

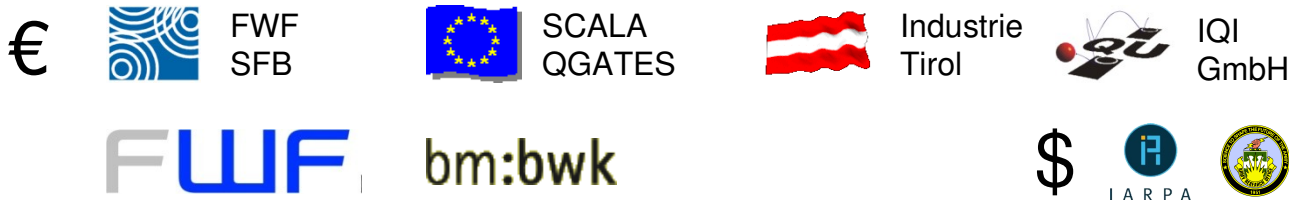
Quantum computing with trapped ions



Hartmut Häffner

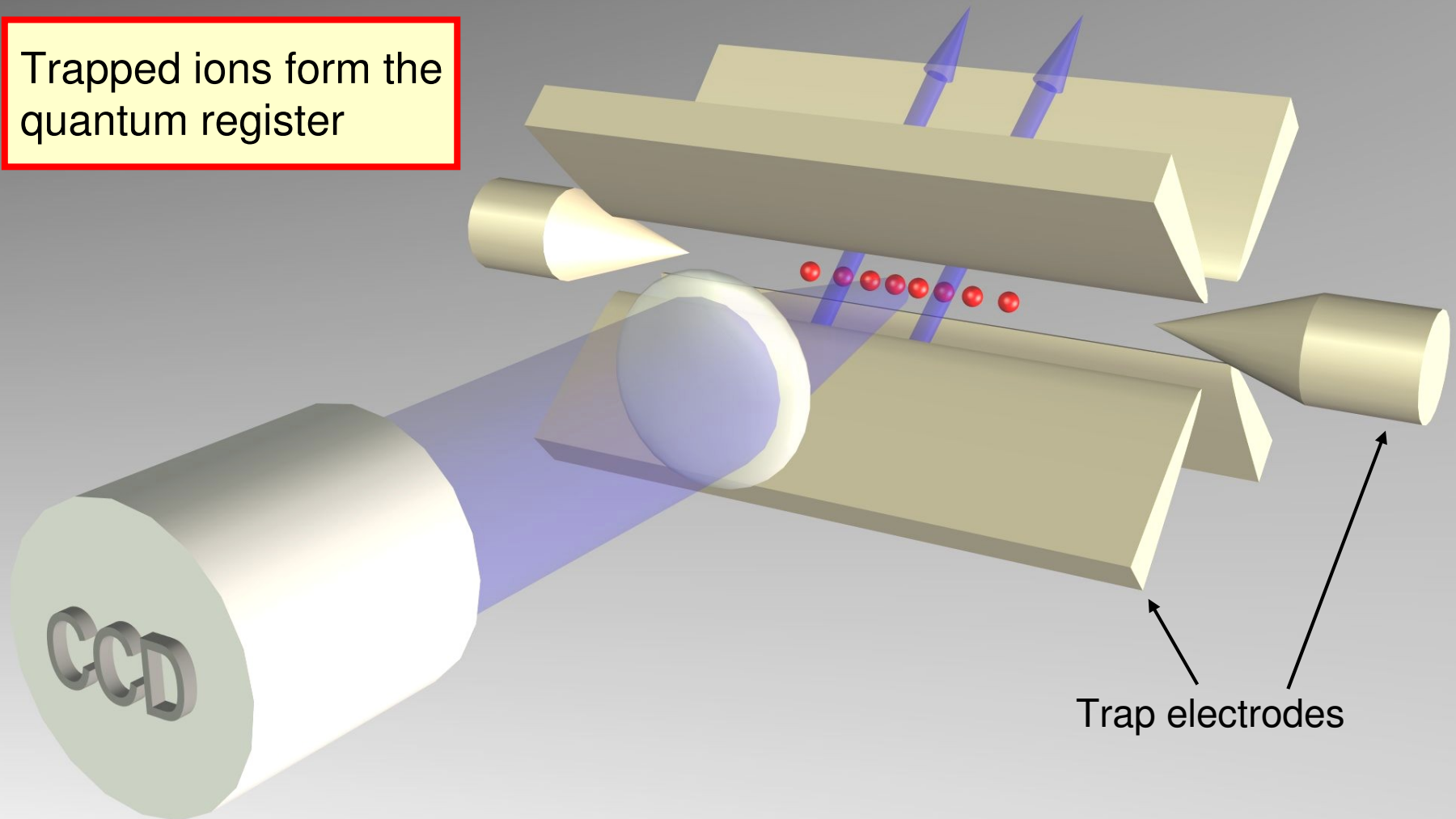
*Institute for Quantum Optics and Quantum Information, Innsbruck and
Department of Physics, University of California, Berkeley, USA*

- Introduction to ion trap quantum computing
- Single ion addressing approach
- Coherent operations with global interactions
- Conclusions

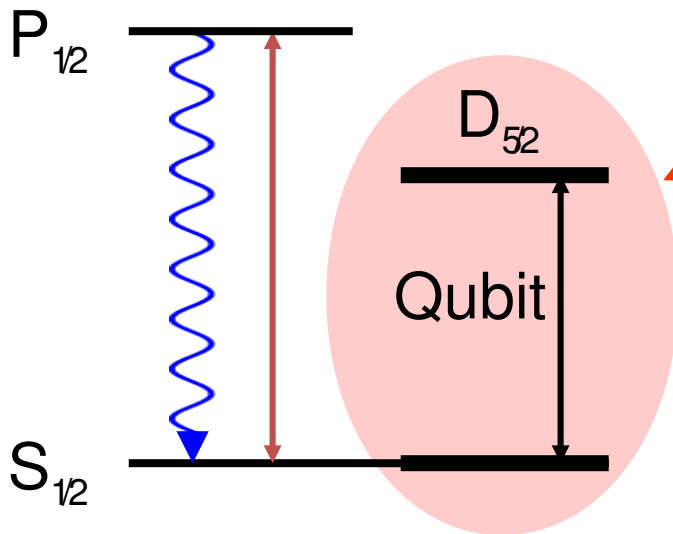
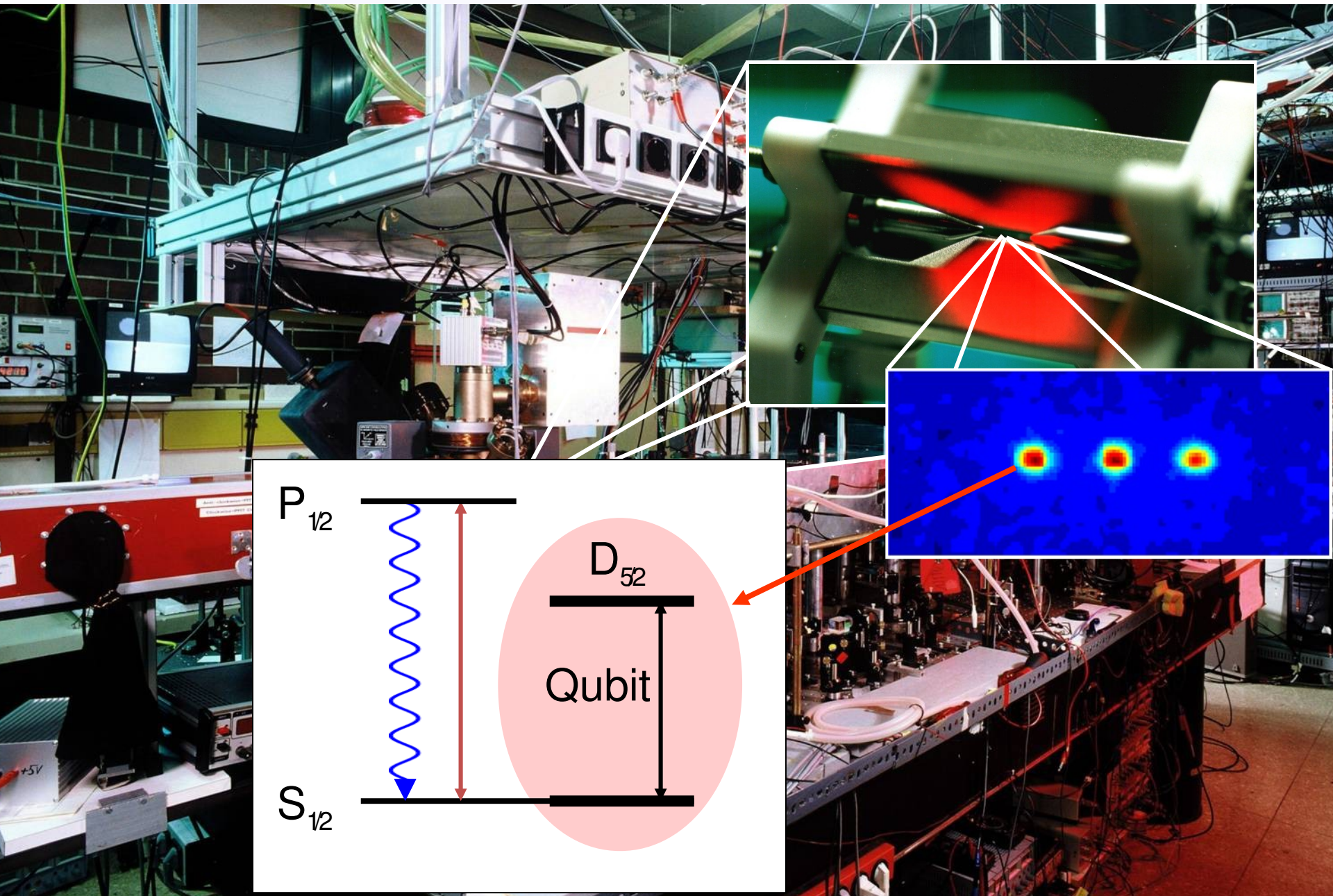


Baltimore, June 5th 2009

Trapped ions form the
quantum register



The hardware



Requirements for quantum computing

Classical computer

- Initialization
- 1-bit operations (NOT)
- 2-bit gates (e.g. NAND)

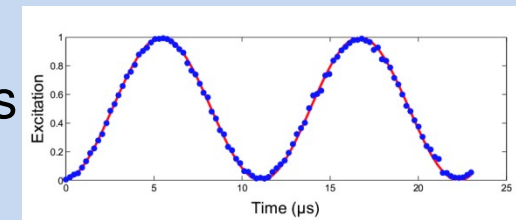
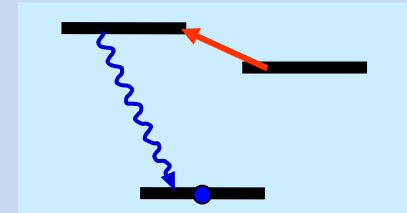
Computational space:

00
01
10
11

- Read out
➡ result

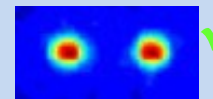
Quantum computer

- Initialization
- 1-qubit rotations
➡ superpositions
- 2-qubit gates (CNOT gate)
➡ entanglement



