Strong, Weak and Finite Element Formulations of 1-D Scalar Problems ME 964: Krishnan Suresh

1. From Strong to Weak

Strong statement:

$$\frac{d}{dx}\left(-a\frac{du}{dx}\right) = f \qquad u(0) = 0; -au_{,x}(1) = -Q \tag{1.1} \label{eq:1.1}$$

To convert into weak form:

1. **Multiply ODE** by a virtual function $u^{v}(x)$ satisfying $u^{v}(0) = 0$:

$$\left(-au_{,x}\right)_{,x}u^{v} = fu^{v} \tag{1.2}$$

2. Shift Derivatives: Recall that

$$(-au_{,x}u^{v})_{,x} \equiv (-au_{,x})_{,x}u^{v} + (-au_{,x})u^{v}_{,x}$$
(1.3)

Thus, replacing LHS

$$\left(-au_{,x}u^{v}\right)_{x} - \left(-au_{,x}\right)u_{,x}^{v} = fu^{v} \tag{1.4}$$

3. Integrate:

$$\int_{0}^{1} \left(-au_{,x}u^{v}\right)_{,x} dx - \int_{0}^{1} \left(-au_{,x}\right)u_{,x}^{v} dx = \int_{0}^{1} fu^{v}(x)dx \tag{1.5}$$

i.e.,

$$\left[-au_{,x}u^{v}\right]_{0}^{1} + \int_{0}^{1} au_{,x}u_{,x}^{v}dx = \int_{0}^{1} fu^{v}(x)dx \tag{1.6}$$

4. Apply boundary conditions

$$-Qu^{v}(1) + \int_{0}^{1} au_{,x}u_{,x}^{v}dx = \int_{0}^{1} fu^{v}(x)dx$$
(1.7)

i.e.,

$$\int_{0}^{1} au_{,x}u_{,x}^{v}dx = \int_{0}^{1} fu^{v}(x)dx + Qu^{v}(1)$$
(1.8)

5. Arrive at weak statement:

Find
$$u(x)$$
, where, $u(0) = 0$, s.t.

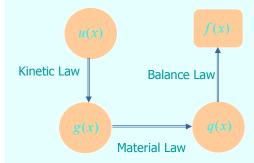
$$\int_{0}^{1} au_{,x}u_{,x}^{v}dx = \int_{0}^{1} fu^{v}(x)dx + Qu^{v}(1)$$
(1.9)

$$\forall u^{v}(x) where, u^{v}(0) = 0$$

2. Graphical Interpretation of Strong & Weak Forms

Strong form:

Strong Form



Kinetic Law: $g(x) = -\frac{du}{dx}$

Fourier Law: $q(x) = -a \frac{du}{dx}$

Balance Law: Heat balance in x

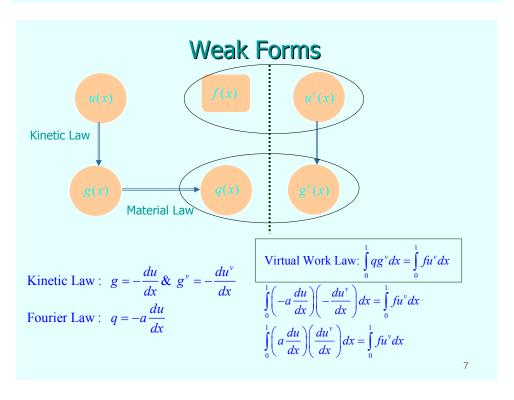
$$q(x+dx) - q(x) = f(x)dx$$

$$\frac{dq}{dx} = f$$

$$\frac{d}{dx}\left(-a\frac{du}{dx}\right) = f$$

Differential Equation: $\frac{d}{dx} \left(-a \frac{du}{dx} \right) = f$

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3. Equivalence

Equation (1.1) is 'equivalent' to Equation (1.9). Consider the example:

$$\frac{d}{dx}\left(-\frac{du}{dx}\right) = 1 \qquad u(0) = 0; u_{,x}(1) = -1 \tag{1.10}$$

i.e., $f=1\ \&\ Q=-1$ The exact solution is $u(x)=-x^2/2$ as one can easily verify. Let us find the same solution via Equation (1.9). Search for a solution of the form $u(x)=Ax+Bx^2$ (satisfies the essential boundary condition). Similarly, let $u^v(x)=A^vx+B^vx^2$. Substituting in Equation (1.9):

Find A & B, s.t.

