

# 7.1

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```
In [1]: # Content from Proakis
# Code 7 2019, Alexander Kain
import numpy as np

from matplotlib import pyplot as plt
%matplotlib inline
plt.rcParams['figure.figsize'] = (10.0, 8.0)

import sympy as sym
sym.init_printing(use_unicode=True)
```

## 1 The Discrete Fourier Transform: Its Properties and Applications

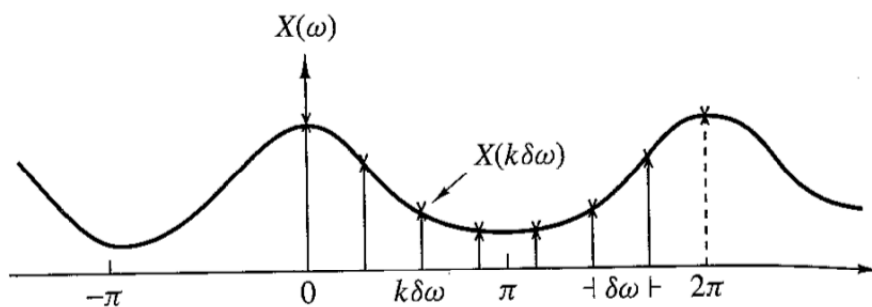
So far, the Fourier Transform  $X(\omega)$  of the sequence  $x[n]$  was a continuous function of frequency. We now consider a discrete frequency-domain representation, the discrete Fourier Transform (DFT)

### 1.1 7.1. Frequency-Domain Sampling: The Discrete Fourier Transform

Suppose we sample

$$X(\omega) = \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n}$$

periodically at  $\omega = 2\pi k/N$ :



Then

$$X(2\pi k/N) = \sum_{n=-\infty}^{\infty} x[n]e^{-j2\pi kn/N}, k = 0, 1, \dots, N-1$$

Let us sub-divide the summation into an infinite number of summations containing  $N$  terms:

$$\begin{aligned} X(2\pi k/N) &= \sum_{l=-\infty}^{\infty} \sum_{n=lN}^{lN+N-1} x[n]e^{-j2\pi kn/N} = \\ &= \sum_{n=lN}^{lN+N-1} \sum_{l=-\infty}^{\infty} x[n]e^{-j2\pi kn/N} = \end{aligned}$$

Let  $n \rightarrow n - lN$ :

$$\begin{aligned} X(2\pi k/N) &= \sum_{n=0}^{N-1} \sum_{l=-\infty}^{\infty} x[n - lN]e^{-j2\pi k(n-lN)/N} = \\ &= \sum_{n=0}^{N-1} \left( \sum_{l=-\infty}^{\infty} x[n - lN] \right) e^{-j2\pi kn/N} \end{aligned}$$