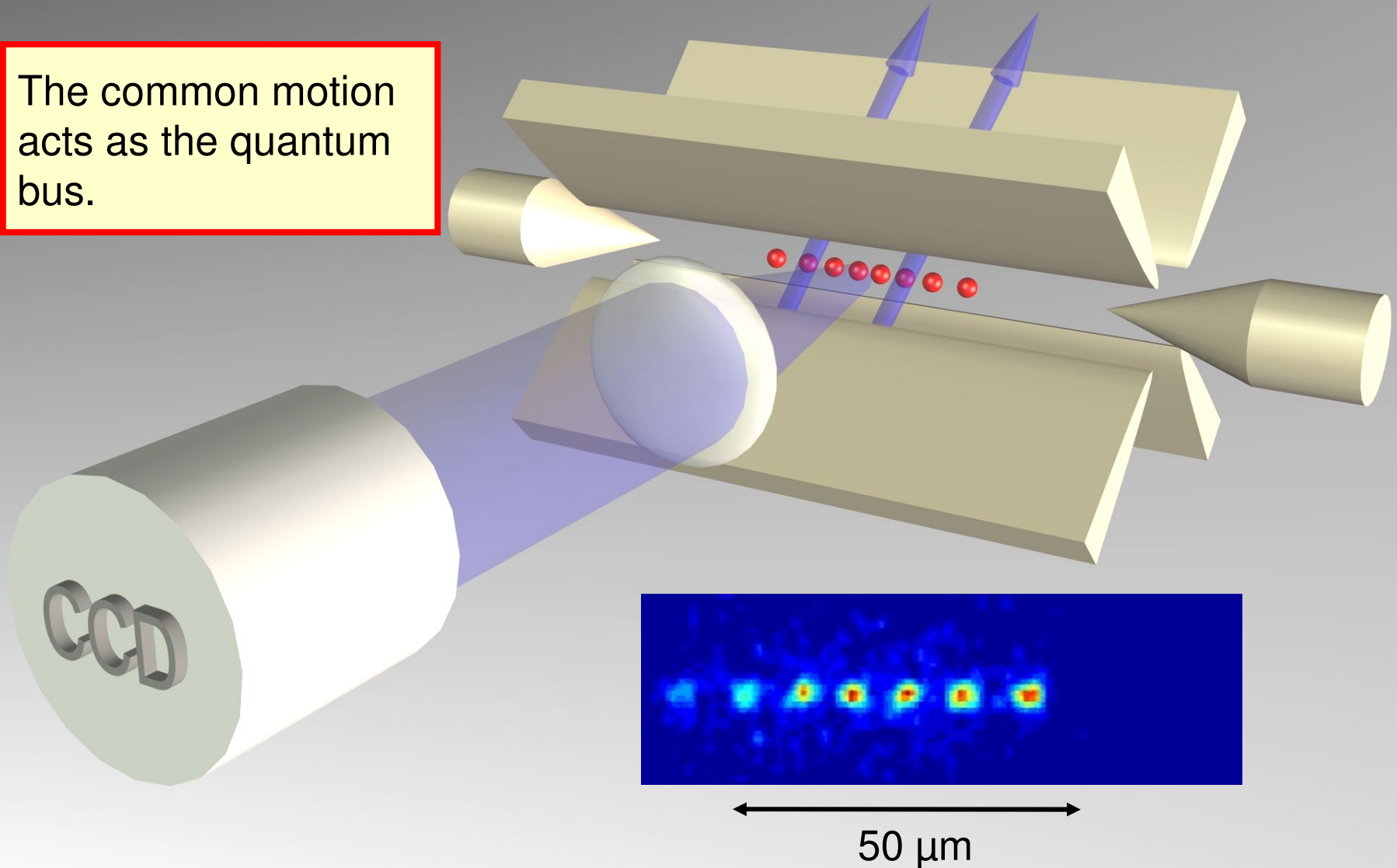
Computational space:

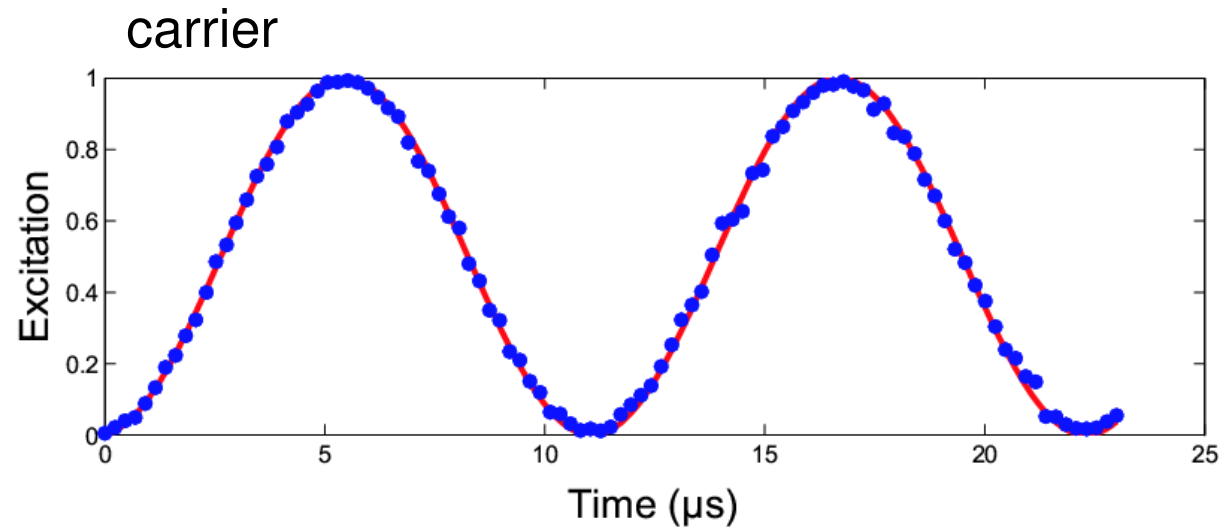
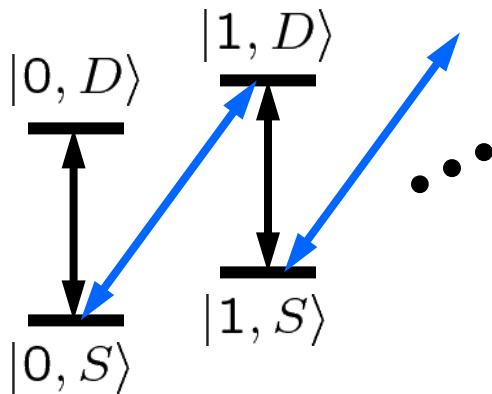
- Read out of qubits
➡ gain of classical information



Having the qubits interact

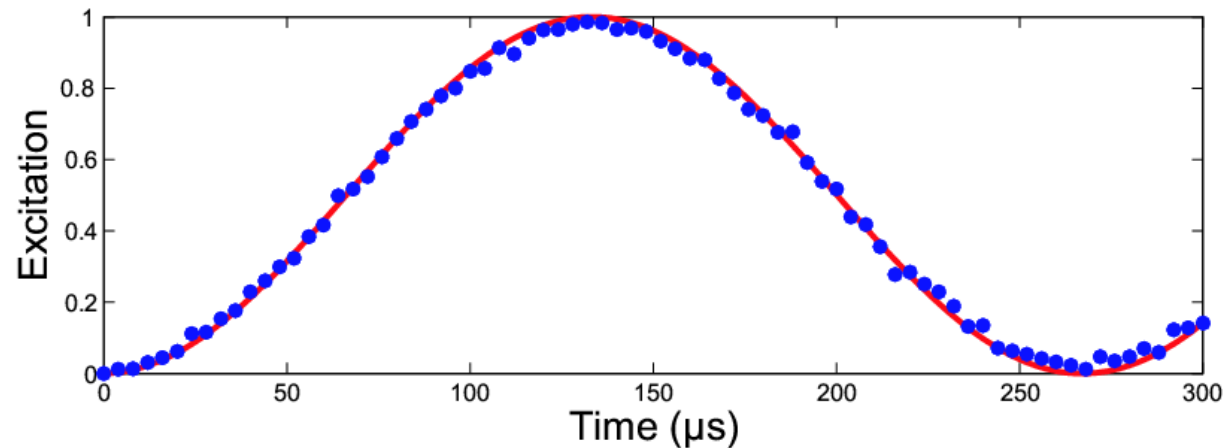
The common motion
acts as the quantum
bus.





carrier and sideband
Rabi oscillations
with Rabi frequencies

$$\Omega, \text{ } \eta\Omega$$



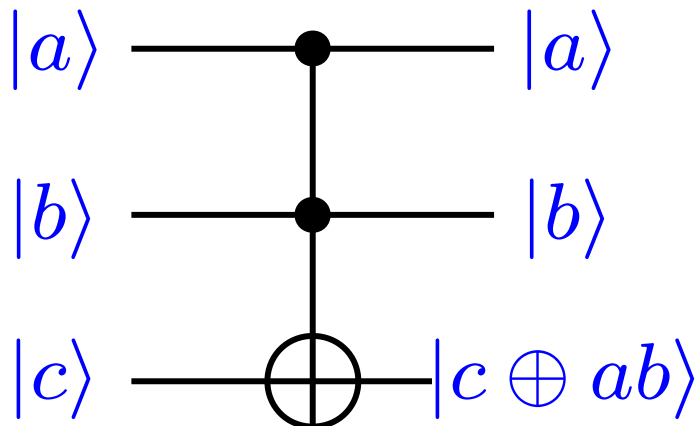
$$\eta = kx_0 \text{ Lamb-Dicke parameter}$$

- Introduction to ion trap quantum information
- **Single ion addressing approach**
- Coherent operations with global interactions
- Conclusions

Toffoli gate (Tommaso Toffoli, 1980):

..... is a universal reversible logic gate, i.e. any reversible circuit can be constructed from Toffoli gates.

also known as the **controlled-controlled-NOT** or **CCNOT**-gate operation

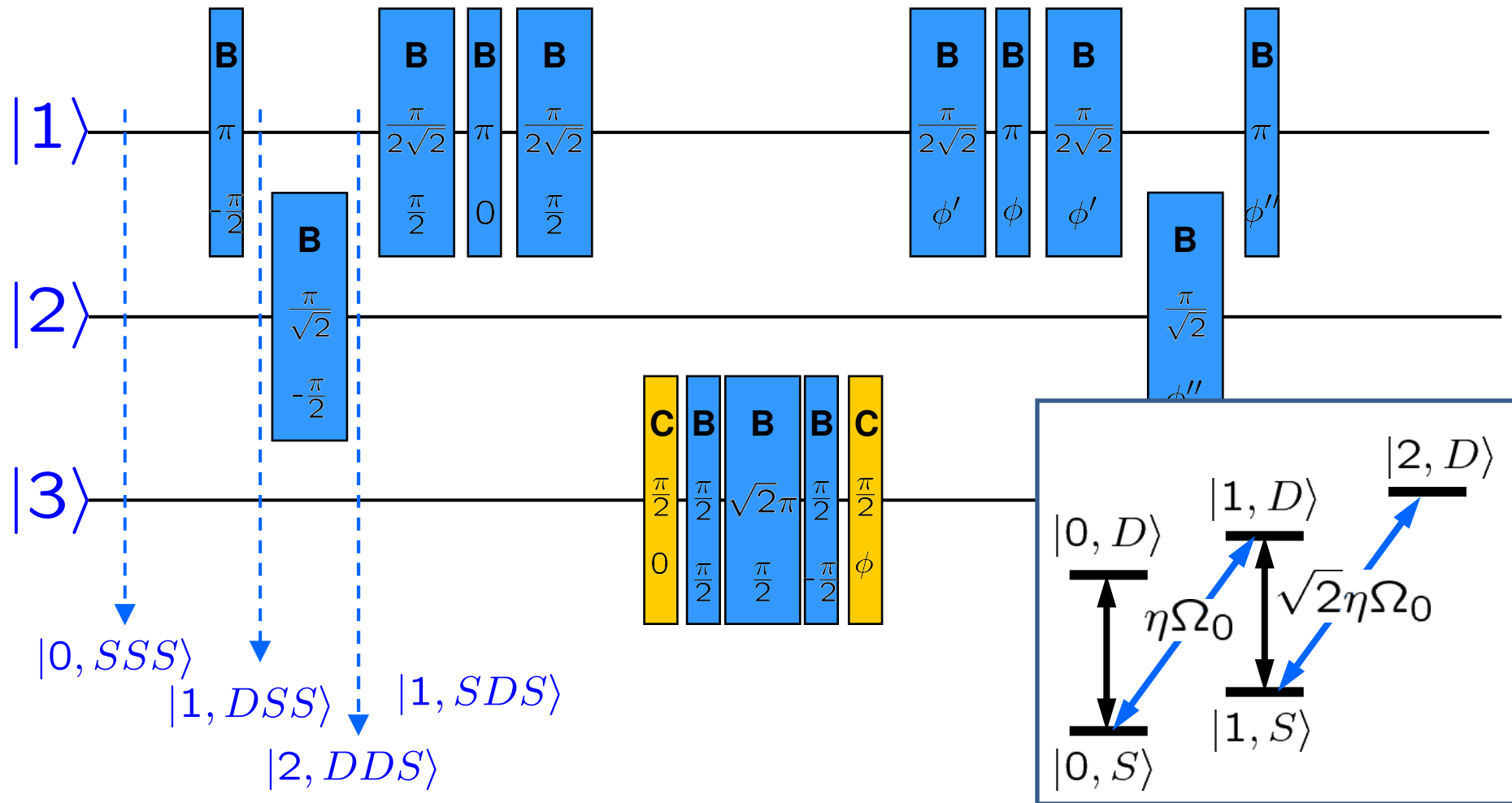


useful, e.g. for error correction

$$= \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & i \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -i & 0 \end{pmatrix}$$

