

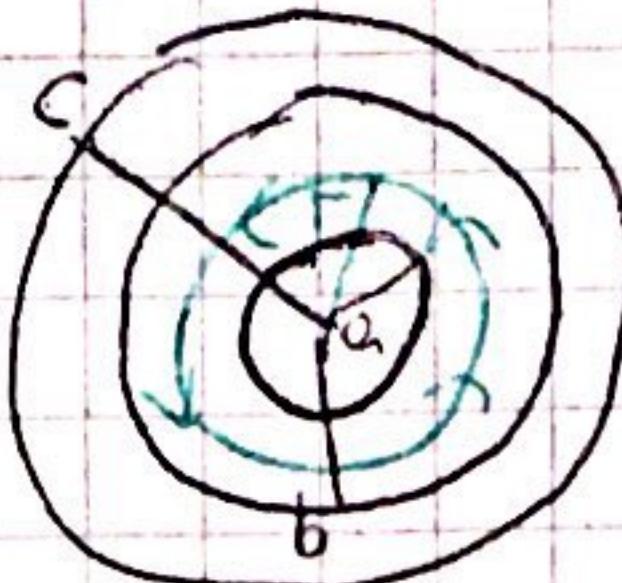
# PRACTICO 6

1)

a)

Caso 1

$$b > r > a$$



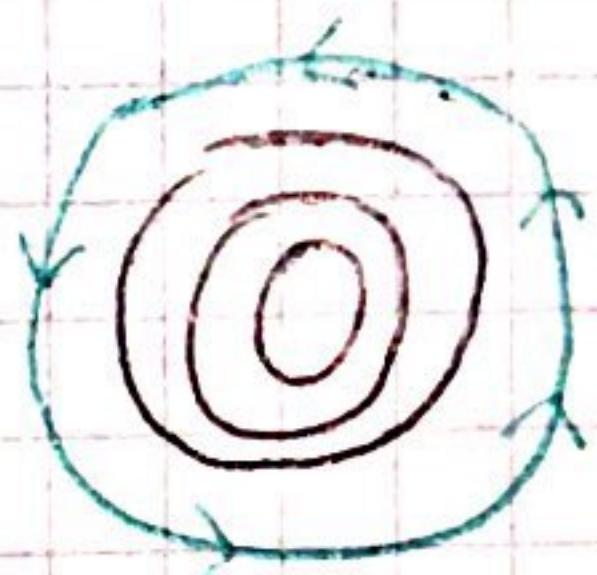
$$d\vec{L} = r d\theta \hat{\theta}, \quad i_{ex} = i$$

$$\mu_0 i_{ex} = \int \vec{B} \cdot d\vec{L} = \int_0^{2\pi} \vec{B} \cdot r d\theta \hat{\theta} = B \int_0^{2\pi} r d\theta \hat{\theta} = B r 2\pi \hat{\theta}$$

