

Mathematical Preliminaries

Elements of Linear Algebra and Calculus

EE4375 - FEM For EE Applications

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Last updated December 27, 2023

In These Slides

- Elements of linear algebra for FDM and FEM
- Elements of calculus of functions in one variable for FDM
- Elements of calculus of functions in several variables for FDM
- Elements of calculus of functions in one variable for FEM
- Elements of calculus of functions in several variables for FEM

FDM: finite difference method - FEM: finite element method

Elements of linear algebra for FDM and FEM

In This Section

- vectors, representation of vectors, transpose of vectors, element-wise operations on vectors, equality of vectors, linear combinations of vectors, magnitude of vectors, inner product of vectors;
- matrices, representation of matrices, transpose of matrices, equality of matrices, linear combinations of matrices;
- matrix-vector multiplication (order of symbols matters)
- linear system solve
- (non-linear system solve and time-stepping)

Elements of linear algebra for FDM and FEM

Vectors (1/2)

- $\mathbf{v} \in \mathbb{R}^n$: n -dimensional column **vector**: $\mathbf{v} = \begin{pmatrix} v_1 \\ \vdots \\ v_n \end{pmatrix} = (v_i)_{1 \leq i \leq n}$
- real-valued elements v_i (unless stated otherwise)
- $\mathbf{v}^T \in \mathbb{R}^n$: **transpose** of \mathbf{v}

n -dimensional row vector: $\mathbf{v}^T = (v_1 \quad \dots \quad v_n)$



- **element-wise operation** on \mathbf{v} :

form new vector by operation on each element of \mathbf{v}

element-wise-operation(\mathbf{v}) = $(\text{operation}(v_i))_{1 \leq i \leq n}$

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Vectors (2/2)

- equal to zero-vector

$\mathbf{v} = \mathbf{0}_n \Leftrightarrow v_i = 0$ for $1 \leq i \leq n$ (all components are zero)

- linear combination of vectors

$\alpha, \beta \in \mathbb{R}$ numbers and $\mathbf{v}, \mathbf{w} \in \mathbb{R}^n$ vectors

$\alpha \mathbf{v} + \beta \mathbf{w} = (\alpha v_i + \beta w_i)_{1 \leq i \leq n}$ (new vector)

- Euclidean norm - magnitude of vector - distance to zero

$\|\mathbf{v}\|_2 = \sqrt{\sum_{i=1}^n v_i^2} \in \mathbb{R}$ (scalar)

- inner product of vectors of same size

$\mathbf{v}, \mathbf{w} \in \mathbb{R}^n$ vectors then $\mathbf{v} \cdot \mathbf{w} = \sum_{i=1}^n v_i w_i \in \mathbb{R}$ (scalar)

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Matrices (1/2)

- $A \in \mathbb{R}^{m \times n}$: m -rows n -columns matrix: $A = \begin{pmatrix} A_{11} & \dots & A_{1n} \\ \vdots & \vdots & \vdots \\ A_{m1} & \dots & A_{mn} \end{pmatrix}$
- real-valued elements A_{ij} (unless stated otherwise)
- $A^T \in \mathbb{R}^{n \times m}$: transpose of A defined as $(A^T)_{ij} = A_{ji}$
- $n \neq m$: A is rectangular - $n = m$: A is square

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Matrices (2/2)

- **equal to zero-matrix**

$$A = \mathbf{0}_{m \times n} \Leftrightarrow A_{ij} = 0 \text{ for } 1 \leq i \leq m \text{ and } 1 \leq j \leq n \text{ (all cmpts zero)}$$

- **linear combination** of matrices of same size

$\alpha, \beta \in \mathbb{R}$ numbers and $A, B \in \mathbb{R}^{m \times n}$ matrices


$$(\alpha A + \beta B)_{ij} = \alpha A_{ij} + \beta B_{ij} \text{ for } 1 \leq i \leq m \text{ and } 1 \leq j \leq n$$

(new matrix)

- **norm** of a matrix: see references

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Matrix-Vector Multiplication

- $\mathbf{v} \in \mathbb{R}^n$ and $A \in \mathbb{R}^{n \times n}$ (A is square $m = n$)
- $\mathbf{w} = A\mathbf{v} \in \mathbb{R}^n$ 
- $w_i = \sum_{j=1}^n A_{ij} v_j$ (i fixed, j summation index)
- exercises: repeat for A rectangular, $\mathbf{v}^T A$ and $\mathbf{v}^T A \mathbf{v}$