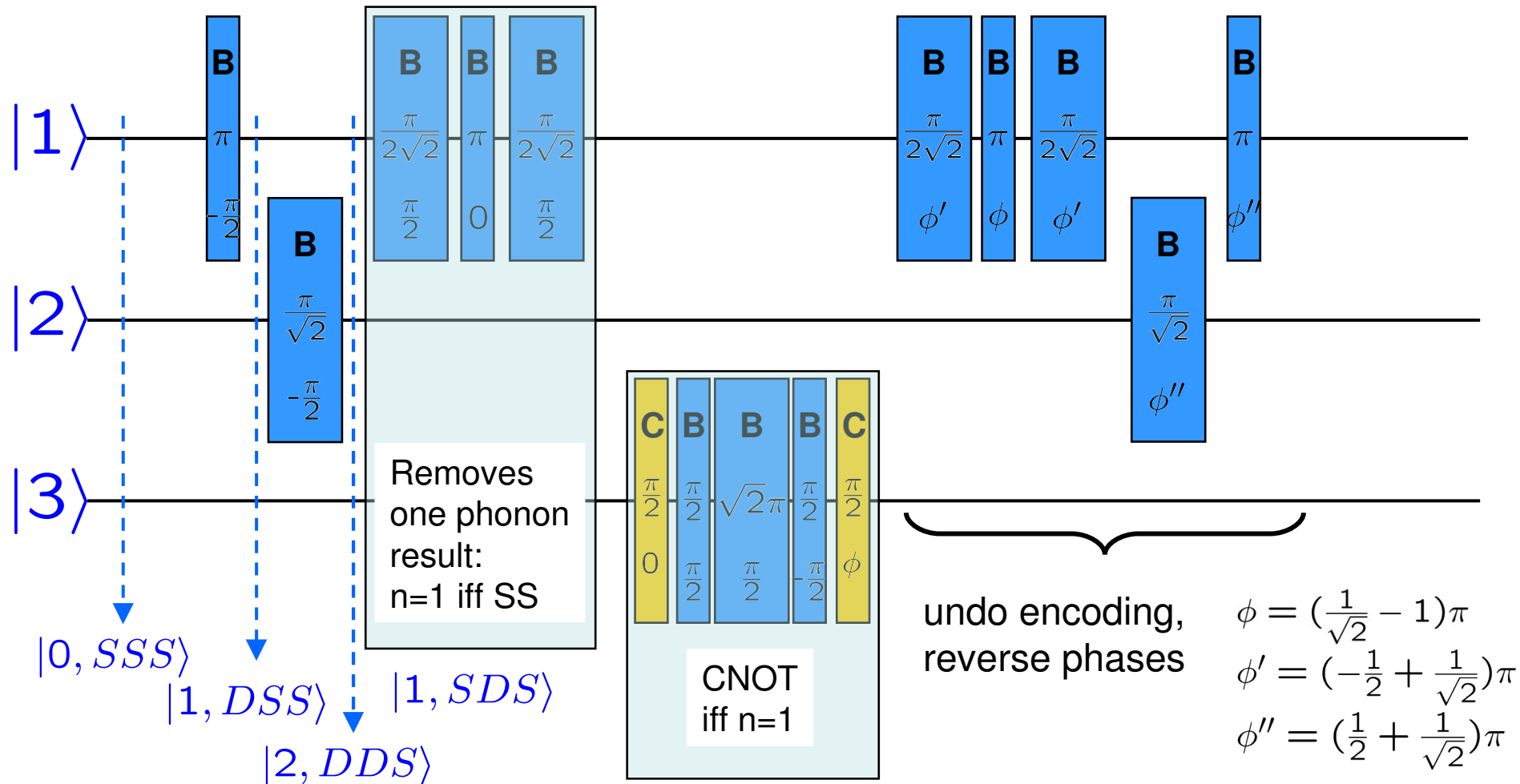
use 2-phonon excitation

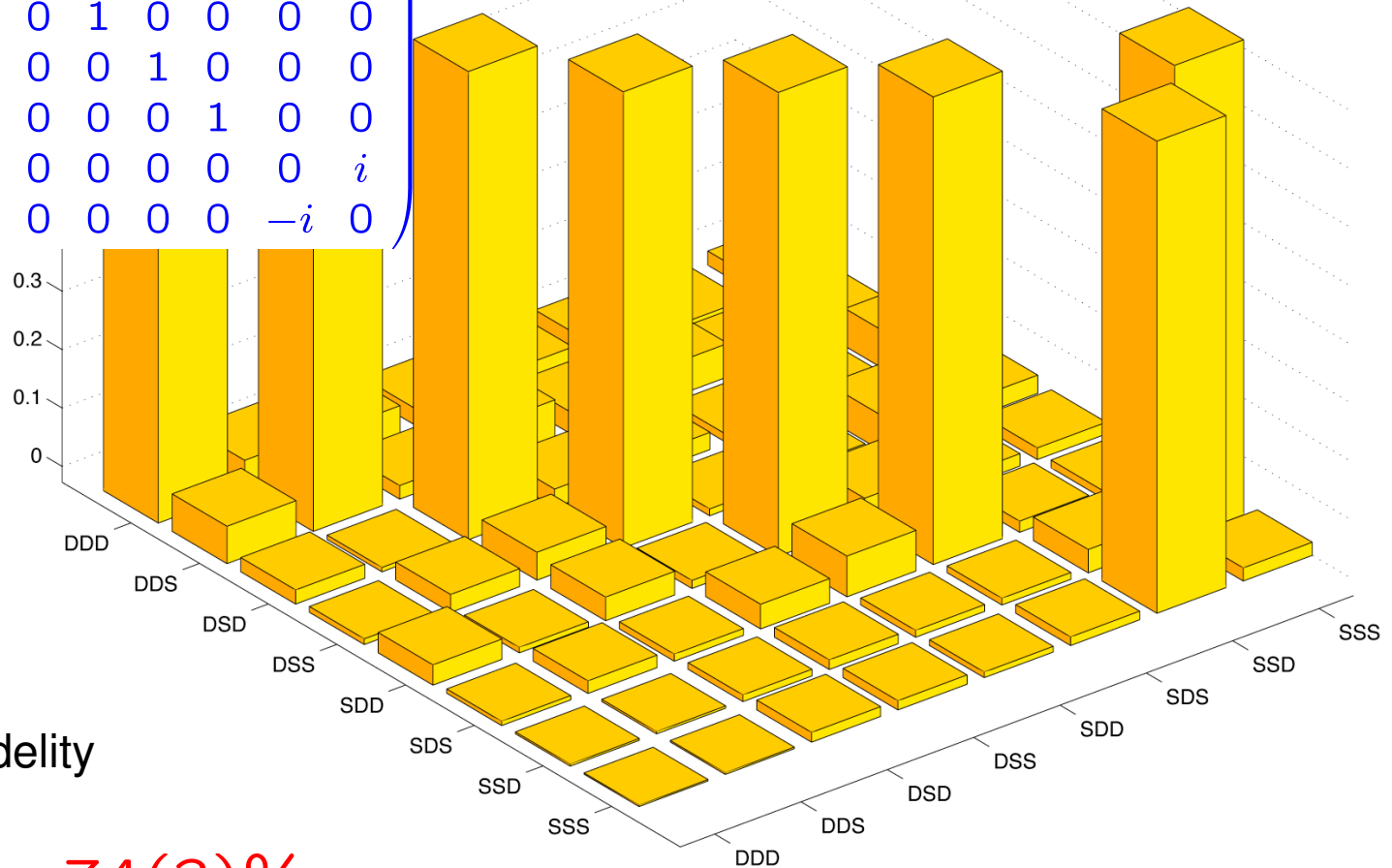
Th. Monz, K. Kim et al., Innsbruck 2008



use 2-phonon excitation

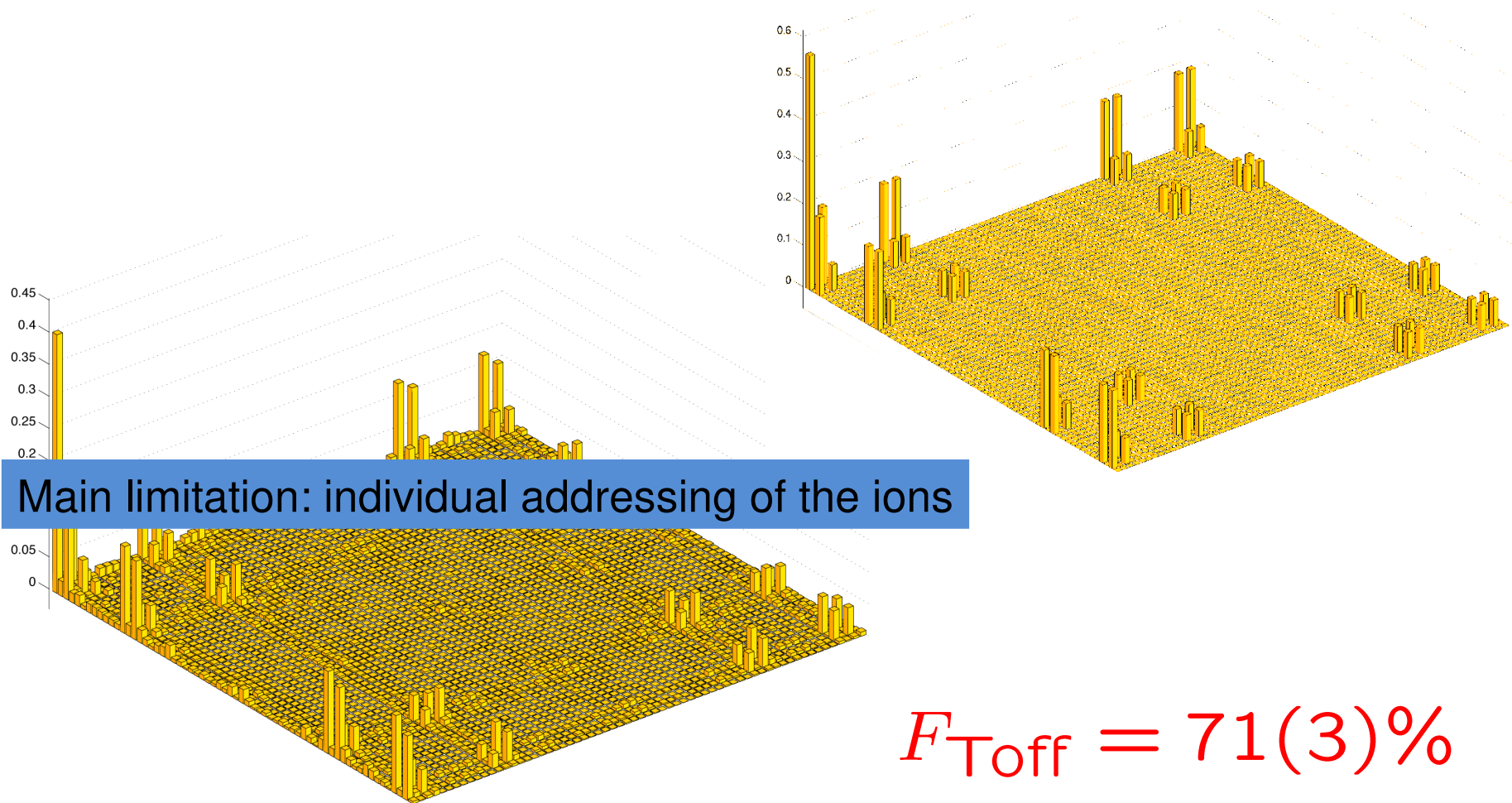
Th. Monz, K. Kim et al., Innsbruck 2008



$$\begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & i \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -i & 0 \end{pmatrix}$$


