

Gradient descent optimization of MPS for Ground state

Estêvão Gomes

Physik-Department E28, Technische Universität München, 85747 Garching, Germany

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In this article, gradient descent is applied to optimize Matrix Product States for finding the ground state of the transverse-field Ising model, leveraging JAX for automatic differentiation. The results are compared with TEBD and DMRG, emphasizing convergence behavior and accuracy across various regimes. The study highlights both the strengths and limitations of gradient-based approaches for quantum many-body systems.

I. INTRODUCTION

The study of quantum many-body systems constitutes a fundamental pillar of modern condensed matter physics, where complex quantum phenomena emerge from the collective behavior of interacting particles. The primary obstacle in studying these systems computationally is the exponential scaling of the Hilbert space with system size, which results in inaccessibly large computational requirements. Tensor networks have become an essential tool for circumventing this exponential bottleneck, enabling efficient large-scale numerical simulations of quantum many-body systems, at the cost of using truncations that limit the entanglement content captured. [1]

To obtain the ground state of a large quantum many-body system, the usage of exact diagonalization methods is often impractical, thus iterative methods such as the Time-Evolving Block Decimation (TEBD) [2] and Density Matrix Renormalization Group (DMRG) [3] are preferred. In this article, direct gradient descent optimization is applied to the Matrix Product State (MPS) [4] representation of the ground state of a transverse field Ising model (TFI), making use of the automatic differentiation capabilities of JAX [5]. This approach is then compared with the results obtained from the TEBD and DMRG methods.

This article is structured as follows: In section II, we provide the theoretical background of the methods used. In section III, we present and discuss the results obtained from our simulations. Finally, we conclude in section IV.

II. THEORETICAL BACKGROUND

A. Matrix Product States

Matrix Product States provide a powerful representation of quantum many-body states, particularly efficient for systems that obey the area law for entanglement entropy, such as the 1D transverse field Ising model, discussed in detail in section II B.

An MPS is expressed as a product of tensor, where each tensor corresponds to a local degree of freedom (spin j_n). For open boundary conditions (OBC), the MPS is given

by

$$|\Psi\rangle = \sum_{j_1, \dots, j_N} M^{[1]j_1} M^{[2]j_2} \dots M^{[N]j_N} |j_1 j_2 \dots j_N\rangle, \quad (1)$$

where each $M^{[n]j_n}$ tensor has shape (χ_{n-1}, d, χ_n) , with d being the local dimension (e.g., $d = 2$ for spin-1/2 systems), and χ_n being the bond dimension at site n . The indices j_n run over the local degrees of freedom, and $\chi_0 = \chi_N = 1$ for OBC.

For periodic boundary conditions (PBC), the MPS is defined similarly, but with the additional condition that the first and last tensors are connected,

$$|\Psi\rangle = \sum_{j_1, \dots, j_N} \text{Tr} \left\{ M^{[1]j_1} M^{[2]j_2} \dots M^{[N]j_N} \right\} |j_1 j_2 \dots j_N\rangle, \quad (2)$$

where $M^{[1]j_1}$ is connected to $M^{[N]j_N}$ through a cyclic permutation of the indices.

B. Transverse Field Ising model

The transverse field Ising model (TFI) is a fundamental model in quantum many-body physics, described for a spin-1/2 system by the Hamiltonian

$$H = -J \sum_{i=1}^L \sigma_i^x \sigma_{i+1}^x - g \sum_{i=1}^L \sigma_i^z, \quad (3)$$

where J is the coupling constant, g is the transverse field strength, and σ_i^x and σ_i^z are the Pauli matrices acting on the i -th spin. The model exhibits a quantum phase transition at a critical value of the transverse field $g_c = J$, where the system transitions from an ordered phase (for $g < g_c$) to a disordered phase (for $g > g_c$).

When using tensor networks, energy expectation values can be computed more efficiently by representing the Hamiltonian in the form of a Matrix Product Operator (MPO), which can be seen schematically in fig. 1.

In this representation, the tensor at each site, $W^{[n]}$, is defined as

$$W^{[n]} = \begin{bmatrix} \mathbb{I} & \sigma^x & -g\sigma^z \\ 0 & 0 & -J\sigma^z \\ 0 & 0 & \mathbb{I} \end{bmatrix}, \quad (4)$$

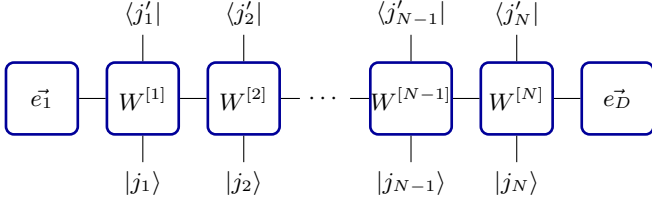


FIG. 1. Schematic representation of an MPO.

with \mathbb{I} being the identity matrix and σ^x, σ^z the Pauli matrices and the vectors \vec{e}_1 and \vec{e}_D are defined as

$$\vec{e}_1 = [1 \ 0 \ 0], \quad (5)$$

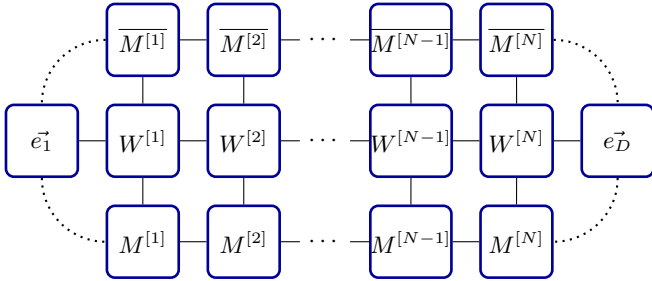
$$\vec{e}_D = [0 \ 0 \ 1]^T. \quad (6)$$

C. Gradient Descent on MPS

To perform gradient descent on an MPS, one first needs to define the loss function to minimize. In this article, we aim to find the ground state of the TFI model, thus a suitable loss to minimize is the normalized energy expectation value,

$$\mathcal{L}(\psi) = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle}, \quad (7)$$

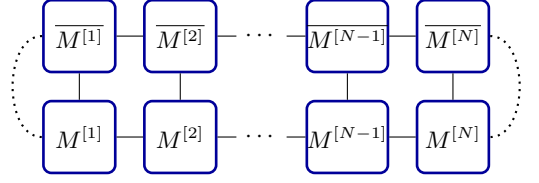
where the energy expectation value is calculated by contracting the MPS with the MPO representation of the Hamiltonian, as shown in fig. 2, and the normalization factor is computed by contracting the MPS with itself, as shown in fig. 3.

FIG. 2. Schematic representation of the energy expectation value, $\langle \psi | H | \psi \rangle$, of an MPS.

To update the MPS tensors $M^{[n]}$ during the optimization process, we compute the gradient of the loss function with respect to each tensor. This can be computed using JAX's automatic differentiation, or using the analytical expression. In section III A, we will compare the results obtained from both methods.

Analytically, the derivative of the loss function with respect to the tensor $M^{[n]}$ is given by

$$\partial_{M^{[n]}} \mathcal{L}(\psi) = \frac{\partial_{M^{[n]}} \langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} - \frac{\langle \psi | H | \psi \rangle \partial_{M^{[n]}} \langle \psi | \psi \rangle}{\langle \psi | \psi \rangle^2} \quad (8)$$

FIG. 3. Schematic representation of the squared norm, $\langle \psi | \psi \rangle$, of an MPS.

with $\partial_{M^{[n]}}$ denoting the derivative with respect to the tensor $M^{[n]}$. This expression can be further simplified by using

$$\partial_{M^{[n]}} \langle \psi | H | \psi \rangle = \langle \partial_{M^{[n]}} \psi | H | \psi \rangle + \langle \psi | H | \partial_{M^{[n]}} \psi \rangle \quad (9)$$

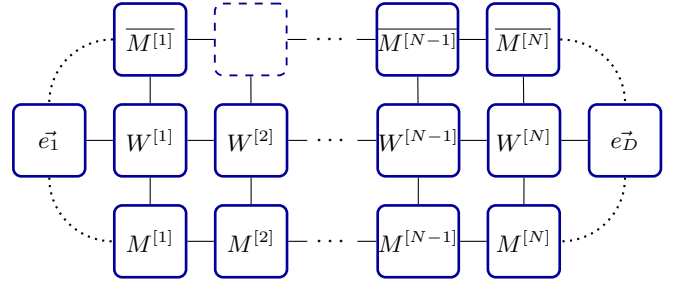
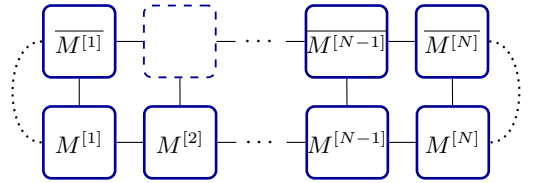
$$= 2\Re(\langle \psi | H | \partial_{M^{[n]}} \psi \rangle), \quad (10)$$

and

$$\partial_{M^{[n]}} \langle \psi | \psi \rangle = \langle \partial_{M^{[n]}} \psi | \psi \rangle + \langle \psi | \partial_{M^{[n]}} \psi \rangle \quad (11)$$

$$= 2\Re(\langle \psi | \partial_{M^{[n]}} \psi \rangle). \quad (12)$$

Equations (10) and (12) can be computed analytically using contractions similar to those used for the energy expectation value, in fig. 2, and normalization factor, in fig. 3, by omitting the contraction with $M^{[n]}$. These can be visualized schematically in figs. 4 and 5, respectively.