$$B_r = \frac{\mu_0 i}{2\pi r} \hat{\theta}$$

Caso 2

$$r > c > b > a$$

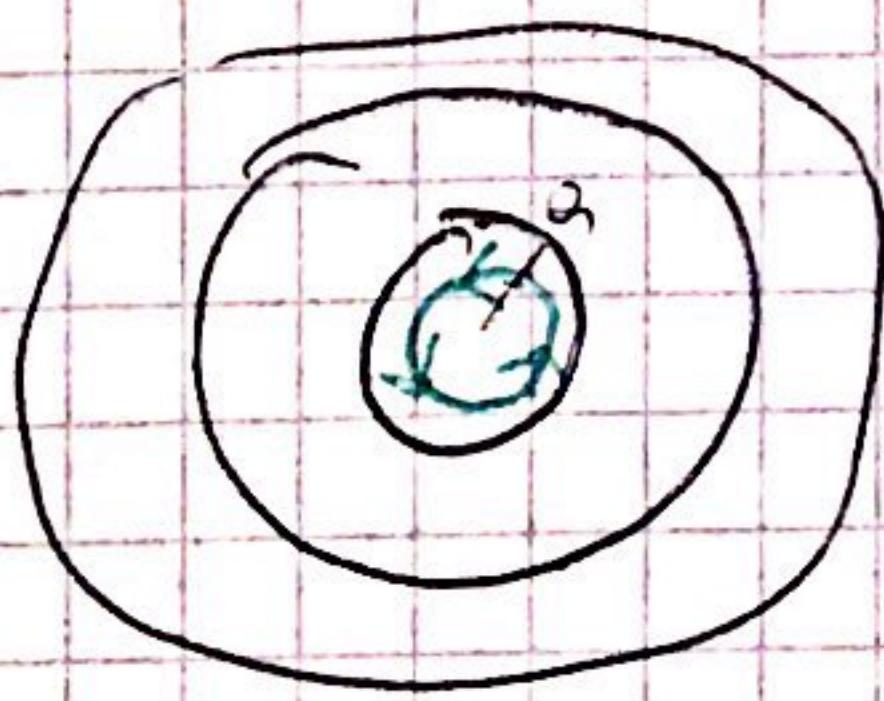


$$i_{ex} = 0 \Rightarrow \vec{B}_2 = 0$$

pues las corrientes  
son iguales pero circulan  
en diferentes sentido

Caso 3

$$r \leq a$$



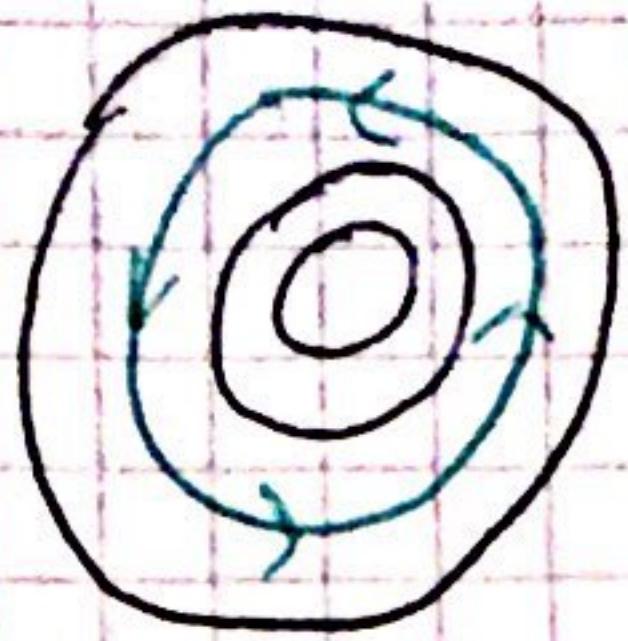
$$i_{ex} = i \cdot \frac{\pi r^2}{\pi a^2} = i \cdot \frac{r^2}{a^2}$$

Calculando  $B$  análogo al Caso 1:

$$\vec{B}_3 = \frac{\mu_0 i \cdot r^2}{2\pi a^2} \hat{\theta}$$

#### Caso 4

$$C > r > b > a$$



$$\begin{aligned} i_{\text{ext}} &= i - \frac{i}{\pi(C^2-b^2)} \cdot \pi(r^2-b^2) \\ &= i - \frac{i(r^2-b^2)}{(C^2-b^2)} \\ &= i \frac{(C^2-b^2 - (r^2-b^2))}{C^2-b^2} \\ &= i \cdot \left( \frac{C^2-r^2}{C^2-b^2} \right) \end{aligned}$$

Entonces:

$$\begin{aligned} \mu_i \left( \frac{C^2-r^2}{C^2-b^2} \right) &= \int_0^{2\pi} B d\theta = B \cdot r \cdot 2\pi \theta \\ &\Rightarrow \vec{B}_4 = \frac{\mu_i}{2\pi r} \left( \frac{C^2-r^2}{C^2-b^2} \right) \hat{\theta} \end{aligned}$$

b)  $(U_B = \int \mu_0 = \int \frac{B^2}{2\mu_0})$

$$U'_B = \underbrace{\int_a^b \int_0^{2\pi} \frac{B_1^2}{2\mu_0} r dr d\theta}_{A} + \underbrace{\int_c^\infty \int_0^{2\pi} \frac{B_2^2}{2\mu_0} r dr d\theta}_{B} + \underbrace{\int_0^a \int_0^{2\pi} \frac{B_3^2}{2\mu_0} r dr d\theta}_{C} + \underbrace{\int_b^c \int_0^{2\pi} \frac{B_4^2}{2\mu_0} r dr d\theta}_{D}$$

es la energía en una sección del cable.

