## HPS/Pl 125: Problem 7

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Complete exercise 2.71 and 2.75 from Nielsen and Chuang.

## 1 2.71

Show that  $Tr(\rho^2) \le 1$  with equality iff  $\rho$  is a pure state.

Begin with the definition of  $\rho$ :  $\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|$ . Then:

$$\rho^{2} = \left(\sum_{i} p_{i} |\psi_{i}\rangle \langle \psi_{i}|\right) \left(\sum_{j} p_{j} |\psi_{j}\rangle \langle \psi_{j}|\right) = \sum_{i,j} p_{i} p_{j} |\psi_{i}\rangle \langle \psi_{i}|\psi_{j}\rangle \langle \psi_{j}|$$

Note from equation 2.157 of the text that in the spectral decomposition of the density matrix, the state-space vectors form an orthogonal basis. Therefore, unless i = j,  $\langle \psi_i | \psi_i \rangle = 0$ . Thus:

$$\rho^2 = \sum_i p_i^2 |\psi_i\rangle \langle \psi_i|$$

Applying the trace and using equation 2.153:

$$\operatorname{Tr}(\rho^2) = \sum_{i} p_i^2$$

In a pure state, the only non-zero  $p_i$  is 1. Therefore,  $\text{Tr}(\rho^2) = 1$  in a pure state. In a mixed state, we have multiple non-zero  $p_i$  which are subject to the condition  $\sum_i p_i = 1$ . Thus each  $p_i < 1$ . If  $p_i < 1$ , then  $p_i^2 < p_i$ . Therefore,  $\sum_i p_i^2 < \sum_i p_i = 1$ . Thus, in the case of a mixed state,  $\text{Tr}(\rho^2) < 1$ . Therefore  $\text{Tr}(\rho^2) = 1$  iff  $\rho$  is a pure state.

## 2 2.75

Find the reduced density operators for each particle in each of the four Bell states:

$$|\beta_{00}\rangle = \frac{|00\rangle + |11\rangle}{\sqrt{2}}, \quad |\beta_{01}\rangle = \frac{|01\rangle + |10\rangle}{\sqrt{2}}$$

$$|\beta_{10}\rangle = \frac{|00\rangle - |11\rangle}{\sqrt{2}}, \quad |\beta_{11}\rangle = \frac{|01\rangle - |10\rangle}{\sqrt{2}}$$

First, find the density operator for each Bell state:

$$\begin{split} \rho_{\beta_{00}} &= |\beta_{00}\rangle \, \langle \beta_{00}| = \frac{|00\rangle \, \langle 00| + |11\rangle \, \langle 00| + |00\rangle \, \langle 11| + |11\rangle \, \langle 11|}{2} \\ \rho_{\beta_{01}} &= |\beta_{01}\rangle \, \langle \beta_{01}| = \frac{|01\rangle \, \langle 01| + |10\rangle \, \langle 01| + |01\rangle \, \langle 10| + |10\rangle \, \langle 10|}{2} \\ \rho_{\beta_{10}} &= |\beta_{10}\rangle \, \langle \beta_{10}| = \frac{|00\rangle \, \langle 00| - |11\rangle \, \langle 00| - |00\rangle \, \langle 11| + |11\rangle \, \langle 11|}{2} \\ \rho_{\beta_{11}} &= |\beta_{11}\rangle \, \langle \beta_{11}| = \frac{|01\rangle \, \langle 01| - |10\rangle \, \langle 01| - |01\rangle \, \langle 10| + |10\rangle \, \langle 10|}{2} \end{split}$$

Apply the reduced density operator (partial trace) to each Bell state for each particle.

$$\rho^1 = \operatorname{Tr}_2(\rho), \quad \rho^2 = \operatorname{Tr}_1(\rho)$$

Note that the following general pattern will emerge:

$$\operatorname{Tr}_{2}(|AB\rangle\langle CD|) = \operatorname{Tr}_{2}(|A\rangle\langle C|\otimes|B\rangle\langle D|) = |A\rangle\langle C|\operatorname{Tr}(|B\rangle\langle D|) = |A\rangle\langle C|\langle D|B\rangle$$

Since the qubit basis is orthogonal,  $\langle D|B\rangle=0$  unless D=B. If D=B, then  $\langle D|B\rangle=1$ . Therefore:

$$\rho_{B_{00}}^{1} = \text{Tr}_{2}(\frac{|00\rangle\langle00| + |11\rangle\langle00| + |00\rangle\langle11| + |11\rangle\langle11|}{2}) = \frac{|0\rangle\langle0| + |1\rangle\langle1|}{2} = \boxed{\frac{I}{2}}$$

Examine the symmetry here. We can see that if we had traced out particle 1 instead of 2, the positions of the outer products in the numerator would've been swapped, i.e  $\frac{|1\rangle\langle 1|+|0\rangle\langle 0|}{2}$ . Therefore

the reduced density operator for particle 2 is the same as for particle 1,  $\left| \frac{I}{2} \right|$ 

$$\rho_{B_{01}}^{1}=\operatorname{Tr}_{2}(\frac{\left|01\right\rangle \left\langle 01\right|+\left|10\right\rangle \left\langle 01\right|+\left|01\right\rangle \left\langle 10\right|+\left|10\right\rangle \left\langle 10\right|}{2})=\frac{\left|0\right\rangle \left\langle 0\right|+\left|1\right\rangle \left\langle 1\right|}{2}=\overline{\left|\frac{I}{2}\right|}$$

An equal symmetry argument applies here as well, so the reduced density operator for particle 2 is the same as for particle 1,  $\frac{I}{2}$ 

$$\rho_{B_{10}}^{1} = \operatorname{Tr}_{2}(\frac{|00\rangle\langle00| - |11\rangle\langle00| - |00\rangle\langle11| + |11\rangle\langle11|}{2}) = \frac{|0\rangle\langle0| + |1\rangle\langle1|}{2} = \boxed{\frac{I}{2}}$$

Symmetry applies here as well, so the reduced density operator for particle 2 is the same as for particle 1,  $\left|\frac{I}{2}\right|$ 

$$\rho_{B_{11}}^{1}=\operatorname{Tr}_{2}(\frac{\left|01\right\rangle \left\langle 01\right|-\left|10\right\rangle \left\langle 01\right|-\left|01\right\rangle \left\langle 10\right|+\left|10\right\rangle \left\langle 10\right|}{2})=\frac{\left|0\right\rangle \left\langle 0\right|+\left|1\right\rangle \left\langle 1\right|}{2}=\overline{\left[\frac{I}{2}\right]}$$

Symmetry applies here as well, so the reduced density operator for particle 2 is the same as for particle 1,  $\left| \frac{I}{2} \right|$