FIG. 4. Schematic representation of the derivative of the energy expectation value, $\partial_{M^{[2]}} \langle \psi | H | \psi \rangle / 2$, of an MPS.FIG. 5. Schematic representation of the derivative of the squared norm, $\partial_{M^{[2]}} \langle \psi | \psi \rangle / 2$, of an MPS.

During the optimization process, we update the tensors $M^{[n]}$ using the ADAM optimizer [6], implemented in the OPTAX library [7], as it offers faster convergence compared to standard gradient descent by leveraging both the first and second moments of the gradients.

D. Time Evolving Block Decimation

The TEBD algorithm is used to obtain the time evolution of a quantum state,

$$|\psi(t)\rangle = U(t)|\psi(0)\rangle, \quad (13)$$

where $U(t) = e^{-itH}$ is the real time evolution operator and $U(\tau) = e^{-\tau H}$ the imaginary time evolution operator. The latter can be used to find the ground state of a quantum system via the relation

$$|\psi_{GS}\rangle = \lim_{\tau \rightarrow \infty} \frac{e^{-\tau H}|\psi_0\rangle}{\|e^{-\tau H}|\psi_0\rangle\|}. \quad (14)$$

Put simply, the TEBD algorithm starts by applying a Suzuki-Trotter decomposition, given in first and second order by

$$e^{(X+Y)\delta} = e^{X\delta}e^{Y\delta} + \mathcal{O}(\delta^2), \quad (15)$$

$$e^{(X+Y)\delta} = e^{X\delta/2}e^{Y\delta}e^{X\delta/2} + \mathcal{O}(\delta^3), \quad (16)$$

with δ a small parameter, to a Hamiltonian decomposed as a sum of two-site operators,

$$H = \sum_{n \text{ odd}} H_{n,n+1} + \sum_{n \text{ even}} H_{n,n+1}, \quad (17)$$

where the local Hamiltonian terms acting on the n -th and $(n+1)$ -th sites, $H_{n,n+1}$, commute for either odd or even n .

This allows us to write the time evolution operator for a small timestep $\delta t \ll 1$ as a product of local operators acting on pairs of neighboring sites, given in first order by

$$e^{-\delta t H} = \prod_{n \text{ odd}} e^{-\delta t H_{n,n+1}} \prod_{n \text{ even}} e^{-\delta t H_{n,n+1}}. \quad (18)$$

The MPS is then updated by applying the local operators to the corresponding tensors, followed by a singular value decomposition to truncate the bond dimension and maintain computational efficiency.

E. Density Matrix Renormalization Group

DMRG is a variational algorithm that iteratively optimizes the MPS representation of a quantum state to minimize the energy expectation value.

A DMRG update consists of minimizing the energy by optimizing the tensors $M^{[n]}$ and $M^{[n+1]}$ in the MPS while keeping the remaining chain fixed. This is done by projecting the Hamiltonian into the reduced Hilbert space spanned by the basis $\{|\alpha_n\rangle \otimes |j_n\rangle \otimes |j_{n+1}\rangle \otimes |\alpha_{n+1}\rangle\}$, with $|\alpha_n\rangle$ and $|\alpha_{n+1}\rangle$ representing the left and right virtual bond states of the MPS, and $|j_n\rangle$, $|j_{n+1}\rangle$ are the physical site indices. The effective Hamiltonian in this

subspace is then diagonalized using an iterative eigensolver such as the Lanczos algorithm [8] to obtain the lowest-energy state.

This two-site update is repeated sequentially for all pairs of neighboring sites in the MPS, starting from the leftmost pair and sweeping rightward to the end, then reversing direction and sweeping back to the left.

On that account, the DMRG algorithm can be seen as a “smart” version of gradient descent, as it performs exact diagonalization in a variational subspace, optimizing one (or two) tensors at a time. Unlike gradient descent, which updates all tensors simultaneously based on local gradient information, this local optimization guarantees that the energy never increases after each step.

III. RESULTS

In this section, we present the results obtained from the simulations of the TFI model using the methods described in section II. We then compare the results obtained from the gradient descent optimization of the MPS with those obtained from the TEBD and DMRG methods.

A. Analytical vs Automatic Differentiation

As a baseline consistency check, a comparison between the gradients computed using JAX’s automatic differentiation and the analytical expression derived in section II C is performed. The mean error, its standard deviation and the maximum error are shown in fig. 6, as well as the average runtime of each gradient computation for the whole optimization.

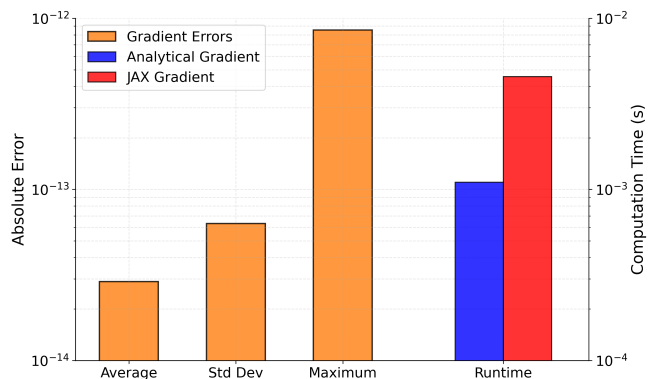


FIG. 6. Error and runtime comparison between analytical and automatic differentiation for the loss in eq. (7).

As expected, the analytical gradients differ from the automatic differentiation results by a value close to machine precision, demonstrating the accuracy of both computations. Moreover, the runtime of the analytical gradients is lower than that of the automatic differentiation,

making it a more efficient choice for large-scale optimizations due to its easy implementation. Nevertheless, the automatic differentiation approach is more flexible and easier to implement for complex models, giving it an edge in terms of usability.

For this project, automatic differentiation could also be used to compute the hessian of the loss, allowing for a faster convergence with the use of Newton's method.

B. Gradient Descent on MPS

To perform gradient descent, the MPS cannot be initialized as a low bond dimension product state such as the all-up state, $|\uparrow \cdots \uparrow\rangle$, as gradient descent cannot increase the bond dimension during optimization. However, simply starting with a high bond dimension product state is not sufficient either: if all tensors are identical, the gradient vanishes and no updates occur.

For this reason, in the remainder of the analysis of the transverse-field Ising model (TFI) with parameters $g > J$, the MPS is initialized as a high bond dimension all-up state with small Gaussian noise added to each tensor.

For the TFI in the regime $g < J$, a more appropriate starting point is the all-right state, $|\rightarrow \cdots \rightarrow\rangle$, with small noise added, as it better approximates the ground state. Furthermore, in the critical case $g = J$, this state remains preferable for initialization. Although the ground state is highly entangled and not well-approximated by any product state, the symmetry of the all-right state is closer to that of the true critical ground state than the all-up state.

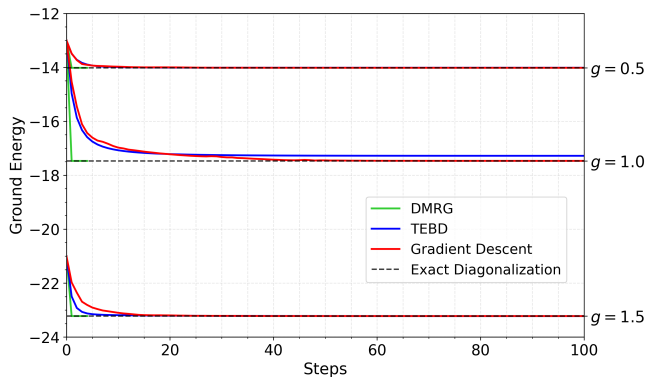


FIG. 7. Comparison of gradient descent with TEBD and DMRG.

Figure 7 shows the results of the optimization of the ground state of a $L = 14$, $J = 1$ TFI for different values of g using gradient descent, compared to the results obtained from the TEBD and DMRG methods. The ground state energy is plotted as a function of the number of the algorithm steps taken to arrive at such state.

For the gradient descent, an exponential learning rate decay was used, starting at 2×10^{-2} and decaying by a factor of 0.5 every 30 steps, as well as a bond dimension of $\chi = 30$.

For the TEBD, a second-order Suzuki-Trotter decomposition was used with an exponential time step decay starting at 10^{-1} and decaying by a factor of 0.1 every 30 steps. This decay was found to give the best results for the $g = 1.5$ case, meaning the TEBD could converge faster for the $g = 1.0$ or $g = 0.5$ by tuning the time step decay accordingly.

For the DMRG, the updates were performed with a maximum bond dimension of 100.

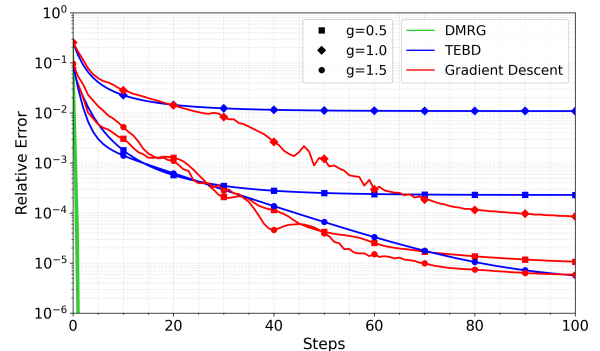


FIG. 8. Comparison of the relative error of the ground state energy obtained from gradient descent, TEBD and DMRG.