$$\int_{0}^{1} (A+2Bx)(A^{v}+2B^{v}x)dx = \int_{0}^{1} (A^{v}x+B^{v}x^{2})dx - (A^{v}+B^{v})$$
(1.11)

 $\forall A^v \& B^v$

i.e.,

$$\int_{0}^{1} (A^{v} + 2B^{v}x)(A + 2Bx)dx = \int_{0}^{1} (A^{v}x + B^{v}x^{2})dx - (A^{v} + B^{v})$$
(1.12)

i.e.,

$$\int_{0}^{1} \left\{ A^{v} \quad B^{v} \right\} \begin{cases} 1 \\ 2x \end{cases} \left\{ 1 \quad 2x \right\} \begin{cases} A \\ B \end{cases} dx = \int_{0}^{1} \left\{ A^{v} \quad B^{v} \right\} \begin{cases} x \\ x^{2} \end{cases} dx - \left\{ A^{v} \quad B^{v} \right\} \begin{cases} 1 \\ 1 \end{cases}$$
 (1.13)

i.e.,

$$\begin{cases}
A^{v} & B^{v}
\end{cases} \begin{cases}
\int_{0}^{1} \begin{cases} 1\\ 2x \end{cases} \begin{cases} 1 & 2x \end{cases} dx
\end{cases} \begin{cases}
A\\ B
\end{cases} = \begin{cases}
A^{v} & B^{v}
\end{cases} \int_{0}^{1} \begin{cases} x\\ x^{2} \end{cases} dx - \begin{cases}
A^{v} & B^{v}
\end{cases} \begin{cases}
1\\ 1
\end{cases}$$
(1.14)

i.e.,

$$\begin{cases}
A^{v} & B^{v} \\
\int_{0}^{1} \begin{bmatrix} 1 & 2x \\ 2x & 4x^{2} \end{bmatrix} dx
\end{cases}
\begin{cases}
A \\
B
\end{cases} = \begin{cases}
A^{v} & B^{v} \\
\int_{0}^{1} \begin{cases} x \\ x^{2} \end{bmatrix} dx - \begin{cases}
A^{v} & B^{v} \\
1 \end{cases}
\end{cases} (1.15)$$

i.e.,

Find A & B, s.t.

$$\left\{ A^{v} \quad B^{v} \right\} \begin{bmatrix} 1 & 1 \\ 1 & 4/3 \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} = \left\{ A^{v} \quad B^{v} \right\} \begin{bmatrix} -1/2 \\ -2/3 \end{bmatrix}
 \tag{1.17}$$

 $\forall A^v \& B^v$

This is true if and only if:

$$\begin{bmatrix} 1 & 1 \\ 1 & 4/3 \end{bmatrix} A B = \begin{cases} -1/2 \\ -2/3 \end{cases}$$
 (1.18)

i.e., A=0, B=-1/2, i.e., $u(x)=-x^2/2$. We obtain the same solution.

4. Why Weak Formulation?

Consider the example:

$$\frac{d}{dx}\left(-\frac{du}{dx}\right) = \frac{-1}{1+x^2} \qquad u(0) = 0; u_{,x}(1) = 0 \tag{1.19}$$

It is difficult to integrate Equation (1.19) to obtain the exact solution. So, let us try an approximate solution: $u(x) = Ax + Bx^2$ where A and B are constants (satisfies the essential boundary condition). Substituting in Equation (1.19), we have:

$$-B = \frac{-1}{1+x^2} \qquad A+2B=0 \tag{1.20}$$

The above suggests that B is not a constant (that's about it!!).

Now consider the weak formulation. As before, let $u(x) = Ax + Bx^2$ and

$$u^{v}(x) = A^{v}x + B^{v}x^{2}$$
. Substituting:

Find A & B, s.t.

$$\left\{ A^{v} \quad B^{v} \right\} \begin{bmatrix} 1 & 1 \\ 1 & 4/3 \end{bmatrix} \begin{Bmatrix} A \\ B \end{Bmatrix} = \int_{0}^{1} (A^{v}x + B^{v}x^{2}) \left(\frac{-1}{1+x^{2}} \right) dx \tag{1.21}$$

 $\forall A^v \& B^v$

Evaluate right hand side numerically, we have:

$$\begin{bmatrix} 1 & 1 \\ 1 & 4/3 \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} = \begin{cases} -0.3469 \\ -0.21439 \end{cases}$$
 (1.22)

i.e.,
$$u(x) \approx -0.75x + 0.4x^2$$

Weak formulations naturally promote computing approximate solutions to challenging problems, and are 'equivalent' to strong forms.

