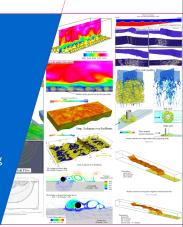


Chapter 4, Part 3: Unsteady Problems

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Outline

Unsteady Problems
Time Advancement
Temporal Schemes in OpenFOAM
Pseudo-transient approach



Advancing the Solution in Time

$$\underbrace{\frac{\partial \rho \phi}{\partial t}}_{\text{temporal derivative}} + \underbrace{\nabla \cdot (\rho \phi \mathbf{u})}_{\text{convection term}} - \underbrace{\nabla \cdot (\Gamma \nabla \phi)}_{\text{diffusion term}} = \underbrace{S_{\phi}}_{\text{source termoral derivative}}$$

- Two basic types of time advancement: Implicit and explicit schemes.
 Properties of the algorithm critically depend on this choice, but both are useful under given circumstances
- ► There is a number of methods, with slightly different properties, *e.g.* fractional step methods,
- ▶ Temporal accuracy depends on the choice of scheme and time step size
- Steady-state simulations
 - If equations are linear, this can be solved directly. We have done this for diffusion and advection-diffusion equations.
 - For non-linear equations, relaxation methods are used. A pseudo time is defined.



Explicit Schemes

- The time derivatives are evaluated using the currently available ϕ and the new ϕ is obtained from the time term
- ► All other terms (advection, diffusion, source, etc.) are evaluated using old time values.
- ► Courant number limit is the major limitation of explicit methods: in each time step, information can only propagate at the order of cell size; otherwise the algorithm is unstable.

$$CFL = \frac{U\Delta t}{\Delta x} < CFL_{max}$$

- Quick and efficient, no additional storage
- Very bad for elliptic behaviour

$$\underbrace{\frac{\partial \rho \phi}{\partial t}}_{\text{emporal derivative}} + \underbrace{\nabla \cdot (\rho \phi \mathbf{u})}_{\text{convection term}} - \underbrace{\nabla \cdot (\Gamma \nabla \phi)}_{\text{diffusion term}} = \underbrace{S_{\phi}}_{\text{source term}}$$



Implicit Schemes

- ► The algorithm is based on the method: each term is expressed in matrix form and the resulting linear system is solved
- ► A new solution takes into account the new values in the complete domain: ideal for elliptic problems
- Implicitness removed the Courant number limitation: we can take larger time-steps (but Δt is still subject to other things such as accuracy requirement)
- Substantial additional storage: matrix coefficients.



Summary of the temporal schemes in the book:

Scheme	Stability	Accuracy	Positive coefficient criterion
Explicit	Conditionally stable	First-order	$\Delta t < \frac{ ho(\Delta x)^2}{2\Gamma}$
Crank-Nicolson	Unconditionally stable	Second-order	$\Delta t < \frac{\rho(\overline{\Delta x})^2}{\Gamma}$
Implicit	Unconditionally stable	First-order	Always positive



From the point of view of temporal discretization, we need to treat both the time derivative term $(\partial/\partial t)$ and the spatial derivative terms (such as gradient, laplacian, divergence, etc.), as well as the source term.

$$\frac{\partial \rho \phi}{\partial t} + \underbrace{\nabla \cdot \left(\rho \phi \mathbf{u}\right)}_{\text{convection term}} - \underbrace{\nabla \cdot \left(\Gamma \nabla \phi\right)}_{\text{diffusion term}} = \underbrace{S_{\phi}}_{\text{source term}}$$

We can write this equation in a more compact form as

$$\frac{\partial \rho \phi}{\partial t} + \underbrace{A \phi}_{\text{spatial derivatives}} = 0$$

where A is an operator including everything else except the temporal derivative.



The first time derivative $(\partial/\partial t)$ term is integrated over a control volume as usual

$$\int_{V} \frac{\partial \rho \phi}{\partial t} dV = \frac{\partial}{\partial t} \int_{V} \rho \phi dV$$

This term is discretized by time differencing schemes using

- new values $\phi^n = \phi(t+\Delta t)$ at the time step we are solving for
- old values $\phi^o = \phi(t)$ that were stored from the previous time step
- ullet old-old values $\phi^{oo}=\phi(t-\Delta t)$ stored from a time step previous to the last



Euler implicit scheme (first order accurate in time)

$$\frac{\partial}{\partial t} \int_{V} \rho \phi dV = \frac{\left(\rho_{P} \phi_{P} V\right)^{n} - \left(\rho_{P} \phi_{P} V\right)^{o}}{\Delta t}$$

Backward differencing scheme (second order accurate in time):

$$\frac{\partial}{\partial t} \int_{V} \rho \phi dV = \frac{3 \left(\rho_{P} \phi_{P} V\right)^{n} - 4 \left(\rho_{P} \phi_{P} V\right)^{o} + \left(\rho_{P} \phi_{P} V\right)^{oo}}{2 \Delta t}$$



For the second time derivative, the only option is Euler scheme:

$$\frac{\partial}{\partial t} \int_{V} \rho \frac{\partial \phi}{\partial t} dV = \frac{\left(\rho_{P} \phi_{P} V\right)^{n} - 2\left(\rho_{P} \phi_{P} V\right)^{o} + \left(\rho_{P} \phi_{P} V\right)^{oo}}{\Delta t^{2}}$$



Next we need to consider how to treat the spatial derivatives in a transient problem. Integrate the whole equation over a control volume and over time, we have

$$\int_{t}^{t+\Delta t} \left[\frac{\partial}{\partial t} \int_{V} \rho \phi \, dV + \int_{V} A \phi \, dV \right] dt = 0$$