Figure 8 shows the relative error of the ground state energy obtained from the 3 methods. From it, it is possible to see the extreme speed at which DMRG converges, taking just one sweep to reach the ground state energy up to machine precision.

It is also possible to notice the worse performance of both TEBD and gradient descent for the critical case $g = 1.0$, where the ground state is highly entangled, leading to a slower and less accurate convergence. It is also possible to see that TEBD starts with a faster convergence than gradient descent, being overtaken by the latter after a few steps, though with an accuracy very close to that of gradient descent for the case $g = 1.5$, where the time step decay was optimized.

IV. CONCLUSION

Overall, the results obtained from simulations of the TFI model using gradient descent on MPS show that this method can be used to obtain accurate ground energies, even for highly entangled states. Nevertheless, DMRG remains the most efficient method for obtaining the ground state of a quantum system, converging to the ground energy in just one sweep, while gradient descent and TEBD require more iterations and careful parameter tuning. Still, gradient-based methods may still be used in systems where DMRG is less suitable, such as those non-Hermitian, with constraints, or when differentiability is needed.

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- [1] J. Hauschild and F. Pollmann, Efficient numerical simulations with Tensor Networks: Tensor Network Python (TeNPy), SciPost Phys. Lect. Notes , 5 (2018).
 - [2] G. Vidal, Efficient simulation of one-dimensional quantum many-body systems, Phys. Rev. Lett. **93**, 040502 (2004).
 - [3] S. R. White, Density matrix formulation for quantum renormalization groups, Phys. Rev. Lett. **69**, 2863 (1992).
 - [4] M. Fannes, B. Nachtergaele, and R. F. Werner, Finitely correlated states on quantum spin chains, Communications in Mathematical Physics **144**, 443 (1992).
 - [5] J. Bradbury, R. Frostig, P. Hawkins, M. J. Johnson, C. Leary, D. Maclaurin, G. Necula, A. Paszke, J. VanderPlas, S. Wanderman-Milne, and Q. Zhang, JAX: composable transformations of Python+NumPy programs, <https://github.com/google/jax> (2018), version 0.3.13.
 - [6] D. P. Kingma and J. L. Ba, Adam: A method for stochastic optimization, 3rd International Conference on Learning Representations, ICLR 2015 - Conference Track Proceedings (2014).
 - [7] DeepMind and Contributors, Optax: A gradient processing and optimization library for JAX, <https://github.com/deepmind/optax> (2020).
 - [8] C. Lanczos, An iteration method for the solution of the eigenvalue problem of linear differential and integral operators, Journal of Research of the National Bureau of Standards **45**, 255 (1950).
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Appendix A: Code listing

1. optimize.py

```

1 import os
2 from time import perf_counter as timer
3
4 import jax.numpy as jnp
5 from jax import jit, grad, config, block_until_ready
6 config.update("jax_enable_x64", True)
7 import optax
8
9 from src.MPS import MPS, init_spinup_MPS, init_spinright_MPS
10 from src.TFI import TFIModel
11 from src.TEBD import TEBD_engine
12 from src.DMRG import DMRGEngine
13 from src.ExactDiag import finite_gs_energy
14
15 import matplotlib.pyplot as plt
16 from matplotlib.lines import Line2D
17 plt.rcParams.update({'font.size': 16})
18
19 images_dir = os.path.join(os.path.dirname(os.path.abspath(__file__)), "images")
20 if not os.path.exists(images_dir):
21     os.makedirs(images_dir)
22
23 L = 14
24 J = 1.0
25 g = 1.5
26
27 theoretical_energies = {
28     0.5: -14.01899646121673726, # ground state energy for g=0.5
29     1.0: -17.47100405473177176, # ground state energy for g=1.0
30     1.5: -23.22295943411735664, # ground state energy for g=1.5
31 }
32
33 energies_DMRG = []
34 energies_TEBD = []
35 energies_GD = []
36
37 theoretical_energies = {}
38
39 g_list = [0.5, 1.0, 1.5]
40 for g in g_list:
41     theoretical_energies[g] = finite_gs_energy(L, J, g)
42
43     if g <= J:

```

```

44     init_MPS = init_spinright_MPS
45 else:
46     init_MPS = init_spinup_MPS
47
48 print(f"Transverse-field Ising model with L={L}, J={J}, g={g}")
49
50 model = TFIModel(L, J, g)
51
52 steps = 100
53 psi = init_MPS(L, 30, noise=True, eps=1e-5)
54
55 print("Chi:", psi.get_chi())
56 print(f"<psi|sigma_z|psi> = {psi.site_expectation_value(model.sigmaz).sum():.8f}")
57 print(f"<psi|sigma_x|psi> = {psi.site_expectation_value(model.sigmax).sum():.8f}")
58 print(f"<psi|psi> = {psi.norm_squared():.8f}")
59 print(f"<psi|H|psi> = {model.energy(psi):.8f}")
60 print(f"<psi|H|psi> = {model.energy_mpo(psi):.8f}")
61
62 #####
63 #                               DMRG                               #
64 #####
65
66 print("\nRunning DMRG...")
67
68 energy_DMRG = [model.energy(psi)]
69 dmrg = DMRGEngine(psi.copy(), model, chi_max=30, eps=1e-12)
70
71 for sweep in range(4):
72     dmrg.sweep()
73     energy = model.energy(dmrg.psi)
74     energy_DMRG.append(energy)
75
76 print(f"DMRG ground state energy: {energy_DMRG[-1]:.5f}")
77
78 energy_DMRG = jnp.array(energy_DMRG)
79 energies_DMRG.append(energy_DMRG)
80
81 #####
82 #                               TEBD                               #
83 #####
84
85 print("\nRunning TEBD...")
86
87 psi_TEBD = init_MPS(L, 2, noise=False)
88 energy_TEBD = [model.energy(psi_TEBD)]
89 scheduler = optax.schedules.exponential_decay(1.e-1, 30, 0.1, staircase=False)
90
91 for step in range(steps):
92     tebd = TEBD_engine(psi_TEBD, model, chi_max=30, eps=1e-10, dt=scheduler(step))
93     energy = tebd.run(1, order=2)
94     energy_TEBD += energy
95
96     if step % 5 == 0:
97         print(f"TEBD ground state energy: {energy[-1]:.5f}")
98
99 energy_TEBD = jnp.array(energy_TEBD)
100 energies_TEBD.append(energy_TEBD)
101
102 #####
103 #                               GD                               #
104 #####
105
106 @jit
107 def loss(psi):
108     """Loss function to minimize, which is the energy expectation value."""
109     return model.energy_mpo(psi) / psi.norm_squared()
110
111 @jit
112 def loss_grad(psi):
113     """Gradient of the loss function."""

```

```

114     energy = model.energy_mpo(psi)
115     norm_squared = psi.norm_squared()
116
117     energy_grad = model.energy_mpo_grad(psi)
118     norm_grad = psi.norm_squared_grad()
119
120     Bs_grad = [energy_grad[i]/ norm_squared - energy * norm_grad[i] / norm_squared**2 for i in
121 range(psi.L)]
122
123     return MPS(Bs_grad, psi.Ss)
124
125 scheduler = optax.schedules.exponential_decay(2.e-2, 30, 0.5, staircase=False)
126 optimizer = optax.adam(learning_rate=scheduler)
127 opt_state = optimizer.init(psi)
128
129 energy_GD = [model.energy(psi)]
130
131 print("\nRunning Gradient Descent...")
132
133 t_grad_analytical = 0
134
135 if g == 1.5:
136     grad_error_avg = 0
137     grad_error_std = 0
138     grad_error_max = 0
139
140     t_grad_jax = 0
141
142     # First run for compilation
143     grad_analytical = block_until_ready(loss_grad(psi))
144     grads = block_until_ready(grad(loss)(psi))
145
146 start = timer()
147 for step in range(steps):
148     t_start_analytical = timer()
149     grad_analytical = block_until_ready(loss_grad(psi))
150     t_grad_analytical += timer() - t_start_analytical
151
152     if g == 1.5:
153         t_start_jax = timer()
154         grads = block_until_ready(grad(loss)(psi))
155         t_grad_jax += timer() - t_start_jax
156
157         for i in range(psi.L):
158             errors = grad_analytical.Bs[i] - grads.Bs[i]
159             grad_error_avg += jnp.mean(jnp.abs(errors))
160             grad_error_std += jnp.std(errors)
161             grad_error_max = max(grad_error_max, jnp.max(jnp.abs(errors)))
162
163     # Apply the optimizer to update the MPS
164     updates, opt_state = optimizer.update(grad_analytical, opt_state, psi)
165     psi = optax.apply_updates(psi, updates)
166
167     # Normalize the MPS (does not change the energy)
168     psi = psi.normalize()
169
170     energy = model.energy_mpo(psi)
171     energy_GD.append(energy)
172
173     if step % 5 == 0:
174         print(f"Step {step:>4}, Loss: {energy:>9.5f}, Learning rate: {scheduler(step):>4.2e}")
175
176 if g == 1.5:
177     grad_error_avg /= steps
178     grad_error_std /= steps
179
180     t_grad_analytical /= steps
181     t_grad_jax /= steps
182
183 energy_GD = jnp.array(energy_GD)

