying glass is a simple converging lens with short focal length that can make an object's image on the retina larger. How does it work? One limitation to the detail you see on an object is the limited angular resolution, which is due to the limited size of the individual receptors. (There is nothing we can do about this.) But the other limitation is that we cannot make the retinal image larger by bringing the object closer than the near point. There is a way to overcome this limitation, and it is used in all optical magnifying devices.

Consider the object and lens of focal length f shown in Figure 4.

The object is placed a focal length f away from the lens. Using the lens equation, we find that a virtual image is formed at infinity as shown. (If you feel uncomfortable

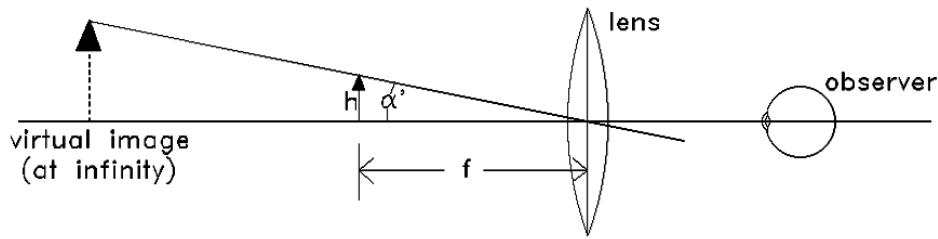


Figure 4: Virtual Image at infinity

thinking of an image at infinity, you might prefer to think of it as simply a very large distance.)

If the rays of light from the ends of the object that pass through the center of the lens are extended to the left, they also pass through the ends of the image. (If this is not obvious, return to Experiment 1-5 and review the rules for ray tracing). This means that the angle $\alpha' = \tan^{-1} h/f$ is the same for both the object and the virtual image. Remember what an image represents - looking at the object through the lens can be represented as looking directly at the image. Since the image is very far away, its size does not appear to change as the eye (to the right of the lens) is moved small distances back and forth.

What actually happened when you put the converging lens one focal length in front of the object? The rays leaving the lens that actually hit the eye (not the ray shown in the figure) will be nearly parallel and therefore the rays appear to come from an infinitely far object unless you are near-sighted, your eye can totally relax.

When we view the image, its size is proportional to $\tan \alpha' = h/f$. The *angular magnification* of the magnifying glass is defined as:

$$M = \frac{\tan \alpha'}{\tan \alpha} = \frac{h/f}{h/25 \text{ cm}} = \frac{25 \text{ cm}}{f \text{ cm}} \quad (1)$$

The **magnification** M represents the ratio of the apparent size of the object when viewed using the magnifying glass to that of the object when viewed directly at a distance of 25 cm.

If f is less than 25 cm, the lens increases the apparent size of the object and permits you to see more detail. Viewing an object with a magnifying glass of focal length 12.5 cm produces a retinal image twice as large as that formed when the object is viewed directly 25 cm away.

If a lens of focal length 25 cm is used as a magnifying glass, the apparent size of the object is the same as that seen when looking directly at the object 25 cm from the eye. Using the lens, however, permits your eye to relax because it views parallel rays,

just as if the object were located very far away.

In practice, one does not always place the object at the focal point of the magnifying glass, but it is at that point that the magnification is greatest. Magnification can also be equivalently defined as the ratio of the height of the image to that of the object. By drawing congruent triangles, it should be clear that this implies that

$$M = \frac{h_i}{h_o} = \frac{-s_i}{s_o}$$

Note that we follow the sign conventions used in the course textbook here - you may wish to review Section 34-7 of Fundamentals of Physics.

2.7 The Microscope

For various reasons, a single lens magnifier (such as a magnifying glass) can only provide good images with relatively small magnifications. To achieve higher magnification, a device known as a compound microscope can be used.

A compound microscope employs two lenses as shown in the next figure. The object to be viewed is placed at a distance s_{o1} , just beyond the focal length, F , of the first lens, called the objective lens. The light collected by the objective lens forms a real (enlarged) image of the object at s_{i1} . (The image is real because a screen placed at s_{i1} would display an actual inverted image of the object.) The second lens, the eyepiece, has a focal length f and acts as a magnifying glass that is used to examine the image s_{i1} . The position of the image formed by the first lens becomes the object for the second lens, and the distance s_{o2} is defined as in Fig. 5. Adjusting the position of the eyepiece adjusts the position of the final image.

As discussed in the preceding section, the observer looking through the eyepiece (magnifier) sees a virtual image s_{i2} of the real image.

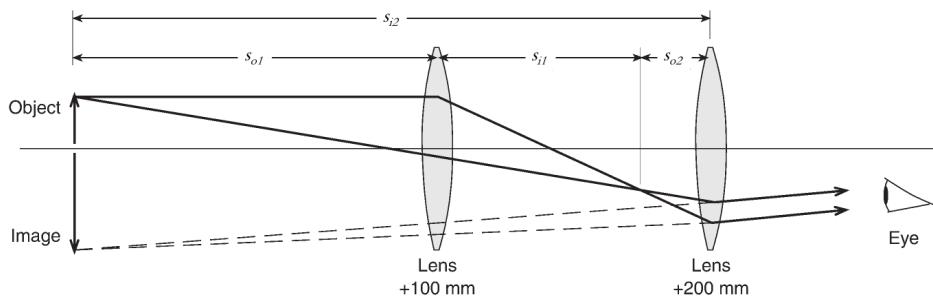


Figure 5: A Compound Microscope

How big does the original object appear? Since magnification represents a multiplicative factor, we simply have to calculate the magnification of the original image I_1 , by the first lens, then multiply that by the magnification of the second lens.

The total magnification of the microscope is then:

$$M = M_{\text{eyepiece}} M_{\text{objective}} = \left(\frac{-s_{i1}}{s_{o1}} \right) \left(\frac{-s_{i2}}{s_{o2}} \right) \quad (2)$$

2.8 The Telescope

The optical system of a refracting telescope has the same elements as those of a compound microscope. In both instruments, the real image formed by an objective lens is viewed through an eyepiece. However, the object of a telescope is distant, the incoming rays are (approximately) parallel and the objective lens forms the first image I_1 , near its focal point F , as shown in the figure below. (By contrast, for the microscope $s' \gg F$.)

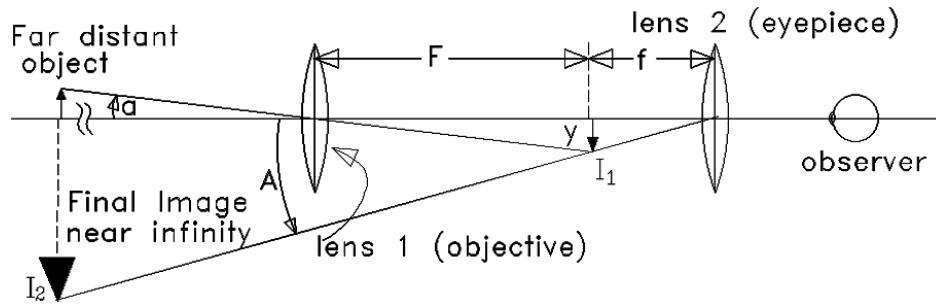


Figure 6: Telescope

The important point here is that although the object is very far away, it still subtends a small angle in our naked eye (otherwise we wouldn't be able to see it). The assumption that it is very far away is to make sure that there will be an image I_1 at the focal point of the objective lens, which again serves as the object for the eyepiece. In order for I_2 – the final virtual image made by the eyepiece – to be near infinity (for relaxed viewing), the eyepiece needs to be located so that I_1 is also near its focal point f . Therefore, the distance between the objective lens and the eyepiece – the length of the telescope – is the sum of the focal lengths $F + f$.

The angular magnification of the telescope is easily derived by noting that α , the very small angle that the distant object would subtend for the unaided eye, is essentially the same as the angle subtended by the image I_1 from the objective lens. Then from the figure above we have:

$$\tan \alpha = \frac{y}{F} \quad (3)$$

and angle A , subtended by I_1 at the eyepiece, we have:

$$\tan A = \frac{y}{f} \quad (4)$$

so:

$$M_{\text{telescope}} = \frac{\tan A}{\tan a} = \frac{y/f}{y/F} \quad (5)$$

$$M_{\text{telescope}} = \frac{F}{f} \quad (6)$$

3 Experiments

3.1 The Near Point

In the first part of the lab, you will measure the near point of each eye. Tape the paper grid pattern onto the screen and place the screen on the optical bench. Place your head so that one eye is at the end of the optical bench. Move the screen as close to your eye as possible, consistent with the requirement that the screen remains in focus. The distance s is the near point for that eye. Repeat for the other eye.

- Is the near point the same for both eyes?
- How do these values for your eyes compare to the standard value of 25 cm?

If the values you obtain are different from the conventional 25 cm value, it should not be surprising; as you compare your values other people in the class, you will probably see a wide variation. So as to make quantitative comparisons of the instruments that we will be using in the following sections, we use a standardized reference near point of 25 cm. We choose this value since the average person cannot get closer to an object than this and still keep it in focus.

3.2 Virtual Images

In this part of the lab, we will study virtual images formed by concave lenses. Concave lenses have negative focal lengths, and are therefore sometimes referred to as “diverging lenses”. You will then use another lens to form a real image of the virtual image. In this way you can identify the location of the virtual image.

Set up the -150 mm lens and light source as in Fig. 7, with the crossed arrow object on the light source pointing towards the lens.

- Look through the lens toward the light source. Describe the image. Is it upright or inverted? Does it appear to be larger or smaller than the object?
- Which do you think is closer to the lens: the image or the object? Why do you think so?

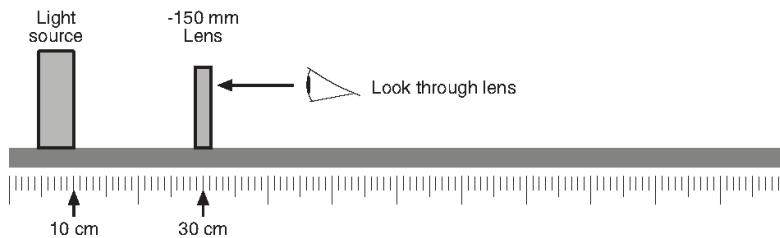


Figure 7: Virtual Image set-up

- Using the lens equation (see Experiment 2-6, and also on the next page) and remembering that the concave lens has a negative focal length, calculate the position of the virtual image.

Place the +200 mm lens anywhere between 50 cm and 80 cm on the optical bench and then move the screen so that the image is focused (Fig. 8).

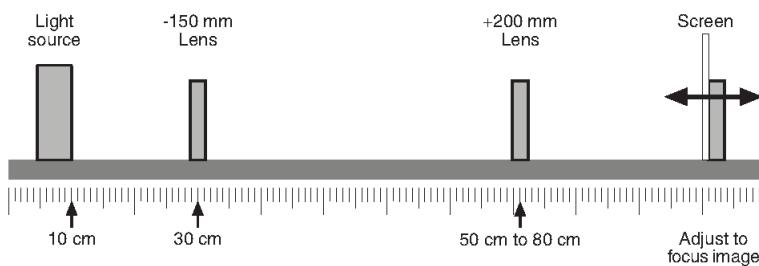


Figure 8: Placing a second lens to form a real image from the virtual one.

- Why are you able to see the image on the screen? What is serving as the object?

In the following step, you will discover the location of the virtual image by replacing it with the light source. Remove the negative lens from the bench, and note that the image on the screen goes out of focus. Slide the light source to a new position so that a clear image is formed on the screen. (Do not move the positive lens or the screen.) Make sure you write down the positions of all components. (Fig. 9)

Your light source is now at the position of the virtual image. Note that it may differ from your prediction made earlier from the thin lens equation. Why is this? (Remember the approximations made in the thin lens equation)

3.3 The Microscope

A microscope magnifies an object that is close to the objective lens. The microscope in this experiment will form an image in the same place as the object (see Figure 10).