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Linear System Formulation

- **given** $\mathbf{f} \in \mathbb{R}^n$ and **given** $A \in \mathbb{R}^{n \times n}$ (A is again square)
- assume A is non-singular
- **unknown** $\mathbf{u} \in \mathbb{R}^n$
- find $\mathbf{u} \in \mathbb{R}^n$ by solving the linear system $\boxed{A\mathbf{u} = \mathbf{f}}$
- find u_1, \dots, u_n such that for all $i \in \{1, \dots, n\}$ we have that

$$\sum_{j=1}^n A_{ij} u_j = f_i$$

- solution \mathbf{u} exists and is unique (by construction)

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Linear System Solve

- $A\mathbf{u} = \mathbf{f} \Leftrightarrow \begin{pmatrix} A_{11} & \dots & A_{1n} \\ \vdots & \vdots & \vdots \\ A_{n1} & \dots & A_{nn} \end{pmatrix} \begin{pmatrix} u_1 \\ \vdots \\ u_n \end{pmatrix} = \begin{pmatrix} f_1 \\ \vdots \\ f_n \end{pmatrix}$
- solve by LU-decomposition of A
- **never** compute A^{-1} ! (add picture here)
- see assignments

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Extensions

- data structures for sparse matrices (few non-zero elements only)
- iterative solution method for large sparse matrices

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Recap

- given $\mathbf{f} \in \mathbb{R}^n$ and given $\mathbf{A} \in \mathbb{R}^{n \times n}$
- find $\mathbf{u} \in \mathbb{R}^n$ such that $\mathbf{A}\mathbf{u} = \mathbf{f}$
- using LU-decomposition of \mathbf{A}

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References

- bachelor courses in linear algebra

Further Reading

- https://en.wikipedia.org/wiki/Numerical_linear_algebra

Elements of calculus of functions in one var - FDM

Throughout This Section

- **given** $x \in \Omega = (0, 1)$ interval or $0 < x < 1$
- $x = 0$ and $x = 1$ left and right boundary points
- $f(x)$: assumed **known** or given function with domain $x \in \Omega$
assume given force or electrical charge distribution
- **given** α number
- $u(x)$: assumed **unknown** with domain $x \in \Omega$
 u is short for unknown - displacement or electrical potential
- $u'(x) = \frac{du}{dx}(x)$ and $u''(x) = \frac{d^2u}{dx^2}(x)$: short hand notation

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Boundary Value Problem for Second Order Differential Equations

- **given** $x \in \Omega = (0, 1)$ with left and right boundary point
- **given** $f(x)$ with domain Ω
- **find** $u(x)$ such that $u''(x) = f(x)$ plus boundary conditions

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Boundary Value Problem for Second Order Differential Equations

- **given** $x \in \Omega = (0, 1)$ with left and right boundary point
- **given**: $f(x)$ given function and α given number
- **find**: $u(x)$ such that

$$-u''(x) = f(x) \text{ for } 0 < x < 1 \text{ (differential equation on } \Omega)$$

$$u(x = 0) = 0 \text{ (Dirichlet boundary condition in } x = 0)$$

$$\frac{du}{dx}(x = 1) = \alpha \text{ (Neumann boundary condition in } x = 1)$$

- differential equation **not** valid in $x = 0$ or $x = 1$
boundary conditions **are** valid in $x = 0$ or $x = 1$

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Exercise

- choose $f(x) = 1$
- **find**: $u(x)$ such that

$$-u''(x) = 1 \text{ for } 0 < x < 1 \text{ (differential equation on } \Omega)$$

$$u(x = 0) = 0 \text{ (Dirichlet boundary condition in } x = 0)$$

$$\frac{du}{dx}(x = 1) = \alpha \text{ (Neumann boundary condition in } x = 1)$$

- solve using **pen and paper** and plot computed solution for various values of α

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Solution (1/2)

- by integration twice and using the boundary conditions

$$u(x) = -\frac{1}{2}x^2 + (1 - \alpha)x$$

- solving by guessing and iteratively improving guess is allowed
- solution can be checking by checking differential equation and boundary conditions
- this examples is will guide remainder of the course

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Solution (2/2)