```

```

183     energies_GD.append(energy_GD)
184
185     print(f"Optimization completed in {timer() - start:.3f} seconds\n")
186
187     #####
188     #                               Plotting                               #
189     #####
190
191     # Plotting energy decay
192
193     fig, ax1 = plt.subplots(figsize=(10, 6))
194
195     for i, energy_DMRG in enumerate(energies_DMRG):
196         ax1.plot(jnp.arange(5), energy_DMRG, label='DMRG' if i == 0 else "", linewidth=2, color="
limegreen")
197     for i, energy_TEBD in enumerate(energies_TEBD):
198         ax1.plot(jnp.arange(steps + 1), energy_TEBD, label='TEBD' if i == 0 else "", linewidth=2, color
="blue")
199     for i, energy_GD in enumerate(energies_GD):
200         ax1.plot(jnp.arange(steps + 1), energy_GD, label='Gradient Descent' if i == 0 else "",
linewidth=2, color="red")
201
202     ax1.set_xlabel("Steps", fontsize=17)
203     ax1.set_ylabel("Ground Energy", fontsize=17)
204     ax1.set_xlim(0, steps)
205     ax1.set_ylim(-24, -12)
206
207     ax2 = ax1.twinx()
208     ax2.set_ylim(-24, -12)
209
210     # Store ticks and labels
211     yticks = []
212     yticklabels = []
213
214     for i, (g, theoretical_energy) in enumerate(theoretical_energies.items()):
215         ax1.axhline(y=theoretical_energy, color='black', linestyle='--', linewidth=1.5, alpha=0.8,
label = "Exact Diagonalization" if i == 0 else "")
216         yticks.append(theoretical_energy)
217         yticklabels.append(fr'$g={g}$')
218
219     # Set the custom ticks and labels
220     ax2.set_yticks(yticks)
221     ax2.set_yticklabels(yticklabels, fontsize=17)
222
223     ax1.minorticks_on()
224     ax1.grid(which='major', axis='y', alpha=0.3, linestyle='--')
225     ax1.grid(which='both', axis='x', alpha=0.3, linestyle='--')
226     ax1.legend(loc='center right', bbox_to_anchor=(0.98, 0.30), fontsize=15)
227
228     # Remove right spine for cleaner look
229     ax1.spines['right'].set_visible(False)
230
231     # Adjust layout
232     plt.tight_layout()
233
234     # Save with high quality
235     plt.savefig(os.path.join(images_dir, "ground_state_optimization.png"), dpi=300, bbox_inches='tight'
, facecolor='white')
236
237
238     # Plotting relative error
239     plt.figure(figsize=(10, 6))
240
241     markers = ["s", "D", "o"]
242     for i, (g, theoretical_energy) in enumerate(theoretical_energies.items()):
243         relative_error_DMRG = jnp.abs((theoretical_energy - energies_DMRG[i]) / theoretical_energy)
244         relative_error_TEBD = jnp.abs((theoretical_energy - energies_TEBD[i]) / theoretical_energy)
245         relative_error_GD = jnp.abs((theoretical_energy - energies_GD[i]) / theoretical_energy)
246

```



```

247 plt.plot(jnp.arange(5), relative_error_DMRG, marker=markers[i], markevery=10, color="limegreen",
248           linewidth=2)
249 plt.plot(jnp.arange(steps + 1), relative_error_TEBD, marker=markers[i], markevery=10, color="
blue", linewidth=2)
250 plt.plot(jnp.arange(steps + 1), relative_error_GD, marker=markers[i], markevery=10, color="red"
, linewidth=2)
251 plt.xlim(0, steps)
252 plt.ylim(1.e-6, 1.e-0)
253 plt.xlabel("Steps", fontsize=17)
254 plt.ylabel("Relative Error", fontsize=17)
255 plt.yscale('log')
256
257 # Custom legend (just measurements)
258 custom_lines = [
259     Line2D([0], [0], color='limegreen'),
260     Line2D([0], [0], color='blue'),
261     Line2D([0], [0], color='red')
262 ]
263 legend1 = plt.legend(custom_lines, ['DMRG', 'TEBD', 'Gradient Descent'], loc='upper right',
264                      fontsize=15)
265 plt.gca().add_artist(legend1) # Add first legend manually
266
267 # Second legend
268 custom_markers = [
269     Line2D([0], [0], color='black', marker='s', linestyle=''),
270     Line2D([0], [0], color='black', marker='D', linestyle=''),
271     Line2D([0], [0], color='black', marker='o', linestyle='')
272 ]
273 legend2 = plt.legend(custom_markers, [f'g={g_list[0]}', f'g={g_list[1]}', f'g={g_list[2]}'], loc='
upper right', fontsize=15, bbox_to_anchor=(0.65, 1.0))
274 plt.gca().add_artist(legend2) # Add second legend manually
275
276 plt.grid()
277 plt.minorticks_on()
278 plt.grid(which='both', alpha=0.3, linestyle='--')
279 plt.savefig(os.path.join(images_dir, "relative_error_ground_state.png"), dpi=300)
280
281 # Plotting gradient comparison
282 fig, ax1 = plt.subplots(figsize=(10, 6))
283
284 # Define colors
285 colors = {
286     'errors': "#FF7F0E", # Modern orange
287     'analytical': "#0000FF", # Deep blue
288     'jax': "#FF0000", # Deep red
289     'edge': "#000000" # Black for edges
290 }
291
292 # Data preparation
293 error_data = [
294     ("Average", jnp.abs(grad_error_avg)),
295     ("Std Dev", jnp.abs(grad_error_std)),
296     ("Maximum", jnp.abs(grad_error_max)),
297 ]
298
299 error_labels = [item[0] for item in error_data]
300 error_vals = [item[1] for item in error_data]
301
302 # Create main error bars
303 X_axis = jnp.arange(len(error_labels))
304 bar_width = 0.55
305
306 bars1 = ax1.bar(X_axis, error_vals, bar_width,
307                 color=colors['errors'], edgecolor=colors['edge'],
308                 linewidth=1.5, alpha=0.8, label="Gradient Errors")
309
310 # Formatting for left y-axis
311 ax1.set_xticks(X_axis)
312 ax1.set_xticklabels(error_labels, fontsize=15)

```

```

312 ax1.set_ylabel("Absolute Error", fontsize=17)
313 ax1.set_yscale("log")
314 ax1.set_ylim(1e-14, 1e-12)
315 ax1.grid(True, which='both', alpha=0.3, linestyle='--')
316 ax1.tick_params(axis='y', labelsize=15)
317
318 # Create second y-axis for runtime
319 ax2 = ax1.twinx()
320 ax2.set_ylabel('Computation Time (s)', fontsize=17)
321 ax2.set_yscale('log')
322 ax2.set_ylim(1e-4, 1e-2)
323 ax2.tick_params(axis='y', labelsize=15)
324
325 # Add runtime bars
326 runtime_x = len(error_labels) + 0.5
327 bar_width_runtime = 0.45
328
329 # Add space for runtime bars
330 ax1.set_xlim(-0.5, runtime_x + 0.8)
331 ax2.set_xlim(-0.5, runtime_x + 0.8)
332
333 bars2 = ax2.bar(runtime_x - bar_width_runtime/2, t_grad_analytical, bar_width_runtime,
334                label="Analytical Gradient", color=colors['analytical'],
335                edgecolor=colors['edge'], linewidth=1.2, alpha=0.8)
336
337 bars3 = ax2.bar(runtime_x + bar_width_runtime/2, t_grad_jax, bar_width_runtime,
338                label="JAX Gradient", color=colors['jax'],
339                edgecolor=colors['edge'], linewidth=1.2, alpha=0.8)
340
341 # Update x-axis labels to include runtime
342 all_labels = error_labels + ["Runtime"]
343 ax1.set_xticks(list(X_axis) + [runtime_x])
344 ax1.set_xticklabels(all_labels, fontsize=15)
345
346 # Create a unified legend in the original position
347 handles1, labels1 = ax1.get_legend_handles_labels()
348 handles2, labels2 = ax2.get_legend_handles_labels()
349 ax1.legend(handles1 + handles2, labels1 + labels2, loc='upper left', fontsize=15)
350
351 # Remove top and right spines for cleaner look
352 ax1.spines['top'].set_visible(False)
353 ax2.spines['top'].set_visible(False)
354 ax1.spines['right'].set_visible(False)
355
356 # Adjust layout
357 plt.tight_layout()
358
359 # Save with high quality
360 plt.savefig(os.path.join(images_dir, "gradient_comparison.png"), dpi=300, bbox_inches='tight',
361            facecolor='white')
362
363 plt.show()

```

2. src/MPS.py

```

1  """Toy code implementing a matrix product state."""
2
3  import jax
4  jax.config.update("jax_enable_x64", True)
5  import jax.numpy as jnp
6  from jax import tree_util, jit
7  from functools import partial
8
9  class MPS:
10     """Class for a matrix product state.
11
12     We index sites with 'i' from 0 to L-1; bond 'i' is left of site 'i'.