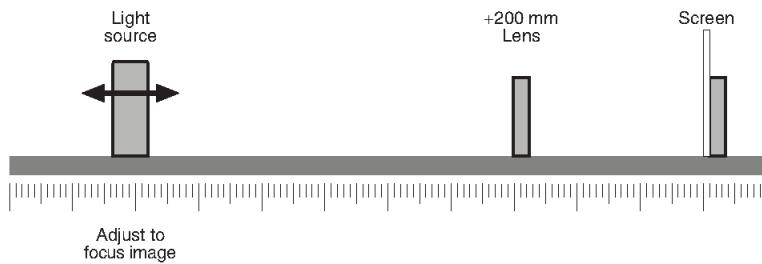


Figure 9: Remove the first lens, then move the light source to the position of where the virtual image was.

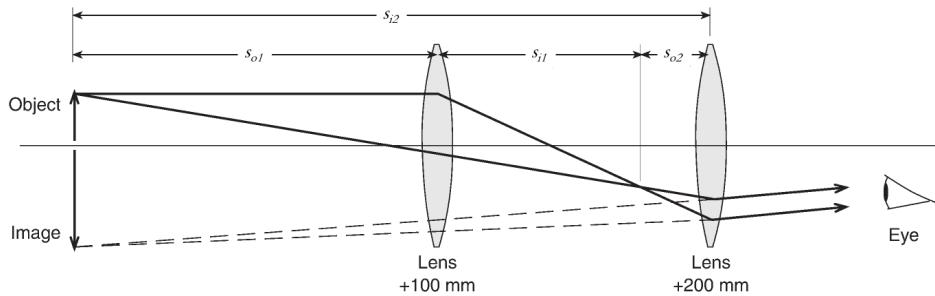


Figure 10: Measuring the Magnification of the Microscope

As discussed earlier, the magnification, M , of a two-lens system is equal to the product of the magnifications of the individual lenses. Magnification is defined to be the negative of the ratio of the image distance to the object distance, as in the lecture course textbook. In a microscope, the eyepiece uses the image of the objective lens as its object. Therefore, the total magnification of the microscope can be given as:

$$M = M_{\text{eyepiece}} M_{\text{objective}} = \left(\frac{-s_{i1}}{s_{o1}} \right) \left(\frac{-s_{i2}}{s_{o2}} \right) \quad (7)$$

In this experiment, we would like to achieve a magnification of **-2** (be careful with sign conventions in the equation above).

Procedure

First, tape the paper grid pattern to the screen to serve as an object. Then, use the 100mm lens as the objective lens and the 200mm lens as the eyepiece, as shown in Fig. 10

The goal here is to create a real image with the objective lens, and then calculate where to place the eyepiece so that the final virtual image is at the same place as the original object.

Set up the Objective lens ($f=100$ mm) 30 cm away from the screen (s_{o1}). Given

this, and the thin lens equation:

$$\frac{1}{f} = \frac{1}{s_o} + \frac{1}{s_i},$$

you can calculate where the image (s_{i1}) from the objective should show up. In order to measure the magnification by eye, we need to place the eyepiece so that the virtual image is formed exactly where the original object is. If we do this, then we can accurately compare the magnified grid to the unmagnified one.

Looking at Figure 10, we can see that the **absolute** value of s_{i2} , the image distance from the eyepiece, is given by:

$$s_{i2} = s_{o2} + s_{i1} + s_{o1}, \quad (8)$$

(Note that with our usual sign convention for object and image distances, s_{i2} is actually negative.) Plugging Equation 8 into Equation 7, we can eliminate s_{i2} . Now, keeping in mind that we would like to achieve a magnification of -2, **calculate where you should place the second lens (s_{o2})**.

Place the eyepiece and look through your microscope. As you move your head from side to side, you should notice that there is little to no parrallax between the lines on the actual object and the lines that you see through the center of the microscope.

- Estimate the magnification by comparing the magnified and unmagnified grids. Quote an error and compare this to the magnification you attempted to get.
- Is the image inverted or upright? How do you explain this?
- Create and fill out a table like the one below in your lab report.

Position of Objective Lens	
Position of Eyepiece Lens	
Position of Screen	
s_{o1}	
s_{i1}	
s_{o2}	
s_{i2}	
Observed Magnification	

3.4 The Telescope

As we have seen before, the telescope works on a similar principle to the microscope, except now, our object distance is effectively infinite.

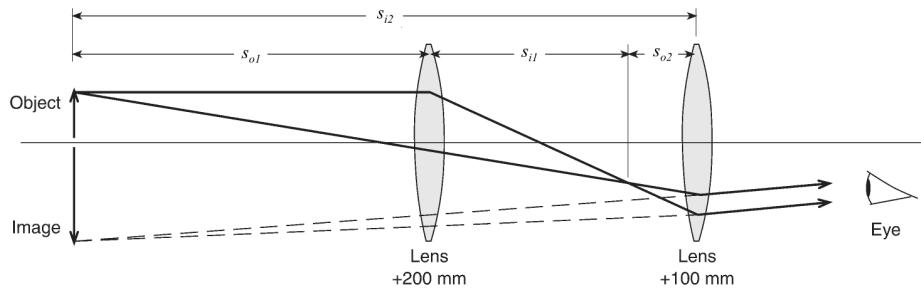


Figure 11: Schematic of a telescope.

We assemble a telescope and measure its magnification. Since the final object is formed near infinity, we can compare far away objects with a magnified image, without worrying too much about where exactly the image is formed. Comparison of these permits an estimate of the magnification.

Using the same 200 mm and 100 mm lenses from before, construct a telescope using your knowledge from reading previous sections. Think carefully about which lens should be the objective and which lens should be the eyepiece.

- Calculate the magnification of your telescope.

Use your telescope to look at a grid that is placed far away (You may have to place your screen off the optical bench to make the “object at infinity” condition valid). Once you have found a clear image looking through the telescope, open your other eye so that you can look simultaneously at the enlarged image with one eye and the unenlarged image with the other eye. If you position the telescope just right, you can see the two grids side by side and estimate the magnification M of your telescope.

- Record your by eye estimate of the telescope magnification. Does it match the theoretical value of M ? If not, why not?
- Is the image inverted? Explain this.
- Carefully draw a ray diagram of the experiment and indicate the type (real or virtual) of each image.
- Create and fill out a table like the one below in your lab report.

Position of Objective Lens	
Position of Eyepiece Lens	
Position of Screen	
s_{o1}	
s_{i1}	
s_{o2}	
s_{i2}	
Observed Magnification	

4 Applications

It may seem strange that we provide yet another example based on the human eye, since you already had that in the last experiment. As with many real-life applications, you are only told part of the story at first (with the complexities left to later). The eye should better be described as a system consisting of two lenses.

The index of refraction of the fluid within the eye is about 1.34 (nearly the same as water). The index of refraction of the solid lens is about 1.44. The difference in the indices of refraction is not large, so the refraction at the interface between them is not strong. Therefore, though the lens provides adjustments in focal length necessary to form the image on the retina at different object distances, the main contribution of focusing the incoming light is refraction at the cornea. This tissue layer separates the eye fluid (or humor) from the air. The difference in indices of refraction between air (1.00) and the aqueous humor (1.34) is much larger and creates substantial refraction. This means that the eye is more precisely treated as a system of two lenses, in which the second lens is variable so as to fine tune the focal length of the system.

Several different optical malfunctions (e.g. farsightedness/nearsightedness) were described in the application part of the previous experiment. These can also arise from malformations of the eyeball. As described before, using an additional lens either in the form of glasses or contact lenses can compensate these.

But a more recent approach is to change the focal length of the first lens (cornea) directly. A high intensity laser beam can be used to evaporate thin layers of the cornea and thus provide the cornea with a new shape. If the curvature of the cornea is increased, the light is focused more (corrects farsightedness) or if you decrease the curvature, the light is focused less (corrects nearsightedness).

5 Lab Preparation Problems

Magnification and Magnifying Glass:

1. The angle subtended by an image with an optical device is $\alpha' = 10^\circ$ and without it, the object subtends $\alpha = 7^\circ$. What is the magnification?
2. What is the magnification of a magnifying glass with a focal length of 10 cm?
3. Assume your eyes have a focal length of 25 cm. What is the advantage for your eyes if you still use a magnifying glass of focal length 25 cm to read a book instead of holding the book at a distance of 25 cm?

Microscope:

4. Assume you have a microscope and measure $s' = 30 \text{ cm}$ and $s = 10 \text{ cm}$. Your second lens has a focal length of 15 cm. What is the magnification?
5. Assume your microscope has a focal length of $F = 10 \text{ cm}$ for the first lens and $f = 10 \text{ cm}$ for the second lens. What is the magnification if you choose $s = 15 \text{ cm}$? What if you choose $s = 12 \text{ cm}$?
6. You build a microscope with a first lens of $F = 5 \text{ cm}$ and a second lens with $f = 10 \text{ cm}$. You separate these two lenses by a distance of 25 cm. What will the magnification of your microscope be?

Hint: You first have to calculate s' and then s !

Telescope:

7. You build a telescope with two lenses. One has a focal length of 5 cm, the other has a focal length of 2 m. In what order do you want to place these two lenses. Which should be the objective lens, which the eyepiece?
8. For the telescope described in 7, what is the magnification?
9. You want to build a telescope that can magnify far away objects in the following way. Everything that appears to be separated at an angle of $1'$ (1/60 of a degree) when viewed without the telescope will appear to be separated by 1° when viewed through the telescope. You have a 5 cm lens for the eyepiece. How do you have to choose the objective lens such that the magnification is as described? How long will your telescope then be? (The length of your telescope is the separation of the two lenses.)

Experiment 2-8

Polarization and Interference

1 Introduction

You have already investigated some wave properties in last semester's Experiment 9. This lab will address three concepts of waves: polarization, interference, and diffraction. We will conduct this experiment with light, but the general concepts apply to many other waves as well.

2 Theory

2.1 Polarization

Polarization occurs only in transverse waves¹. By definition, the oscillation is perpendicular to the direction of wave propagation in a transverse wave. Since in three-dimensional space there are three independent dimensions, there are two possible perpendicular directions needed to describe any oscillatory motion perpendicular to the direction of motion. If you think about a transverse wave in a string, the oscillation could be up-and-down or left-and-right, while the wave front propagates forward along the string. The two independent and perpendicular directions are called the two polarization states. (Or in coordinates, if a wave propagates along the z -axis, there are two polarizations perpendicular to the wave motion—one along the x -axis and one along the y -axis.) Polarization cannot occur in longitudinal waves because the direction of oscillation has only one possibility, along the direction of motion.

You may argue that you can make oscillations happen on a string in more than just the two directions described above—and you are right! For example, you can make oscillations diagonally from the upper left to the lower right or you can even make circular oscillations by moving your hand continuously in a circle, producing a corkscrew pattern along the string². But no matter which wave you make, the oscillations can always be represented as two superimposed waves with displacements along the two mutually perpendicular polarization axes. Any vector can always be described by decomposing it into its x - and y - components. Here, the relevant vector is the oscillating displacement from the z -axis. The amplitudes of the resolved x and y components indicate how much of the wave is in each polarization state.

¹Examples of transverse waves are oscillations on a string, water waves, or the electromagnetic radiation we treat here. Sometimes you may also hear the term polarization used in the context of transverse- and longitudinal polarization.

²This wave is sometimes called left- or right- circular polarization.

For electromagnetic radiation, the oscillating quantities are the electric (and magnetic) fields. These undulate perpendicular to the direction of the wave. Like waves in a string, electromagnetic waves are transverse waves and have polarization properties. The polarization state describes the axis along which the electric field in the wave oscillates. Visible light is electromagnetic radiation in a particular frequency-wavelength regime, with large frequency of oscillation ($\sim 10^{15}$ Hz) and small wavelength ($\sim 0.5 \times 10^{-6}$ m).