$$\begin{aligned} A &= \int_a^b \int_0^{2\pi} \frac{\mu_0 i^2}{4\pi^2 r^2} \cdot \frac{1}{2\mu_0} \cdot r dr d\theta = \int_a^b \int_0^{2\pi} \frac{\mu_0 i^2}{8\pi^2 r} dr d\theta = \frac{\mu_0 i^2}{8\pi^2} \int_a^b \int_0^{2\pi} \frac{1}{r} dr d\theta \\ &= \frac{\mu_0 i^2}{8\pi^2} \cdot \ln\left(\frac{b}{a}\right) \cdot 2\pi = \ln\left(\frac{b}{a}\right) \frac{\mu_0 i^2}{4\pi} \end{aligned}$$

B  
= 0

$$\begin{aligned}
 C &= \int_0^a \int_0^{2\pi} \frac{\mu_0 i^2 r^2}{4\pi^2 a^4} \cdot \frac{1}{2\mu_0} \cdot r dr d\theta = \int_0^a \int_0^{2\pi} \frac{\mu_0 i^2}{8\pi^2 a^4} \cdot r^3 dr d\theta = \frac{\mu_0 i^2}{8\pi^2 a^4} \int_0^a r^3 dr \int_0^{2\pi} d\theta \\
 &= \frac{\mu_0 i^2}{8\pi^2 a^4} \cdot \frac{a^4}{4} \cdot 2\pi = \frac{\mu_0 i^2}{16\pi}.
 \end{aligned}$$

$$\begin{aligned}
 D &= \int_b^c \int_0^{2\pi} \frac{B_4^2}{2\mu_0} \cdot r dr d\theta = \int_b^c \int_0^{2\pi} \frac{\mu_0 i^2}{4\pi^2 r^2} \left( \frac{C^2 - r^2}{C^2 - b^2} \right)^2 \cdot \frac{1}{2\mu_0} \cdot r dr d\theta \\
 &= \int_b^c \int_0^{2\pi} \frac{\mu_0 i^2}{8\pi^2 r} \cdot \frac{1}{(C^2 - b^2)^2} \cdot (C^2 - r^2)^2 dr d\theta \\
 &= \frac{\mu_0 i^2}{8\pi^2 (C^2 - b^2)^2} \int_b^c \frac{(C^2 - r^2)^2}{r} dr \int_0^{2\pi} d\theta
 \end{aligned}$$

$$P \cdot 2\pi = \frac{\mu_0 i^2}{4\pi (C^2 - b^2)^2}$$

$$\begin{aligned}
 &= P \cdot 2\pi \cdot \int_b^c \frac{C^4}{r} dr - \int_b^c \frac{2Cr^2}{r} dr + \int_b^c \frac{r^4}{r} dr \\
 &= P \cdot 2\pi \left[ C^4 \ln \left( \frac{c}{b} \right) - 2C^2(C-b) + \frac{C^4 - b^4}{4} \right]
 \end{aligned}$$

$$\begin{aligned}
 \text{Por lo tanto} \\
 C_B &= \frac{\mu_0 i^2}{4\pi} \left[ h \left( \frac{b}{a} \right) + \frac{1}{4} + \frac{C^4 \ln \left( \frac{c}{b} \right) - 2C^2(C-b) + \frac{C^4 - b^4}{4}}{(C^2 - b^2)} \right] \\
 &= \frac{i^2}{2} \cdot \frac{\mu_0}{2\pi} \cdot S
 \end{aligned}$$