```

```

13 We *assume* that the state is in right-canonical form.
14
15 Parameters
16 -----
17 Bs, Ss:
18     Same as attributes.
19
20 Attributes
21 -----
22 Bs : list of jnp.Array[ndim=3]
23     The 'matrices' in right-canonical form, one for each physical site.
24     Each 'B[i]' has legs (virtual left, physical, virtual right), in short 'vL i vR'
25 Ss : list of jnp.Array[ndim=1]
26     The Schmidt values at each of the bonds, 'Ss[i]' is left of 'Bs[i]'.
27 L : int
28     Number of sites.
29 """
30
31 def __init__(self, Bs, Ss):
32     self.Bs = Bs
33     self.Ss = Ss #stop_gradient(s) for s in Ss]
34     self.L = len(Bs)
35
36 @jit
37 def copy(self):
38     # return tree_util.tree_map(jnp.array, self)
39     return MPS([jnp.array(B) for B in self.Bs], [jnp.array(S) for S in self.Ss])
40
41 @jit
42 def norm_squared(self):
43     """Calculate the norm squared of the MPS, which is the overlap with itself."""
44     return overlap(self, self).real
45
46 @jit
47 def norm_squared_grad(self):
48     """
49     Compute the gradient of the norm squared with respect to the MPS tensors.
50     Args:
51         psi: MPS object
52     Returns:
53         Gradient of the norm squared with respect to the MPS tensors.
54     """
55     grad = []
56
57     contr_left = jnp.ones((1, 1), dtype=self.Bs[0].dtype)
58     contr_right = jnp.ones((1, 1), dtype=self.Bs[0].dtype)
59
60     left_blocks = [contr_left]
61     right_blocks = [contr_right]
62
63     for n in range(self.L - 1):
64         M_bra = self.Bs[n].conj() # vL* i* vR*
65         M_ket = self.Bs[n]         # vL i vR
66
67         contr_left = jnp.tensordot(contr_left, M_bra, [0, 0]) # [vR*] vR, [vL*]
68         i* vR*
69         contr_left = jnp.tensordot(contr_left, M_ket, axes=([0, 1], [0, 1])) # [vR] [i*] vR*,
70         [vL] [i] vR
71
72         left_blocks.append(contr_left)
73
74     for n in reversed(range(1, self.L)):
75         M_bra = self.Bs[n].conj() # vL* i* vR*
76         M_ket = self.Bs[n]         # vL i vR
77
78         contr_right = jnp.tensordot(M_bra, contr_right, [2, 0]) # vL* i* [vR*],
79         [vL*] vL
80         contr_right = jnp.tensordot(contr_right, M_ket, axes=([1, 2], [1, 2])) # vL* [i*] [vL
81         ], vL [i] [vR]
82

```

```

79         right_blocks.append(contr_right)
80
81     for n in range(self.L):
82         M_ket = self.Bs[n] # vL i vR
83         contr_left = left_blocks[n] # vR* vR
84         contr_right = right_blocks[self.L - n - 1] # vL* vL
85
86         grad_n = jnp.tensordot(contr_left, M_ket, [1, 0]) # vR* [vR], [vL] i vR
87         grad_n = jnp.tensordot(grad_n, contr_right, [2, 1]) # vR* i [vR], vL* [vL]
88
89         grad.append(2*grad_n)
90
91     assert all(grad[i].shape == self.Bs[i].shape for i in range(self.L))
92     return grad
93
94 @jit
95 def normalize(self):
96     center = self.L // 2
97     psi = self.copy()
98     psi.Bs[center] = psi.Bs[center] / jnp.sqrt(psi.norm_squared())
99     return psi
100
101 @jit
102 def canonicalize(self):
103     """
104     Return a new MPS in right-canonical form by a left-to-right SVD sweep.
105     Optionally truncate to chi_max and discard singular values < eps.
106     Does not modify self.
107     """
108     psi = self.copy() # Create a copy to avoid modifying self
109     chis = psi.get_chi()
110     for i in range(psi.L - 1):
111         chivC = chis[i]
112         j = i + 1
113         theta = psi.get_theta2(i) # vL i j vR
114
115         Ai, Sj, Bj = split_theta(theta, chivC)
116
117         # put back into MPS
118         Gi = jnp.tensordot(jnp.diag(psi.Ss[i]**(-1)), Ai, axes=[1, 0]) # vL [vL*], [vL] i vC
119         psi.Bs[i] = jnp.tensordot(Gi, jnp.diag(Sj), axes=[2, 0]) # vL i [vC], [vC] vC
120         psi.Ss[j] = Sj #jax.lax.stop_gradient(Sj) # vC
121         psi.Bs[j] = Bj # vC j vR
122     return psi
123
124 @partial(jit, static_argnames=["i"])
125 def get_theta1(self, i):
126     """Calculate effective single-site wave function on sites i in mixed canonical form.
127
128     The returned array has legs 'vL, i, vR' (as one of the Bs)."""
129     return jnp.tensordot(jnp.diag(self.Ss[i]), self.Bs[i], [1, 0]) # vL [vL'], [vL] i vR
130
131 @partial(jit, static_argnames=["i"])
132 def get_theta2(self, i):
133     """Calculate effective two-site wave function on sites i,j=(i+1) in mixed canonical form.
134
135     The returned array has legs 'vL, i, j, vR'."""
136     j = i + 1
137     return jnp.tensordot(self.get_theta1(i), self.Bs[j], [2, 0]) # vL i [vR], [vL] j vR
138
139 def get_chi(self):
140     """Return bond dimensions."""
141     return [self.Bs[i].shape[2] for i in range(self.L - 1)]
142
143 @jit
144 def site_expectation_value(self, op):
145     """Calculate expectation values of a local operator at each site."""
146     result = []
147     for i in range(self.L):
148         theta = self.get_theta1(i) # vL i vR

```

```

149         op_theta = jnp.tensordot(op, theta, axes=[1, 1]) # i [i*], vL [i] vR
150         result.append(jnp.tensordot(theta.conj(), op_theta, [[0, 1, 2], [1, 0, 2]]))
151         # [vL*] [i*] [vR*], [i] [vL] [vR]
152     return jnp.real(jnp.array(result))
153
154 @jit
155 def bond_expectation_value(self, op):
156     """Calculate expectation values of a local operator at each bond."""
157     result = []
158     for i in range(self.L - 1):
159         theta = self.get_theta2(i) # vL i j vR
160         op_theta = jnp.tensordot(op[i], theta, axes=[[2, 3], [1, 2]])
161         # i j [i*] [j*], vL [i] [j] vR
162         result.append(jnp.tensordot(theta.conj(), op_theta, [[0, 1, 2, 3], [2, 0, 1, 3]]))
163         # [vL*] [i*] [j*] [vR*], [i] [j] [vL] [vR]
164     return jnp.real(jnp.array(result))
165
166 @jit
167 def entanglement_entropy(self):
168     """Return the (von-Neumann) entanglement entropy for a bipartition at any of the bonds."""
169     result = []
170     for i in range(1, self.L):
171         S = jnp.array(self.Ss[i])
172         S = S[S > 1e-30] # 0*log(0) should give 0; avoid warnings or NaN by discarding small S
173         S2 = S * S
174         assert abs(jnp.linalg.norm(S) - 1.) < 1.e-14
175         result.append(-jnp.sum(S2 * jnp.log(S2)))
176     return jnp.array(result)
177
178 def _tree_flatten(self):
179     children = (self.Bs, self.Ss) # arrays / dynamic values
180     aux_data = {} # static values
181     return (children, aux_data)
182
183 @classmethod
184 def _tree_unflatten(cls, aux_data, children):
185     Bs, Ss = children
186     return cls(Bs, Ss)
187
188 tree_util.register_pytree_node(MPS,
189                                MPS._tree_flatten,
190                                MPS._tree_unflatten)
191
192 @jit
193 def overlap(mps_bra, mps_ket):
194     """
195     Compute the overlap <mps_bra|mps_ket> for two MPS in right-canonical form.
196     Both should be lists of tensors of the same length.
197     """
198     L = len(mps_bra.Bs)
199     contr = jnp.ones((1, 1), dtype=mps_bra.Bs[0].dtype)
200     for n in range(L):
201         M_bra = mps_bra.Bs[n].conj() # vL* i* vR*
202         M_ket = mps_ket.Bs[n] # vL i vR
203
204         contr = jnp.tensordot(contr, M_ket, axes=(1, 0)) # vR* [vR], [vL] j vR
205         contr = jnp.tensordot(M_bra, contr, axes=[[0, 1], [0, 1]]) # [vL*] [j*] vR*, [vR*] [j] vR
206     assert contr.shape == (1, 1)
207     return contr[0, 0]
208
209 def init_spinup_MPS(L: int, chi_max: int, noise: bool = False, eps: float = 1e-4, key=None) -> MPS:
210     """
211     Create an all-up spin MPS with maximum bond dimension chi_max.
212     Optionally add small noise to each tensor.
213
214     Args:
215         L: int, number of sites
216         chi_max: int, maximum bond dimension (at the center)
217         key: jax.random.PRNGKey or None
218         noise: bool, whether to add noise