5. Finite Element Solutions of Weak Formulation

Consider the model problem:

Find u(x), where, u(0) = 0, s.t.

$$\int_{0}^{1} au_{,x}u_{,x}^{v}dx = \int_{0}^{1} fu^{v}(x)dx + Qu^{v}(1)$$
(1.23)

$$\forall u^v(x) where, u^v(0) = 0$$

Break up the domain into finite elements; thus:

$$\sum_{e} \int_{x_{i}}^{x_{i+1}} a u_{,x} u_{,x}^{v} dx = \sum_{e} \int_{x_{i}}^{x_{i+1}} f u^{v} dx + Q u^{v}(1)$$
(1.24)

Define linear shape functions:

$$L_1(\xi) = \left(\frac{1-\xi}{2}\right); L_2(\xi) = \left(\frac{1+\xi}{2}\right), -1 \le \xi \le 1$$
 (1.25)

Now, let (over a single element):

$$x = L_1(\xi)x_i + L_2(\xi)x_{i+1}$$

$$u = L_1(\xi)u_i + L_2(\xi)u_{i+1}$$

$$u^v = L_1(\xi)u_i^v + L_2(\xi)u_{i+1}^v$$
(1.26)

Note that:

and

$$\frac{dx}{d\xi} = L_{1,\xi} x_i + L_{2,\xi} x_{i+1} = \frac{x_{i+1} - x_i}{2} = \frac{h}{2}$$

$$u_{,x} = \frac{2}{h} \left(L_{1,\xi} u_i + L_{2,\xi} u_{i+1} \right) = \frac{2}{h} \left\{ L_{1,\xi} \quad L_{2,\xi} \right\} \begin{cases} u_i \\ u_{i+1} \end{cases}$$

$$u_{,x}^v = \frac{2}{h} \left(L_{1,\xi} u_i^v + L_{2,\xi} u_{,i+1}^v \right) = \frac{2}{h} \left\{ u_i^v \quad u_{,i+1}^v \right\} \begin{cases} L_{1,\xi} \\ L_{2,\xi} \end{cases}$$
(1.27)

Considering integration over a single element:L

$$\int_{x_{i}}^{x_{i+1}} a u_{,x} u_{,x}^{v} dx = \int_{-1}^{1} a \frac{2}{h} \left\{ u_{i}^{v} - u_{,i+1}^{v} \right\} \left\{ L_{1,\xi} \atop L_{2,\xi} \right\} \frac{2}{h} \left\{ L_{1,\xi} - L_{2,\xi} \right\} \left\{ u_{i} \atop u_{i+1} \right\} \frac{h}{2} d\xi$$

$$= \frac{2}{h} \left\{ u_{i}^{v} - u_{,i+1}^{v} \right\} \left(\int_{-1}^{1} a \begin{bmatrix} L_{1,\xi} L_{1,\xi} - L_{1,\xi} L_{2,\xi} \\ L_{2,\xi} L_{1,\xi} - L_{2,\xi} L_{2,\xi} \end{bmatrix} d\xi \right) \left\{ u_{i} \atop u_{i+1} \right\} \tag{1.28}$$

 $\int\limits_{\cdot}^{x_{i+1}}fu^{v}dx=\int\limits_{\cdot}^{1}f\left\{ u_{i}^{v}\quad u_{,i+1}^{v}\right\} \begin{cases} L_{1}\\L_{2}\end{cases} \frac{h}{2}d\xi$

Thus, from Equation (1.24) and Equation (1.28), we have:

$$\sum_{e} \frac{2}{h} \left\{ u_{i}^{v} \quad u_{,i+1}^{v} \right\} \left(\int_{-1}^{1} a \begin{bmatrix} L_{1,\xi} L_{1,\xi} & L_{1,\xi} L_{2,\xi} \\ L_{2,\xi} L_{1,\xi} & L_{2,\xi} L_{2,\xi} \end{bmatrix} d\xi \right) \left\{ u_{i} \\ u_{i+1} \right\} \\
= \sum_{e} \int_{-1}^{1} f \left\{ u_{i}^{v} \quad u_{,i+1}^{v} \right\} \left\{ L_{1} \\ L_{2} \right\} \frac{h}{2} d\xi + Q u^{v} (1) \tag{1.29}$$