2.2 Polarizer

Usually light is produced (e.g. incandescent light from a candle or a light bulb) by hot electrons and atoms, which are moving or are oriented in random directions, so that the wave is **unpolarized**. This means that the electric field associated with the wave oscillates in random directions (though always at right angles to the direction of the propagation of the light). So there is no favored direction for polarization. It is as if you were moving your hand in random directions at the end of a string. Note that this is very different from the cases of waves in a string described above. Even in circular polarization the oscillations were not random.

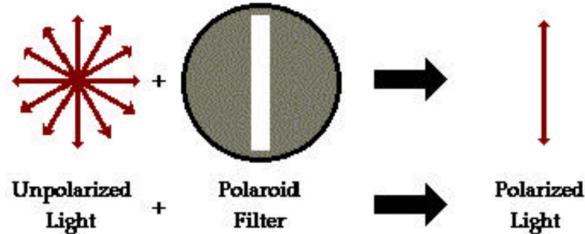


Figure 1: Polarization Filter and Creation of Polarized Light

A polarization filter, shown in Figure 1, is a device with a specified direction, called the polarization axis. The filter absorbs all the wave's incoming electric field perpendicular to the polarization axis and transmits the electric field parallel to the polarization axis. So no matter what the polarization of the light shining into a polarizer might be, the light coming out of the polarizer is always polarized in the direction specified by the polarization axis³. An unpolarized wave just prior to the polarizer has an electric field (**E**) pointing randomly in all directions perpendicular to the wave direction, as shown in Figure 2. The polarized wave after the polarizer has only half the energy content, and the **E**-field is oscillating solely along the polarization axis. This electric field has the same value as the component along that axis just before the polarizer.

³Such light is called linearly polarized light.

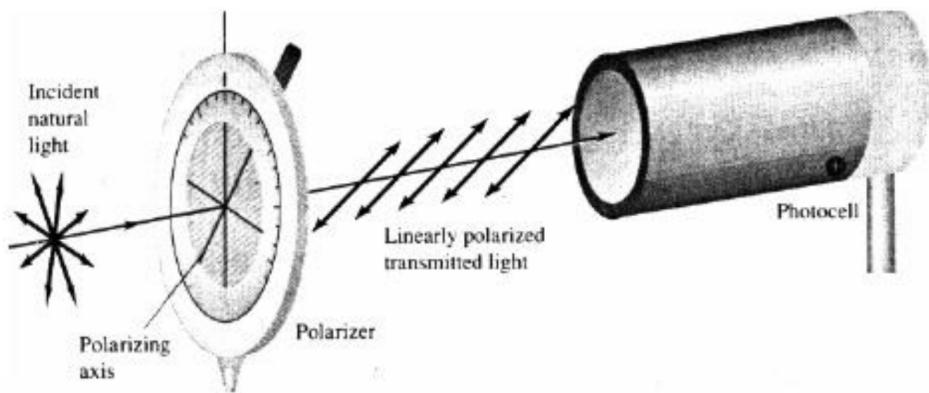


Figure 2: Detection of Polarized Light

2.3 Interference—The Double Slit

As you've learned in class, Young's double slit experiment is a demonstration of interference and diffraction. A double slit, as illustrated in Figure 3, permits what remains of an incoming wave (from the left) to travel to a distant screen (to the right) along two different paths with different lengths. The light waves from the two slits interfere, resulting in an interference pattern of bright (constructive) and dim (destructive) patches as viewed on the screen. Let's see how the double slit works using Figure 3.

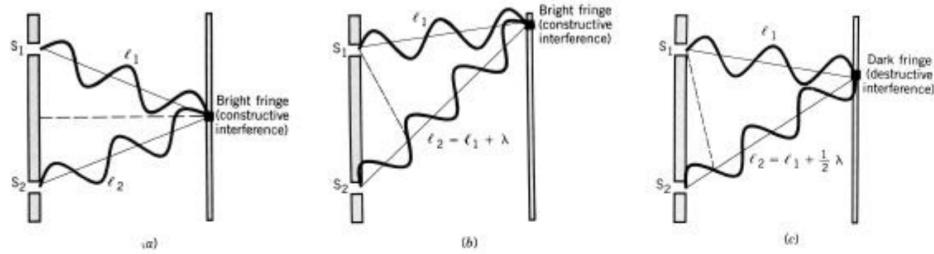


Figure 3: Young's Double Slit Experiment

We shine light from the left onto the double slit, which allows two light waves to propagate from the two slits. Straight-ahead they will always be in phase since they travel the same distance to the screen. But when the two waves propagate at an angle θ , they cover different distances to reach a specific point on the screen. (We assume that the two rays are parallel, a good approximation since the screen is effectively infinitely far away compared to the distance between the two slits.)

As you can convince yourself by applying the rules of trigonometry to the triangle in Figure 4, the difference in path (Δl) between the two rays is, for slit separation, d ,

$$\Delta l = d \sin \theta \quad (1)$$

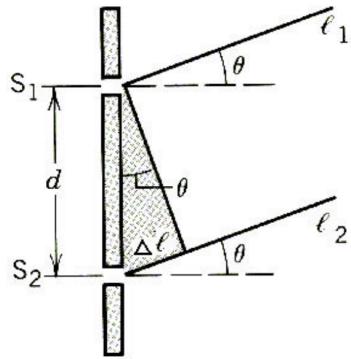


Figure 4: Geometry of the Double Slit

So what quantitatively determines a point of maximum intensity on the screen? Maximum intensity occurs if the two waves are precisely in phase at that point on the screen. This follows from what you learned about constructive and destructive interference in the lab on standing waves. As you saw, two waves can add up constructively or destructively. In the case of maximal constructive interference, the waves combine at the screen so that a crest on one wave coincides with a crest on the other wave and a trough coincides with a trough. In the case of maximal destructive interference, the waves combine so that a crest on one wave coincides with a trough on the other wave.

The waves will be in phase when they meet at the screen if the path difference between the two waves is an exact number of wavelengths, m , where m is an integer called the order of the maximum. This requirement, expressed mathematically, is

$$\Delta l = m\lambda \quad (2)$$

Note that if $m = 0$, there is no difference in the path lengths, resulting in a maximum straight ahead, $\theta = 0$, as shown in Figure 5. Subsequent maxima occur for subsequent positive and negative integers, $m = \pm 1, \pm 2, \pm 3, \dots$

Combining the two equations for Δl , we get for the specific angles, θ_m , that give maximal constructive interference:

$$d \sin \theta_m = m\lambda \quad (3)$$

If we know the order m of a particular maximum, we can determine the wavelength of the light incident on the double slit by measuring the angle of the maximum on the screen. Conversely, we can determine the order of a particular maximum, m , from the wavelength of the light and the angle.

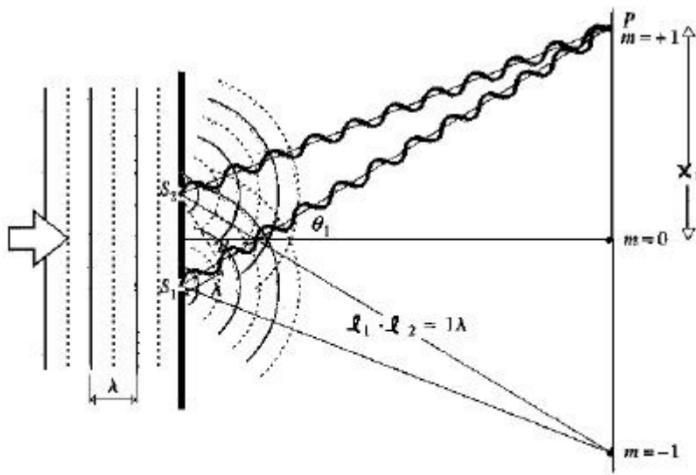


Figure 5: Interference of Light Coming from Two Small Slits

2.4 Intensity of the Double-Slit Maxima

You will soon note that the intensities for subsequent maxima are different. What causes this? Just as the double slit creates an interference pattern, each single slit creates a diffraction pattern due to interference from waves coming from different parts of a single slit as described in section 39-6 of the lecture course text⁴. This occurs because each individual slit has a finite width. The overall (more realistic) pattern shown in figure 6 is a superposition of the single slit and double slit interference patterns. The narrow spikes (of interest for this experiment) are due to the double slit interference and the envelope, or large-scale pattern, is due to the single slit diffraction caused by each slit.

2.5 Light from a Laser

In the double slit experiment we use a laser as the light source. Laser light differs from incandescent or fluorescent light in two important ways. First, lasers emit intense light of a single wavelength or frequency (monochromatic). Second, laser light is coherent, meaning that all light waves are in phase.

Let's illustrate the differences between coherent and incoherent light through a simple familiar example. Contrast a crowd of people cramming the city streets during rush hour with a regiment of soldiers at a parade. In the midtown crowd, even if everybody were to step at the same frequency (which they don't!), each pedestrian's steps are independent of what others are doing. But soldiers at a parade move in step, or in phase, with a similar frequency, taking their steps at exactly the same time. Laser

⁴Sears, Zemansky, and Young, **College Physics 7th edition**, Addison-Wesley, p. 903-906.

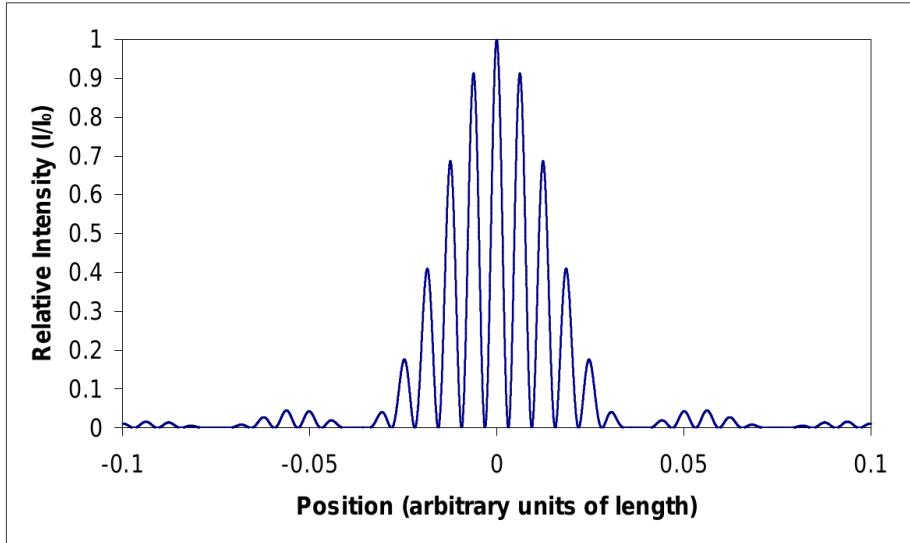


Figure 6: Intensity Pattern of the Double Slit. The Many Peaks of the Narrow Spikes Correspond to the Maxima in $d \sin \theta_m = m\lambda$

light is like the (coherent) soldiers at a parade, whereas light from a light bulb is like the (incoherent) civilian crowd at Times Square.

Although incandescent light sources⁵ can produce diffraction and interference patterns, a laser is better suited to illustrate coherent phenomena.

3 Experiments

3.1 Polarization

Suppose a wave has linearly polarized light characterized by an electric field vector, \mathbf{E}_0 , at an angle, ϕ , relative to the direction of a polarizer. The electric field that is transmitted through the polarizer is $\mathbf{E} = \mathbf{E}_0 \cos \phi$, the component that was not absorbed. Since the energy intensity in the wave is proportional to \mathbf{E}^2 , then the intensity passing through two polarizers will be proportional to the square of the cosine of the angle between the two polarizing axes.⁶

$$I = I_{\max} \cos^2 \phi \quad (4)$$

In this experiment, we use an incandescent unpolarized light source and then polarize it by first passing it through a polarizer. The electric field vector \mathbf{E}_0 will then

⁵This is easily visible in the colors produced by soap bubbles.

⁶This relation is sometimes referred to as Malus' Law, after Etienne Malus, who also discovered the polarization of reflected light in the early nineteenth century.

point along the same direction as the polarizing axis of the first polarizer. As the now polarized light falls on the second polarizer, only the component of \mathbf{E}_0 along the polarizing axis of the second polarizer will make it through unaffected.