```

```

219     eps: float, noise amplitude (if noise=True)
220 Returns:
221     mps: list of jnp.ndarray, each of shape (left, 2, right)
222 """
223 # Compute the staircase bond dimensions
224 chi = [min(2**min(i, L-i), chi_max) for i in range(L+1)]
225 shapes = [(chi[i], 2, chi[i+1]) for i in range(L)]
226 Bs = []
227 for left, phys, right in shapes:
228     tensor = jnp.zeros((left, phys, right), dtype=jnp.float64)
229     tensor = tensor.at[0, 0, 0].set(1.0)
230     if noise:
231         import jax.random as jr
232         if key is None:
233             key = jr.PRNGKey(42)
234             subkey, key = jr.split(key)
235             tensor += eps * jr.normal(subkey, shape=(left, phys, right), dtype=jnp.float64)
236     Bs.append(tensor)
237 Ss = [jnp.pad(jnp.ones([1], jnp.float64), (0, chi[i]-1)) for i in range(L)]
238 mps = MPS(Bs, Ss)
239 mps = mps.canonicalize() # Canonicalize to ensure the noise is properly incorporated
240 # mps = mps.normalize() # Not needed, as canonicalization already normalizes the MPS
241 return mps
242
243 def init_spinright_MPS(L: int, chi_max: int, noise: bool = False, eps: float = 1e-4, key=None) ->
MPS:
244 """
245 Create an all-right spin MPS with maximum bond dimension chi_max.
246 Optionally add small noise to each tensor.
247
248 Args:
249     L: int, number of sites
250     chi_max: int, maximum bond dimension (at the center)
251     key: jax.random.PRNGKey or None
252     noise: bool, whether to add noise
253     eps: float, noise amplitude (if noise=True)
254 Returns:
255     mps: list of jnp.ndarray, each of shape (left, 2, right)
256 """
257 # Compute the staircase bond dimensions
258 chi = [min(2**min(i, L-i), chi_max) for i in range(L+1)]
259 shapes = [(chi[i], 2, chi[i+1]) for i in range(L)]
260 Bs = []
261 for left, phys, right in shapes:
262     tensor = jnp.zeros((left, phys, right), dtype=jnp.float64)
263     tensor = tensor.at[0, 0, 0].set(0.5**0.5)
264     tensor = tensor.at[0, 1, 0].set(0.5**0.5)
265     if noise:
266         import jax.random as jr
267         if key is None:
268             key = jr.PRNGKey(42)
269             subkey, key = jr.split(key)
270             tensor += eps * jr.normal(subkey, shape=(left, phys, right), dtype=jnp.float64)
271     Bs.append(tensor)
272 Ss = [jnp.pad(jnp.ones([1], jnp.float64), (0, chi[i]-1)) for i in range(L)]
273 mps = MPS(Bs, Ss)
274 mps = mps.canonicalize() # Canonicalize to ensure the noise is properly incorporated
275 # mps = mps.normalize() # Not needed, as canonicalization already normalizes the MPS
276 return mps
277
278 @partial(jit, static_argnames=["chivC"])
279 def split_theta(theta, chivC):
280     """Split a two-site wave function in mixed canonical form.
281
282     Split a two-site wave function as follows::
283         vL --(theta)-- vR      =>    vL --(A)--diag(S)--(B)-- vR
284             |   |                |           |
285             i   j                i           j
286
287 Parameters

```

```

288 -----
289 theta : jnp.Array[ndim=4]
290     Two-site wave function in mixed canonical form, with legs 'vL, i, j, vR'.
291 chivC : int
292     Maximum number of singular values to keep
293
294 Returns
295 -----
296 A : jnp.Array[ndim=3]
297     Left-canonical matrix on site i, with legs 'vL, i, vC'
298 S : jnp.Array[ndim=1]
299     Singular/Schmidt values.
300 B : jnp.Array[ndim=3]
301     Right-canonical matrix on site j, with legs 'vC, j, vR'
302 """
303 chivL, dL, dR, chivR = theta.shape
304 theta = jnp.reshape(theta, [chivL * dL, dR * chivR])
305
306 X, Y, Z = jnp.linalg.svd(theta, full_matrices=False) # returns Y sorted in descending order
307
308 # truncate
309 X, Y, Z = X[:, :chivC], Y[:chivC], Z[:chivC, :]
310
311 Y = jnp.maximum(Y, 1e-12) # avoid division by zero
312
313 # renormalize
314 S = Y / jnp.linalg.norm(Y) # == Y/sqrt(sum(Y**2))
315
316 # split legs of X and Z
317 A = jnp.reshape(X, [chivL, dL, chivC])
318 B = jnp.reshape(Z, [chivC, dR, chivR])
319 return A, S, B
320
321 def split_truncate_theta(theta, chi_max, eps):
322     """Split and truncate a two-site wave function in mixed canonical form.
323
324     Split a two-site wave function as follows::
325         vL --(theta)-- vR      =>   vL --(A)--diag(S)--(B)-- vR
326             |   |                |           |
327             i   j                i           j
328
329     Afterwards, truncate in the new leg (labeled 'vC').
330
331 Parameters
332 -----
333 theta : jnp.Array[ndim=4]
334     Two-site wave function in mixed canonical form, with legs 'vL, i, j, vR'.
335 chi_max : int
336     Maximum number of singular values to keep
337 eps : float
338     Discard any singular values smaller than that.
339
340 Returns
341 -----
342 A : jnp.Array[ndim=3]
343     Left-canonical matrix on site i, with legs 'vL, i, vC'
344 S : jnp.Array[ndim=1]
345     Singular/Schmidt values.
346 B : jnp.Array[ndim=3]
347     Right-canonical matrix on site j, with legs 'vC, j, vR'
348 """
349 chivL, dL, dR, chivR = theta.shape
350 theta = jnp.reshape(theta, [chivL * dL, dR * chivR])
351
352 X, Y, Z = jnp.linalg.svd(theta, full_matrices=False) # returns Y sorted in descending order
353
354 # truncate
355 chivC = min(chi_max, jnp.sum(Y > eps))
356 X, Y, Z = X[:, :chivC], Y[:chivC], Z[:chivC, :]
357

```

```

358 # renormalize
359 S = Y / jnp.linalg.norm(Y) # == Y/sqrt(sum(Y**2))
360
361 # split legs of X and Z
362 A = jnp.reshape(X, [chivL, dL, chivC])
363 B = jnp.reshape(Z, [chivC, dR, chivR])
364 return A, S, B

```

3. src/TFI.py

```

1  """Toy code implementing the transverse-field ising model."""
2
3  import jax
4  jax.config.update("jax_enable_x64", True)
5  import jax.numpy as jnp
6  from jax import tree_util, jit
7
8  class TFIModel:
9      """Class generating the Hamiltonian of the transverse-field Ising model.
10
11      The Hamiltonian reads
12      .. math ::
13          H = - J \sum_i \sigma^x_i \sigma^x_{i+1} - g \sum_i \sigma^z_i
14
15      Parameters
16      -----
17      L : int
18          Number of sites.
19      J, g : float
20          Coupling parameters of the above defined Hamiltonian.
21
22      Attributes
23      -----
24      L : int
25          Number of sites.
26      d : int
27          Local dimension (=2 for spin-1/2 of the transverse field ising model)
28      sigmax, sigmay, sigmaz, id :
29          Local operators, namely the Pauli matrices and identity.
30      H_bonds : list of jnp.Array[ndim=4]
31          The Hamiltonian written in terms of local 2-site operators, 'H = sum_i H_bonds[i]'.
32          Each 'H_bonds[i]' has (physical) legs (i out, (i+1) out, i in, (i+1) in),
33          in short 'i j i* j*'.
34      """
35
36      def __init__(self, L, J, g):
37          self.L, self.d = L, 2
38          self.J, self.g = J, g
39          self.sigmax = jnp.array([[0., 1.], [1., 0.]])
40          self.sigmay = jnp.array([[0., -1j], [1j, 0.]])
41          self.sigmaz = jnp.array([[1., 0.], [0., -1.]])
42          self.id = jnp.eye(2)
43          self._H_bonds = None
44          self._H_mpo = None
45
46      def _init_H_bonds(self):
47          """Initialize 'H_bonds' hamiltonian. Called by H_bonds."""
48          sx, sz, id = self.sigmax, self.sigmaz, self.id
49          d = self.d
50          H_list = []
51          for i in range(self.L - 1):
52              gL = gR = 0.5 * self.g
53              if i == 0: # first bond
54                  gL = self.g
55              if i + 1 == self.L - 1: # last bond
56                  gR = self.g
57              H_bond = -self.J * jnp.kron(sx, sx) - gL * jnp.kron(sz, id) - gR * jnp.kron(id, sz)

```



```

58         # H_bond has legs 'i, j, i*, j*'
59         H_list.append(jnp.reshape(H_bond, [d, d, d, d]))
60     self._H_bonds = H_list
61
62     def _init_H_mpo(self):
63         """Initialize the MPO representation of the Hamiltonian."""
64         W = jnp.zeros((3, 3, self.d, self.d))
65
66         W = W.at[0, 0].set(self.id)
67         W = W.at[0, 1].set(self.sigmax)
68         W = W.at[0, 2].set(-self.g * self.sigmaz)
69         W = W.at[1, 2].set(-self.J * self.sigmax)
70         W = W.at[2, 2].set(self.id)
71
72         self._H_mpo = [W.copy() for _ in range(self.L)]
73
74     @property
75     def H_bonds(self):
76         """Return the Hamiltonian bonds."""
77         if self._H_bonds is None:
78             self._init_H_bonds()
79         return self._H_bonds
80
81     @property
82     def H_mpo(self):
83         """Return the MPO representation of the Hamiltonian."""
84         if self._H_mpo is None:
85             self._init_H_mpo()
86         return self._H_mpo
87
88     @jit
89     def energy(self, psi):
90         """Evaluate energy  $E = \langle \psi | H | \psi \rangle$  for the given MPS."""
91         assert psi.L == self.L
92         return jnp.sum(psi.bond_expectation_value(self.H_bonds))
93
94     @jit
95     def energy_mpo(self, psi):
96         """
97         Compute the expectation value  $\langle \text{mps\_bra} | \text{MPO} | \text{mps\_ket} \rangle$  for two MPS and an MPO.
98         All should be lists of tensors of the same length.
99         Args:
100             mps_bra: MPS object (bra, conjugated)
101             mps_ket: MPS object (ket)
102             mpo: list of MPO tensors (one per site)
103         Returns:
104             Scalar energy expectation value
105         """
106         assert psi.L == self.L
107         left_vec = jnp.array([1, 0, 0], dtype=psi.Bs[0].dtype)
108         contr = left_vec.reshape(1, 3, 1)
109
110         for n in range(self.L):
111             M_bra = psi.Bs[n].conj() # vL* i* vR*
112             M_ket = psi.Bs[n]        # vL i vR
113             W = self.H_mpo[n]        # wL wR i* i
114
115             contr = jnp.tensordot(contr, M_bra, [0, 0]) # [vR*] wR vR, [vL*] i* vR*
116             contr = jnp.tensordot(contr, W, axes=([0, 2], [0, 2])) # [wR] vR [i*] vR*, [wL] wR
117             # [i*] i
118             contr = jnp.tensordot(contr, M_ket, axes=([0, 3], [0, 1])) # [vR] vR* wR [i], [vL] [i]
119             # vR
120
121         assert contr.shape == (1, 3, 1)
122         return contr[0, 2, 0] # right_vec = jnp.array([0, 0, 1], dtype=psi.Bs[0].dtype)
123
124     @jit
125     def energy_mpo_grad(self, psi):
126         """
127         Compute the gradient of the energy expectation value with respect to the MPS tensors.