The discretization of the temporal derivative and operator A is independent. For example, we can use Euler implicit scheme for the temporal,

$$\int_{t}^{t+\Delta t} \left[\frac{\partial}{\partial t} \int_{V} \rho \phi \, dV \right] dt = \frac{\left(\rho_{P} \phi_{P} V \right)^{n} - \left(\rho_{P} \phi_{P} V \right)^{o}}{\Delta t} \Delta t$$



For the operator A, it can be expressed as

$$\int_{t}^{t+\Delta t} \left[\int_{V} A\phi \, dV \right] dt = \int_{t}^{t+\Delta t} A^{*} dt$$

where A^* represent the spatial discretization of the operator. The time integral on the right hand side of the equation can be discretized in different ways:

▶ Euler implicit: first order accurate, bounded, unconditionally stable

$$\int_{t}^{t+\Delta t} A^* dt = A^* \phi^n \Delta t$$

Explicit: first order accurate, unstable if the Courant number $C_o = \frac{U_f \cdot \mathbf{d}}{|\mathbf{d}|^2 \delta t} > 1$

$$\int_{t}^{t+\Delta t} A^* dt = A^* \phi^o \Delta t$$

 Crank Nicolson: second order accurate, unconditionally stable, not guaranteed boundedness

$$\int_t^{t+\Delta t} A^* dt = A^* \left(\frac{\phi^n + \phi^o}{2} \right) \Delta t$$



The first time derivative $(\partial/\partial t)$ terms are specified in the *ddtSchemes* sub-dictionary. The discretisation scheme for each term can be selected from those listed in the following table.

Scheme	Description	
Euler	First order, bounded, implicit	
localEuler	Local-time step, first order, bounded, implicit	
CrankNicolson Ψ	Second order, bounded, implicit	
backward	Second order, implicit	
steadyState	Does not solve for time derivatives	

There is an off-centering coefficient Ψ with the CrankNicolson scheme that blends it with the Euler scheme.

- $\Psi = 1$ corresponds to pure CrankNicolson
- $\mathbf{V} = \mathbf{0}$ corresponds to pure Euler

The blending coefficient can help to improve stability in cases where pure CrankNicolson are unstable.

Second time derivative $(\partial^2/\partial t^2)$ terms are specified in the *d2dt2Schemes* sub-dictionary. Only the Euler scheme is available for *d2dt2Schemes*.



In OpenFOAM[®], the temporal discretization of spatial operator A is actually controlled by the implementation of the spatial derivative in the equation. For example, the transient diffusion equation

$$\frac{\partial \phi}{\partial t} = \Gamma \nabla^2 \phi$$

An Euler implicit treatment would be

where the *laplacian* member function in the *fvm* namespace in OpenFOAM[®] treats the Laplacian term implicitly (Euler implicit). However, you can specify how to discretize the *fvm::ddt(phi)* term differently in the *fvSchemes* file.



An explicit treatment would be

```
solve( fvm::ddt(phi) == Gamma*fvc::laplacian(phi) )
```

where the *laplacian* member function in the *fvc* namespace in OpenFOAM® treats the Laplacian term explicitly.



```
The Crank Nicolson scheme treatment would be solve( fvm::ddt(phi) == Gamma*0.5*(fvm::laplacian(phi) + fvc::laplacian(phi)))
```

where it uses the *laplacian* member functions in both the \emph{fvm} and \emph{fvc} namespaces in OpenFOAM® .

For Crank Nicolson scheme, you also have the choice of directly specify the *ddtScheme* as *CrankNicolson* in the *fvSchemes* file.

Here the number "1.0" following the keyword *CrankNicolson* means it is a pure *CrankNicolson* scheme. If the number is "0", then it becomes a pure *Euler* implicit scheme. Any number in between "0" and "1.0" means a blending between *Euler* and *CrankNicolson*.



Pseudo-transient approach for steady state problem

In previous chapters, we have focused on steady state physics, i.e., the $\partial \phi/\partial t$ term does not show up in the governing equation.

However, steady state solution can also be simulated as the final converged solution of a transient simulation. We call this approach the "pseudo-transient" approach.

The pseudo-transient approach is attractive for circumstances where for example nonlinearity presents in the problem. It will pose great difficulty to the solver if don't use the pseudo-transient approach.



Pseudo-transient approach for steady state problem

Closely related to this topic is the under-relaxation which controls the change of ϕ , i.e.,

$$\phi_P^n = \phi_P^{n-1} + \alpha \left(\phi_P^{n*} - \phi_P^{n-1} \right)$$

where α is the relaxation factor.

- ho α < 1: under-relaxation. This will slow down the convergence but increase the stability.
- $\alpha = 1$: no relaxation.
- ho lpha > 1: over-relaxation. This can be used to accelerate the convergence but it will also decrease the stability.



Pseudo-transient approach for steady state problem

More on the relaxation:

- Unfortunately, the specification of the relaxation factor α might be problem-dependent.It needs user's experience
- ▶ The general purpose of under-relaxation is to suppress oscillations
- OpenFOAM[®] comes with some default relaxation factor values which are usually recommended

```
relaxationFactors
{
    p      0.3;
    U      0.7;
    k      0.7;
    omega     0.7;
}
```

 If you write your own code, you need to relax the solution variable explicitly. For example, in the solver simpleFoam, the velocity and pressure are relaxed explicitly by

```
UEqn.relax()
p.relax()
```



Other schemes

 $\mathsf{OpenFOAM}^\circledR$ has temporal schemes to solve ordinary differential equations (ODEs):

- source code located at: src/ODE/ODESolvers
- ► Fifth-order Cash-Karp Runge-Kutta for non-stiff systems: RK
- ► Fourth-order semi-implicit Runge-Kutta scheme of Kaps, Rentrop and Rosenbrock for stiff systems: KRR
- ► Semi-Implicit Bulirsh-Stoer: SIBS



Questions?