The term in parentheses

$$x_p[n] = \sum_{l=-\infty}^{\infty} x[n - lN]$$

is periodic with fundamental period  $N$ . Therefore, it can be written as a Fourier Series:

$$x_p[n] = \sum_{k=0}^{N-1} c_k e^{j2\pi kn/N}, n = 0, 1, \dots, N-1$$

with Fourier coefficients

$$c_k = \frac{1}{N} \sum_{n=0}^{N-1} x_p[n] e^{-j2\pi kn/N}, k = 0, 1, \dots, N-1$$

Comparing this to the formula for  $X(2\pi k/N)$ , we find that

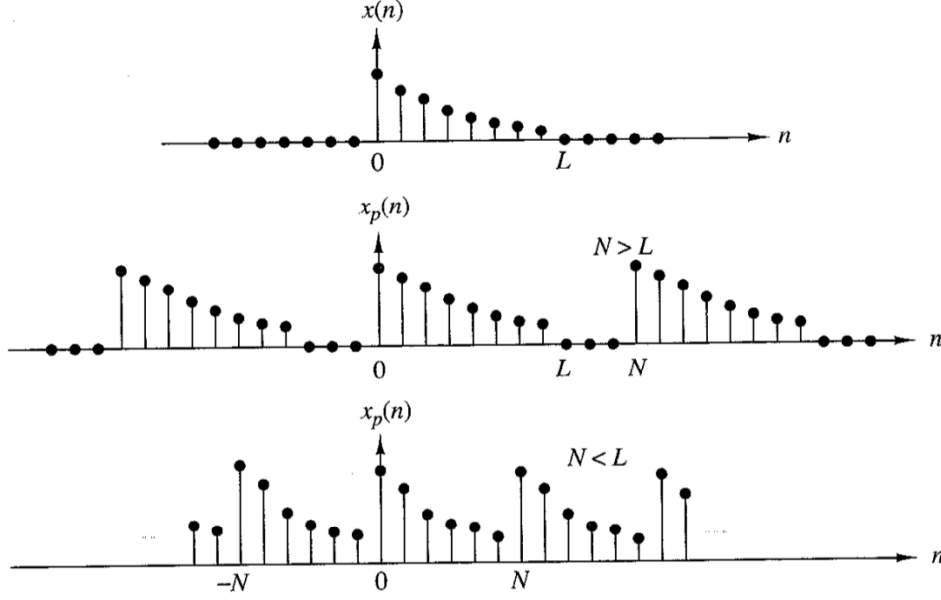
$$c_k = \frac{1}{N} X(2\pi k/N)$$

and

$$x_p[n] = \frac{1}{N} \sum_{k=0}^{N-1} X(2\pi k/N) e^{j2\pi kn/N}, n = 0, 1, \dots, N-1$$

This relationship provides the reconstruction of the periodic signal  $x_p[n]$  from the samples of the spectrum  $X(\omega)$ .

However, we would like to also recover  $x[n]$  and  $X(\omega)$ . The former can be recovered from  $x_p[n]$  if there is no *time-domain aliasing*, that is,  $x[n]$  is time-limited to less than the period  $N$ .



If there is no time-domain aliasing, i.e.

$$x[n] = \begin{cases} x_p[n], & 0 \leq n \leq N-1 \\ 0, & \text{elsewhere} \end{cases}$$

it is possible to express  $X(\omega)$  in terms of its samples. Since  $x[n]$  is time-limited we can use appropriate integral boundaries

$$x[n] = \frac{1}{N} \sum_{k=0}^{N-1} X(2\pi k/N) e^{j2\pi kn/N}, n = 0, 1, \dots, N-1$$

Now taking the Fourier Transform