- in example $f(x) = 1$
- for more general $f(x) = \text{whatever}(x)$ proceed as follows
- integrate using pen-and-paper: $u'(x) = \int^x f(\eta) d\eta + C_1$
(here η is the integration variable)
- thus by replacing x by ξ : $u'(\xi) = \int^\xi f(\eta) d\eta + C_1$
- integrate once more

$$u(x) = \int^x u'(\xi) d\xi = \int^x \int^\xi f(\eta) d\eta d\xi + C_1 x + C_2$$

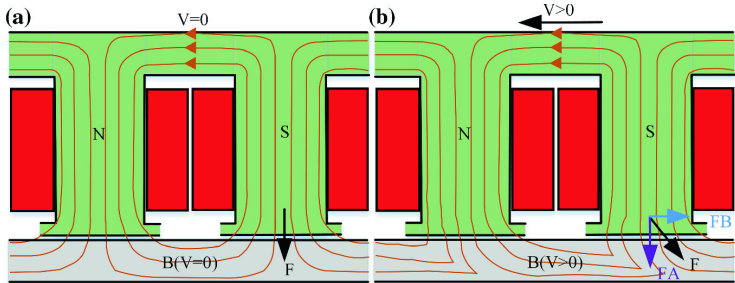
with C_1 and C_2 from the boundary conditions

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Example of Rotor/Stator Motion

(a) $-u''(x) = f(x)$

(b) $u''(x) + \beta u'(x) = f(x)$



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Solve Using Available Tools

- **find**: $u(x)$ such that

$$-u''(x) = 1 \text{ plus boundary conditions}$$

- solve symbolically using function `dsolve()` in **Sympy.jl**
- solve numerically using collocation method or shooting method implemented in `BVPproblem` in **DifferentialEquations.jl**
- **important**: know that methods exists and how to apply
- **outside scope of course**: details of the methods being methods

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More Involved Models in 1D

- variable diffusion $c(x)$: $-\frac{d}{dx} \left[c(x) \frac{du(x)}{dx} \right] = f(x) + bc$
- induced current effects: $u''(x) + \beta u(x) = f(x) + b.c.$
- rotor/stator motion: $u''(x) + \beta u'(x) = f(x)$
- ferromagnetic saturation: $-\frac{d}{dx} \left[c[u'(x)] \frac{du(x)}{dx} \right] = f(x) + bc$
- lots more ...
- no worries now - see later

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From here on

Goodbye Analytics

Hello Numerics

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First Order Derivative and Its Finite Difference Approximation

- **definition:**
$$\frac{d u}{d x}(x = x_i) = \lim_{\Delta x \rightarrow 0} \frac{u(x_i + \Delta x) - u(x_i)}{\Delta x}$$
- **forward** difference:
$$\frac{d u}{d x}(x = x_i) \approx \frac{u(x_i + \Delta x) - u(x_i)}{\Delta x}$$
- **backward** difference:
$$\frac{d u}{d x}(x = x_i) \approx \frac{u(x_i) - u(x_i - \Delta x)}{\Delta x}$$
- **central** difference:
$$\frac{d u}{d x}(x = x_i) \approx \frac{u(x_i + \Delta x) - u(x_i - \Delta x)}{2 \Delta x}$$

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Second Order Derivative and Its Finite Difference Approximation

- set $v(x) = \frac{d u}{d x}$ then $\frac{d v}{d x} = \frac{d^2 u}{d x^2}$
- use half steps: $v_{i+1/2} \approx v(x = x_{i+1/2}) = \frac{d u}{d x}(x_{i+1/2}) \approx \frac{u_{i+1} - u_i}{\Delta x}$
- similar to obtain $v_{i-1/2}$
- combine to obtain
$$\frac{d^2 u}{d x^2}(x = x_i) \approx \frac{v_{i+1/2} - v_{i-1/2}}{\Delta x} \approx \frac{u_{i-1} - 2u_i + u_{i+1}}{\Delta x^2}$$
- accuracy of second order accurate approximation using Taylor series expansions ($\mathcal{O}(\Delta x^2)$)

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Accuracy of Three-Point Finite Difference Approximation

- three-point finite difference scheme

$$\frac{d^2 u}{dx^2}(x = x_i) \approx \frac{u_{i-1} - 2u_i + u_{i+1}}{\Delta x^2} = \frac{u_{i-1} - 2u_i + u_{i+1}}{h^2}$$