As figure 7 illustrates, an electric field strength of only $\mathbf{E}_0 \cos \phi$ will make it through the second polarizer, so that the new intensity will be given by Equation 4 above. A human eye or a photometer can detect this change in light intensity.

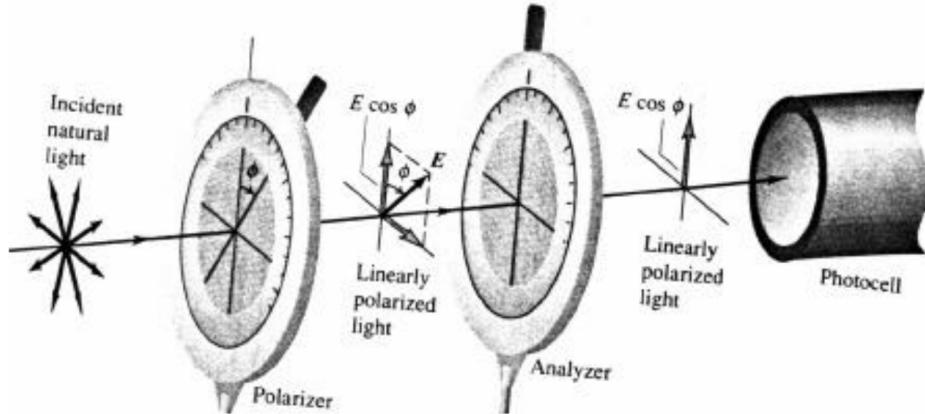


Figure 7: Creation and Detection of Polarized Light

3.2 Double Slit

In this part of the experiment we use a laser instead of incandescent light source. Laser light shines onto a double slit, which produces an interference pattern of bright and dark lines on the wall or a piece of paper beyond the double slit.

We know that the angle, θ_m , at which the m^{th} order bright spot appears is given by

$$d \sin \theta_m = m\lambda \quad (5)$$

For small θ_m (in radians), the approximation $\sin \theta \approx \tan \theta \approx \theta$ is valid⁷. So on a screen far away from the double slit, we observe maxima located close to the optical axis (where the 0^{th} order maximum hits the screen) approximately given by

$$\sin \theta_m \approx \tan \theta_m = \frac{x_m}{L} \quad (6)$$

where x_m is the distance between the m^{th} order maximum and the 0^{th} order maximum, and L is the distance between the double slit and the screen. We then find

$$\frac{x_m}{L} = m \frac{\lambda}{d} \quad (7)$$

⁷See for yourself by punching $\sin 0.1$ and $\tan 0.1$ into your calculator (using radiant measure)!

or for adjacent maxima

$$\frac{\Delta x}{L} = \frac{\lambda}{d} \quad (8)$$

with Δx equal to the distance between the maxima. Since Δx and L are easy to measure, we can straightforwardly determine d (or λ) given λ (or d).

4 Specifics of the Experiment

SAFETY NOTE:

*Although the laser is of low intensity for most situations, it can be dangerous in certain circumstances unless used carefully. In particular, **do not** use the laser in such a way that it can shine into any person's eye (The warning label on the laser states "Do not stare into Beam!"). When you are not actually using the laser, turn it off or close the beam shutter at the front of the laser.*

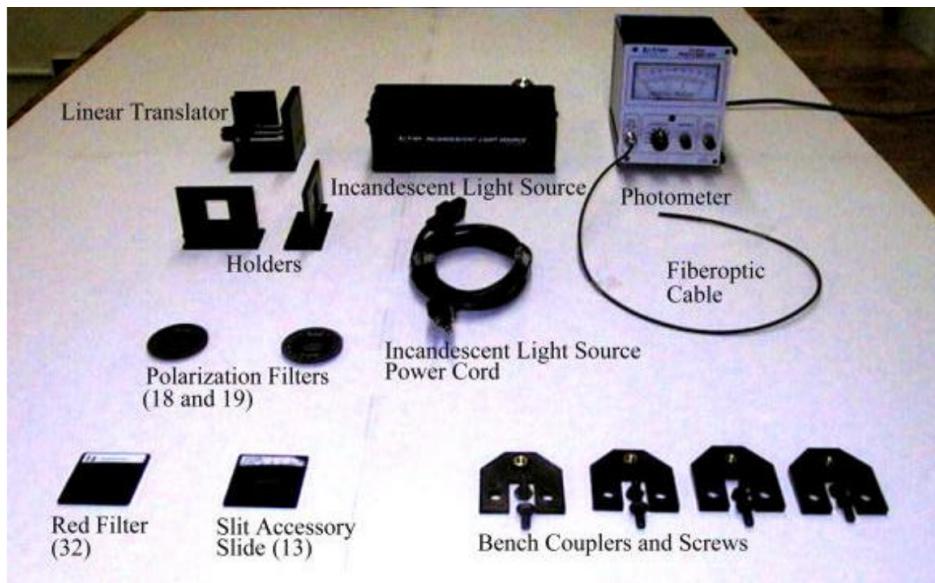


Figure 8: Equipment Components

5 Equipment

Linear Translator: A holder for the fiberoptic cable that can be moved transversely in small increments, allowing a positioning with better than 0.1 mm precision.

Fiberoptic Cable: Cable made from glass, quartz, plastic or a similar material, which carries light waves. The cable uses total internal reflection at the walls to provide very high efficiency (or low losses) between the cable ends.

Photometer: Device to measure the total power contained in light impinging on it. It uses the photoelectric effect to convert the photons to electrons and various electron amplification techniques to produce a measurable electric current.

5.1 Polarization

- Make sure only the dim incandescent ceiling lights in the room are on.
- The linear translator should be placed at the far end of the track opposite the laser. Make sure that the linear translator is in the middle position (at 2.5 cm).
- Place the fiber-optic cable in the hole of the linear translator.
- Place the incandescent light source about 10-20 cm from the linear translator such that the light shines towards the linear translator.
- Place the polarizers (numbers 18 and 19) in the holders.
- Rotate the polarizers such that the white mark on the holder aligns with the 0° angle reading on the polarizer.
- Place the two polarizers between the incandescent light source and the linear translator.
- Put the photometer to the lowest sensitivity (the 1000 setting) and use the zero adjust knob to make sure that the pointer is at the 0 position when the incandescent light source is turned off.
- Turn the incandescent light source back on and adjust the sensitivity of the photometer such that the needle is at the highest position without being at the maximum reading. The value of the sensitivity corresponds to the maximum value on the analog scale.
- Record this measurement as the intensity for 0° difference between the two polarizing axes.
- Measure the intensity for angles between 0° and 90° in increments of 10 degrees. If the readings become too small you may have to switch the photometer to a setting with a higher sensitivity.
- Plot the relative intensity (I divided by I_0 , the intensity at 0°) vs. the \cos^2 of the angle between the axes of the polarizers.
- What is the general shape of your graph? What does this indicate? Did your line or curve pass through the origin?

- What reading do you obtain when the polarizers are at an angle of 90° relative to each other? (This is called the “noise” of your measuring device.) What reading would you expect if there was no noise?
- What is the signal to noise ratio? (The “signal” is the reading at 0° .) What does this tell you?
- If you had a lot of background light, how could you reduce the influence it would have on your results (by changing the setting or by changing your data analysis)?
- If you had aligned the 90° mark instead of the 0° mark of both polarizers with the white mark on the holder would your results have been any different? What relative intensity would you expect if the angle between the two polarizers was 180° ? 270° ?
- What are the main sources of error? Which do you think contributes most?

5.2 Double Slit

- Before you start this section, be sure you have read the laser safety note above.
- Your TA will demonstrate the interference pattern by projecting it on the wall. Note down your observations!
- Remove the incandescent light source, the polarizers, and the holders from the track, and the fiberoptic cable from the linear translator.
- The laser should already be mounted on the laser holder at the far end of the track. Rotate the dial on the linear translator so that it is in the middle position.
- Turn on the laser. It should propagate through the center of the hole in the linear translator, producing a bright red dot on the wall. If the laser is not properly aligned, alter the position of the laser by adjusting *only the back legs* on the laser holder. Make sure you avoid lifting the laser and shining it in your classmates’ eyes.
- Place the red filter (number 32) onto the linear translator. Carefully slide the fiberoptic cable back into the hole in the linear translator so that the front of the cable sits right against the red filter.
- Place the slide with the slits (number 13) in a holder and place the holder right in front of the laser.
- Adjust the slide position in the holder so that the laser propagates through slit pattern A.

- Note the interference pattern on the front of the linear translator.
- Adjust the photometer sensitivity to an appropriate setting.
- Turn the knob on the linear translator. This should cause the intensity reading on the photometer to change.
- Record the position and intensity for each maximum until the linear translator has moved as far as possible, starting from the 0 position of the linear translator. As you turn the knob on the linear translator take care not to shift the position of the linear translator on the bench.
- Measure and record the distance L from the slits to the tip of the fiberoptic cable. Where is the 0th order maximum located? Why might it be located somewhere other than the middle position of the linear translator?
- Plot the position of the maxima x_m vs. the order m . Determine the slope of your graph and use it to calculate the wavelength of the laser light. The slit separation for slit A is 0.250 mm.
- Construct a graph of the relative intensity, I/I_0 , vs. x_m . Explain the shape of your graph.
- Would the maxima be more spread out or closer together if the slit separation decreased?
- Remove the slit slide and its holder from the track and the fiberoptic cable from the linear translator. Rotate the dial on the linear translator so that it is in the middle position.
- Place the unknown slit apparatus on the track approximately 25 cm from the front of the laser so that the laser propagates through the slit pattern marked with a pink dot. You should see a horizontal interference pattern on the linear translator; the light in the center of the interference pattern should propagate through the hole in the linear translator. Note: If the laser does not propagate through the appropriate slit, slide the unknown slit apparatus towards and away from the laser until the laser strikes the slit pattern at the appropriate place. If the light from the pattern does not propagate through the center of the linear translator, alter the position of the laser by adjusting *only the back legs* on the laser holder.
- Again measure the distance between consecutive maxima as above. Also, measure the distance L from the slits to the tip of the fiberoptic cable. Use this data and the wavelength of the laser you determined previously to determine the slit separation.

- What is the purpose of using the red filter?
- Give the main sources of error.

6 Applications

Some chemical molecules have a property called chirality. This means that the mirror image of the molecule is not identical to the original. For example, no matter how good you might be at 3D puzzles, you cannot arrange images of the molecule and of its mirror image to superimpose precisely. A famous example is the DNA double helix, a highly complex molecule, which carries all genetic information. The molecule always twists in one direction, while its mirror image always twists in the opposite direction. Many such chiral molecules are optically active, and rotate the polarization-axis of transmitted or reflected polarized light⁸. If you shine polarized light into such a sample, you will see that the light leaving the sample has its polarization axis rotated a few degrees clockwise or counterclockwise, depending on the substance.

This effect can be used to determine the concentration of chiral molecules in a solution. For example, you can determine the glucose concentration in a blood specimen in a non-chemical way. The calculation of the concentration is particularly simple, since the change in angle of the polarization axis depends only on the concentration of the chiral substance in the solution, the distance the light travels through the sample, and a constant specific for the substance. Calculating the glucose concentration, therefore, requires only a measurement of the rotation in polarization vector and known constants – a simple task for a rudimentary computer.

⁸For instance, the glucose dextrose turns the polarization axis to the right: “dexeter” is Latin for “right”.

