

0.1 The use of Monte Carlo methods to solve the linear Boltzmann equation.

The transport equation in an infinite medium with its corresponding deterministic collisional component can be expressed as:

$$\partial_t u(x, t, \mathbf{v}) + \mathbf{v} \cdot \nabla u(x, t, \mathbf{v}) + v \sigma_t(x, t, \mathbf{v}) u(x, t, \mathbf{v}) = v \sigma_s(x, t, \mathbf{v}) \int P(x, t, \mathbf{v}, \mathbf{v}') u(x, t, \mathbf{v}') d\mathbf{v}' \quad (1)$$

Where

$$\sigma_s(x, t, \mathbf{v}) = \int \sigma_s(x, t, \mathbf{v}, \mathbf{v}') d\mathbf{v}', \quad P(x, t, \mathbf{v}, \mathbf{v}') = \frac{\sigma_s(x, t, \mathbf{v}, \mathbf{v}')}{\sigma_s(x, t, \mathbf{v})}$$

The approach involves a series of variable changes. The initial step involves re-expressing the transport equation (1) with respect to a characteristic $x + \mathbf{v}t$. As a result, it transforms into:

$$\partial_s u(x + \mathbf{v}s, s, \mathbf{v}) = -v \sigma_t(x + \mathbf{v}s, s, \mathbf{v}) u(x + \mathbf{v}s, s, \mathbf{v}) + v \sigma_s(x + \mathbf{v}s, s, \mathbf{v}) \int P(x + \mathbf{v}s, s, \mathbf{v}, \mathbf{v}') u(x + \mathbf{v}s, s, \mathbf{v}') d\mathbf{v}' \quad (2)$$

Let's multiply both sides of the equation by:

$$e^{\int_0^s v \sigma_t(x + \mathbf{v}\alpha, \alpha, \mathbf{v}) d\alpha}$$

Following that, we obtain

$$\partial_s [u(x + \mathbf{v}s, s, \mathbf{v}) e^{\int_0^s v \sigma_t(x + \mathbf{v}\alpha, \alpha, \mathbf{v}) d\alpha}] = e^{\int_0^s v \sigma_t(x + \mathbf{v}\alpha, \alpha, \mathbf{v}) d\alpha} v \sigma_s(x + \mathbf{v}s, s, \mathbf{v}) \int P(x + \mathbf{v}s, s, \mathbf{v}, \mathbf{v}') u(x + \mathbf{v}s, s, \mathbf{v}') d\mathbf{v}'$$

We get after integrating the equation (0.1) between $[0, t]$:

$$u(x + \mathbf{v}t, t, \mathbf{v}) = u_0(x, \mathbf{v}) \exp \left(- \int_0^t v \sigma_t(x + \mathbf{v}\alpha, \alpha, \mathbf{v}) d\alpha \right) + \int_0^t \int v \sigma_s(x + \mathbf{v}s, s, \mathbf{v}) u(x + \mathbf{v}s, s, \mathbf{v}') e^{-\int_s^t v \sigma_t(x + \mathbf{v}\alpha, \alpha, \mathbf{v}) d\alpha} P(x + \mathbf{v}s, s, \mathbf{v}, \mathbf{v}') d\mathbf{v}' ds \quad (3)$$

We obtain:

$$u(x, t, \mathbf{v}) = u_0(x - \mathbf{v}t, \mathbf{v}) \exp \left(- \int_0^t v \sigma_t(x - \mathbf{v}(t - \alpha), \alpha, \mathbf{v}) d\alpha \right) + \int_0^t \int v \sigma_s(x - \mathbf{v}(t - s), s, \mathbf{v}) u(x - \mathbf{v}(t - s), s, \mathbf{v}') e^{-\int_s^t v \sigma_t(x - \mathbf{v}(t - \alpha), \alpha, \mathbf{v}) d\alpha} P(x - \mathbf{v}(t - s), s, \mathbf{v}, \mathbf{v}') d\mathbf{v}' ds \quad (4)$$

We also have:

$$\exp \left(- \int_0^t v \sigma_t(x - \mathbf{v}(t - \alpha), \alpha, \mathbf{v}) d\alpha \right) = \exp \left(- \int_0^t v \sigma_t(x - \mathbf{v}\alpha, t - \alpha, \mathbf{v}) d\alpha \right) = \int_t^\infty v \sigma_t(x - \mathbf{v}s, t - s, \mathbf{v}) \exp \left(- \int_0^s v \sigma_t(x - \mathbf{v}\alpha, t - \alpha, \mathbf{v}) d\alpha \right) ds \quad (5)$$

Then the integral representation of 1 is provided by:

$$\begin{aligned}
u(x, t, \mathbf{v}) &= \int_t^\infty u_0(x - \mathbf{v}t, \mathbf{v}) v \sigma_t(x - \mathbf{v}s, t - s, \mathbf{v}) \exp\left(-\int_0^s v \sigma_t(x - \mathbf{v}\alpha, t - \alpha, \mathbf{v}) d\alpha\right) ds \\
&+ \int_0^t \int v \sigma_s(x - \mathbf{v}(t - s), s, \mathbf{v}) u(x - \mathbf{v}(t - s), s, \mathbf{v}') e^{-\int_s^t v \sigma_t(x - \mathbf{v}(t - \alpha), \mathbf{v}) d\alpha} P(x - \mathbf{v}(t - s), s, \mathbf{v}, \mathbf{v}') d\mathbf{v}' ds
\end{aligned} \tag{6}$$

Developing a Monte Carlo scheme involves introducing a group of random variables along with their associated probability measure to express 6 as an expectation. The selection of this set of random variables is not unique, resulting in various Monte Carlo schemes with distinct properties.