```

```

126     Args:
127         psi: MPS object
128     Returns:
129         Gradient of the energy expectation value with respect to the MPS tensors.
130     """
131     assert psi.L == self.L
132
133     grad = []
134
135     left_vec = jnp.array([1, 0, 0], dtype=psi.Bs[0].dtype)
136     right_vec = jnp.array([0, 0, 1], dtype=psi.Bs[0].dtype)
137     contr_left = left_vec.reshape(1, 3, 1)
138     contr_right = right_vec.reshape(1, 3, 1)
139
140     left_blocks = [contr_left]
141     right_blocks = [contr_right]
142
143     for n in range(self.L - 1):
144         M_bra = psi.Bs[n].conj() # vL* i* vR*
145         M_ket = psi.Bs[n]         # vL i vR
146         W = self.H_mpo[n]         # wL wR i* i
147
148         contr_left = jnp.tensordot(contr_left, M_bra, [0, 0]) # [vR*] wR vR, [
vL*] i* vR*
149         contr_left = jnp.tensordot(contr_left, W, axes=([0, 2], [0, 2])) # [wR] vR [i*] vR
*, [wL] wR [i*] i
150         contr_left = jnp.tensordot(contr_left, M_ket, axes=([0, 3], [0, 1])) # [vR] vR* wR [i
], [vL] [i] vR
151
152         left_blocks.append(contr_left)
153
154     for n in reversed(range(1, self.L)):
155         M_bra = psi.Bs[n].conj() # vL* i* vR*
156         M_ket = psi.Bs[n]         # vL i vR
157         W = self.H_mpo[n]         # wL wR i* i
158
159         contr_right = jnp.tensordot(M_bra, contr_right, [2, 0]) # vL* i* [vR*],
[vL*] wL vL
160         contr_right = jnp.tensordot(contr_right, W, axes=([1, 2], [2, 1])) # vL* [i*] [wL]
vL, wL [wR] [i*] i
161         contr_right = jnp.tensordot(contr_right, M_ket, axes=([1, 3], [2, 1])) # vL* [vL] wL [
i], vL [i] [vR]
162
163         right_blocks.append(contr_right)
164
165     for n in range(self.L):
166         M_ket = psi.Bs[n] # vL i vR
167         W = self.H_mpo[n] # wL wR i* i
168         contr_left = left_blocks[n] # vR* wR vR
169         contr_right = right_blocks[self.L - n - 1] # vL* wL vL
170
171         grad_n = jnp.tensordot(contr_left, M_ket, [2, 0]) # vR* wR [vR], [vL]
i vR
172         grad_n = jnp.tensordot(grad_n, W, axes=([1, 2], [0, 3])) # vR* [wR] [i] vR, [
wL] wR i* [i]
173         grad_n = jnp.tensordot(grad_n, contr_right, axes=([1, 2], [2, 1])) # vR* [vR] [wR] i*,
vL* [wL] [vL]
174
175         grad.append(2*grad_n)
176
177     assert all(grad[i].shape == psi.Bs[i].shape for i in range(self.L))
178     return grad
179
180     def _tree_flatten(self):
181         children = (self._H_bonds,) # arrays / dynamic values
182         aux_data = {
183             "L": self.L,
184             "J": self.J,
185             "g": self.g,
186         } # static values

```

```

187         return (children, aux_data)
188
189     @classmethod
190     def _tree_unflatten(cls, aux_data, children):
191         obj = cls(aux_data["L"], aux_data["J"], aux_data["g"])
192         obj._H_bonds = children
193         return obj
194
195 tree_util.register_pytree_node(TFIModel,
196                                TFIModel._tree_flatten,
197                                TFIModel._tree_unflatten)

```

4. src/ExactDiag.py

```

1  """Provides exact ground state energies for the transverse field ising model for comparison.
2
3  The Hamiltonian reads
4  .. math ::
5      H = - J \sum_{i} \sigma^x_i \sigma^x_{i+1} - g \sum_{i} \sigma^z_i
6  """
7  import numpy as np
8  import scipy.sparse as sparse
9  import warnings
10
11 def finite_gs_energy(L, J, g):
12     """For comparison: obtain ground state energy from exact diagonalization.
13
14     Exponentially expensive in L, only works for small enough 'L' <~ 20.
15     """
16     if L >= 20:
17         warnings.warn("Large L: Exact diagonalization might take a long time!")
18     # get single site operators
19     sx = sparse.csr_matrix(np.array([[0., 1.], [1., 0.]])
20     sz = sparse.csr_matrix(np.array([[1., 0.], [0., -1.]])
21     id = sparse.csr_matrix(np.eye(2))
22     sx_list = [] # sx_list[i] = kron([id, id, ..., id, sx, id, .... id])
23     sz_list = []
24     for i_site in range(L):
25         x_ops = [id] * L
26         z_ops = [id] * L
27         x_ops[i_site] = sx
28         z_ops[i_site] = sz
29         X = x_ops[0]
30         Z = z_ops[0]
31         for j in range(1, L):
32             X = sparse.kron(X, x_ops[j], 'csr')
33             Z = sparse.kron(Z, z_ops[j], 'csr')
34         sx_list.append(X)
35         sz_list.append(Z)
36     H_xx = sparse.csr_matrix((2*L, 2*L))
37     H_z = sparse.csr_matrix((2*L, 2*L))
38     for i in range(L - 1):
39         H_xx = H_xx + sx_list[i] * sx_list[(i + 1) % L]
40     for i in range(L):
41         H_z = H_z + sz_list[i]
42     H = -J * H_xx - g * H_z
43     E, V = sparse.linalg.eigsh(H, k=1, which='SA', return_eigenvectors=True)
44     return E[0]

```

5. src/TEBD.py

```

1  """Toy code implementing the time evolving block decimation (TEBD)."""
2
3  import numpy as np

```

```

4 from scipy.linalg import expm
5 from src.MPS import split_truncate_theta
6
7 class TEBD_engine:
8     def __init__(self, psi, model, chi_max, eps, dt):
9         self.psi = psi
10        self.model = model
11        self.chi_max = chi_max
12        self.eps = eps
13        self.dt = dt
14        self._U_bonds_dt = None
15        self._U_bonds_half_dt = None
16
17    def _init_U_bonds(self, dt):
18        """Given a model, calculate 'U_bonds[i] = expm(-dt*model.H_bonds[i])'.'.
19
20        Each local operator has legs (i out, (i+1) out, i in, (i+1) in), in short 'i j i* j*'.
21        Note that no imaginary 'i' is included, thus real 'dt' means imaginary time evolution!
22        """
23        H_bonds = self.model.H_bonds
24        d = H_bonds[0].shape[0]
25        U_bonds = []
26        for H in H_bonds:
27            H = np.reshape(H, [d * d, d * d])
28            U = expm(-dt * H)
29            U_bonds.append(np.reshape(U, [d, d, d, d]))
30        return U_bonds
31
32    @property
33    def U_bonds_dt(self):
34        """Return the U_bonds for the full time step."""
35        if self._U_bonds_dt is None:
36            self._U_bonds_dt = self._init_U_bonds(self.dt)
37        return self._U_bonds_dt
38
39    @property
40    def U_bonds_half_dt(self):
41        """Return the U_bonds for the half time step."""
42        if self._U_bonds_half_dt is None:
43            self._U_bonds_half_dt = self._init_U_bonds(self.dt / 2)
44        return self._U_bonds_half_dt
45
46    def update_bond(self, i, bond="dt"):
47        """Apply 'U_bond' acting on i,j=(i+1) to 'psi'."""
48        if bond == "dt":
49            U_bond = self.U_bonds_dt[i]
50        elif bond == "half_dt" or bond == "dt/2":
51            U_bond = self.U_bonds_half_dt[i]
52        else:
53            raise ValueError("bond must be 'dt', 'half_dt' or 'dt/2'")
54        j = i + 1
55        # construct theta matrix
56        theta = self.psi.get_theta2(i) # vL i j vR
57        # apply U
58        Utheta = np.tensordot(U_bond, theta, axes=([2, 3], [1, 2])) # i j [i*] [j*], vL [i] [j] vR
59        Utheta = np.transpose(Utheta, [2, 0, 1, 3]) # vL i j vR
60        # split and truncate
61        Ai, Sj, Bj = split_truncate_theta(Utheta, self.chi_max, self.eps)
62        # put back into MPS
63        Gi = np.tensordot(np.diag(self.psi.Ss[i]**(-1)), Ai, axes=[1, 0]) # vL [vL*], [vL] i vC
64        self.psi.Bs[i] = np.tensordot(Gi, np.diag(Sj), axes=[2, 0]) # vL i [vC], [vC] vC
65        self.psi.Ss[j] = Sj # vC
66        self.psi.Bs[j] = Bj # vC j vR
67
68    def run(self, steps, order=2):
69        """Evolve the state 'psi' for 'N_steps' time steps with TEBD.
70        The state psi is modified in place."""
71        energy_TEBD = []
72        Nbonds = self.psi.L - 1
73