7 Lab Preparation Examples

Polarization:

1. You look at the sky through a polarizer and as you turn the polarizer you see that the sky appears darker or lighter, depending on the position of the polarizer. What does this observation tell you?
2. You look at a light source through a polarizer. As you turn the polarizer you see no change in intensity. Is the light emitted by this source polarized?
3. Unpolarized light passes through two polarizers whose polarization axes are aligned. You note down the observed intensity as I_0 . Now you turn the polarization axis of one of the polarizers by 60° . What intensity do you observe now?
4. Now you add, after the second polarizer, a third polarizer. The third has a polarization axis of 30° relative to the second polarizer. What intensity do you observe now?
5. Does the result depend on which of the two polarizers you rotate in question 3?
6. Would your answer to question 5 change if the light source produced polarized light? Explain!
7. You have two polarizers with an angle of 90° between their polarization axes. Adding a third polarizer between these two polarizers will typically produce light after the final polarizer. At what angle (relative to the first polarizer's axis) would you observe the maximum intensity after the third polarizer?

The Double-Slit:

While performing a double slit experiment, with $L = 1.00 \pm 0.01$ m, $d = 0.5 \pm 0.1$ mm, you obtain the following locations for maxima:

order	-4	-3	-2	-1	0	1	2	3	4
Position (mm)	-4.1	-3.2	-2.0	-0.9	0.0	1.3	2.1	3.2	3.9

8. Calculate x for all of the consecutive maxima and use the $2/3$ estimate to determine the uncertainty of x . Now determine the wavelength of the light (including uncertainty).
9. Using the data above, determine x from a plot of position vs. order. Given the wavelength as $\lambda = 600$ nm and $L = 1.00 \pm 0.01$ m, what is the spacing between the two slits?

8 Appendix for Experts: How a Laser works (Not required!)

Quantum mechanics tells us that light is produced when electrons in an atom make a transition from a higher to a lower state, or energy level. Whenever an electron makes a transition between states, the difference in energy is emitted (or absorbed) as a discrete energy package called a photon⁹, which carries (with other photons) the electromagnetic wave. The transition from a higher to a lower state can happen in two different ways. First, it can happen randomly, i.e. the electron just falls down spontaneously. This is what happens when a candle produces light. Electrons that make such random transitions result in emitted photons which are incoherent, or out of phase with one another.

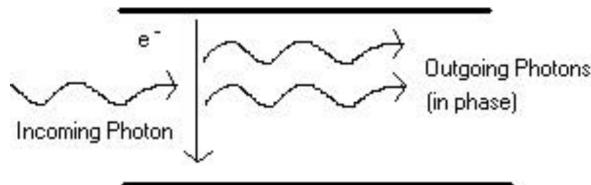


Figure 9: Process of Stimulated Emission

The second way this process can happen is called stimulated emission. If you already have a field of photons that are in phase and have the same frequency, and that frequency corresponds exactly to the transition energy, the emission is no longer random. When an electron falls to the lower energy level, a photon is emitted precisely in phase with the already existing photons, as shown in Figure 9.¹⁰ In our analogy of soldiers marching on parade, new entrants to the parade are required to start walking in step with the others (or they would miss out on the parade).

For stimulated emission to work, the upper atomic state must be full of electrons and the lower one must be almost empty. This is called an inversion, since atoms (at ordinary temperatures) tend to be in their lowest energy state. The inversion is actually achieved by pumping the electrons from the lower state into a third transient state, with an energy even higher than the second laser state, so that the electrons fall

⁹See the experiment on the Photoelectric Effect.

¹⁰Photons are of a particle type called bosons. Such particles are highly “social” and like to be in exactly the same energy state and the same phase with their friends. This contrasts with another particle type, called fermions, which hate being in the same state as their associates (Pauli Exclusion Principle). Electrons are fermions, and this antisocial property accounts for the unique state of each electron in an atom. If electrons were bosons, all electrons in an atom would be in the ground state and we would not have chemistry!

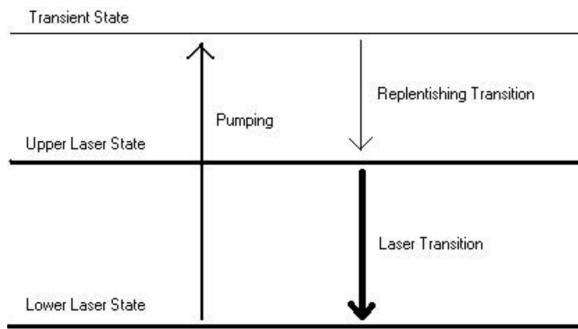


Figure 10: The basic three - level scheme for laser operation.

down to replenish the upper laser state. So real lasers use at least three energy levels; efficient lasers usually use even more¹¹.

Finally, the setup is placed in a cavity with a mirror at one end and a semi-transparent¹² mirror at the other end. This increases the efficiency of the laser by maintaining a high density of laser light within the cavity. The little leakage emitted from the semitransparent mirror is actually what we use as the laser beam.

¹¹Semiconductor lasers (or laser diodes) work somewhat differently, but the underlying concepts of stimulated emission and inversion are the same.

¹²These semi-transparent mirrors are still highly reflective. They reflect about 99% of incident light to keep most of the light inside the cavity.

Experiment 2-9

The Spectrum of the Hydrogen Atom

1 Introduction

In this experiment we will observe the discrete light spectrum one observes from a gas discharge lamp. We will see that the spectrum consists of a collection of sharp, single colored lines. We will be able to measure the wavelength of the light emitted quite precisely, usually better than 1 part in a thousand. Therefore it is crucial to make all calculations to 5 significant figures.

2 Theory

2.1 The Spectrum of the Hydrogen Atom

According to the classical theory of electro magnetism, atoms should radiate continuously and the electrons should fall into the nuclei within a short timespan. Obviously this is not the case since we have stable atoms. Therefore the classical description seems to be wrong.

Niels Bohr suggested an alternative theory for atoms, suggesting that the electrons can only exist in certain energy states and can only jump between these discreet states discontinuously. This discontinuity explains why we see a collection of sharp lines in a gas discharge lamp, and not a continuous spectrum as in a usual light bulb.

Bohr was able to derive the following formula for the energy of the energy levels only using natural constants

$$E_n = -\frac{2\pi^2 me^4 k_e^2}{n^2 h^2} = -\frac{me^4}{h^2} \left(\frac{1}{8\varepsilon_0^2} \right) \frac{1}{n^2} \quad (1)$$

Therefore the energy one gets for a transition from an initial level n_i to a final level n_f is $\Delta E = E_{n_i} - E_{n_f}$.

This energy is then set free in the form of light and one can relate it to the wavelength via

$$\frac{1}{\lambda} = \frac{f}{c} = \frac{\Delta E}{hc} = \frac{E_{n_i} - E_{n_f}}{hc} = R \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right) \quad (2)$$

where R is called the Rydberg constant, which is given by

$$R = \frac{2\pi^2 me^4 k_e^2}{h^3 c} = 1.0974 \times 10^7 \text{ m}^{-1} \quad (3)$$

2.2 Resolving a Light Spectrum with a Grating

A grating is a collection of small parallel slits. In our case the slits are so fine that we cannot see them any more. A grating has the nice property that light of different wavelength gets diffracted in different directions. Therefore one can resolve a continuous light spectrum (containing light of various wavelengths) into its components.

How does a grating do this trick? To understand this it is sufficient if we look at only two slits with a distance d between them. (d is also called the lattice constant)

We look at two light rays coming from the two slits. If the two rays travel perpendicular to the two slits the two waves will always be in phase, since they travel the same distance. But if the two rays propagate at an angle θ to the normal the two rays will have to travel different distances to reach the eye. (we always assume that the two rays are parallel, which is a good approximation since the eye is almost infinitely far away compared to the distance between the two slits.)

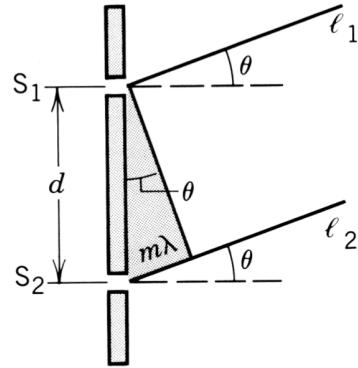


Figure 1: Geometry of the Double Slit

As you may convince yourself by looking at the graph, the difference in path between the two rays is

$$\Delta l = d \sin \theta \quad (4)$$

As in the lab about standing waves we can have different possibilities how these two rays add up. They can add up to form a node or they can add up to form an anti node (or something in between). To form a node every mountain shall be made lower and every valley shall be made higher. Therefore each hill from one wave should meet a valley from the other ray and vice versa. To get an antinode and therefore a

maximum in the total motion of the wave every hill should meet a hill and every valley should meet a valley.

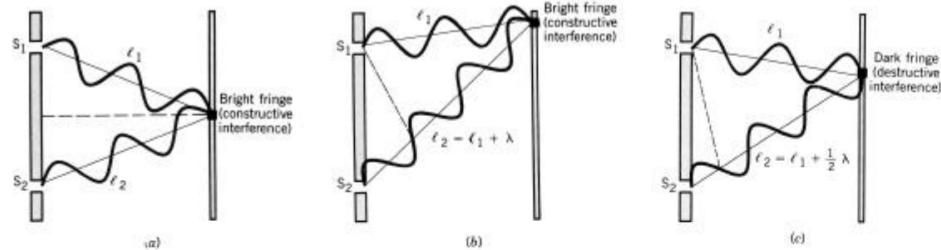


Figure 2: Young's Double Slit

How do we now achieve maximum intensity (node)? We always get this if the two waves are in phase which means that there fits an exact number of wavelength m in the difference between the two waves, i.e.

$$\Delta l = m\lambda \quad (5)$$

m is also called the order. If we now put the two equations for Δl together, we get

$$d \sin \theta = m\lambda \quad (6)$$

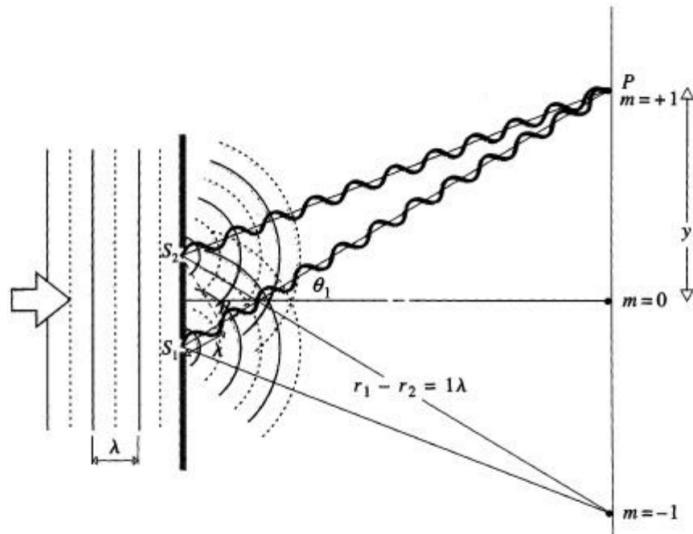


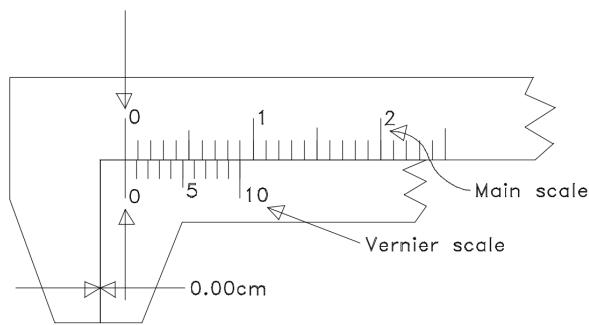
Figure 3: Interference of Light Coming from Two Small Slits

This tells us that given the order m we can determine the wavelength of the light we see through the grating by reading off the angle. This also works the other way

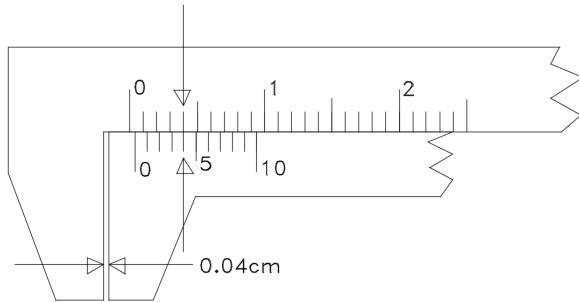
around. (We can also get m by count up the orders starting with 0 at an angle of 0 degrees, and then count how often a particular color repeats as you go to higher or lower angles.)