Técnica autoinducción energía

Q)

$$U_B = l \cdot U_B = l \cdot \frac{1}{2} i^2 \cdot S \frac{l}{2\pi} \stackrel{\text{A}}{\downarrow} = \frac{1}{2} \cdot L \cdot i^2$$

$\boxed{A \cdot S = L}$

②

Usando la aproximación de solenoide infinito, el campo dentro es uniforme y fuera es cero.

$$U_B = \int_0^r \int_0^{2\pi} \frac{B^2}{2\mu_0} \cdot r \, d\theta \, dr = \frac{B^2}{2\mu_0} \int_0^r r \, dr \int_0^{2\pi} d\theta$$

es uniforme

$$= \frac{B^2}{2\mu_0} \cdot \frac{r^2}{2} \cdot 2\pi = \frac{B^2 r^2 \pi}{2\mu_0} = \frac{0,4^2 \cdot 0,45^2 \cdot \pi}{2 \cdot 4\pi \cdot 10^{-7}} \cdot \frac{T \cdot m^2 \cdot A}{T \cdot m}$$

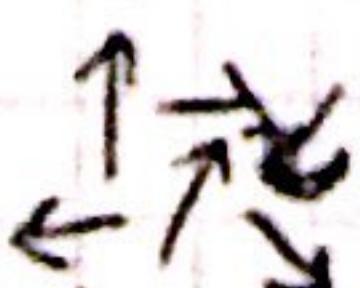
$$= \frac{0,4^2 \cdot 0,45^2}{8 \cdot 10^{-7}} \cdot T \cdot m \cdot A = 40500 \text{ Kg} \cdot \frac{m}{s^2}$$

Por lo tanto:

$$U_B = U_B \cdot l = 40500 \cdot 2,2 \text{ Kg} \cdot \frac{m}{s^2} = 89100 \text{ J}$$

③

Los materiales ferromagnéticos tienen cada átomo con momentos dipolares magnéticos pero no alineados, es decir que están así:



Cuando el material se relaciona con un campo, dichos momentos dipóles buscan alinearse. La magnetización de saturación, es aquella en la que todos los momentos logran alinearse. Entonces:

$$M_s = \mu_0 \cdot \frac{\text{cantidad de átomos}}{\text{cm}^3}$$

densidad de átomos:  $7,8 \frac{\text{g}}{\text{cm}^3}$

en 1 Kg hay  $1,08 \cdot 10^{25}$  átomos de hierro  $\Rightarrow$  cantidad de átomos  $= \frac{8,42 \cdot 10^{22}}{\text{cm}^3}$

$$M_s = (2,2 \cdot 9,3 \cdot 10^{-24}) \cdot 8,42 \cdot 10^{22} \frac{\text{T}}{\text{A}} \cdot \frac{1}{\text{cm}^3}$$

$$= 1,7 \frac{\text{T}}{\text{cm}^3} = 1,7 \cdot 10^6 \frac{\text{T}}{\text{m}^3}$$

4

$$N = 60 \text{ m}$$

$$\vec{B} = k_m \cdot \vec{B}_0$$

$$i = 5 \text{ A}$$

campo del solenoide con  
un hierro dentro

campo del solenoide solo

$$B_0 = \mu_0 \cdot i \cdot N = \mu_0 \cdot 5 \text{ A} \frac{60}{\text{m}}$$

$$\downarrow$$

campo del solenoide solo  $= \mu_0 \cdot 300 \cdot \frac{\text{A}}{\text{m}} = 4\pi \cdot 10^{-7} \cdot 300 \text{ T} \cdot \frac{\text{m}}{\text{A}} \cdot \frac{\text{A}}{\text{m}} = 0,00037 \text{ T}$

$$B = k_m \cdot B_0 = 5000 \mu_0 \cdot 300 \frac{\text{A}}{\text{m}} = 15 \cdot 10^5 \mu_0 \cdot \frac{\text{A}}{\text{m}} = 18,5 \cdot 10^{-2} \text{ T}$$