- Taylor expansion of u_{i+1} and u_{i-1} around x_i

$$\begin{aligned} u_{i+1} &= u(x_{i+1}) = u_i + h u'_i + \frac{h^2}{2} u''_i + \frac{h^3}{6} u'''_i + \frac{h^4}{24} u''''_i \\ u_{i-1} &= u(x_{i-1}) = u_i - h u'_i + \frac{h^2}{2} u''_i - \frac{h^3}{6} u'''_i + \frac{h^4}{24} u''''_i \end{aligned}$$

- add Taylor expansion, subtract $2u_i$ and divide by h^2 to obtain

$$\frac{u_{i-1} - 2u_i + u_{i+1}}{h^2} = u''_i + \frac{h^2}{12} u''''_i$$

- three-point scheme is thus **second order** accurate. Reducing meshwidth h by 2, reduces error in finite difference approximation by 4.

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Accuracy of Three-Point Finite Difference Approx. (cont'd)

- we obtained earlier that

$$\frac{u_{i-1} - 2u_i + u_{i+1}}{h^2} = u_i'' + h^2/12 u_i''''$$

- error zero and thus scheme exact in case that $u''''(x) = 0$
- three-point difference scheme is thus exact for first, second and third order polynomials

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Accuracy of Two-Point Finite Difference Approximation

- forward and backward difference scheme for u'_i are **first** order accurate
(exercise: show this result using Taylor expansions)
- central difference scheme for u'_i is **second** order accurate
(exercise: show this result using Taylor expansions)
- use ghost points and **central** scheme to approximate Neumann boundary conditions
(details to be added)
(exercise: show this result using Taylor expansions)

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Finite Difference Method in Short

- start from $-u''(x) = f(x)$ + boundary conditions
- approximate differences using finite differences
- form linear system for unknown in grid points
- details coming soon

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Extensions

- differential equations closer to practice
- more accurate finite difference approximation
- fast solution
- embed in digital twin for electrical machine or high voltage line

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References

- bachelor courses in calculus

Further Reading

- https://en.wikipedia.org/wiki/Finite_difference

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Recap

- boundary value problem for the Poisson equation for $u(x)$
- analytical solution methods
- finite difference approximation

$$-\frac{d^2 u}{dx^2}(x = x_i) = \frac{-u_{i-1} + 2u_i - u_{i+1}}{\Delta x^2} + \mathcal{O}(\Delta x^2)$$

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In This Section

- functions in two variables $u(x, y)$ and $f(x, y)$
- analytical solution method for the Laplace and Poisson equation using separation of variables
- finite difference approximation of the derivative of functions in two variables

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Domain of Computation

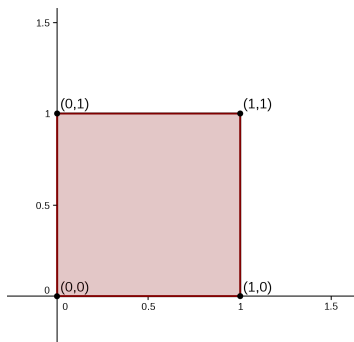
Before

Unit Interval: $\Omega = (0, 1)$



Here

Unit Square: $\Omega = (0, 1)^2$



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Domain of Computation

- **Before:** $u(x)$ and $f(x)$ for $x \in \Omega = (0, 1)$
- **Here:** $u(x, y)$ and $f(x, y)$ for $(x, y) \in \Omega = (0, 1)^2$
- **Goal:** formulate $u''(x, y) = f(x, y)$
- **However:** what are $u'(x, y)$ and $u''(x, y)$?

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Gradient (1/2): Definition

- unit (basic) vectors $\mathbf{i} = (1, 0) \in \mathbb{R}^2$ and $\mathbf{j} = (0, 1) \in \mathbb{R}^2$
- gradient of $u(x, y)$ using $\nabla = (\partial_x, \partial_y)$:

$$\text{grad } u(x, y) = \nabla u = \left(\frac{\partial u}{\partial x}, \frac{\partial u}{\partial y} \right) = \frac{\partial u}{\partial x} \mathbf{i} + \frac{\partial u}{\partial y} \mathbf{j}$$

- $-\nabla u$: direction of largest increase in u - direction of diffusion

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Gradient (2/2): Example

- electrical potential: $\phi(x, y)$

- electrical field vector: $\mathbf{E}(x, y) = -\nabla\phi(x, y)$



- electrical field components: $E_x(x, y) = -\frac{\partial\phi}{\partial x}$ and $E_y(x, y) = -\frac{\partial\phi}{\partial y}$

- or $\mathbf{E}(x, y) = E_x(x, y)\mathbf{i} + E_y(x, y)\mathbf{j} = -\frac{\partial\phi}{\partial x}\mathbf{i} - \frac{\partial\phi}{\partial y}\mathbf{j}$

- how about magnetic flux $\mathbf{B}(x, y)$?