$$F_{\text{Toff}} = 74(3)\%$$

$|\chi|$ - matrix for ideal TOFFOLI gate operation

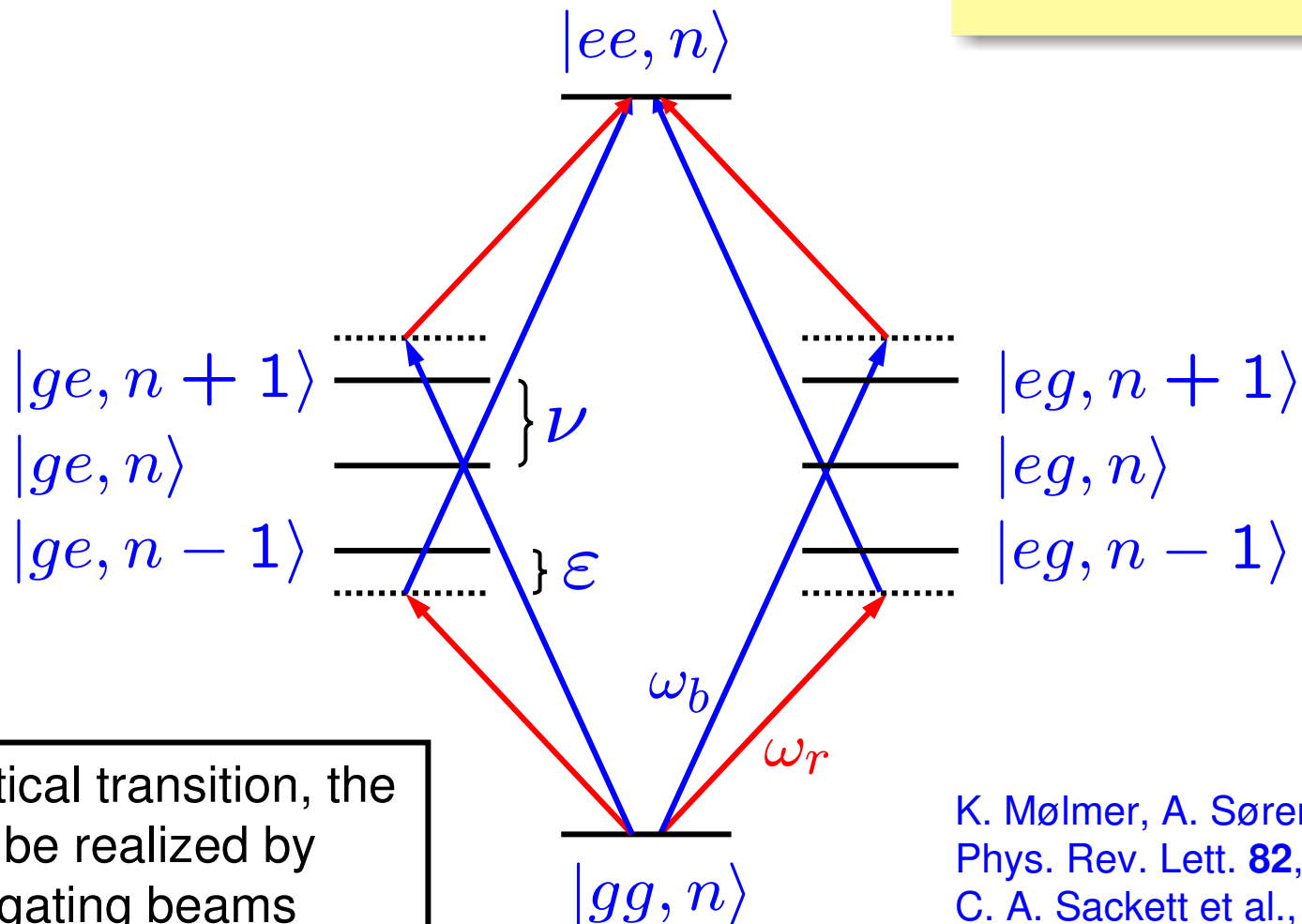


- Introduction to ion trap quantum information
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$$|gg\rangle \rightarrow |ee\rangle, \quad |ge\rangle \rightarrow |eg\rangle$$

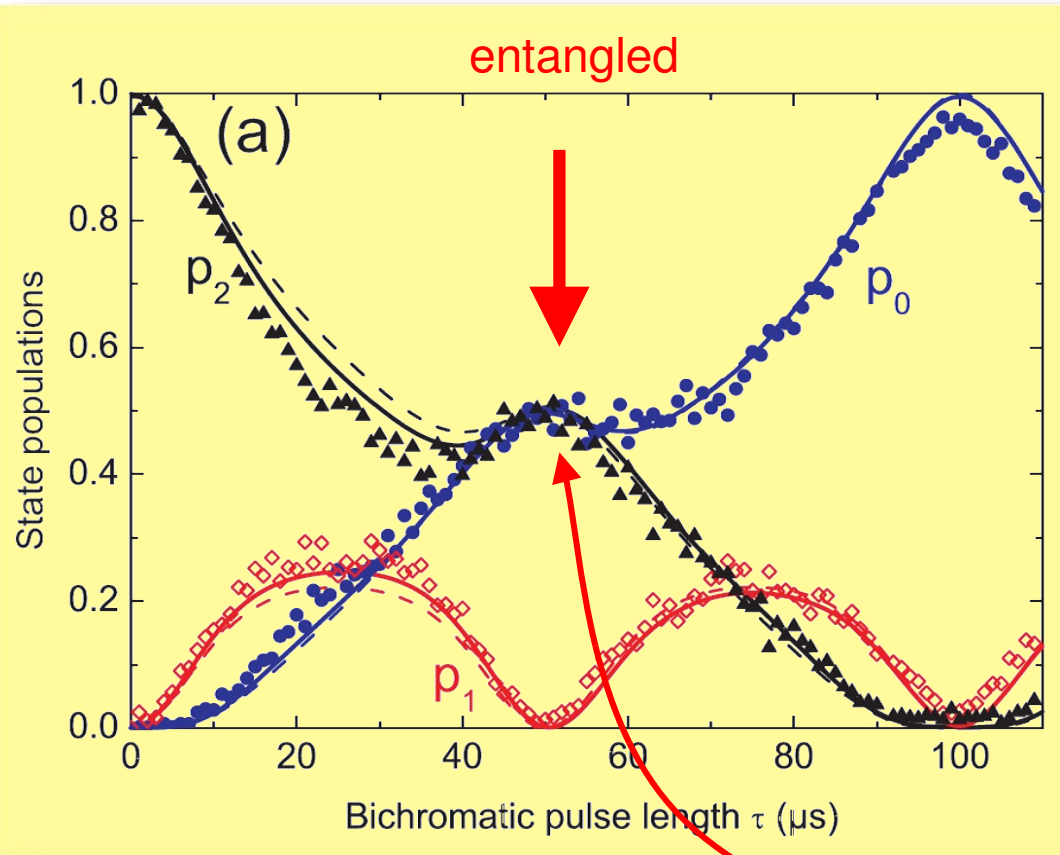
$$\omega_b = \omega_0 + (\nu - \varepsilon)$$

$$\omega_r = \omega_0 - (\nu - \varepsilon)$$



on an optical transition, the gate can be realized by co-propagating beams

K. Mølmer, A. Sørensen,
Phys. Rev. Lett. **82**, 1971 (1999)
C. A. Sackett et al.,
Nature **404**, 256 (2000)

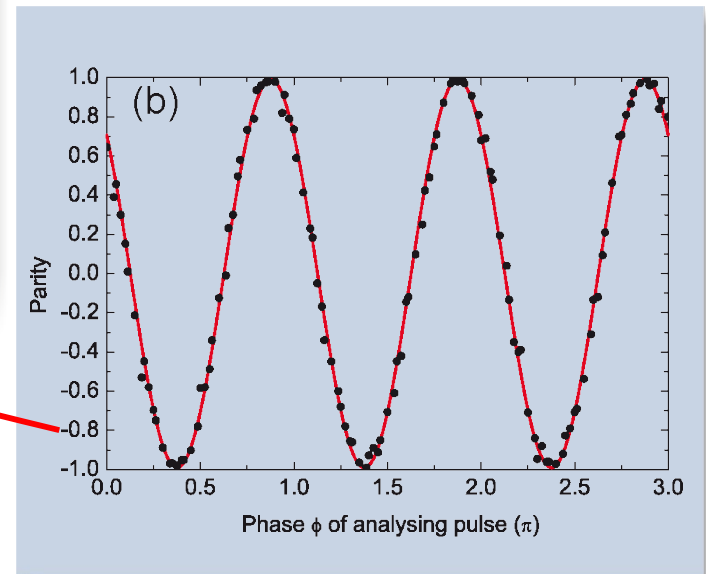


gate duration $51\mu\text{s}$
average fidelity

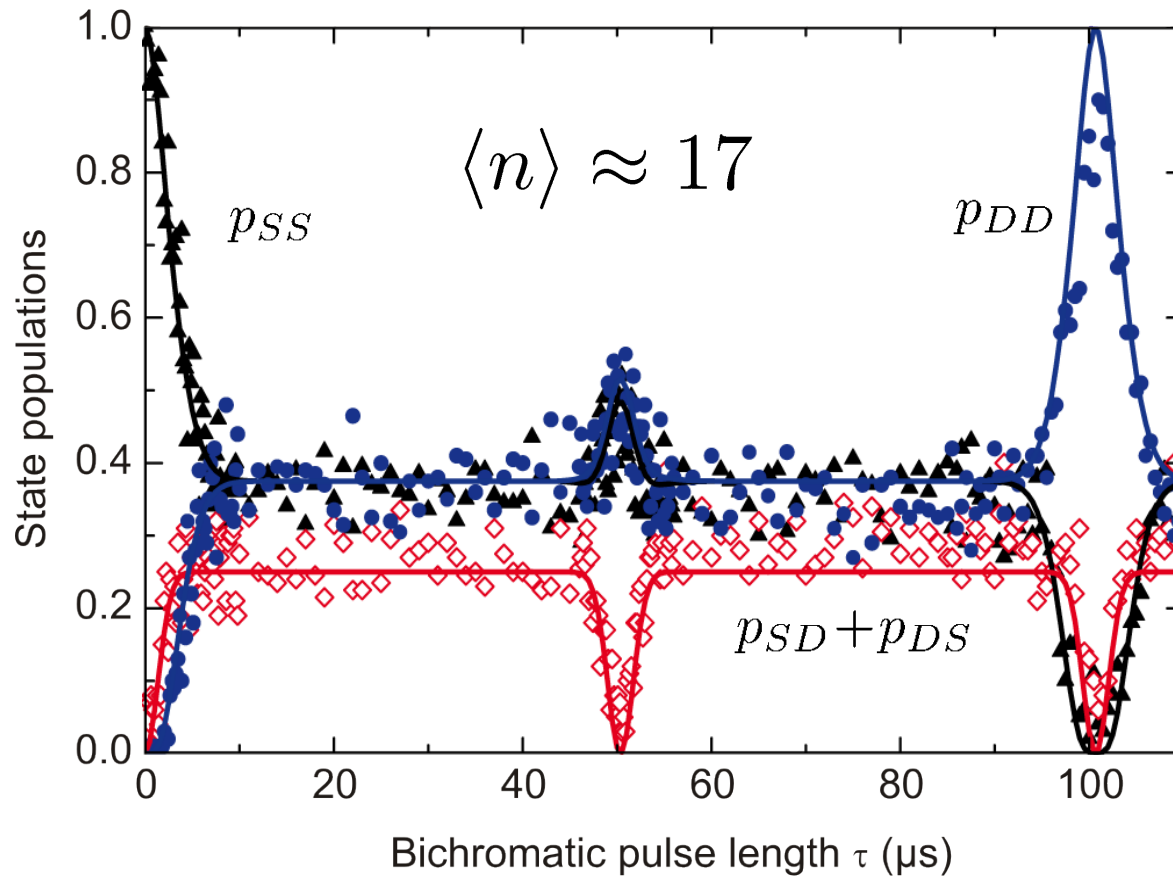
$$F_{\text{MS}} = 99.3(0.2)\%$$

J. Benhelm, G. Kirchmair, C. Roos,
Nature Physics **4** 463 (2008).

Theory: C. Roos,
New Journal of Physics **10**,
013002 (2008).