$$\begin{aligned} X(\omega) &= \sum_{n=0}^{N-1} \left( \frac{1}{N} \sum_{k=0}^{N-1} X(2\pi k/N) e^{j2\pi kn/N} \right) e^{-j\omega n} = \\ &= \sum_{k=0}^{N-1} X(2\pi k/N) \left( \frac{1}{N} \sum_{n=0}^{N-1} e^{-j(\omega - 2\pi k/N)n} \right) \end{aligned}$$

The inner term in parentheses represents the interpolation function (shifted by  $2\pi k/N$  in frequency). Let us define the unshifted version

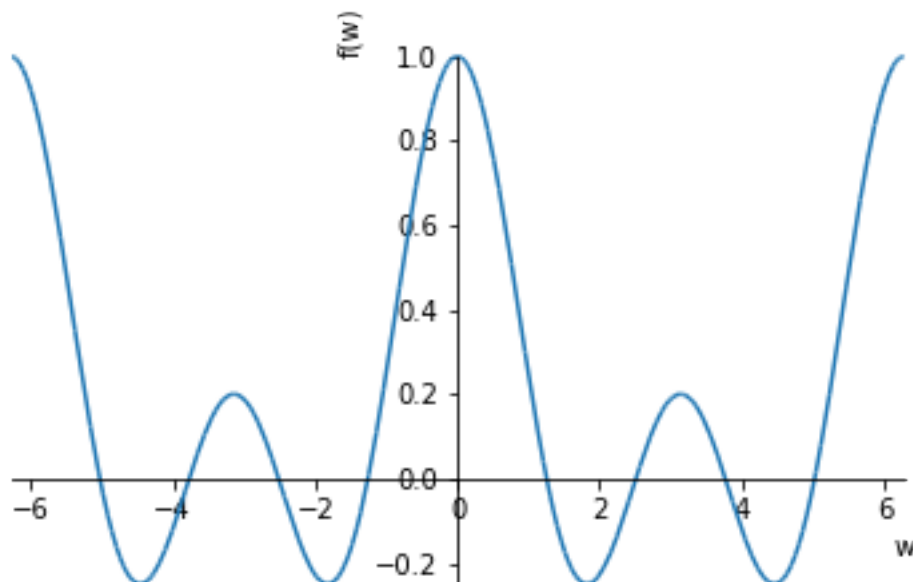
$$\begin{aligned} P(\omega) &= \frac{1}{N} \sum_{n=0}^{N-1} e^{-j\omega n} = \frac{1}{N} \frac{1 - e^{-j\omega N}}{1 - e^{-j\omega}} \\ &= \frac{\sin(\omega N/2)}{N \sin(\omega/2)} e^{-j\omega(N-1)/2} \end{aligned}$$

This is the periodic counterpart to the familiar sinc function. Then

$$X(\omega) = \sum_{k=0}^{N-1} X(2\pi k/N) P(\omega - 2\pi k/N)$$

```
In [2]: # visualize part of interpolation function P
w = sym.symbols('w', real=True)
N, k = sym.symbols('N k', integer=True)
P = sym.sin(w * N / 2) / (N * sym.sin(w / 2))

sym.plot(P.subs({N: 5}), (w, -2 * sym.pi, 2 * sym.pi));
```



$P(\omega)$  has the property that  $P(0) = 1$  and

$$P\left(\frac{2\pi k}{N}\right) = 0, k = 1, 2, \dots, N-1$$

Consequently, the interpolation function gives exactly the sample values  $X(2\pi j/N)$  for  $\omega = 2\pi k/N$ ; for all other frequencies the result is a linear combination of the original spectral samples.

