

## Lecture 15: Self-Adjointness

- Recall from math 54 that if  $A = A^T$ ,  $A$  is called a symmetric matrix and we may apply the Spectral Theorem to diagonalize  $A$ :

e.g.)  $A = \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix}$  has  $A = A^T$

and  $A$  has eigenvalues that are roots of

$$(2-x)^2 - 1 = x^2 - 4x + 3 \quad \text{i.e. } x=1, x=3.$$

The eigenvectors are  $\begin{bmatrix} 1 \\ 1 \end{bmatrix}, \begin{bmatrix} 1 \\ -1 \end{bmatrix}$  giving

$$A = \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 3 \end{bmatrix} \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix}$$

- Another way to view symmetry: on  $\mathbb{R}^n$   
 $\langle Ax, y \rangle = (Ax) \cdot (y) = x \cdot (A^T y) = \langle x, A^T y \rangle$   
so  $A = A^T$  says  $\langle Ax, y \rangle = \langle x, Ay \rangle$

- This is a property we may generalize:  
for a linear map  $A: H \rightarrow H$ , we will  
focus on self-adjoint cases where  $\langle Ax, y \rangle = \langle x, Ay \rangle$   
for all  $x, y \in H$ .

**Lemma 7.11** Suppose that  $\Omega \subseteq \mathbb{R}^n$  is a bdd. domain with  $C^1$  bdy. If  $u, v \in C^2(\bar{\Omega})$  and both satisfy either homogeneous Dirichlet or Neumann bdy conditions, on  $\partial\Omega$ , then

$$\langle \Delta u, v \rangle = \langle u, \Delta v \rangle.$$

**Prf:**  $\int_{\Omega} (\Delta u) \bar{v} - u (\Delta \bar{v}) dx = \int_{\partial\Omega} \left[ \bar{v} \frac{\partial u}{\partial \eta} - u \frac{\partial \bar{v}}{\partial \eta} \right] dx$

and Dirichlet  $v|_{\partial\Omega} = 0 = u|_{\partial\Omega}$

Neumann  $\frac{\partial v}{\partial \eta}|_{\partial\Omega} = 0 = \frac{\partial u}{\partial \eta}|_{\partial\Omega}$

give RHS = 0  $\square$

• Thus, under appropriate space restrictions and Boundary conditions, the Laplacian ~~is~~ acts self-adjoint (Spaces/domains to make it self-adjoint are more nuanced).

**Lemma 7.12** Suppose  $\{\lambda_j\}$  is a sequence of eigenvalues of  $-\Delta$  on a bounded domain  $U \subset \mathbb{R}^n$ , with eigenvectors in  $C^2(\bar{U})$  subject to a self-adjoint b.c. condition.

Then,  $\lambda_j \in \mathbb{R}$  and, after a possible rearrangement, the eigenvectors form an orthonormal sequence in  $L^2(U)$ .

Furthermore,  $\lambda_j > 0$  for Dirichlet conditions and  $\lambda_j \geq 0$  for Neumann.

**[Pf]** Suppose we have a sequence  $\{\phi_j\} \in C^2(U)$  satisfying  $-\Delta \phi_j = \lambda_j \phi_j$ . Notice  $\langle \Delta \phi_j, \phi_j \rangle = \langle \phi_j, \Delta \phi_j \rangle \Rightarrow (\lambda_j - \bar{\lambda}_j) \|\phi_j\|^2 = 0$  so  $\lambda_j \in \mathbb{R}$ .

Then,  $\langle \phi_j, \Delta \phi_k \rangle = \langle \Delta \phi_j, \phi_k \rangle$

$\Rightarrow -\lambda_j \langle \phi_j, \phi_k \rangle = -\bar{\lambda}_k \langle \phi_j, \phi_k \rangle$

and for  $\phi_j, \phi_k$  ~~for~~  $\lambda_j \neq \lambda_k \Rightarrow \langle \phi_j, \phi_k \rangle = 0$ .

If some  $\lambda_j$ 's are equal, we use Gram-Schmidt to obtain orthonormal eigenvectors.

Normalizing so  $\|\phi_j\|_2 = 1$

$$\begin{aligned} \lambda_j &= \langle -\Delta \phi_j, \phi_j \rangle = \int_{\Omega} (-\Delta \phi_j)(\bar{\phi}_j) dx \\ &= \int_{\Omega} |\nabla \phi_j|^2 dx - \underbrace{\int_{\partial \Omega} \bar{\phi}_j \frac{\partial \phi_j}{\partial \eta} dS}_{\text{vanishes in Dirichlet/Neumann conditions.}} \end{aligned}$$

If  $\lambda_j = 0$ ,  $\nabla \phi_j = 0$  gives  $\phi_j$  constant. In the Dirichlet case,  $\phi_j = 0$  is trivial and not an eigenvector.  $\square$

ex.) In lecture 10, we found a set of eigenfunctions for  $-\Delta$  on  $D \subseteq \mathbb{R}^2$ , with Dirichlet B.C

$$\phi_{k,m}(r, \theta) = e^{ik\theta} \bar{J}_k(j_{k,m}r) \quad k \in \mathbb{Z}, m \in \mathbb{N}$$

Corresponding to eigenvalue  $j_{k,m}^2$   
 (no  $j_{k,m}^2$  was repeated, as it was the  $\theta$  of  $J_{k+\frac{1}{2}}$ , except possibly  $j_{k,m}^2 = j_{-k,m}^2$  b/c  $\bar{J}_{-k} = -\bar{J}_k$ )

•) The orthogonality condition is

$$\langle \phi_{k,m}, \phi_{k',m'} \rangle_{L^2} = \int_0^1 \int_0^{2\pi} \phi_{k,m}(r, \theta) \phi_{k',m'}(r, \theta) r d\theta dr$$

$$= \int_0^1 \int_0^{2\pi} e^{i(k-k')\theta} \bar{J}_k(j_{k,m}r) \bar{J}_{k'}(j_{k',m'}r) d\theta dr$$

$$\text{if } k \neq k', \int_0^{2\pi} e^{i(k-k')\theta} d\theta = 0 \text{ so this}$$

whole integral is 0.

if  $k = k'$

$$\begin{aligned} \langle \phi_{k,m}, \phi_{k',m'} \rangle &= 2\pi \int_0^1 r \bar{J}_k(j_{k,m}r) \bar{J}_k(j_{k',m'}r) dr \\ &= 0 \text{ for } m \neq m' \end{aligned}$$

by the orthog. of eigenfunctions.

This occurs because of oscillations in  $J_k$