Gate operation after Doppler cooling

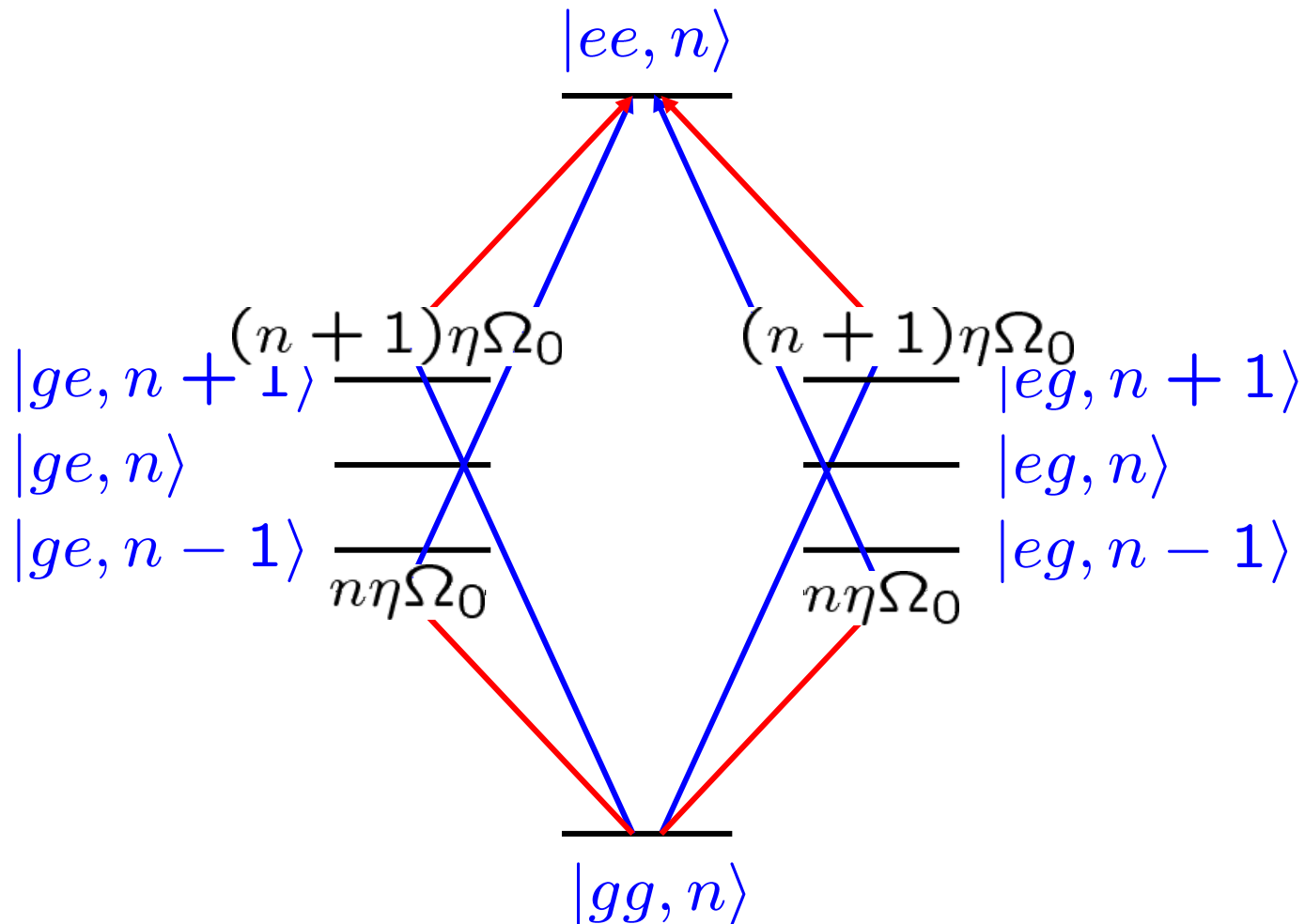


Bell state:
 $\Psi = |SS\rangle + i|DD\rangle$

Fidelity :

$$F = 96.1(5) \%$$

Gate operation \approx independent of motional state !



Arbitrary quantum gates

N ions

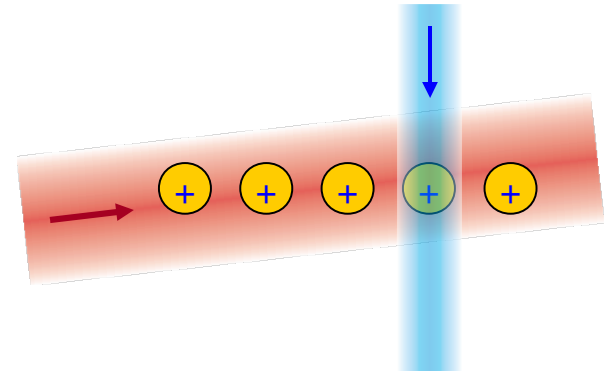
Basic set of operations:

$$H_i \in \{S_y^2, S_y, \underbrace{\sigma_z^{(1)}, \sigma_z^{(2)} \dots \sigma_z^{(N)}}\}$$

Mølmer-Sørensen gate

individual light shift gates

global single atom spin flips



- favorable ion addressing by light shifts ($\sim \Omega^2$)
- no interferometric stability between beams required

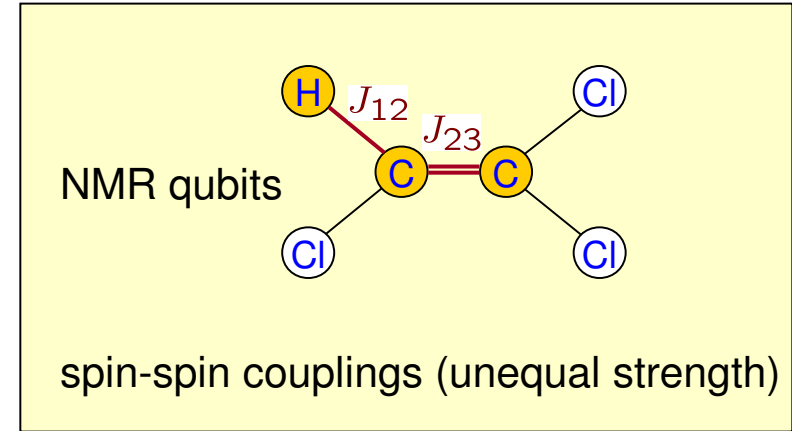
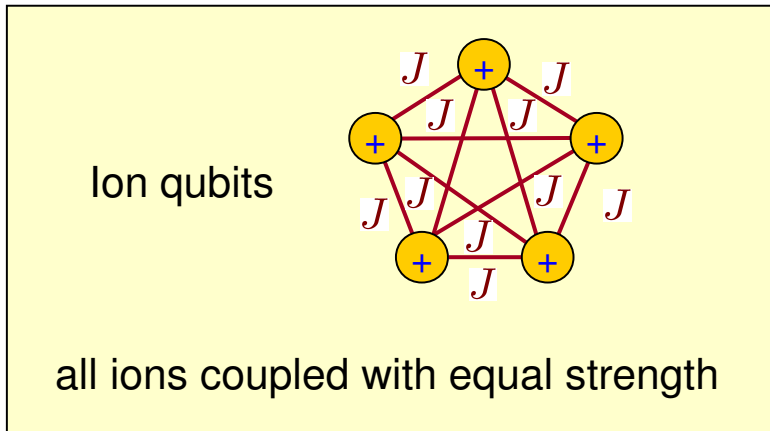
$H_i, H_j \rightarrow [H_i, H_j]$ generate Lie algebra \mathcal{L} with $\dim \mathcal{L} = 4^N$

Arbitrary unitary operations can be achieved !

...but how ?

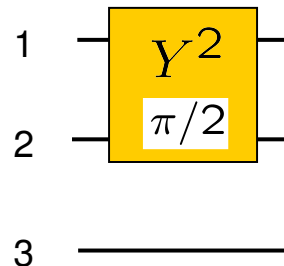
Arbitrary quantum gates

Similarity with NMR systems:

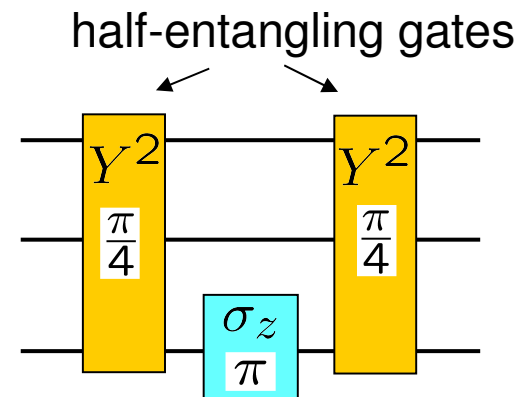


Refocussing of unwanted interactions:

Entangling gate
between ion 1 and 2



=



light shift pulse : phase shift by π

Arbitrary quantum gates

Quantum optimal control:

$$H(t) = \sum_{k=1}^n \alpha_k(t) H_k$$

Find $\{\alpha_k(t), k = 1 \dots n\}$ such that $U_{gate} \stackrel{!}{=} \mathcal{T} \int_0^\tau dt e^{-\frac{i}{\hbar} \sum_k \alpha_k(t) H_k}$

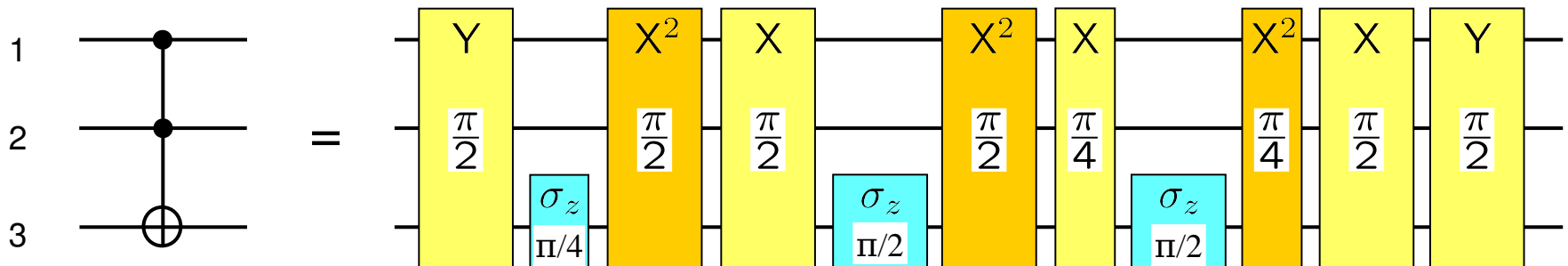
Gradient ascent algorithm: [N. Khaneja et al., J. Magn. Res. 172, 296 \(2005\).](#)

Modification of search algorithm: [V. Nebendahl et al., PRA \(2009\)](#)

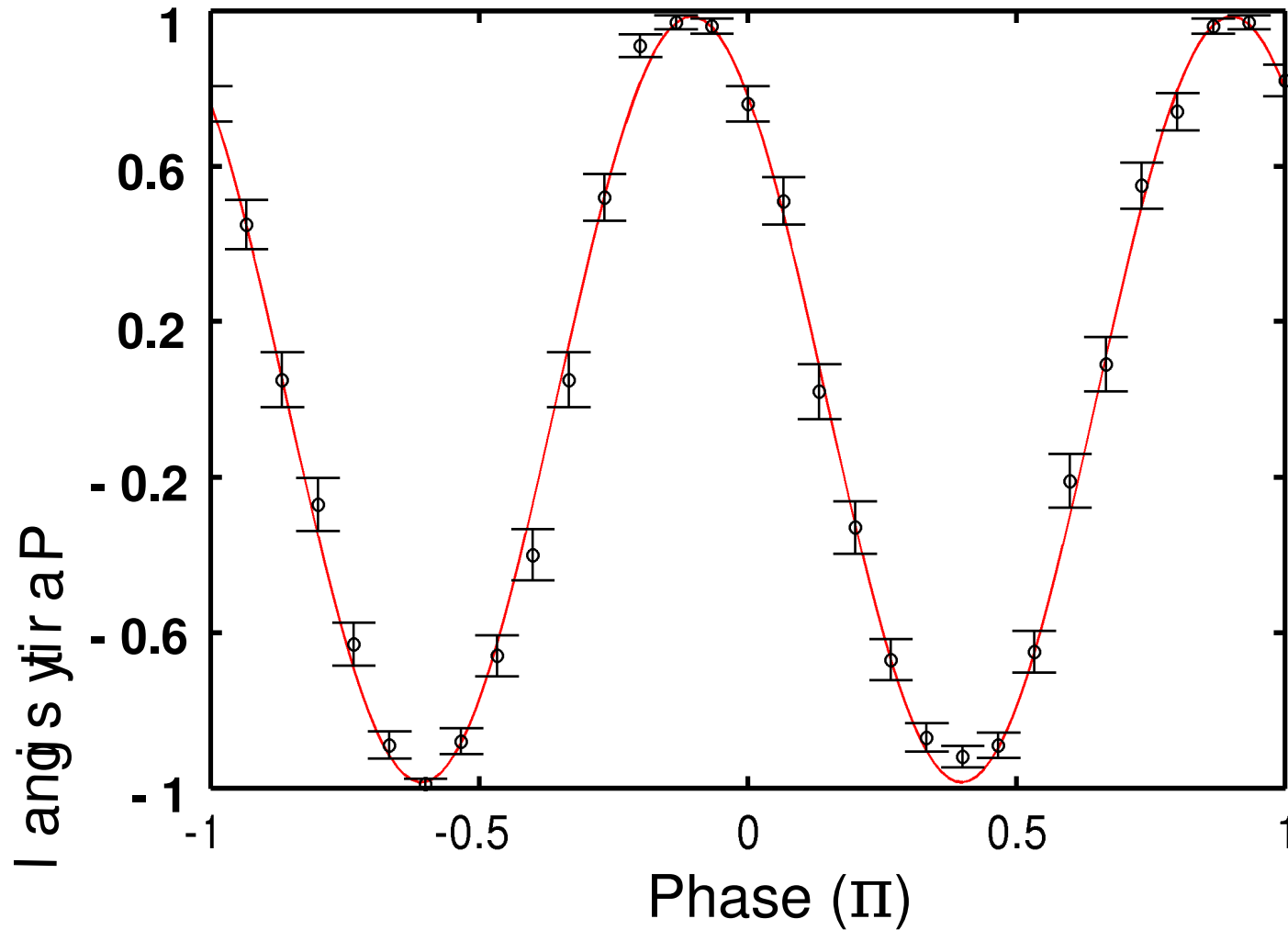
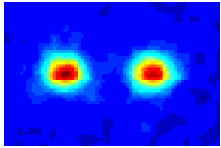
No simultaneous application of several Hamiltonians !