Por otro lado:

$$M = \frac{B - B_0}{\mu_0} = 1,47 \cdot 10^6 \frac{\text{A}}{\text{m}}$$

$B = B_0 + M \cdot \mu_0$ , donde fue notado un campo externo  $B_0 \Rightarrow B = M \cdot \mu_0$

NOTA  $\Rightarrow B = 1,7 \cdot 10^6 \frac{\text{T}}{\text{m}^3} \cdot 4\pi \cdot 10^{-7} \cdot \frac{\text{T} \cdot \text{m}}{\text{A}} = 2,13 \frac{\text{T}}{\text{m}^2 \text{A}}$

5

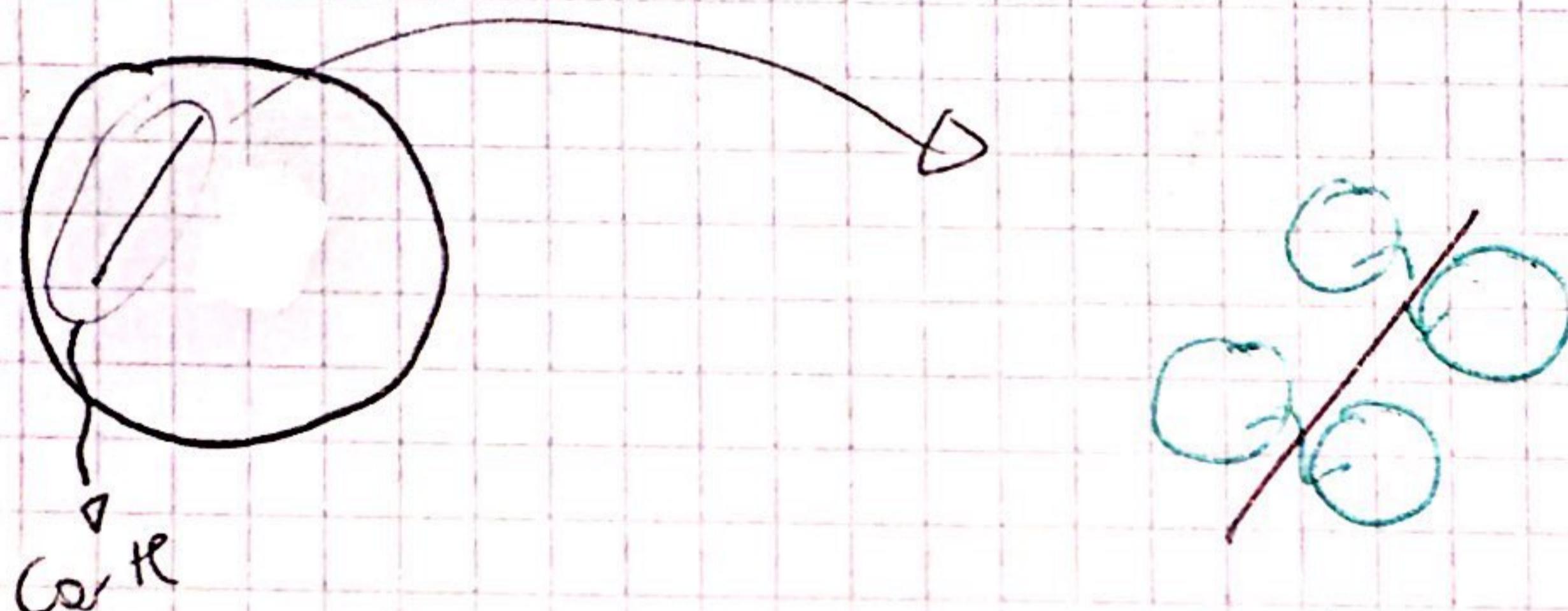
Q)  $\vec{B} = \vec{B}_0 \cdot k_m \Rightarrow \frac{\vec{B}}{\vec{B}_0} = k_m$  Pero no existe ningún  $\vec{B}_0$