2.3 How to read a Vernier Scale

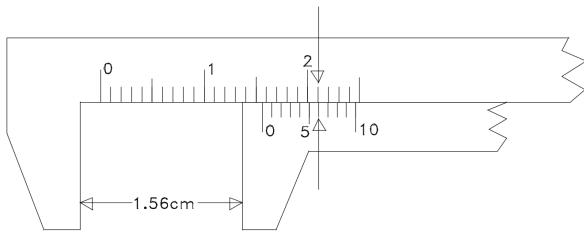
With a usual ruler you can read distances up to a mm resolution. But if you would like to have a measuring device to read with a 0.1 mm resolution you have a problem. In principle you can scratch such a fine scale in the ruler, but in practice you can easily imagine that due to the width of the marks themselves such a scale would be almost impossible to read. That is why we use a trick and introduce a Vernier scale. A Vernier scale works in the following way:



- a) If the zero points of the vernier and main scales are aligned, the first vernier division is $1/10$ of a main scale division short of a mark on the main scale; the 2nd vernier division is $2/10$ of a main scale division short of the next mark on the main scale, etc., and the 10th vernier mark coincides with a main scale mark.



- b) If the vernier scale is now moved to the right until the 4th vernier division is lined up with the nearest main scale division to its right, the distance the 0 point of the vernier scale has moved past the 0 point on the main scale is $4/10$ of a main scale division. Thus the vernier scale gives the fraction of a main scale division that the zero point of the vernier has moved beyond any main scale mark.



- c) In this last example, the zero line of the vernier scale lies between 1.5 cm and 1.6 cm on the main scale. The fraction of a division from 1.5 cm to the zero line can be determined as follows: Line number 6 of the vernier scale coincides with a line on the main scale, so the zero line on the vernier has moved $6/10$ of a main scale division away from the main scale line 1.5. Hence, the reading, to the nearest hundredth, is 1.56. The maximum uncertainty is now ± 0.01 cm, corresponding to the smallest division on the vernier scale. The use of the vernier has reduced the maximum uncertainty of the caliper measurement from ± 0.1 cm (without the vernier) to ± 0.01 cm.

In this experiment we will not use a simple Vernier scale since we do not measure distances, but we will use an angular Vernier scale to measure sub-divisions of angles. If you look at the main scale you will see that the scale is divided in half-degree steps. (You have slightly longer marks for the full degrees and shorter marks for the half degrees.) For angles the next smaller unit is not $1/10$ th of a degree but minutes, which is $1/60$ th of a degree. But since we can already read the scale up to $1/2$ degree we need not have a Vernier scale with 60 devision but with 30 divisions.

So you first use the 0 mark on the Vernier scale to read off the angle in degrees and if you are in between 0 and 30 minutes or 30 and 60 minutes. Then you look which of the marks on the Vernier scale exactly matches a mark on the other scale. So e.g. if the 0 mark on the Vernier scale is between 20.5 and 21 degrees and the 13th mark on the Vernier scale matches a mark on the degree scale we would have a reading of $20 + 1/2 + 13/60$ degrees or 20 degrees and 43 minutes.

3 Experiments

3.1 Adjusting the Spectrometer

The first part of the experiment will basically include the procedure to set up the equipment. This should be done with as much care as possible. Only then we will be able to measure the wavelength on the limit of our apparatus. If you don't set up the spectrometer correctly you will get systematical errors, which screw up your data.

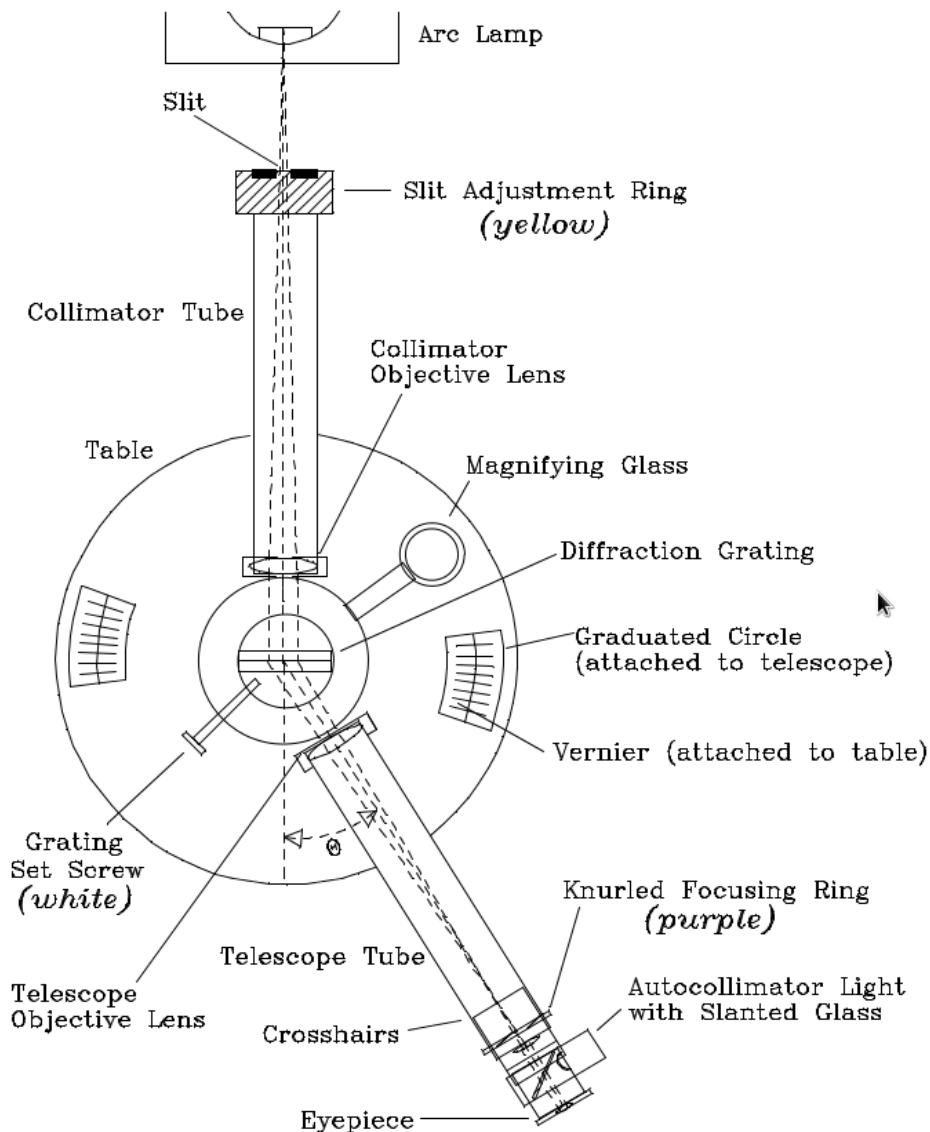


Figure 4: Schematic of the Spectrometer

3.2 Obtaining the Lattice Constant

After adjusting the spectrometer we will measure the yellow line of a Helium discharge lamp. Since we know that the wavelength of this light is $\lambda = 5.8756 \times 10^{-7}$ m, we can determine the lattice constant of the grating quite accurately. Even though the lattice has written on it 600 lines/mm (which is only an approximate value anyway), we want to get the lattice constant with 5 relevant digits and not just 3, and this means that we have to measure it!

3.3 Measuring the Spectrum of Hydrogen Atoms

With the lattice constant determined in the previous part we now measure the wavelength of the light emitted from the Hydrogen discharge lamp.

4 Step by Step List

4.1 Adjusting the Spectrometer

- Take the grating out off the holder and close the green knob.
- Rotate the yellow knob such that the slit is about half open.
- Look through the eyepiece and turn the purple focusing ring until you see a sharp image of the slit.
- Loosen the red knob and move the telescope tube until the cross hair is in the middle of the slit. Tighten red knob.
- Now open the green knob and turn the tabletop such that the 0 mark from the Vernier scale with the magnifying glass is lined with either the 180 or the 360 from the outer scale. (Always use only this Vernier scale and don't switch to the other one in between)
- Close the green knob (and don't open it again for the rest of the experiment!).
- Now you can fine adjust the relative position of the inner and outer scale by turning the blue knob. Make sure that the lineup between the 0 on the Vernier and the 180/360 is done as careful as possible. (For some of the spectrometers there is a small mark to the left of the 0 mark on the Vernier scale. Make sure that you line up the 0 mark and not this extra mark with the 180/360.)
- Now put the grating in the holder such that it is perpendicular to the telescope tube-collimator tube line. Close the white screw to lock the grating.

4.2 Obtaining the Lattice Constant

- Switch on the Helium lamp and line the spectrometer up such that you can see the slit well illuminated by the lamp as you look through the spectrometer.
- Put the black cardboard over the front end of your collimator tube and you can use the black piece of cloth over the spectrometer to block light from the surrounding. Also this part and most important the next one should be performed

in the dark, therefore switch off the light and use the lamps provided (they are supposed to be that dim) to read the scale.

- Open the red knob and move the telescope tube to the left until the crosshair is in the center of the yellow line. (You should first see a few blue and green lines, then the isolated yellow line and then red lines. The yellow line should be somewhere around 20 degrees)
- Note down the angle in degrees and minutes at which you see the 1st order of the yellow line. Do the same on the right side and average these two numbers.
- Use the average and plug it into the grating equation ($m = 1$) to determine the lattice constant d (at least 5 relevant figures).
- How many lines per mm does this lattice constant correspond to?
- Can you also see the 2nd order and 3rd order yellow lines on either side?

4.3 Measuring the Spectrum of Hydrogen Atoms

- Now switch off the helium lamp and set up the spectrometer for the hydrogen lamp.
- The light purple line you see in the middle is the 0th order and not yet one of the lines we are going to measure.
- There are 4 visible lines for the Hydrogen spectrum. One red (furthest out), one greenish blue, one purple blue and one dark purple. The dark purple line is very faint and you may not be able to see it, so don't despair.
- Measure the angles of the 4 lines on both sides and average the angles.
- How many orders do you see on either side? (look e.g. for the red line and count how often it appears as you go further out)
- Do the higher orders overlap? (i.e. does a new order start before the old one ended?)

4.4 Analyzing the Data

- Put together formulas (2) and (6) and solve the resulting formula for n_i .
- Use your data and the value for R from above and $n_f = 2$ to determine n_i .
- n_i should be an integer number labeling from which initial atomic shell the electrons in the hydrogen atoms fell to the 2nd atomic shell.

- Are the results integers (or close to them)?
- Are the numbers (integers) the one they are supposed to be?
- For which color did the electrons jump from the lowest shell? Explain why you could have predicted that anyway!
- Within how many % have you been able to measure the data? Compare that to other experiments you have performed in this course!
- Discuss briefly the factors that determine the uncertainty in your measurements. Which of them are random, which systematic?

5 Lab Preparation Examples

Hydrogen Spectrum:

1. What is the wavelength of emitted if an electron in an hydrogen atom makes a transition from $n_i = 3$ to $n_f = 1$?
2. For a hydrogen atom give all the transitions that fall within the visible range of the light spectrum (i.e. between $\lambda_{\min} = 400 \text{ nm}$ and $\lambda_{\max} = 800 \text{ nm}$). Give the transitions in the following form: A transition from energy level $n_i = 3$ to $n_f = 1 : 3 \rightarrow 1$.

Hint: It may be simplest to calculate the energies $\lambda_{\min} = 400 \text{ nm}$ and $\lambda_{\max} = 800 \text{ nm}$ correspond to. Then make a list of E_n and choose the differences between the energy level that fit in the calculated range.

3. What wavelengths do the transitions in the previous problem correspond to?
4. If you look at the light of an object at high temperature (e.g. a normal light bulb), you will find that it usually emits a continuous spectrum without any gaps or lines in it. (This is due to the effect that the atoms in a solid (or plasma) are so close together that they disturb each other so strongly that the initially sharp transition lines of the atoms get so broadened that they appear as a continuous spectrum.)