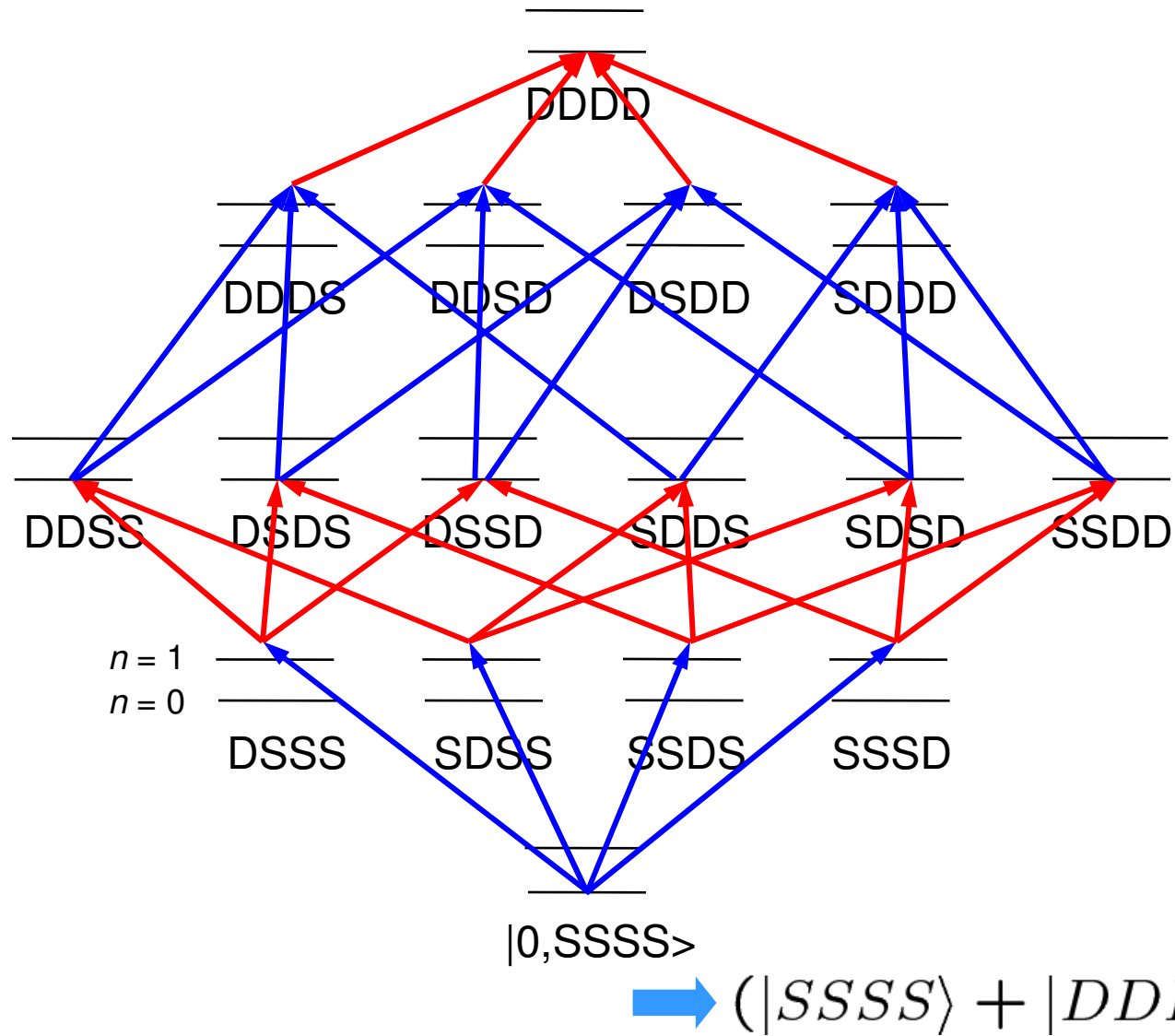
→ Sequence of pulses with variable length

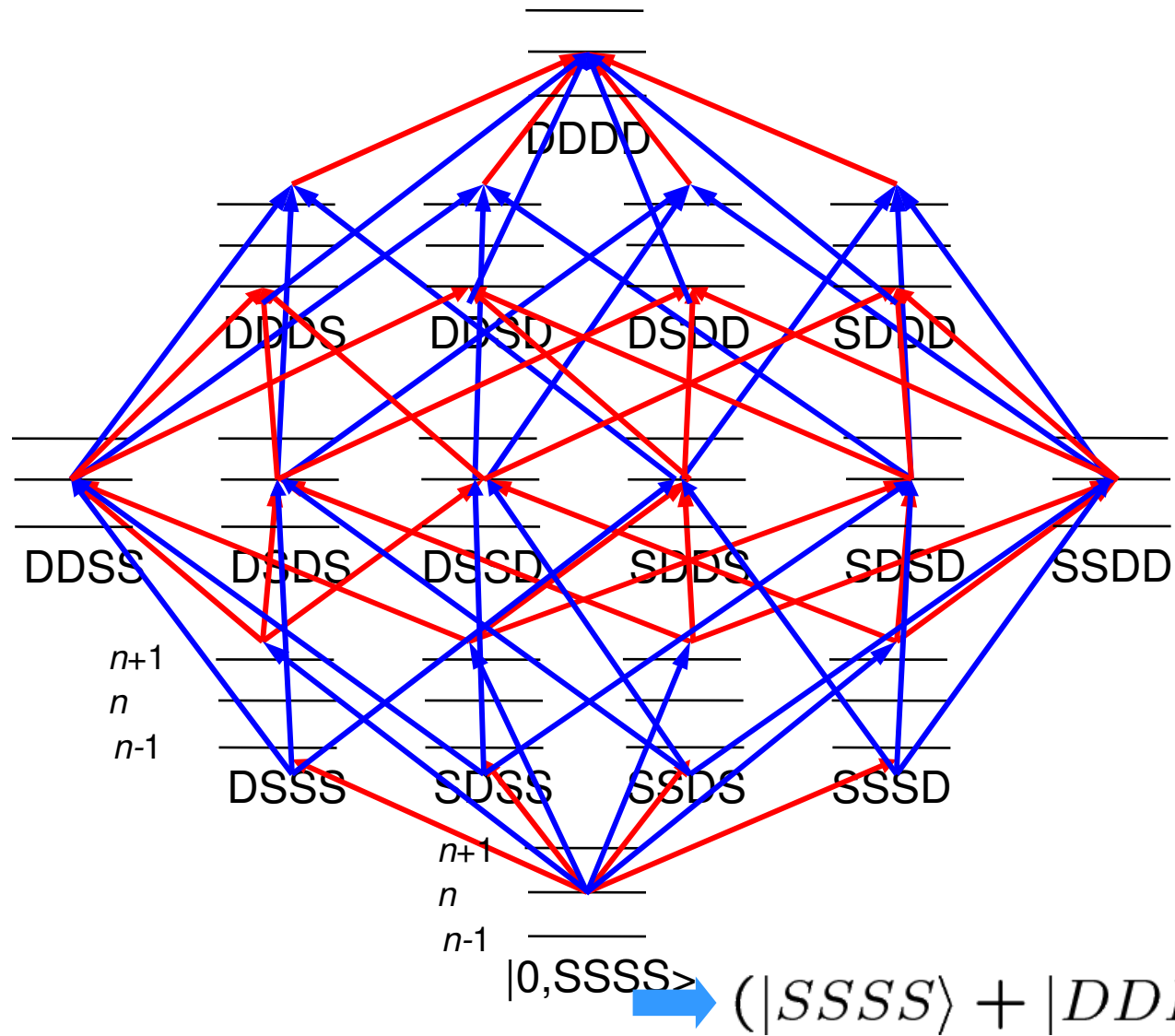
Example: quantum Toffoli gate



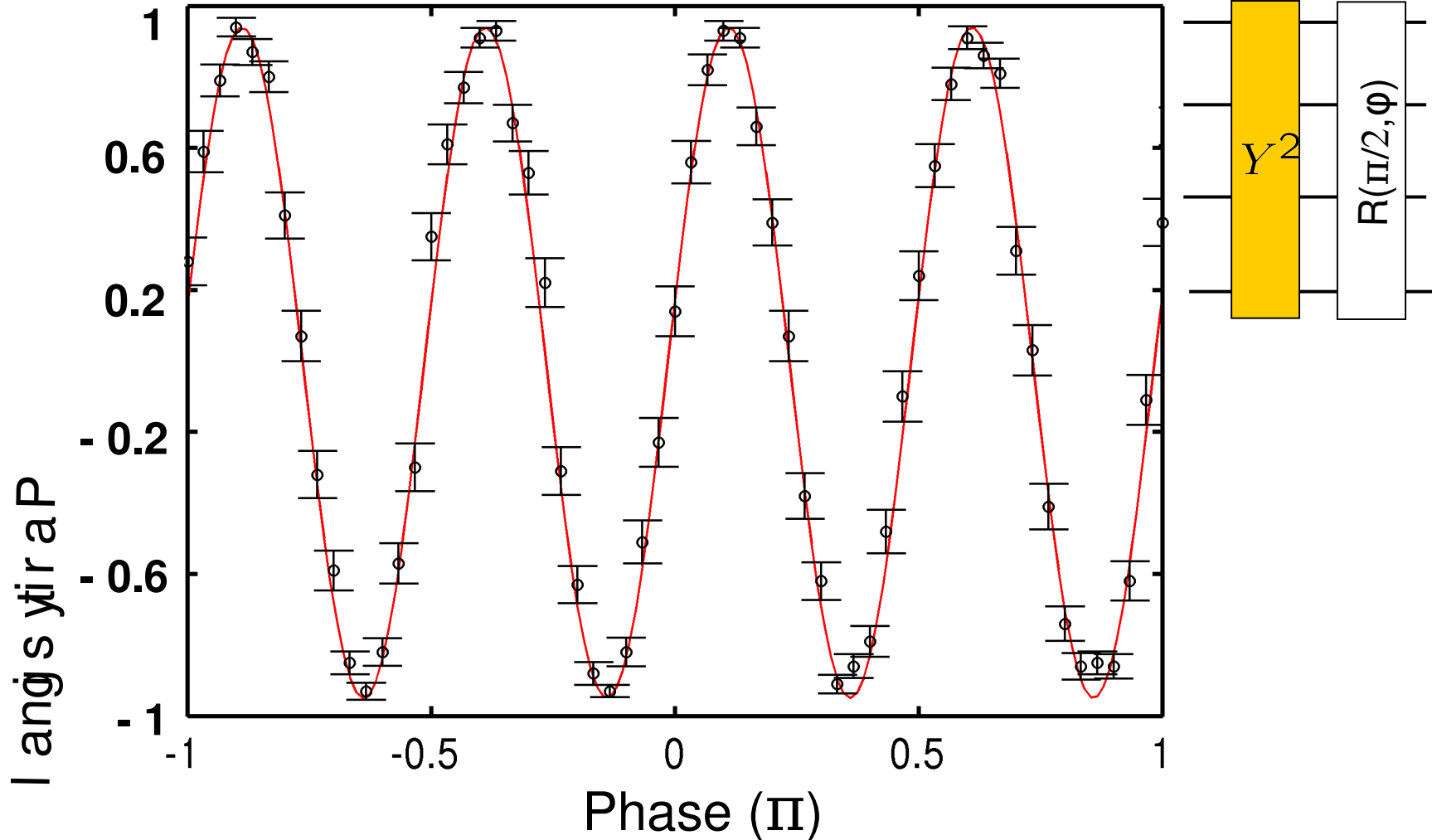
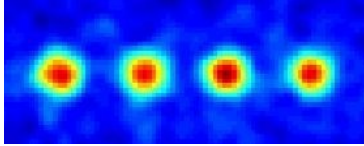
Two-ion GHZ-state





