## Anomalous Hall effect

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The Hamiltonian of the problem is

$$\hat{H} = \frac{\hbar^2 k^2}{2m} + \lambda (k_x \sigma_y - k_y \sigma_x) - \Delta \sigma_z \tag{1}$$

It can be easily diagonalized, the energies are

$$E_k^{(1,2)} = \frac{\hbar^2 k^2}{2m} \pm \sqrt{\lambda^2 k^2 + \Delta^2}$$
 (2)

It is convenient to introduce  $\epsilon_k = \sqrt{\lambda^2 k^2 + \Delta^2}$ .

The eigenvectors are, as usual,

$$u_{\uparrow} = \begin{pmatrix} \cos\frac{\theta}{2} \\ \sin\frac{\theta}{2}e^{i\phi} \end{pmatrix}$$

$$u_{\downarrow} = \begin{pmatrix} -\sin\frac{\theta}{2}e^{-i\phi} \\ \cos\frac{\theta}{2} \end{pmatrix}$$
(3)

 $\operatorname{Here}$ 

$$\cos \theta = -\frac{\Delta}{\epsilon_k}$$

$$e^{i\phi} = \frac{-k_y + ik_x}{k}$$
(4)

The Berry curvature can be evaluated from the gauge—invariant formula. For the upper band,

$$\Omega_z = \Omega_{xy} = -\frac{i\lambda^2}{4\epsilon_k^2} \left[ \langle u_{\uparrow} | \sigma_y | u_{\downarrow} \rangle \langle u_{\downarrow} | \sigma_x | u_{\uparrow} \rangle - \langle u_{\uparrow} | \sigma_x | u_{\downarrow} \rangle \langle u_{\downarrow} | \sigma_y | u_{\uparrow} \rangle \right]$$

$$\Omega_x = \Omega_y = 0$$
(5)

The matrix elements:

$$\langle u_{\uparrow} | \sigma_x | u_{\downarrow} \rangle = \cos \frac{\theta^2}{2} - e^{-2i\phi} \sin \frac{\theta}{2}$$

$$\langle u_{\downarrow} | \sigma_y | u_{\uparrow} \rangle = i(\cos \frac{\theta^2}{2} + e^{2i\phi} \sin \frac{\theta}{2})$$
(6)

So,

$$\Omega_z^{\text{upper}} = -\frac{\lambda^2}{2\epsilon_k^2} \cos \theta = \frac{\lambda^2 \Delta}{2\epsilon_k^3}$$
(7)

For lower band, we obtain similarly

$$\Omega_z^{\text{lower}} = -\frac{\lambda^2 \Delta}{2\epsilon_k^3} \tag{8}$$

Now let's find the Hall conductivity. Let the electric field be (E,0,0). Then the equations of motion for the wave packet read

$$\dot{x} = \frac{1}{\hbar} \frac{\partial E_k}{\partial k_x} - \dot{y}\Omega_z$$

$$\dot{y} = \frac{1}{\hbar} \frac{\partial E_k}{\partial k_y} + \dot{x}\Omega_z$$

$$\hbar \dot{k_x} = -eE$$

$$\hbar \dot{k_y} = 0$$
(9)