But if you now look at the spectrum of the sun (or another star), you will find a continuous spectrum that has a number of sharp black lines, so no (or almost no) light of that frequency reaches you; the light of that frequency is missing in the spectrum. These black lines are at exactly the same positions where you would see the colored lines in emission spectra (that is what we do in lab) of the elements present on the sun. For example since hydrogen is present on the sun we will find black lines at the wavelength calculated in 3.

Can you give a short explanation of how we get these inverse spectra (compared to the ones we see in the lab from emission spectra) from the sun?

Spectrum with a Grating:

5. What is the lattice constant d if the lattice has 400 lines per mm?
6. If you have a grating with 600 lines per mm at what angle would you observe the $m = 1$ maximum for a wavelength of 600 nm?
7. If you have $d = 10^{-1} \text{ m}$ and $\lambda = 1 \text{ cm}$. at what angles will you see the maxima for $m = 1, 2, 3$?

Experiment 2-10

Absorption of Beta and Gamma Rays

Objective: To study the behavior of gamma and beta rays passing through matter; to measure the range of beta-particles from a given source and hence to determine the endpoint energy of decay; to determine the absorption coefficient in lead of the gamma radiation from a given source.

Safety Note: The radioactive sources used in this experiment are extremely weak and therefore can be handled without any danger of overexposure to radiation. However, it is always prudent not to handle any radioactive source more than is necessary.

1 Discussion

All nuclei heavier than lead (and many isotopes of lighter nuclei) have a finite probability of decaying spontaneously into another nucleus plus one or more lighter particles. One of the decay products may be an alpha-particle (two protons and two neutrons—the stable nucleus of a helium atom). Alternatively, a nucleus with more neutrons than it can maintain in stability may decay by emission of an electron from the nucleus (beta-decay) which corresponds to the conversion of a neutron to a proton. These electrons may emerge with a kinetic energy of up to several MeV.

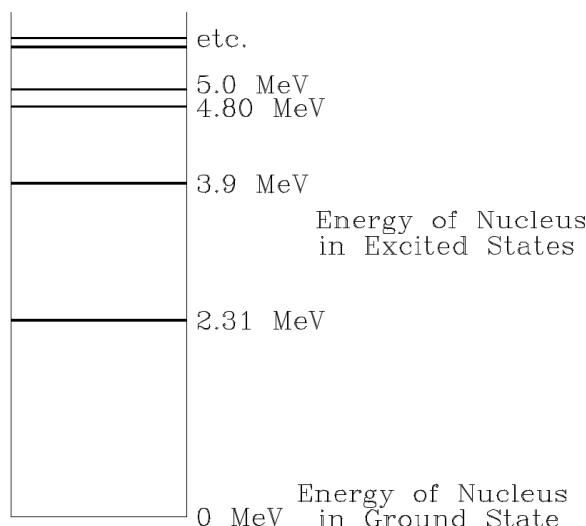


Figure 1: Typical Nuclear Energy Level Diagram.

After alpha or beta decay, the residual nucleus may be left in an excited state. In this case, a transition to a state of lower energy of the same nucleus will occur almost immediately with the emission of a photon (gamma ray). The spectrum of photons emitted from the various excited states of a nucleus will have discrete frequencies ν , corresponding to transitions $\Delta E = h\nu$, between discrete energy levels of the nucleus. The gamma ray spectrum from an excited nucleus is thus **analogous** to the line spectrum of visible radiation from an atom due to excited electrons, with the notable difference that the MeV energy changes of the nucleus are approximately 10^6 times as large as energy changes in transitions between atomic states (where $\Delta E_{\text{atomic}} \approx$ several eV). See Figure 1.

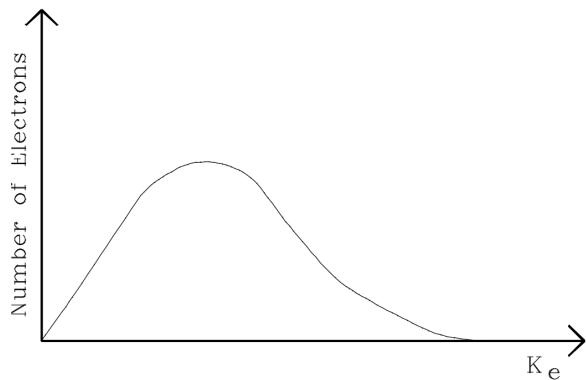


Figure 2: Beta Ray (Electron) Energy Spectrum

In early experiments on beta-decay, it was observed that each decay was not a simple one in which an electron and the recoil nucleus came off with equal and opposite momentum. The electrons, in fact, were emitted with a continuous spectrum of energies (Figure 2). It was subsequently suggested by Pauli and Fermi that in each decay, another particle of zero mass and charge, called the neutrino, was emitted. Experimental verification of the neutrino has been obtained by observation of its rare interaction with matter.

2 Detection of Charged Particles (the Geiger Counter)

When an energetic charged particle traverses matter, it will give electrostatic impulses to nearby electrons and thus ionize atoms of the material. Most methods for detecting nuclear particles rely on observing the results of this ionization. In a photographic emulsion, a trail of ionized grains will show up when the film is developed. A trail of droplets or bubbles forms along the wake of ionization left by a charged particle passing through a cloud or bubble chamber. In a scintillation counter, ionized molecules will very rapidly radiate visible light. A charged electroscope will be discharged by ion

pairs created in air. In the **geiger counter**, which will be used as a detector in this experiment, the ionization produced by a charged particle causes a violent electrical discharge.

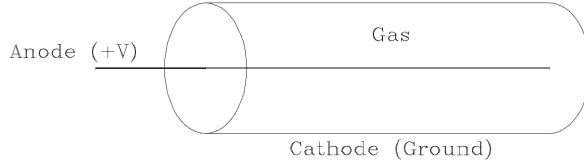


Figure 3: The Geiger Counter.

As shown in Figure 3, the counter consists of a metal cylinder (cathode) with insulating ends supporting a fine axial wire (anode). When a charged particle ionizes the gas (e.g., argon) in the tube, electrons will move toward the positively charged anode wire, accelerated more and more by the rapid increase in electric field near the wire. When an electron acquires kinetic energy greater than the ionization energy of the gas molecules, it can create by collision a new ion and electron, which in turn can accelerate and create another ion, etc., thus initiating an avalanche of charge. The process, called the **Townsend avalanche**, is possible only if the voltage maintained between anode and cathode is sufficiently high.

The basic counter circuit, shown in Figure 4, supplies a positive high voltage of up to 900 volts to the center wire. When an avalanche occurs, current flows through R , the counter side of R drops in potential, and this negative pulse is fed through C to a stage of amplification and then to a scaling device.

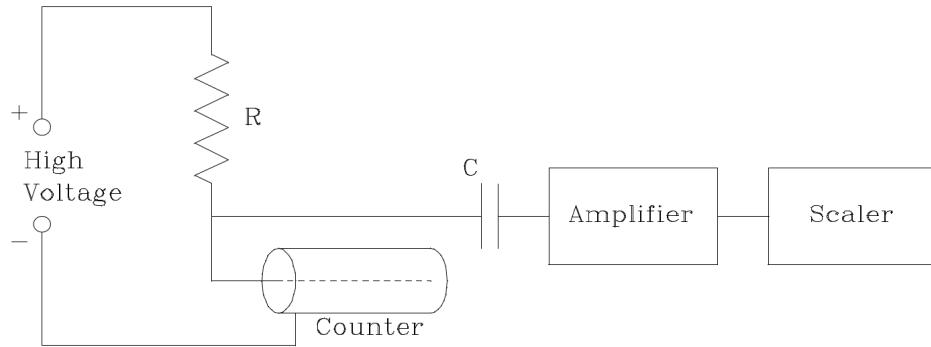


Figure 4: Geiger Counter Circuit.

For a large enough value of R , the voltage drop across R would be sufficient to stop the discharge after a count has been recorded. This method of quenching the discharge, however, has the disadvantage of creating a long time constant, $\tau = RC$, for reestablishing voltage across the counter, i.e., of creating a long “dead time” during which the passage of another particle cannot be detected. The quenching can be

achieved for a lower value of R by adding halogen vapor to the gas to help absorb ion pairs. Nevertheless, the resultant dead time is substantial and must be taken into account when the particle flux is high.

3 Energy Loss and Range of Beta Particles

Because of its ionizing action (Figure 5), a *charged*, incident particle in matter will continuously lose kinetic energy, and the particle will subsequently come to rest after traversing a path length called its **range**. *For a particle of known charge and mass, there will be a unique range associated with each incident energy.* A formula can be theoretically deduced for the rate of energy loss—and hence the range—of a particle (of known mass, charge, and initial velocity) in a particular “stopping” material (of known electron density and ionization potential).

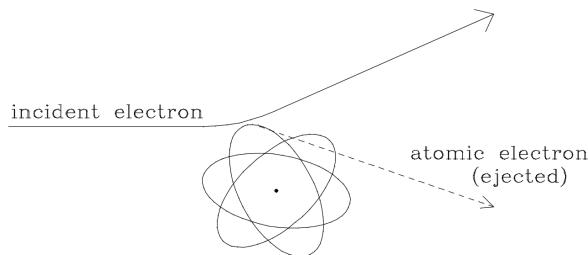


Figure 5: Ionizing Action of Incident Electron.

In each interaction with atomic electrons, however, an incident electron may be scattered through small or relatively large angles, and as it traverses the material it may follow a rather tortuous, winding path (especially at low energies). Therefore, the actual path of the electron may be considerably longer than the observed distance that it penetrates into the material. For this reason, the incident electron range is not sharply determined, and the theoretical calculation is of limited usefulness for electrons of less than 1 MeV energy.

In this experiment, therefore, we will use an approximate empirical relationship between range and energy for low energy electrons:

$$r = \frac{0.412 \text{ g/cm}^2}{\rho} E^{1.29} \quad (1)$$

where r is in cm, E is in MeV, and ρ is the density of the stopping material in g/cm^3 . Note that the density of Aluminum is 2.702 g/cm^3 .

(This result is described in: L. Katz and A. S. Penfold, “Range-Energy Relations for Electrons and the Determination of Beta-Ray End-Point Energies by Absorption,” Revs.Modern Phys. **24**, 1 (1952).)

4 Absorption of Gamma Rays

Gamma rays, or high-energy photons, can interact with matter by three distinct processes:

- 1) Compton Scattering: This refers to a photon-electron collision in which the energy lost by the scattered photon is given to the recoil electron.
- 2) Photoelectric Effect: The photon is absorbed by the atom as a whole, releasing an electron with kinetic energy equal to $E_\gamma - E_b$, where E_γ is the photon energy and E_b is the relatively small binding energy of the electron in the shell from which it is released.
- 3) Pair Production: If the photon has energy greater than 1.02 MeV, it can create an electron-positron pair in the neighborhood of a nucleus. The radiative source used in this experiment does **not** emit photons with energy greater than 1 MeV.

The probability of each of the three processes taking place in a given thickness of material depends on the energy of the photon and the atomic structure of the material. The *total* probability for interaction of photons in Pb, i.e., the sum of the probabilities of the three processes, varies with photon energy as indicated in Figure 6. The ordinate plotted on the graph is μ , the **total linear absorption coefficient** in units of cm^{-1} . It is defined by the equation:

$$\frac{dN}{dx} = \mu N \quad (2)$$

where N is the number of incident photons and dN is the number removed from the beam (i.e. absorbed) in an absorber of thickness dx (in cm). Note that dN and dx are the calculus equivalents of infinitesimally small values of ΔN and Δx , respectively. As in any process where the rate of decrease is proportional to the number present (such as the discharge of a capacitor), the solution of this differential equation is:

$$N(x) = N_0 e^{-\mu x} \quad (3)$$

where $N(x)$ is the number of photons passing through x cm of absorber and $N_0 = N(x)$ at $x = 0$, and e is the base of natural logarithms.