### 1.1.1 7.1.2. The Discrete Fourier Transform

In general, the equally-spaced frequency samples  $X(2\pi k/N), k = 0, 1, \dots, N-1$  do *not* uniquely represent the original sequence  $x[n]$ ; instead they correspond to a periodic sequence  $x_p[n]$  of period  $N$  where  $x_p[n]$  may be an aliased version of  $x[n]$

$$x_p[n] = \sum_{l=-\infty}^{\infty} x[n - lN]$$

When the sequence  $x[n]$  has a finite duration of length  $L \leq N$ , then  $x_p[n]$  is simply a periodic repetition of  $x[n]$ , where

$$x_p[n] = \begin{cases} x[n], & 0 \leq n \leq L-1 \\ 0, & L \leq n \leq N-1 \end{cases}$$

In this case the frequency samples *uniquely* represent the finite-duration sequence  $x[n]$ , and the Fourier Transform is

$$X(\omega) = \sum_{n=0}^{L-1} x[n]e^{-j\omega n}$$

When we sample this we obtain

$$X[k] \equiv X(2\pi k/N) = \sum_{n=0}^{L-1} x[n]e^{-j2\pi kn/N}$$

and without loss of generality we can increase the upper index in the sum to  $N - 1$ :

$$X[k] = \sum_{n=0}^{N-1} x[n]e^{-j2\pi kn/N}, k = 0, 1, 2, \dots, N - 1$$

which is the *discrete Fourier Transform* (DFT). The inverse DFT (IDFT) is

$$x[n] = \frac{1}{N} \sum_{k=0}^{N-1} X[k]e^{j2\pi kn/N}, n = 0, 1, \dots, N - 1$$

**Example 7.1.2** A finite-duration sequence of length  $L$  is given as

$$x[n] = \begin{cases} 1, & 0 \leq n \leq L - 1 \\ 0, & \text{otherwise} \end{cases}$$

Determine the DFT of this sequence for  $N \geq L$ :

