

# Quantum Mechanics B (Physics 212B) Winter 2019

## Worksheet 1 – Solutions

### Problems

#### 1. Bogliubov Transformation

In our last discussion we solved a new Hamiltonian by defining a set of transformed creation/annihilation operators which satisfy the same algebra  $[A, A^\dagger] = \mathbb{1}$

More generally consider  $\hat{b} = \hat{a} \cosh \eta + \hat{a}^\dagger \sinh \eta$

- (a) Show that  $[\hat{b}, \hat{b}^\dagger] = \mathbb{1}$   
 $[\hat{b}, \hat{b}^\dagger] = \cosh^2 \eta [a, a^\dagger] - \sinh^2 \eta [a, a^\dagger] = 1[a, a^\dagger] = \mathbb{1}$

- (b) Show that  $\hat{b} = U \hat{a} U^\dagger$  for  $U = e^{\frac{\eta}{2}(\hat{a}\hat{a} - \hat{a}^\dagger \hat{a}^\dagger)}$

Time to bust out BCH:  $U \hat{a} U^\dagger = \hat{a} + [A, \hat{a}] + \frac{1}{2}[A, [A, \hat{a}]] + \dots$

Where  $A \equiv \frac{\eta}{2}(\hat{a}\hat{a} - \hat{a}^\dagger \hat{a}^\dagger)$ . Note  $[a^\dagger a^\dagger, a] = -2a^\dagger$  and  $[aa, a^\dagger] = 2a$

So  $[A, \hat{a}] = \eta a^\dagger$  and  $[A, [A, \hat{a}]] = \eta [A, a^\dagger] = \eta^2 a$

So everything decomposes into odd terms with  $a^\dagger$  and even terms  $a$ . All plus signs. This gives the form of  $b$

Now consider the Hamiltonian

$$\hat{H} = \omega \hat{a}^\dagger \hat{a} + \frac{V}{2}(\hat{a}\hat{a} + \hat{a}^\dagger \hat{a}^\dagger) \quad (1)$$

- (c) Diagonalize the Hamiltonian (1) using the  $\hat{b}$  operators for suitably chosen  $\eta$

We should look for an operator of the form  $H = \Omega b^\dagger b + F$  for some constant  $F$

$$b^\dagger b = (\cosh^2 \eta + \sinh^2 \eta) a^\dagger a + \sinh^2 \eta + \cosh \eta \sinh \eta (aa + a^\dagger a^\dagger)$$

$$\text{So } \Omega \cosh(2\eta) = \omega \text{ and } \Omega \sinh(2\eta) = V \implies V = \omega \tanh(2\eta)$$

It also must be that  $\Omega \sinh^2 \eta + F = 0$ ; solving for  $\Omega$  and  $F$  independently is just algebra.  $\Omega = \omega \cosh(2\eta) - V \sinh(2\eta)$

The spectrum is simple now though!  $E_n = \Omega n + F$

- (d) Show there is a limit on  $V$  for which this Hamiltonian makes physical sense

$$\Omega \text{ must be positive. } \omega \cosh(2\eta) > V \sinh(2\eta) \implies V < \frac{\omega}{\tanh(2\eta)} = \frac{\omega^2}{V}$$

Thus  $V < \omega$

#### 2. Fermionic Harmonic Oscillator

Let's discuss a system with no simple classical analog. Consider "fermionic" operators  $c$  and  $c^\dagger$  which obey the following anti-commutation relation:

$$\hat{c}^2 = 0 = (\hat{c}^\dagger)^2 \quad \{\hat{c}, \hat{c}^\dagger\} \equiv c c^\dagger + c^\dagger c = \mathbb{1} \quad (2)$$

- (a) Assuming the Hilbert space cannot be empty, show there exists a state  $|0\rangle$  for which  $c|0\rangle = 0$ . Similarly, show there must exist a non-vanishing state  $|1\rangle \equiv c^\dagger|0\rangle$ . By assumption there must exist a state  $|\psi\rangle$ . Either  $c|\psi\rangle = 0$ , in which case  $|\psi\rangle \equiv |0\rangle$  and we're done, or  $c|\psi\rangle \neq 0$ . In the latter case  $|0\rangle \propto c|\psi\rangle$  as  $c^2|\psi\rangle = 0$  by fermionic algebra. To show  $c^\dagger|0\rangle \neq 0$  we consider acting  $c$  on it.  $cc^\dagger|0\rangle = (1 - c^\dagger c)|0\rangle = |0\rangle$  from algebra. This cannot happen if  $c^\dagger|0\rangle$  vanished.
- (b) Write explicit matrix representations of the operators  $c$  and  $c^\dagger$  in the  $|0\rangle, |1\rangle$  basis.  $c = \sigma^- = X - \mathbf{i}Y$  and  $c^\dagger = \sigma^+ = X + \mathbf{i}Y$  where  $X$  and  $Y$  are Pauli matrices. From Pauli algebra they obey the fermionic statistics.
- (c) Show that  $\hat{d} = \hat{c} \cos \theta + \hat{c}^\dagger \sin \theta$  is the analogous Bogliobuv transform. Similar but the anticommutation means  $\{d, d^\dagger\} = \cos^2 \theta \{\hat{c}, \hat{c}^\dagger\} + \sin^2 \theta \{\hat{c}, \hat{c}^\dagger\} = \mathbb{1}$ .
- (d) Define the Hamiltonian in analogy to the bosonic SHO as:

$$H = \frac{\omega}{2}(c^\dagger c - cc^\dagger) \quad (3)$$

What are it's eigenstates and eigenvalues?

It's easy to see that it's diagonal in the  $|0\rangle, |1\rangle$  basis which eigenvalues  $E_\pm = \pm \frac{\omega}{2}$  with  $|0\rangle$  as the groundstate.

The operator  $c^\dagger c$  is again a number operator and labels 0, 1 are fermion number.