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Divergence (1/3): Recap on Inner Product

- **inner product** denoted by \cdot

- performs a **contraction**

suppose that $\mathbf{a} = (a_x, a_y)$ and $\mathbf{b} = (b_x, b_y)$, then

$$\mathbf{a} \cdot \mathbf{b} = (a_x, a_y) \cdot (b_x, b_y) = a_x b_x + a_y b_y \in \mathbb{R} \text{ (single leg only)}$$

- example: electrical field $\mathbf{E}(x, y)$ - electrical displacement $\mathbf{D}(x, y)$

$$\mathbf{E}(x, y) \cdot \mathbf{D}(x, y) \text{ (electrostatic energy density)}$$

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Divergence (2/3): Definition

- vector field

$$\mathbf{F}(x, y) = P(x, y)\mathbf{i} + Q(x, y)\mathbf{j} \text{ (vector)}$$

- examples: $\mathbf{F}(x, y) = \nabla u(x, y)$ or $\mathbf{F}(x, y) = \mathbf{E}(x, y)$
- divergence of the vector field $\mathbf{F}(x, y)$

$$\operatorname{div} \mathbf{F}(x, y) = \nabla \cdot \mathbf{F}(x, y) = \frac{\partial P(x, y)}{\partial x} + \frac{\partial Q(x, y)}{\partial y} \text{ (scalar)}$$

- interpretation: $\nabla \cdot \mathbf{F}(x, y)$: change in F over small volume

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Divergence (3/3): Example

- electrical charge density $\rho(x, y)$ (scalar)

- electrical displacement

$$\mathbf{D}(x, y) = (D_x(x, y), D_y(x, y)) = D_x(x, y)\mathbf{i} + D_y(x, y)\mathbf{j} \text{ (vector)}$$

- Gauss's Law: $\nabla \cdot \mathbf{D}(x, y) = \rho(x, y)$ or

$$\frac{\partial D_x}{\partial x} + \frac{\partial D_y}{\partial y} = \rho \text{ (scalar)}$$

- how about Ampère's Law?

**Stay
Tuned!**

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Laplacian (1/2): Definition

- if $u(x, y)$ scalar function, then $\nabla u(x, y)$ is a vector function
- if \mathbf{F} vector function, then $\nabla \cdot \mathbf{F}$ is a scalar function
- put $\mathbf{F} = \nabla u(x, y)$, then

$$\nabla \cdot \mathbf{F} = \nabla \cdot \nabla u(x, y) = \nabla \cdot \left(\frac{\partial u}{\partial x} \mathbf{i} + \frac{\partial u}{\partial y} \mathbf{j} \right) = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}$$

- **Laplacian** - double derivative for $u(x, y)$

$$u'' \Rightarrow \Delta u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}$$

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Laplacian (2/2): Example in Electrostatics

- $\mathbf{x} = (x, y) \in \Omega$ - $f(\mathbf{x})$: **known/given** electrical charge density (possibly denoted by $\rho(\mathbf{x})$ elsewhere)
- **unknown/desired** electrical potential $u(\mathbf{x})$ (possibly denoted by $\phi(\mathbf{x})$ elsewhere)
- $\mathbf{E} = -\nabla u(\mathbf{x})$ and $\mathbf{D} = \epsilon \mathbf{E}$
- Gauss's Law:
 $\text{div } \mathbf{D} = f(\mathbf{x}) \Leftrightarrow -\text{div } (\epsilon \mathbf{E}) = f(\mathbf{x}) \Leftrightarrow -\text{div } (\text{grad } u(\mathbf{x})) = f(\mathbf{x})/\epsilon$
- Poisson equation for the electric potential:
(assume $\epsilon = 1$ - beware of the minus sign!) $-\Delta u(\mathbf{x}) = f(\mathbf{x})$
- similar for **magnetostatics**: stay tuned