5 Procedure

5.1 Adjustments and Measurement of Errors in Counting

5.1.1 High Voltage Variations

Every Geiger tube that is in good working order has a plateau region in which its counting rate is relatively insensitive to changes in the high voltage supply. This region

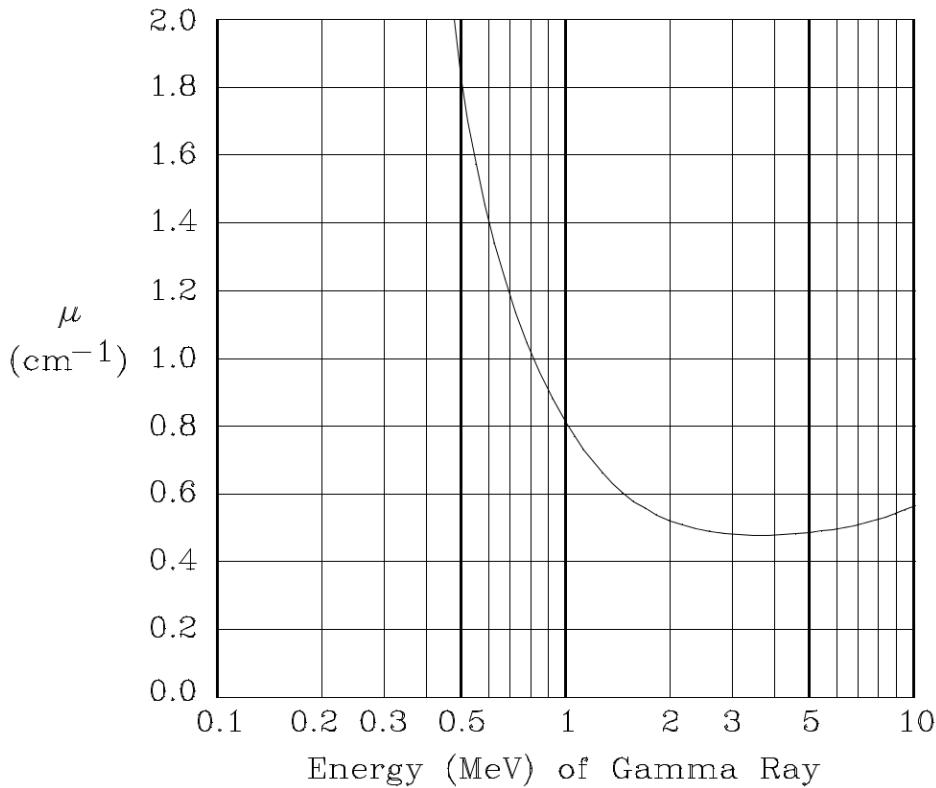


Figure 6: Linear Absorption Coefficient for gamma rays in lead as a function of energy.

can be found by putting a source under the tube and increasing the high voltage from its lowest value until the tube just begins to count. One can then take 15 second counts while raising the voltage in 50 volt steps. The curve of counting rate vs. voltage levels off, i.e., the count rate rises by less than 10% for a 100-volt increase. Do not raise the voltage further as this may damage the tube due to continuous breakdown. Set the high voltage to a value on this plateau for the remainder of the experiment. If this procedure is followed correctly, high voltage variations may be ignored as a source of error.

5.1.2 Statistical Accuracy

Particles decay randomly in time from a radioactive source (over a period short compared to the half-life). The probability distribution for measuring a given number of these counts in a given time interval is an almost bell shaped curve (a Poisson Distribution) centered around N_0 , the most probable value. The distribution will have a standard deviation about the peak of $\pm\sqrt{N_0}$. If N counts are measured in an interval, the best estimate of the error is $\pm\sqrt{N}$.

Note that the magnitude of the statistical error—your uncertainty in the measurement—

increases significantly for trials involving a very small number of counts. For a high counting rate in a given interval, say, 900 counts in one minute, the estimated error will be 30 counts per minute or 3.3% error. However, for a much lower counting rate in the same interval, say 25 counts, the error of ± 5 counts per minute amounts to a 20% statistical error. In order to achieve the same precision as in the first case, it would be necessary to collect 900 counts—in other words, to take a 36-minute measurement. While such a long measurement is impractical for this lab, you should keep in mind the relationship between a small number of counts and higher errors, and do your best to minimize these errors by taking longer measurements when necessary.

5.1.3 Measuring Background Radiation

In order to make accurate counting measurements of the sources, it is necessary to know the counting rate due to natural background radiation (mostly cosmic rays coming through the earth's atmosphere). Additionally, there will be some excess counts due to the Cs-137 gamma sources nearby in the room, whose gamma rays can pass through the side of the detector. At a distance of 30 cm, for example, the Cs-137 source contributes roughly as many counts as the natural background radiation (doubling the distance would reduce its contribution to one-fourth this level). It's best to try to minimize these secondary effects—by keeping your detector far from others' sources, and by shielding your own Cesium source with the lead sheets when not in use), but it's even more important to try to keep these background effects *constant*. If all of your data is shifted by roughly the same constant amount, then it is possible to isolate the results you're interested in by subtracting out this constant background.

5.2 Range of Beta Particles

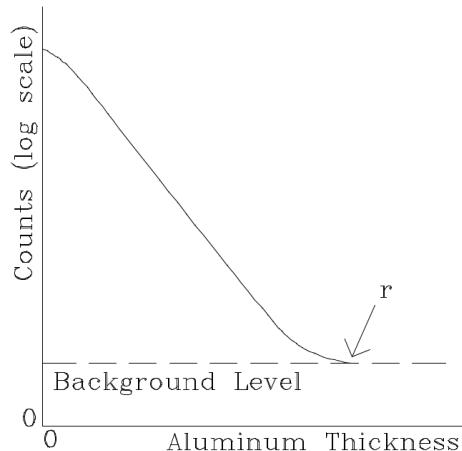


Figure 7: Beta Absorption Curve

Thallium 204 beta decays to Lead 204 with a half-life of 3.9 years. The range of the most energetic of the decay electrons can be determined by placing aluminum foil absorbers between the source and the geiger counter.

A typical absorption curve is shown in Figure 7. The maximum range r is the point where the absorption curve meets the background. You should start by making a careful measurement of background, and you should repeat this measurement after taking the absorption curve to check for constancy.

Place the Thallium source on the second shelf below the detector, as shown in Figure 8, to maximize the number of counts while leaving enough room to stack aluminum absorbers. Begin taking measurements for the absorption curve, adding aluminum foils until the counting rate reaches the background level. (*Note that the thickness of the aluminum absorbers is marked on the foil in **mils**, or thousandths of an inch, not millimeters.*)

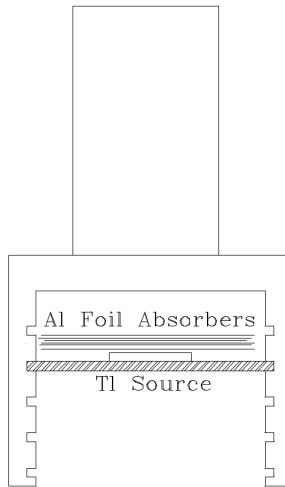


Figure 8: Beta Absorption

Make a table of the results, including the background level, and include estimates of the statistical error in each measurement. Note that you *should not* subtract background from your data for this experiment. Plot the results on the semi-log paper provided (if your measurements are of different durations, you will need to scale them to a constant duration which is convenient for graphing).

Determine the approximate value of the maximum range from the graph, and use Eq. (1) to compare your result with the value of 0.765 MeV for the maximum beta energy for Thallium 204 as measured in a magnetic spectrometer.

5.3 Absorption of Gamma Rays

The source for this part of the experiment is Cesium 137, which decays with a half-life of 30 years to Barium 137 with the emission of a 0.52 MeV beta ray (Figure 9). The resulting Barium is in an excited state and decays by emitting a 0.662 MeV gamma ray almost instantaneously. Lead absorbers are used for the gamma absorption study. They are thick enough(0.062 inches) so that one absorber will stop all the betas.

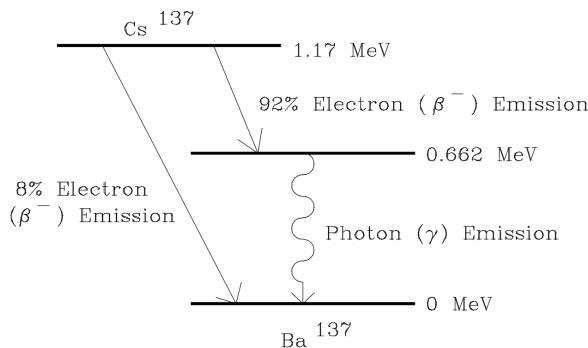


Figure 9: Cesium 137 Decay Scheme

The gamma rays are detected by means of the same geiger counter used in the previous part. Note that the efficiency of the geiger counter for detecting photons is much less than for detecting the beta rays, since it depends on a collision of the photon with the gas or wall of the counter, resulting in the emission of an electron, which in turn initiates the discharge.

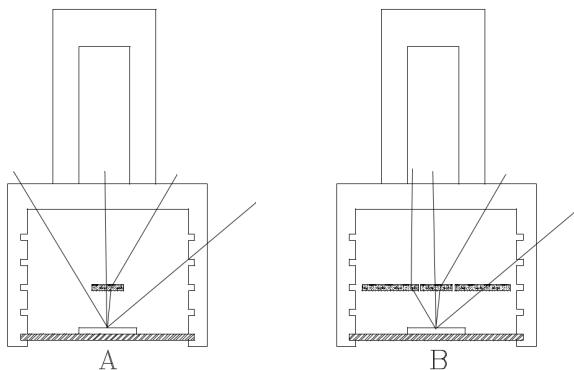


Figure 10: Scattering Effects: In (A), the gamma ray on the left passes outside of the detector tube. In (B), it can be seen how increasing the area of lead absorber can cause the same gamma ray to scatter *into* the detector.

Figure 10 represents one of the effects of gamma ray scattering in the lead sheets. Increasing the *area* of lead through which the gamma rays pass tends to *increase*, rather

than decrease, the number of counts one measures, since gamma rays which otherwise would not have entered the detector may now be scattered into it.

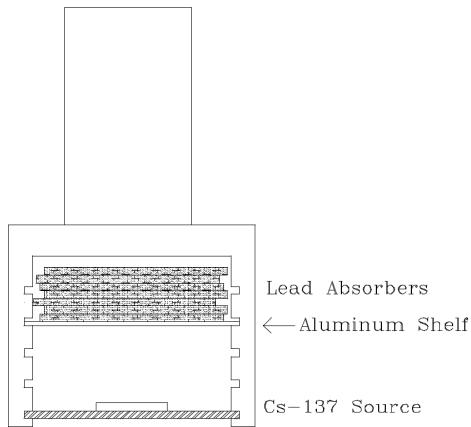


Figure 11: Absorption of Gamma Rays Set-Up.

The arrangement shown in Figure 11 is designed to reduce the effect of such scattering. By keeping the lead sheets high above the source, one reduces the excess area exposed to gamma rays, and one reduces the effective difference in area between the top and bottom sheets as well.

Take measurements for an absorption curve. Note how this differs fundamentally from beta absorption: there is no maximum range for gamma rays passing through lead; rather, one expects to lose a fixed *fraction* of the remaining gamma rays passing through each successive layer of lead absorber. For this reason, the theoretical absorption curve never intersects the background curve. When the absorption curve is plotted on semi-log graph paper with *background subtracted*, the exponential decay curve appears as a straight line.

Make a table of the data and estimated errors. Plot the data on semi-log graph paper as in Part II.

Using a straight line fitted to the data, find the absorption coefficient μ by choosing two points on the line a distance x apart, substituting their y values for N_0 and $N(x)$ in Eq. (3) and solving for μ . Use this value and Figure 6 to estimate the energy of the gamma rays, and compare with the accepted value. What factors limit the precision of this estimate?