$$X[k] = \sum_{n=0}^{N-1} x[n]e^{-j2\pi kn/N}, k = 0, 1, 2, \dots, N - 1$$

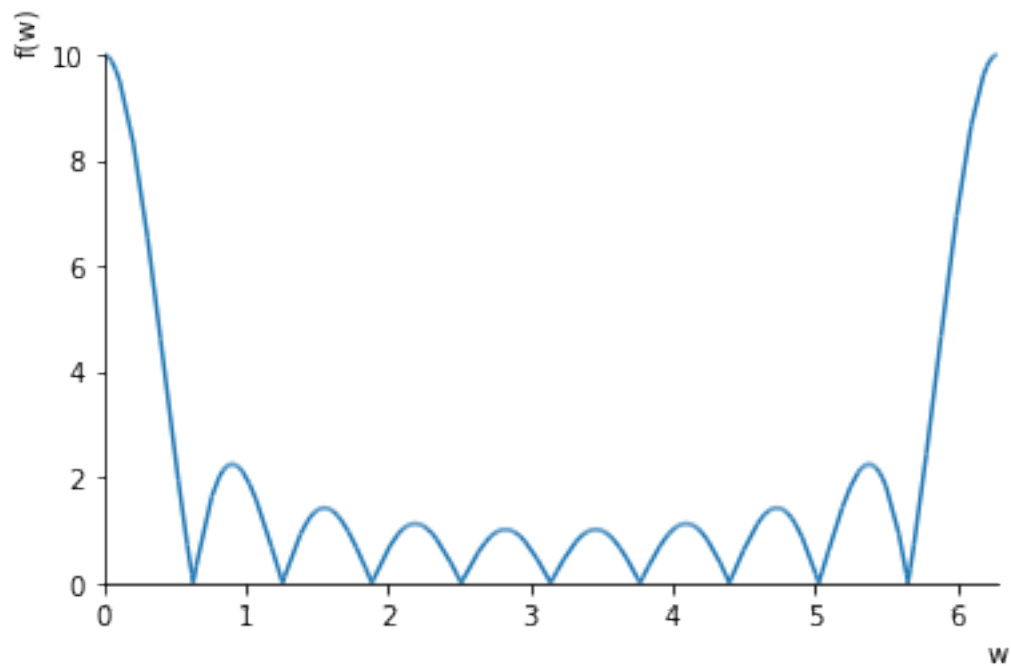
The Fourier Transform is

$$\begin{aligned} X(\omega) &= \sum_{n=0}^{L-1} x[n]e^{-j\omega n} \\ &= \sum_{n=0}^{L-1} e^{-j\omega n} = \frac{1 - e^{-j\omega L}}{1 - e^{-j\omega}} \end{aligned}$$

In [3]: # code here

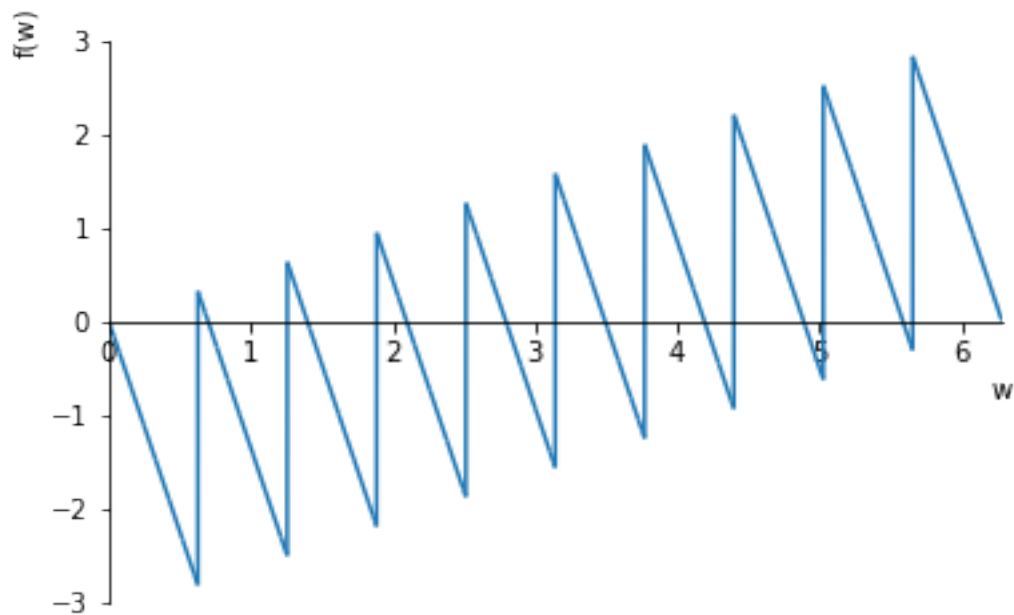
```
w = sym.symbols('w', real=True)
L, N, k = sym.symbols('L N k', integer=True)
X = (1 - sym.exp(-1j * w * L)) / (1 - sym.exp(-1j * w))

sym.plot(sym.Abs(X.subs({L: 10})), (w, 0, 2 * sym.pi));
```



Clearly, at  $X(0) = L$ .

In [4]: `sym.plot(sym.arg(X.subs({L: 10})), (w, 0, 2 * sym.pi));`



The DFT is  $X(\omega)$  evaluated at  $\omega = 2\pi k/N$ :

```
In [5]: X.subs({w: 2 * sym.pi * k / N})
```

Out[5]:

$$\frac{1 - e^{-\frac{2.0i\pi Lk}{N}}}{1 - e^{-\frac{2.0i\pi k}{N}}}$$

If  $N = L$  then:

```
In [6]: X.subs({w: 2 * sym.pi * k / L})
```

Out[6]:

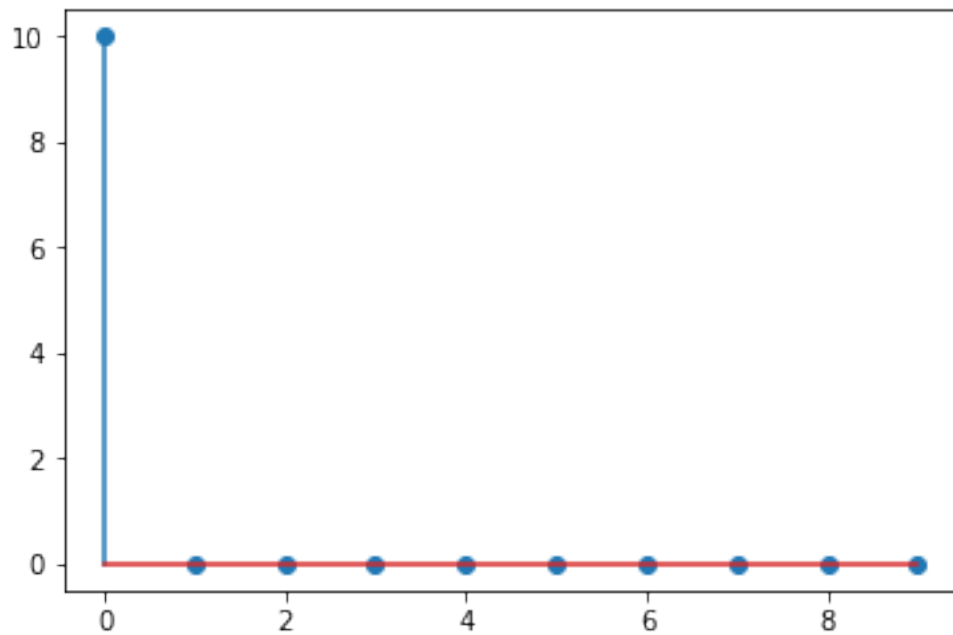
$$\frac{1 - e^{-2.0i\pi k}}{1 - e^{-\frac{2.0i\pi k}{L}}}$$

which is zero for any  $k = 1, 2, \dots, L - 1$ , in other words

$$X[k] = \begin{cases} L, & k = 0 \\ 0, & k = 1, 2, \dots, L - 1 \end{cases}$$

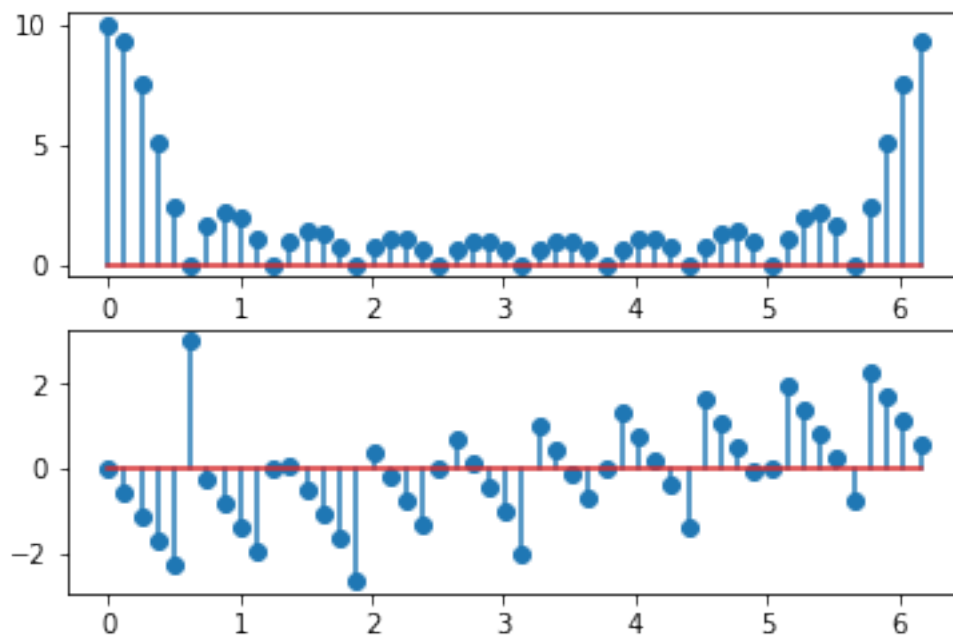
```
In [7]: # calculate this numerically as opposed to analytically
from numpy.fft import fft # fft is a fast *implementation* of the dft
```

```
L = 10