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Boundary conditions (1/4)

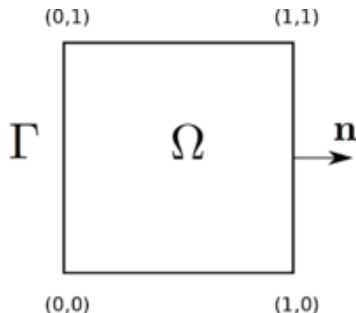
- **Before:** $u''(x) = f(x)$
- **Now:** $u''(x, y) = f(x, y)$ actually means $\Delta u = f(x, y)$
- what about the boundary conditions?
- **Before:** $\Omega = (0, 1)$ has boundary $x = 0$ and $x = 1$
- **Now:** $\Omega = (0, 1)^2$ has boundary Γ

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Boundary conditions (2/4)

$\Omega = (0,1)^2$: Unit Square

Γ boundary of Ω - \mathbf{n} : outward normal on Γ



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Boundary conditions (3/4)

- **direction** or vector: $\mathbf{s} = (s_x, s_y) = s_x \mathbf{i} + s_y \mathbf{j}$
- **directional derivative** of $u(x, y)$ in the direction of \mathbf{s}

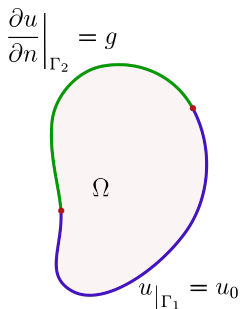
$$\frac{\partial u}{\partial \mathbf{s}} = \underbrace{\nabla u \cdot \mathbf{s}}_{\text{innerproduct}} = s_x \frac{\partial u}{\partial x} + s_y \frac{\partial u}{\partial y} \text{ (scalar)}$$

- choose $\mathbf{s} = \mathbf{n}$ or $\mathbf{s} = \mathbf{t}$: normal or tangential derivative
- example: electromagnetic force/torque computation using the Maxwell stress tensor

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Boundary conditions (4/4)

Two Types of Boundary Conditions: $\Gamma = \Gamma_D \cup \Gamma_N$



- Dirichlet condition on Γ_D

fix u

(equivalent of $x = 0$ in 1D)

- Neumann condition on Γ_N

fix $\frac{\partial u}{\partial n}$

(equivalent of $x = 1$ in 1D)

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Boundary Value Problem for Second Order Differential Equations

- **given** $(x, y) \in \Omega = (0, 1)^2$ with $\Gamma = \Gamma_D \cup \Gamma_N$ the boundary of Ω
- **given**: $f(x, y)$ given function and α given number
- **find**: $u(x, y)$ such that
 - $\Delta u(x, y) = f(x, y)$ for $0 < x, y < 1$ (differential equation on Ω)
 - $u = 0$ (Dirichlet boundary condition on Γ_D)
 - $\frac{\partial u}{\partial n} = \alpha$ (Neumann boundary condition on Γ_N)
- differential equation **invalid** on Γ - boundary conditions **valid** on Γ

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Exercise: Verify Solution

- suppose **given**: $u(x, y) = x(x-1)y(y-1)$
- then $\frac{\partial u}{\partial x} = (2x-1)y(y-1)$ and $\frac{\partial^2 u}{\partial x^2} = 2y(y-1)$
- thus $-\Delta u(x, y) = -2x(x-1) - 2y(y-1)$
- furthermore $u(x, y) = 0$ on Γ
- thus $u(x, y)$ solves

$$\begin{aligned} -\Delta u(x, y) &= -2x(x-1) - 2y(y-1) \text{ on } \Omega \\ u &= 0 \text{ on } \Gamma \end{aligned}$$

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Construction of Analytical Solution

- **separation of variables** $u(x, y) = X(x) Y(y)$
- $X(x)$ and $Y(y)$ solved separately as 1D problems
- see References section for details
- **in this course**: not applied analytically
- **in this course**: applied as cheat-sheet to solve numerically