The element stiffness matrix and element forcing vector are:

$$K_{e} = \frac{h}{2} \left(\int_{-1}^{1} a \begin{bmatrix} L_{1,x} L_{1,x} & L_{1,x} L_{2,x} \\ L_{2,x} L_{1,x} & L_{2,x} L_{2,x} \end{bmatrix} d\xi \right)$$

$$f_{e} = \int_{-1}^{1} f \begin{Bmatrix} L_{1} \\ L_{2} \end{Bmatrix} \frac{h}{2} d\xi$$
(1.30)

After assembly, and eliminating the virtual variables, we have:

$$Ku = f (1.31)$$

where

$$K = \sum_{e} K_{e}$$

$$f = \sum_{e} f_{e} + f_{Q}$$
(1.32)

The essential (Dirichlet) boundary conditions are imposed via the Lagrange multiplier method:

$$\overline{K}\overline{u} = \overline{f} \tag{1.33}$$

where

$$\overline{K} = \begin{bmatrix} K & C^T \\ C & 0 \end{bmatrix}, \overline{u} = \begin{bmatrix} u \\ \lambda \end{bmatrix}, \overline{f} = \begin{bmatrix} f \\ b \end{bmatrix}$$
 (1.34)

where C & b capture the boundary condition at x=0. The variables λ are the Lagrange multipliers that contain useful information about the sensitivity of the essential boundary conditions.

6. Non-Linear Problems

Equations (1.31) and (1.33) is valid for both linear and non-linear problems. For linear problems, the equation is solved once since the stiffness matrix and forcing vector are independent of u. On the other hand, for non-linear problems, either the stiffness matrix or the forcing vector or both is dependent on u. So, Equation (1.31) reads:

$$[K(u)]u = f(u) \tag{1.35}$$

where:

$$K(u) = \sum_{e} K_{e}(u)$$

$$f(u) = \sum_{e} f_{e}(u) + f_{Q}$$
(1.36)

and

$$K_{e}(u) = \frac{2}{h} \left[\int_{-1}^{1} a(u) \begin{bmatrix} L_{1,\xi} L_{1,\xi} & L_{1,\xi} L_{2,\xi} \\ L_{2,\xi} L_{1,\xi} & L_{2,\xi} L_{2,\xi} \end{bmatrix} d\xi \right]$$

$$f_{e}(u) = \int_{-1}^{1} f(u) \begin{cases} L_{1} \\ L_{2} \end{cases} \frac{h}{2} d\xi$$
(1.37)

while Equation (1.33) reads

$$\overline{K}(u)\overline{u} = \overline{f}(u) \tag{1.38}$$

i.e., a non-linear problem.

We will consider two different methods of solving Equation (1.38):

- 1. Picard iteration
- 2. Newton Raphson iteration

6.1 Picard Iteration

Theorems (for Picard's Method)

Theorem 1: (Banach's Fixed Point Theorem) If $\varphi(x): \Re^n \to \Re^n$ is a continuous function such that:

$$d(\varphi(x), \varphi(y)) < d(x, y), \forall x, y \tag{1.39}$$

where d(.,.) is an appropriate metric defined in \Re^n , then the iteration:

$$x^{n+1} = \varphi(x^n) \tag{1.40}$$

will converge (to a fixed point of $\varphi(x)$).

Theorem 2: If $\varphi(x)$ is a differentiable function in a range [a,b], then $\varphi(x)$ has a unique fixed point if $|\varphi_{,x}(x)| < 1$ for all $x \in [a,b]$.

Proof: Mean Value Theorem.

From Equation (1.38), $\varphi(\overline{u}) : \left| \overline{K}(\overline{u}) \right|^{-1} \overline{f}(\overline{u})$

Suppose:

$$\left\| \left[\overline{K}(\overline{u}) \right]^{-1} \overline{f}(\overline{u}) - \left[\overline{K}(\overline{v}) \right]^{-1} \overline{f}(\overline{u}) \right\| < \left\| \overline{u} - \overline{v} \right\|, \forall u, v$$
(1.41)

then the Picard iteration for Equation (1.35) will proceed as follows:

- 1. Guess u^0 (say 0). Set n = 0.
- 2. Assemble $\left[\overline{K}(\overline{u}^n)\right] \& \overline{f}(\overline{u}^n)$
- 3. Compute $u^{n+1} = \left[\overline{K}(\overline{u}^n) \right]^{-1} \overline{f}(\overline{u}^n)$
- 4. Stop when $\left\|u^{^{n+1}}-u^{^{n}}\right\|<\varepsilon$, else go to step 2.

Of course, establishing Equation (1.41) is not easy ... a case-by-case analysis at best! We will later consider stiffening versus softening non-linear problems, and study the convergence of Picard iteration.