x = np.ones(L)
plt.stem(abs(fft(x))); # N = L
```



Although the  $L$ -point DFT is sufficient to uniquely represent the sequence  $x[n]$  in the frequency domain, one can see that it does not provide sufficient detail to give a good picture of the underlying spectral characteristics of  $x[n]$ . We can interpolate  $X(\omega)$  at more closely-spaced frequencies by increasing  $N$ , which is equivalent to appending  $N - L$  zeros to the sequence, also known as *zero padding*:

```
In [11]: N = 50 # try 50 or 100
w = np.arange(N) * 2 * np.pi / N
plt.subplot(2,1,1)
plt.stem(w, abs(fft(x, N))); # second argument specifies the desired # of points
plt.subplot(2,1,2)
plt.stem(w, np.angle(fft(x, N))); # note what happens when the magnitude is zero
```



### 1.1.2 7.1.4. Relationship of the DFT to Other Transforms

**Relationship to the Fourier series coefficients of a period sequence** A periodic sequence can be represented in the form

$$x_p[n] = \sum_{k=0}^{N-1} c_k e^{j2\pi nk/N}$$

where the Fourier coefficients are

$$c_k = \frac{1}{N} \sum_{n=0}^{N-1} x_p[n] e^{-j2\pi nk/N}, k = 0, 1, \dots, N-1$$



In comparison, we see that this has the identical form of a DFT; if we let  $x[n] = x_p[n]$ ,  $n = 0, 1, \dots, N-1$ , then the DFT is

$$X[k] = Nc_k$$

In conclusion, the  $N$ -point DFT provides the line spectrum of the periodic extension of a finite signal  $x[n]$ .