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More Involved Models in 2D

- as before, replace $u(x)$ by $u(x, y)$

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From here on

Goodbye Analytics

Hello Numerics

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Mesh Generation

- equidistant mesh in both x and y -direction (insert figure here)
- mesh spacing: $\Delta x = \Delta y = h = 1/N$
- $x_i = (i - 1)h$ and $y_j = (j - 1)h$ for $1 \leq i, j \leq N + 1$
- **interior nodes:** (x_i, y_j) for $2 \leq i, j \leq N$
- **east boundary** where $x = 0$: fix $i = 1$
similar for other 3 boundaries

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Finite Difference Discretization (1/2)

- apply finite difference discretization in both x and y direction
- **as before by varying x** and keeping y fixed

$$\frac{\partial^2 u}{\partial x^2}(x = x_i, y = y_j) \approx \frac{u_{i-1,j} - 2u_{i,j} + u_{i+1,j}}{h^2} \quad (\leftrightarrow)$$

- **likewise in y -direction** and keeping x fixed

$$\frac{\partial^2 u}{\partial y^2}(x = x_i, y = y_j) \approx \frac{u_{i,j-1} - 2u_{i,j} + u_{i,j+1}}{h^2} \quad (\updownarrow)$$

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Finite Difference Discretization (2/2)

- **add contributions** in x and y -direction and multiply with -1

$$-\Delta u(x = x_i, y = y_j) \approx \frac{-u_{i-1,j} - u_{i,j-1} + 4u_{i,j} - u_{i+1,j} - u_{i,j+1}}{h^2}$$

- treat **boundary conditions**
- formulate system of linear equations for $u_{i,j}$
- solve for $u_{i,j}$

Elements of calculus of functions - two vars - FDM

References

- bachelor courses in calculus

Further Reading

- https://en.wikipedia.org/wiki/Separation_of_variables
- https://en.wikipedia.org/wiki/Finite_difference
- https://en.wikipedia.org/wiki/Kronecker_sum_of_discrete_Laplacians

Elements of calculus of functions - two vars - FDM

Recap

- boundary value problem for the Poisson equation for $u(x, y)$
- analytical solution methods
- finite difference approximation

$$-\Delta u(x = x_i, y = y_j) \approx \frac{-u_{i-1,j} - u_{i,j-1} + 4u_{i,j} - u_{i+1,j} - u_{i,j+1}}{h^2}$$

Elements of calculus of functions in one var: FEM

Strong and Weak Equality of Functions

- strong equality of functions on Ω
- weak equality of functions on Ω

Elements of calculus of functions in one var: FEM

Derivative: Integration by Parts in one variable (1/2)

- **fundamental theorem of calculus:**

assume $0 < x < 1$ or $x \in \Omega = (0, 1)$ and $F'(x) = \frac{dF(x)}{dx}$

$$\int_{\Omega} F'(x) dx = \int_0^1 F'(x) dx = [F(x)]_0^1 = F(1) - F(0)$$

- **choose** $F(x) = \frac{du}{dx} v(x)$ (motivation later)

observe $u(x)$ has prime and $v(x)$ has no prime

- then $F'(x) = \frac{d}{dx} \left(\frac{du}{dx} v(x) \right) = u''(x) v(x) + u'(x) v'(x)$

Elements of calculus of functions in one var: FEM

Derivative: Integration by Parts in one variable (2/2)

- subsequently from $\int_0^1 F'(x) dx = [F(x)]_0^1$

$$\int_0^1 [u''(x) v(x) + u'(x) v'(x)] dx = [u'(x) v(x)]_0^1$$

- after rearranging terms

$$\boxed{- \int_0^1 u''(x) v(x) dx = \int_0^1 u'(x) v'(x) dx - [u'(x) v(x)]_0^1}$$

LHS: $u(x)$ has double prime and $v(x)$ has no prime

RHS: both $u(x)$ and $v(x)$ have one prime - additional fudge term
(used later)

Elements of calculus of functions in one var: FEM

Integration: Quadrature by Trapezoidal Rule

- trapezoidal rule: $0 \leq a < b \leq 1$

$$\int_a^b g(x) dx \approx \frac{b-a}{2} [g(a) + g(b)]$$

- Simpson rule: $0 \leq a < b \leq 1$

$$\int_a^b g(x) dx \approx \frac{b-a}{6} \left[g(a) + 4g\left(\frac{a+b}{2}\right) + g(b) \right]$$