6.2 Newton-Raphson Iteration

Theorem (for N-R Method): If h(x) is a second order differentiable in [a,b], and if $h(x^*) = 0$ where $x^* = [a,b]$, and $h_x(x^*) = 0$ then the sequence:

$$x^{n+1} = x^n - h(x^n) / h_{,x}(x^n)$$
(1.42)

will converge to x^*

Proof: See popular text-books on numerical methods.

Derivation of the Tangent Matrix

The non-linear equation that we must solve is:

$$R(u) = [K(u)]u - f = 0 (1.43)$$

Applying Equation (1.42):

$$u^{n+1} = u^n - [T(u^n)]^{-1} R(u^n)$$
(1.44)

where $T(u^n)$ is the tangent or Jacobian matrix given by:

$$T_{ij} = \frac{\partial R_i}{\partial u_j} = \frac{\partial \left(\sum_m K_{im} u_m\right)}{\partial u_j}$$
(1.45)

Considering a single element (for linear shape functions m = 1, 2):

$$T_{ij}^{e} = \frac{\partial \left(\sum_{m=1}^{2} K_{im}^{e} u_{m}\right)}{\partial u_{j}} = \sum_{m=1}^{2} \frac{\partial \left(K_{im}^{e} u_{m}\right)}{\partial u_{j}} = \sum_{m=1}^{2} \frac{\partial K_{im}^{e}}{\partial u_{j}} u_{m} + \sum_{m=1}^{2} K_{im}^{e} \frac{\partial u_{m}}{\partial u_{j}}$$
(1.46)

Considering the definition of element stiffness matrix in Equation (1.30), we have:

$$\frac{\partial K_{im}^e}{\partial u_j} = \frac{\partial}{\partial u_j} \left(\frac{h}{2} \int_{-1}^1 a(x, u) L_{i,x} L_{m,x} d\xi \right) = \left(\frac{h}{2} \int_{-1}^1 a_{,u}(x, u) \frac{\partial u}{\partial u_j} L_{i,x} L_{m,x} d\xi \right)$$
(1.47)

Further, since $u = L_1u_1 + L_2u_2$, we have:

$$\frac{\partial K_{im}^e}{\partial u_j} = \left(\frac{h}{2} \int_{-1}^1 a_{,u}(x,u) L_j L_{i,x} L_{m,x} d\xi\right) \tag{1.48}$$

The second term in Equation (1.46):

$$\sum_{m=1}^{2} K_{im}^{e} \frac{\partial u_{m}}{\partial u_{i}} = K_{ij}^{e} \tag{1.49}$$

Thus, substituting Equation (1.48) and Equation (1.49) in Equation (1.46), we have:

$$T_{ij}^{e} = \sum_{m=1}^{2} \left(\frac{h}{2} \int_{-1}^{1} a_{,u}(x,u) L_{j} L_{i,x} L_{m,x} d\xi \right) u_{m} + K_{ij}^{e}$$
(1.50)

$$T_{ij}^{e} = \left(\frac{h}{2} \int_{-1}^{1} a_{,u}(x,u) L_{j} L_{i,x} d\xi\right) \sum_{m=1}^{2} L_{m,x} u_{m} + K_{ij}^{e}$$
(1.51)

Thus, Equation (1.51) is used to assemble the tangent matrix.

Applying zero boundary conditions on Δu^n , we have:

$$\begin{bmatrix} T(u^n) & C^T \\ C & 0 \end{bmatrix} \begin{bmatrix} \Delta u^n \\ \zeta \end{bmatrix} = \begin{bmatrix} R(u^n) \\ 0 \end{bmatrix}$$
 (1.52)

Thus, we must assemble both the stiffness matrix and the tangent matrix $T(u^n)$ defined via Equation in each step.

Thus, the Newton-Raphson iteration will proceed as follows:

- 1. Guess u^0 (say 0). Set n = 0.
- 2. Assemble $[K(u^n)] \& f(u^n)$
- 3. Assemble/ compute $\left[T(u^n)\right] \& R(u^n)$
- 4. Compute $\Delta \overline{u}^{n+1} = \left[\overline{T}(\overline{u}^n)\right]^{-1} \overline{R}(\overline{u}^n)$ per Equation (1.52)
- 5. Compute $u^{n+1} = u^n + \Delta u^{n+1}$
- 6. Stop when $\left\|u^{n+1}-u^{n}\right\|<\varepsilon$, else go to step 2.