**Relationship to the Fourier transform of an aperiodic sequence** We have already shown that given a periodic finite sequence  $x[n]$  with Fourier transform  $X(\omega)$ , if we sample the spectrum

$$X[k] = X(\omega)|_{\omega=2\pi k/N} = \sum_{n=-\infty}^{\infty} x[n]e^{-j2\pi nk/N}, k = 0, 1, \dots, N-1$$

then the  $X[k]$  are the DFT coefficients of the periodic sequence

$$x_p[n] = \sum_{l=-\infty}^{\infty} x[n - lN]$$

The finite-duration sequence

$$\hat{x}[n] = \begin{cases} x_p[n], & 0 \leq n \leq N-1 \\ 0, & \text{otherwise} \end{cases}$$

is distinct from the original sequence  $x[n]$ , unless  $x[n]$  has finite duration of length  $L \leq N$ , in which case  $\hat{x}[n] = x[n]$ , and the IDFT of  $X[k]$  will return the original sequence  $x[n]$ .

**Relationship to the z-transform** Consider a sequence  $x[n]$  with z-transform

$$X(z) = \sum_{n=-\infty}^{\infty} x[n]z^{-n}$$

with a ROC that includes the unit circle. If  $X(z)$  is sampled we obtain

$$\begin{aligned} X[k] &= X(z)|_{z=e^{j2\pi k/N}}, k = 0, 1, \dots, N-1 \\ &= \sum_{n=-\infty}^{\infty} x[n]e^{-j2\pi nk/N} \end{aligned}$$

If the sequence  $x[n]$  has finite duration of length  $\leq N$ , the sequence can be recovered from its  $N$ -point DFT; thus its z-transform is uniquely determined by its  $N$ -point DFT as follows:

$$\begin{aligned} X(z) &= \sum_{n=0}^{N-1} x[n]z^{-n} \\ &= \sum_{n=0}^{N-1} \left( \frac{1}{N} \sum_{k=0}^{N-1} X[k]e^{j2\pi kn/N} \right) z^{-n} \\ &= \frac{1}{N} \sum_{k=0}^{N-1} X[k] \sum_{n=0}^{N-1} e^{j2\pi kn/N} z^{-n} \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{N} \sum_{k=0}^{N-1} X[k] \sum_{n=0}^{N-1} \left( e^{j2\pi k/N} z^{-1} \right)^n \\
&= \frac{1}{N} \sum_{k=0}^{N-1} X[k] \frac{1 - (e^{j2\pi k/N} z^{-1})^N}{1 - e^{j2\pi k/N} z^{-1}} \\
&= \frac{1 - z^{-N}}{N} \sum_{k=0}^{N-1} \frac{X[k]}{1 - e^{j2\pi k/N} z^{-1}}
\end{aligned}$$

When this expression is evaluated on the unit circle:

$$\begin{aligned}
X(\omega) &= \frac{1 - e^{-j\omega N}}{N} \sum_{k=0}^{N-1} \frac{X[k]}{1 - e^{j2\pi k/N} e^{-j\omega}} \\
&= \frac{1 - e^{-j\omega N}}{N} \sum_{k=0}^{N-1} \frac{X[k]}{1 - e^{-j(\omega - 2\pi k/N)}}
\end{aligned}$$

it gives an expression for the Fourier Transform in terms of a polynomial (Lagrange) interpolation formula; this is reducable to the interpolation formula seen earlier.