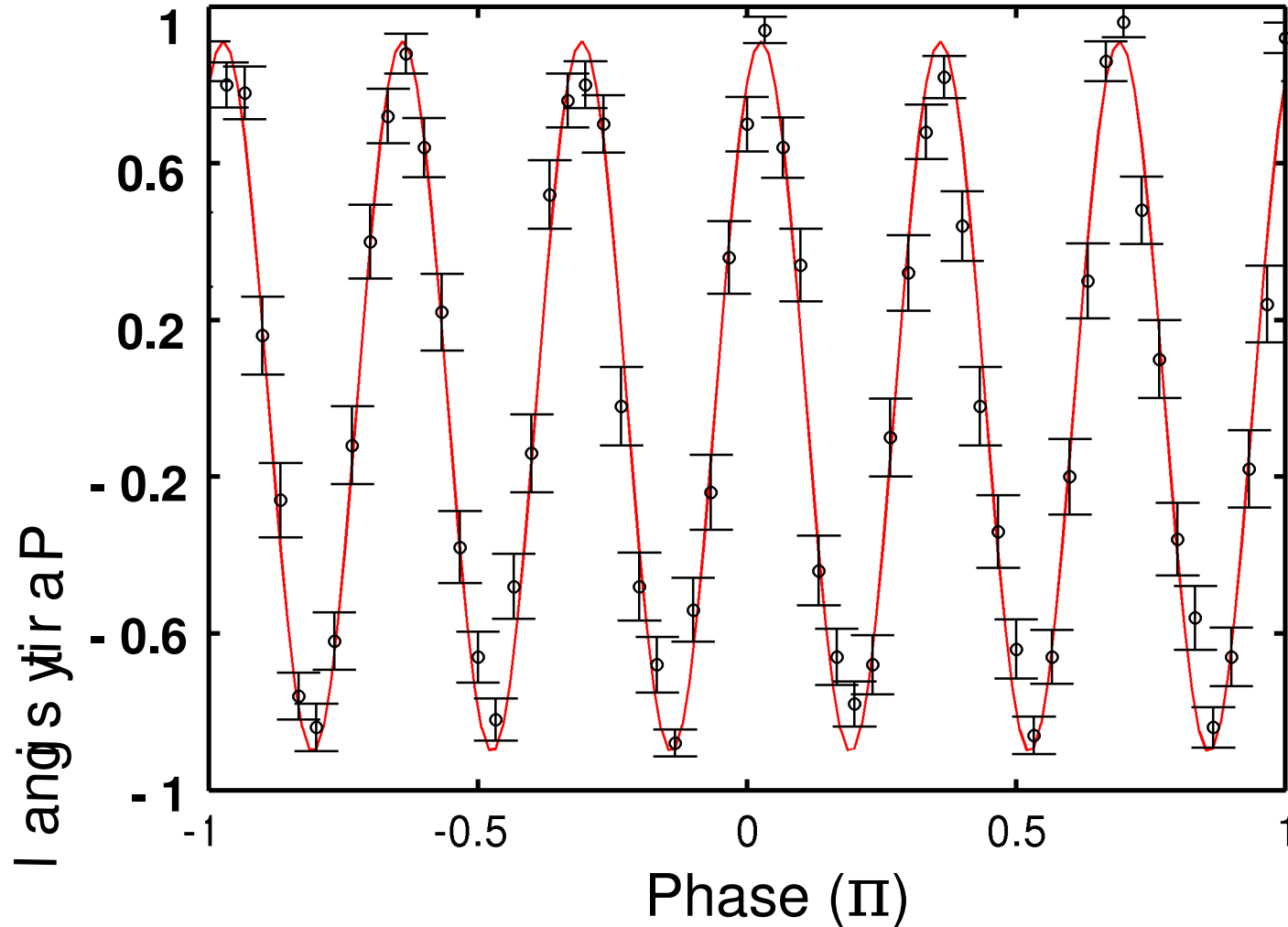
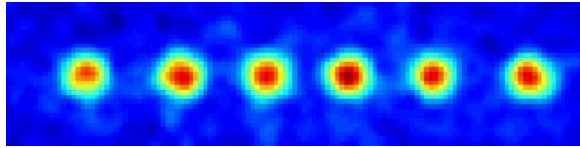


Four-ion GHZ state



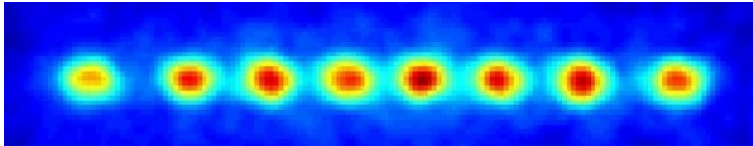
See also: Leibfried *et al.*, Nature **438**, 639 (2005)

Six-ion GHZ-state

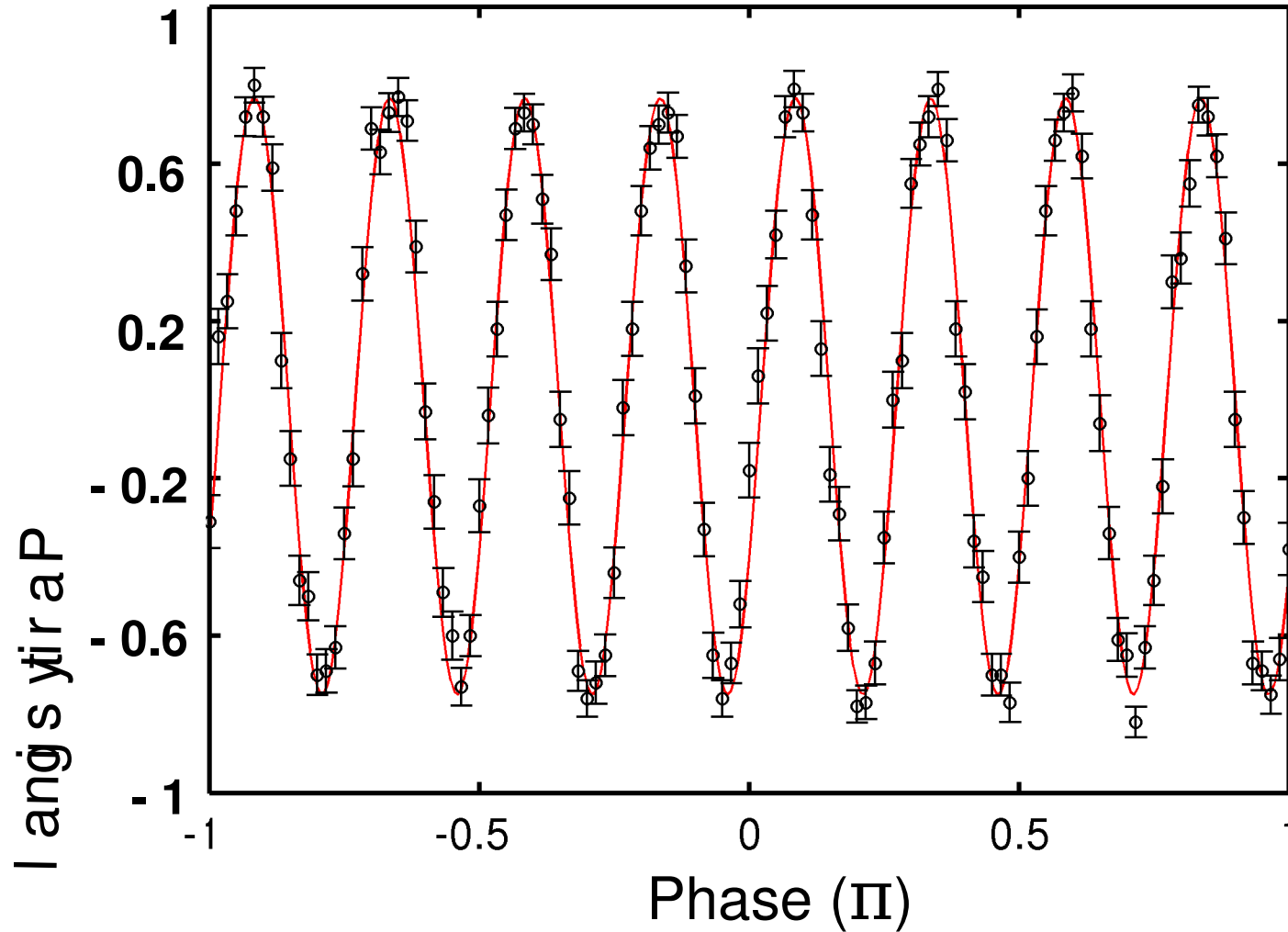


See also: Leibfried *et al.*, Nature **438**, 639 (2005)

Eight-ion GHZ state



T. Monz, P. Schindler, J. Barreiro, M. Hennrich

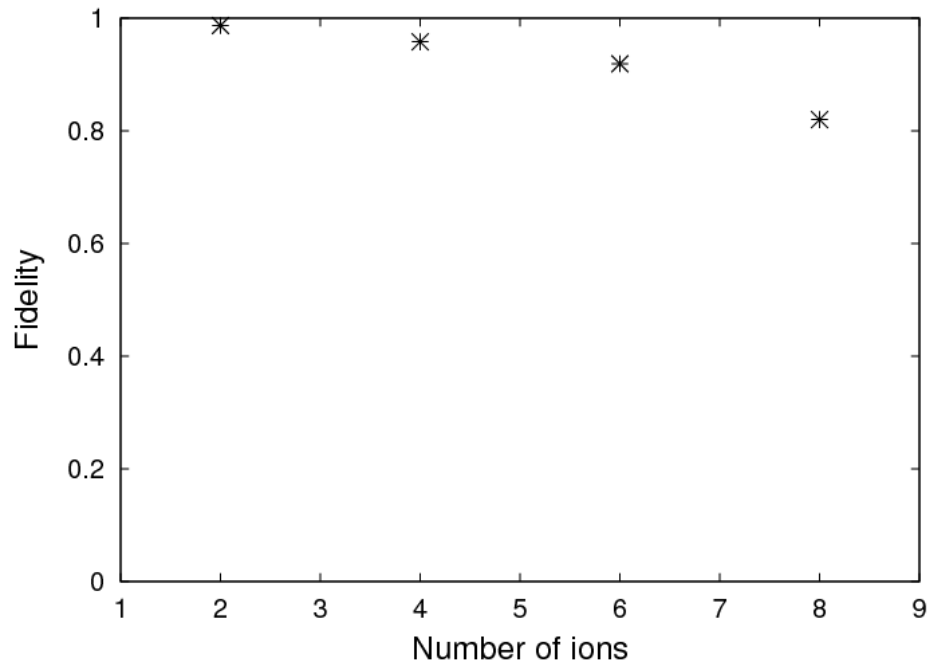


# ions	Fidelity	Witness
2	0.99 (1)	-0.97 (1)
4	0.96 (2)	-0.89 (1)
6	0.92 (3)	-0.81 (3)
8	0.82 (3)	-0.52 (3)

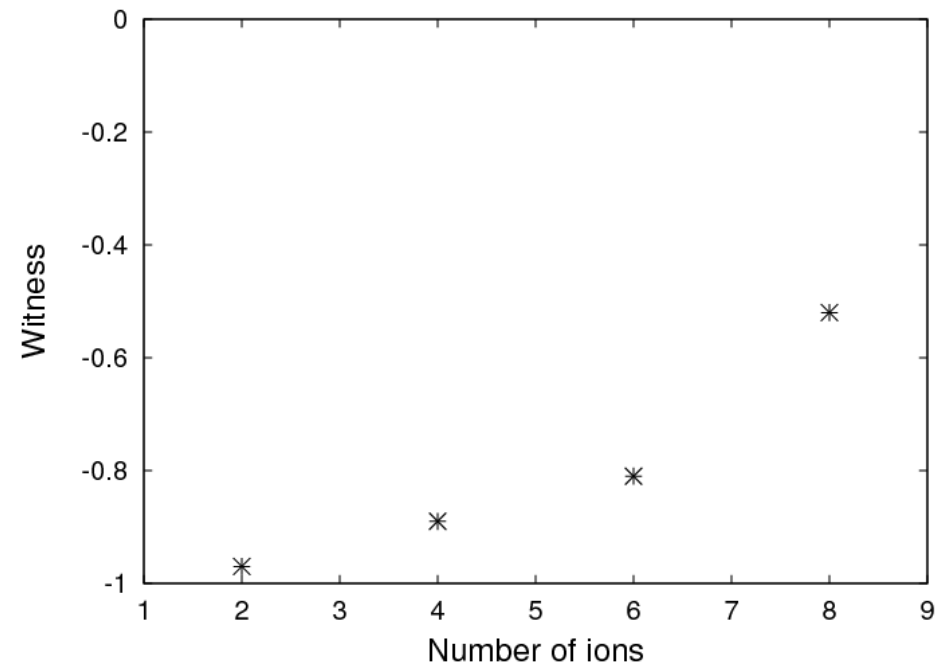
$$F = (P_{SS\dots S} + P_{DD\dots D} + 2 \text{ Contrast})/2$$