Elements of calculus of functions in one var: FEM

Function Spaces

- $\Omega = (0, 1)$
- $V(\Omega)$
- $V_0(\Omega)$

Elements of calculus of functions in one var: FEM

Fourier Analysis

- Fourier analysis
- basis of function space

Elements of calculus of functions in one var: FEM

Finite Dimensional Approximation of a Function Space

- set of basis functions

Elements of calculus of functions in one var - FEM

Extensions

- to be decided

Elements of calculus of functions in one var - FEM

References

- bachelor courses in calculus

Further Reading

- `https://en.wikipedia.org/wiki/Integration_by_parts`
- `https://en.wikipedia.org/wiki/Trapezoidal_rule`

Elements of calculus of functions in one var - FEM

Recap

- integration by parts in 1D

Elements of calculus of functions in n vars: FEM

Two Small Exercises

- suppose that $\mathbf{F} = v \nabla u$
- Exercise 1: compute $\nabla \cdot \mathbf{F} = \nabla \cdot (v \nabla u)$
- Exercise 2: compute $\mathbf{F} \cdot \mathbf{n} = (v \nabla u) \cdot \mathbf{n}$

Elements of calculus of functions in n vars: FEM

Two Small Exercises

- suppose that $\mathbf{F} = v \nabla u$
- Ex 1: solution $\nabla \cdot \mathbf{F} = \nabla u \cdot \nabla v + v \Delta u$
- Ex 2: solution $\mathbf{F} \cdot \mathbf{n} = (v \nabla u) \cdot \mathbf{n} = v (\nabla u \cdot \mathbf{n}) = v \frac{\partial u}{\partial n}$
(by definition of $\frac{\partial u}{\partial n}$)

Elements of calculus of functions in n vars: FEM

Derivative: Integration by Parts in two variables (1/2)

- **Gauss Integration Theorem:** $(x, y) \in \Omega$ with boundary Γ

$$\int_{\Omega} \nabla \cdot \mathbf{F} d\Omega = \int_{\Gamma} \mathbf{F} \cdot \mathbf{n} ds$$

- **choose:** $\mathbf{F} = v \nabla u = (v \frac{\partial u}{\partial x}, v \frac{\partial u}{\partial y})$

Elements of calculus of functions in n vars: FEM

Derivative: Integration by Parts in two variables (1/2)

- then

$$\begin{aligned}\int_{\Omega} \nabla \cdot \mathbf{F} d\Omega &= \int_{\Omega} \nabla \cdot (v \nabla u) d\Omega = \int_{\Omega} [\nabla u \cdot \nabla v + v \Delta u] d\Omega \\ &= \int_{\Gamma} \mathbf{F} \cdot \mathbf{n} ds = \int_{\Gamma} \frac{\partial u}{\partial n} v ds\end{aligned}$$

- after rearranging terms:

$$\int_{\Omega} (-\Delta u) v d\Omega = \int_{\Omega} \nabla u \cdot \nabla v - \int_{\Gamma} \frac{\partial u}{\partial n} v ds$$

Elements of calculus of functions in n vars: FEM

Derivative: Integration by Parts in two variables (2/2)

- integration by parts formula becomes

$$\int_{\Omega} (-\Delta u) v \, d\Omega = \int_{\Omega} \nabla u \cdot \nabla v - \int_{\Gamma} \frac{\partial u}{\partial n} v \, ds$$

- observe that as before:

LHS: double derivatives in u - no derivatives on v

RHS: first order derivatives on both u and v - additional term on the boundary

Elements of calculus of functions in n vars: FEM

Integration: Quadrature by Trapezoidal Rule

- trapezoidal rule: \mathbf{e} triangle with vertices \mathbf{x}_1 , \mathbf{x}_2 and \mathbf{x}_3

$$\int_{\mathbf{e}} g(x, y) d\mathbf{e} \approx \frac{\text{area}(\mathbf{e})}{3} [g(\mathbf{x}_1) + g(\mathbf{x}_2) + g(\mathbf{x}_3)]$$

- more accurate rules exist (Gauss quadrature)

Elements of calculus of functions in n vars: FEM

Example in Magnetostatics

- curl

-

-

Elements of calculus of functions in n vars - FEM

Extensions

- to be decided

Elements of calculus of functions in n vars - FEM

References

- bachelor courses in calculus

Further Reading



Elements of calculus of functions in n vars - FEM

Recap

- integration by parts in 2D