```

```

74     if order == 1:
75         for n in range(steps):
76             for k in [0, 1]: # even, odd
77                 for i_bond in range(k, Nbonds, 2):
78                     self.update_bond(i_bond, "dt")
79                     energy_TEBD.append(self.model.energy(self.psi))
80             return energy_TEBD
81
82     elif order == 2:
83         for i_bond in range(0, Nbonds, 2): # even bonds
84             self.update_bond(i_bond, "half_dt")
85         for i_bond in range(1, Nbonds, 2): # odd bonds
86             self.update_bond(i_bond, "dt")
87         energy_TEBD.append(self.model.energy(self.psi))
88         for n in range(steps - 1):
89             for k in [0, 1]: # even and odd
90                 for i_bond in range(k, Nbonds, 2):
91                     self.update_bond(i_bond, "dt")
92                     energy_TEBD.append(self.model.energy(self.psi))
93         for i_bond in range(0, Nbonds, 2): # even bonds
94             self.update_bond(i_bond, "half_dt")
95         energy_TEBD.pop()
96         energy_TEBD.append(self.model.energy(self.psi))
97         return energy_TEBD
98
99     else:
100         raise ValueError("order must be 1 or 2")

```

6. src/DMRG.py

```

1  """Toy code implementing the density-matrix renormalization group (DMRG)."""
2
3  import numpy as np
4  from src.MPS import split_truncate_theta
5  import scipy.sparse
6  import scipy.sparse.linalg._eigen.arpack as arp
7
8
9  class HEffective(scipy.sparse.linalg.LinearOperator):
10     """Class for the effective Hamiltonian.
11
12     To be diagonalized in 'DMRGEngine.update_bond'. Looks like this::
13
14         .--vL*          vR*--.
15         |          i*    j*    |
16         |          |      |      |
17         (LP)---(W1)---(W2)----(RP)
18         |          |      |      |
19         |          i      j      |
20         .--vL          vR--.
21     """
22
23     def __init__(self, LP, RP, W1, W2):
24         self.LP = LP # vL wL* vL*
25         self.RP = RP # vR* wR* vR
26         self.W1 = W1 # wL wC i i*
27         self.W2 = W2 # wC wR j j*
28         chi1, chi2 = LP.shape[0], RP.shape[2]
29         d1, d2 = W1.shape[2], W2.shape[2]
30         self.theta_shape = (chi1, d1, d2, chi2) # vL i j vR
31         self.shape = (chi1 * d1 * d2 * chi2, chi1 * d1 * d2 * chi2)
32         self.dtype = W1.dtype
33
34     def _matvec(self, theta):
35         """calculate |theta> = H_eff |theta>"""
36         x = np.reshape(theta, self.theta_shape) # vL i j vR
37         x = np.tensordot(self.LP, x, axes=(2, 0)) # vL wL* [vL*], [vL] i j vR

```

```

38     x = np.tensordot(x, self.W1, axes=([1, 2], [0, 3])) # vL [wL*] [i] j vR, [wL] wC i [i*]
39     x = np.tensordot(x, self.W2, axes=([3, 1], [0, 3])) # vL [j] vR [wC] i, [wC] wR j [j*]
40     x = np.tensordot(x, self.RP, axes=([1, 3], [0, 1])) # vL [vR] i [wR] j, [vR*] [wR*] vR
41     x = np.reshape(x, self.shape[0])
42     return x
43
44
45 class DMRGEngine(object):
46     """DMRG algorithm, implemented as class holding the necessary data.
47
48     Parameters
49     -----
50     psi, model, chi_max, eps:
51         See attributes
52
53     Attributes
54     -----
55     psi : MPS
56         The current ground-state (approximation).
57     model :
58         The model of which the groundstate is to be calculated.
59     chi_max, eps:
60         Truncation parameters, see :func:`a_mps.split_truncate_theta`.
61     LPs, RPs : list of np.Array[ndim=3]
62         Left and right parts ("environments") of the effective Hamiltonian.
63         'LPs[i]' is the contraction of all parts left of site 'i' in the network '<psi|H|psi>',
64         and similar 'RPs[i]' for all parts right of site 'i'.
65         Each 'LPs[i]' has legs 'vL wL* vL*', 'RPs[i]' has legs 'vR* wR* vR'
66     """
67
68     def __init__(self, psi, model, chi_max=100, eps=1.e-12):
69         assert psi.L == model.L # ensure compatibility
70         self.H_mpo = model.H_mpo
71         self.psi = psi
72         self.LPs = [None] * psi.L
73         self.RPs = [None] * psi.L
74         self.chi_max = chi_max
75         self.eps = eps
76         # initialize left and right environment
77         D = self.H_mpo[0].shape[0]
78         chi = psi.Bs[0].shape[0]
79         LP = np.zeros([chi, D, chi], dtype="float") # vL wL* vL*
80         RP = np.zeros([chi, D, chi], dtype="float") # vR* wR* vR
81         LP[:, 0, :] = np.eye(chi)
82         RP[:, D - 1, :] = np.eye(chi)
83         self.LPs[0] = LP
84         self.RPs[-1] = RP
85         # initialize necessary RPs
86         for i in range(psi.L - 1, 1, -1):
87             self.update_RP(i)
88
89     def sweep(self):
90         # sweep from left to right
91         for i in range(self.psi.L - 2):
92             self.update_bond(i)
93         # sweep from right to left
94         for i in range(self.psi.L - 2, 0, -1):
95             self.update_bond(i)
96
97     def update_bond(self, i):
98         j = i + 1
99         # get effective Hamiltonian
100         Heff = HEffective(self.LPs[i], self.RPs[j], self.H_mpo[i], self.H_mpo[j])
101         # Diagonalize Heff, find ground state 'theta'
102         theta0 = np.reshape(self.psi.get_theta2(i), [Heff.shape[0]]) # initial guess
103         e, v = arp.eigsh(Heff, k=1, which='SA', return_eigenvectors=True, v0=theta0)
104         theta = np.reshape(v[:, 0], Heff.theta_shape)
105         # split and truncate
106         Ai, Sj, Bj = split_truncate_theta(theta, self.chi_max, self.eps)
107         # put back into MPS

```

```

108 Gi = np.tensordot(np.diag(self.psi.Ss[i]**(-1)), Ai, axes=[1, 0]) # vL [vL*], [vL] i vC
109 self.psi.Bs[i] = np.tensordot(Gi, np.diag(Sj), axes=[2, 0]) # vL i [vC], [vC*] vC
110 self.psi.Ss[j] = Sj # vC
111 self.psi.Bs[j] = Bj # vC j vR
112 self.update_LP(i)
113 self.update_RP(j)
114
115 def update_RP(self, i):
116     """Calculate RP right of site 'i-1' from RP right of site 'i'."""
117     j = i - 1
118     RP = self.RPs[i] # vR* wR* vR
119     B = self.psi.Bs[i] # vL i vR
120     Bc = B.conj() # vL* i* vR*
121     W = self.H_mpo[i] # wL wR i i*
122     RP = np.tensordot(B, RP, axes=[2, 0]) # vL i [vR], [vR*] wR* vR
123     RP = np.tensordot(RP, W, axes=[[1, 2], [3, 1]]) # vL [i] [wR*] vR, wL [wR] i [i*]
124     RP = np.tensordot(RP, Bc, axes=[[1, 3], [2, 1]]) # vL [vR] wL [i], vL* [i*] [vR*]
125     self.RPs[j] = RP # vL wL vL* (== vR* wR* vR on site i-1)
126
127 def update_LP(self, i):
128     """Calculate LP left of site 'i+1' from LP left of site 'i'."""
129     j = i + 1
130     LP = self.LPs[i] # vL wL vL*
131     B = self.psi.Bs[i] # vL i vR
132     G = np.tensordot(B, np.diag(self.psi.Ss[j]**(-1)), axes=[2, 0]) # vL i [vR], [vR*] vR
133     A = np.tensordot(np.diag(self.psi.Ss[i]), G, axes=[1, 0]) # vL [vL*], [vL] i vR
134     Ac = A.conj() # vL* i* vR*
135     W = self.H_mpo[i] # wL wR i i*
136     LP = np.tensordot(LP, A, axes=[2, 0]) # vL wL* [vL*], [vL] i vR
137     LP = np.tensordot(W, LP, axes=[[0, 3], [1, 2]]) # [wL] wR i [i*], vL [wL*] [i] vR
138     LP = np.tensordot(Ac, LP, axes=[[0, 1], [2, 1]]) # [vL*] [i*] vR*, wR [i] [vL] vR
139     self.LPs[j] = LP # vR* wR vR (== vL wL* vL* on site i+1)

```