$$W = 1 - 4 \text{ Contrast}$$

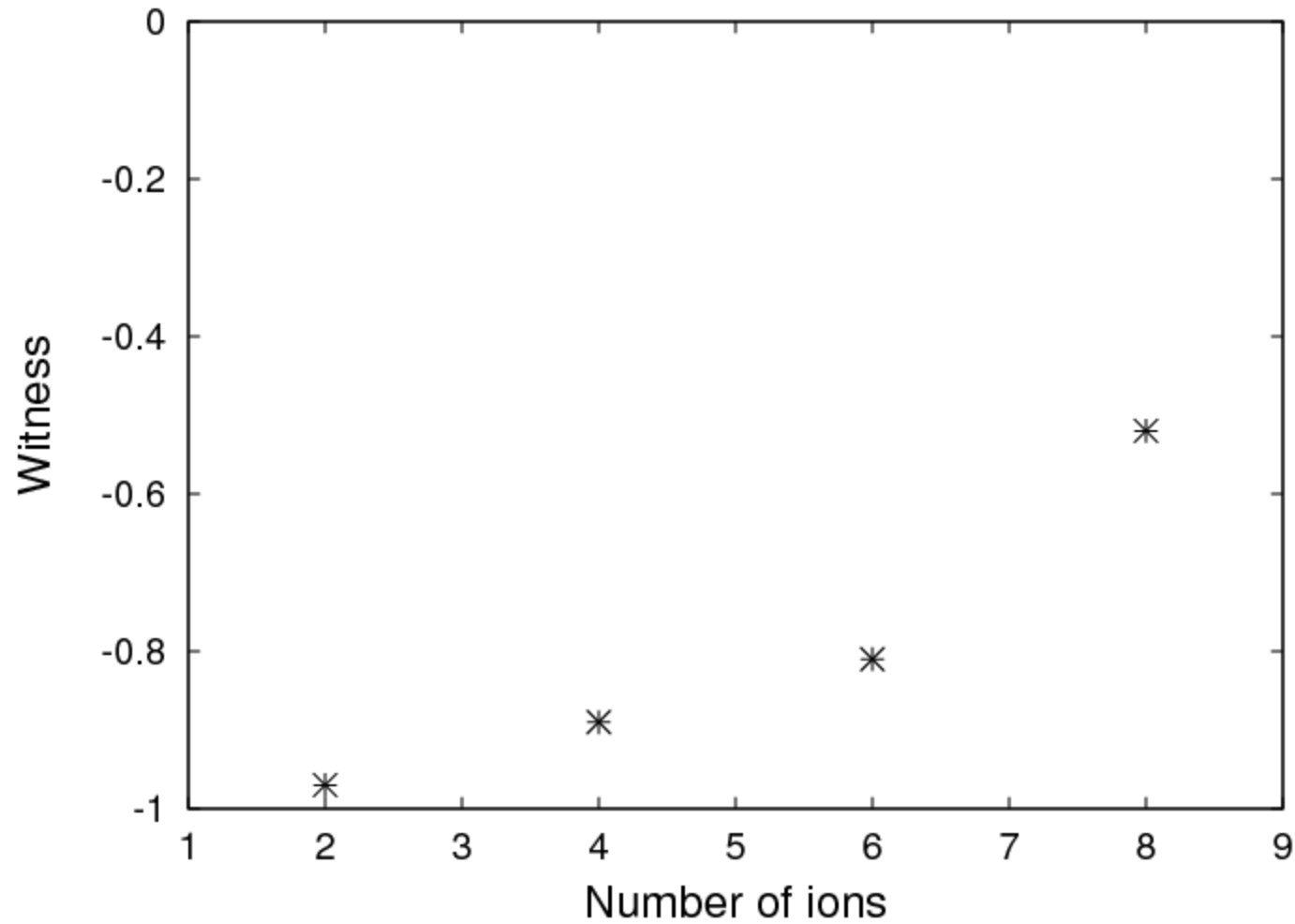
Fidelity



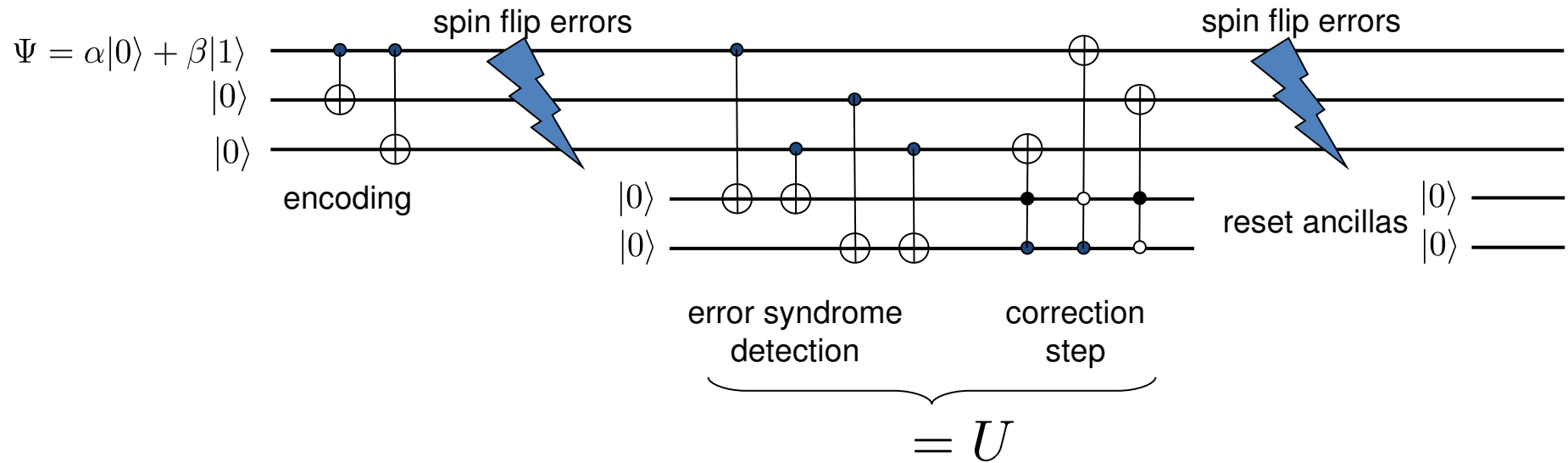
Expectation value of the Witness



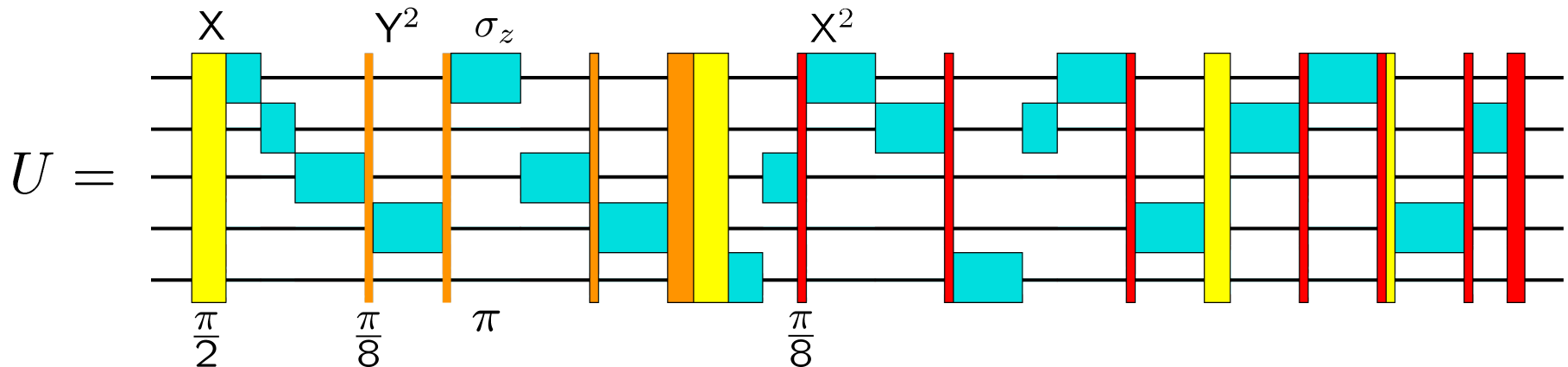
Entanglement



Example: quantum error correction: 3 qubits encode a logical qubit (protection against spin flips)

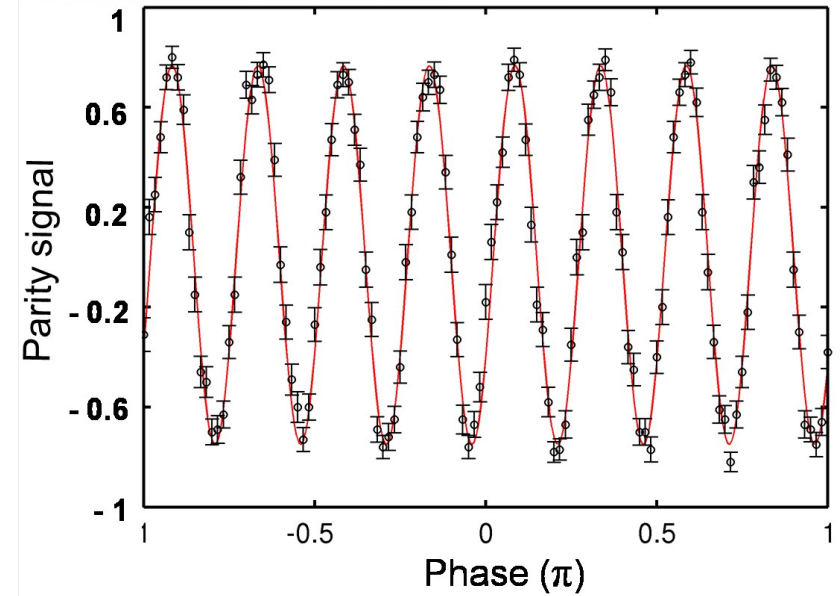
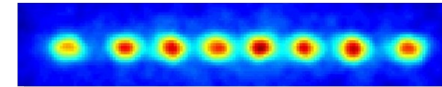
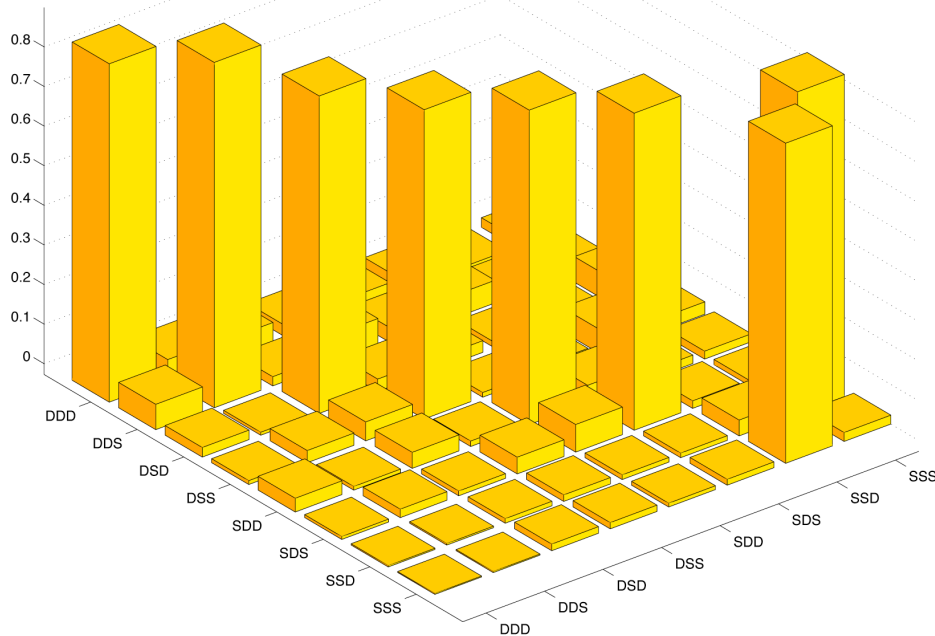


Implementation : 34 laser pulses (11 entangling pulses)

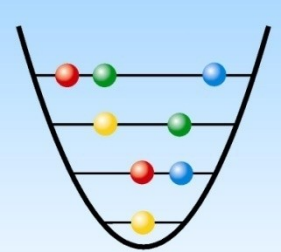


Conclusions

Toffoli gate



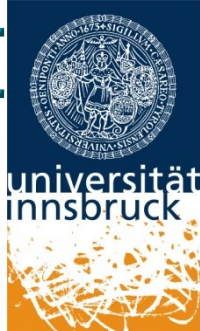
GHZ-states



AG Quantenoptik
und Spektroskopie



The Innsbruck ion trap group



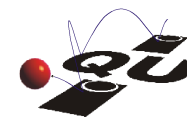
FWF
SFB



CONQUEST
SCALA



Industrie
Tirol



IQI
GmbH

FWF | **bm:bwk**