Y  
el cuerpo sin el material magnético

comportamiento

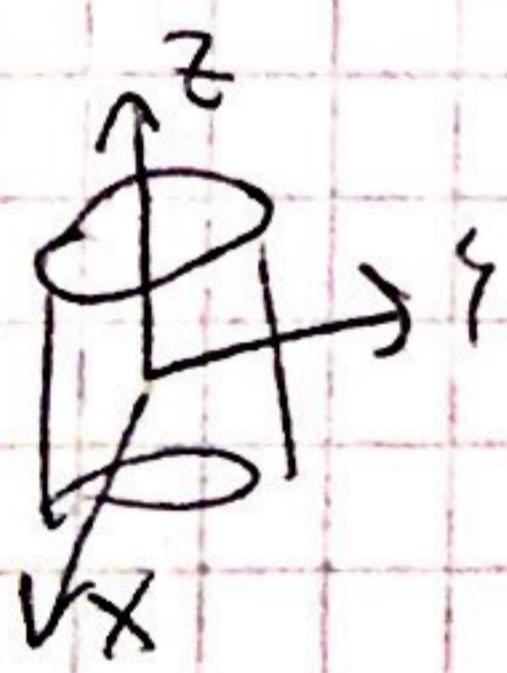
c)

Si agarro un imán y le hago cualquier corte:



Las corrientes de cada imanito se van a "cancelar" excepto en el borde.

Por lo tanto podemos pensarlo como un único imán grande donde la corriente de magnetización circula únicamente cerca en los bordes.

 $I_a$ 

$$d\vec{I} = R \cdot d\theta \cdot \hat{\theta}$$

$$\oint \vec{B}(\vec{z}) = \frac{\mu_0 \cdot i}{4\pi} \cdot \int_{\frac{\pi}{2}}^{2\pi} \frac{R \cdot d\theta \cdot \hat{\theta} \times (C\hat{z} - (z\hat{z} + R\hat{r}))}{((C-z)^2 + R^2)^{\frac{3}{2}}} d\theta$$

$$di = \frac{i}{a} \cdot dz$$

$$\vec{B}(z) = \int_{-\frac{a}{2}}^{\frac{a}{2}} \int_{0}^{2\pi} \dots$$

NOTA

Entonces:

$$\begin{aligned}
 B(\vec{z}) &= \frac{\mu_0 i}{4\pi a} \int_{-\frac{Q}{2}}^{\frac{Q}{2}} \int_0^{2\pi} dz \cdot \delta \\
 &= \frac{\mu_0 i}{4\pi a} \int_{-\frac{Q}{2}}^{\frac{Q}{2}} \int_0^{2\pi} \frac{R d\theta (c-z) r^2 + R^2 d\theta z^2}{((c-z)^2 + R^2)^{\frac{3}{2}}} dz \\
 &= \frac{\mu_0 i}{4\pi a} \int_{-\frac{Q}{2}}^{\frac{Q}{2}} \int_0^{2\pi} \frac{R(c-z) d\theta dz}{((c-z)^2 + R^2)^{\frac{3}{2}}} + \frac{\mu_0 i}{4\pi a} \int_{-\frac{Q}{2}}^{\frac{Q}{2}} \int_0^{2\pi} \frac{R^2 d\theta dz}{((c-z)^2 + R^2)^{\frac{3}{2}}} \\
 &= \frac{\mu_0 i}{4\pi a} \left[ \frac{4\pi R}{\sqrt{4c^2 - 4cq + q^2 + 4R^2}} - \frac{4\pi R}{\sqrt{4c^2 + 4cq + q^2 + 4R^2}} \right] \\
 &\quad + \frac{\mu_0 i}{4\pi a} \left[ \frac{4\pi c + 2\pi q}{\sqrt{4R^2 + (2c+q)^2}} - \frac{4\pi c - 2\pi q}{\sqrt{4R^2 + (2c-q)^2}} \right]
 \end{aligned}$$

$$\begin{cases}
 \hat{r} = R d\theta (c-z) \\
 \hat{\theta} = 0 \\
 \hat{z} = R d\theta
 \end{cases}$$