

MASTER THESIS

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### **Abstract**

We consider  $\mathbb{R}$  as well as  $\dim_{\mathcal{S}}, \dim_{\mathcal{H}} \subseteq \dim_{\mathcal{W}}$ .

# Chapter 1

## Fractal Dimensions and the Einstein Relation

In this introductory chapter, we wish to briefly expose the ingredients of the Einstein relation - the Hausdorff dimension  $\dim_{\mathcal{H}}$ , the spectral dimension  $\dim_{\mathcal{S}}$ , and the walk dimension  $\dim_{\mathcal{W}}$  - and state some of their properties.

### 1.1 Hausdorff measure and Hausdorff dimension

Although the concepts of Hausdorff measure and dimension are well-known, we give the definitions in the interest of completeness. In what follows, let  $(E, d)$  be a metric space.

**Definition 1.1** (Hausdorff outer measure). For fixed  $s \geq 0$ , any subset  $S \subseteq E$  and any  $\delta > 0$ , let

$$\mathcal{H}_{\delta}^s(S) := \inf \left\{ \sum_{i \in I} (\text{diam } U_i)^s : |I| \leq \aleph_0, S \subseteq \bigcup_{i \in I} U_i \subseteq E, \text{diam } U_i \leq \delta \right\},$$

i.e. the infimum is taken over all countable coverings of  $S$  with diameter at most  $\delta$ . The  $s$ -dimensional Hausdorff outer measure of  $S$  is now defined to be

$$\mathcal{H}^s(S) := \lim_{\delta \searrow 0} \mathcal{H}_{\delta}^s(S). \quad (1.1)$$

Observe that the limit in (1.1) exists or equals  $\infty$ , since  $\mathcal{H}_{\delta}^s(S)$  is monotonically nonincreasing in  $\delta$ , yet bounded from below by 0. Furthermore, it can be shown that  $\mathcal{H}^s$  defines a metric outer measure on  $E$ , and thus restricts to a measure on a  $\sigma$ -algebra containing the Borel  $\sigma$ -algebra  $\mathcal{B}(E)$  (cf. [Mat99, p.54ff]). By definition, the obtained measure then is the  $s$ -dimensional Hausdorff measure which we will denote by  $\mathcal{H}^s$  as well. Note that for  $\mathcal{H}^s$  to be a Radon measure, i.e. locally finite and inner regular,  $\mathcal{H}^s(E) < \infty$  is sufficient.

In the special case of  $(E, d)$  being an Euclidean space, Hausdorff measures interpolate between the usual Lebesgue measures  $\lambda^n$ : For  $s = 0$ , we have simply  $\mathcal{H}^0(S) = \#S$ , whereas for any integer  $n > 0$ , it can be shown that there exists a constant  $c_n > 0$  depending only on  $n$  such that

$\mathcal{H}^n = c_n \lambda^n$ , where the constant evaluates to the volume of the unit ball.

Since exponential functions are monotonically increasing, the Hausdorff measures' dependence on  $s$  exhibits the same behaviour for a fixed set  $S$ . At the same time, simple estimates yield that if  $\mathcal{H}^s(S)$  is finite for some  $s$ , it vanishes for all  $s' < s$ , and conversely, if  $\mathcal{H}^s(S) > 0$ , then  $\mathcal{H}^{s'}(S) = \infty$  for all  $s' > s$ . Therefore, there exists precisely one real number  $s$  where  $\mathcal{H}^s(S)$  jumps from 0 to  $\infty$  (by possibly attaining any value of  $[0, \infty]$  there). This motivates the following definition of Hausdorff dimension:

**Definition 1.2.** The Hausdorff dimension  $\dim_{\mathcal{H}}(S)$  of  $S \subseteq E$  is defined as

$$\dim_{\mathcal{H}}(S) := \inf\{s \geq 0 : \mathcal{H}^s(S) > 0\}.$$

Due to the above discussion, we have the following equalities:

$$\begin{aligned} \dim_{\mathcal{H}}(S) &= \inf\{s \geq 0 : \mathcal{H}^s(S) > 0\} = \inf\{s \geq 0 : \mathcal{H}^s(S) = \infty\} \\ &= \sup\{s \geq 0 : \mathcal{H}^s(S) = 0\} = \sup\{s \geq 0 : \mathcal{H}^s(S) < \infty\}, \end{aligned}$$

providing some alternative characterisations of the Hausdorff dimension.

We further collect some important facts. To this end, let  $S, S'$  and  $S_1, S_2, \dots$  be subsets of  $E$  as before. Then, the following properties hold (cf. [Fal07, p.32f] for a discussion in the Euclidean setting; however all arguments adapt to our more general situation without complication):

**Monotonicity.** If  $S \subseteq S'$  then  $\dim_{\mathcal{H}}(S) \leq \dim_{\mathcal{H}}(S')$ .

**Countable Stability.** For a sequence  $(S_n)_{n \geq 1}$ , we have the equality

$$\dim_{\mathcal{H}}\left(\bigcup_{n \geq 1} S_n\right) = \sup_{n \geq 1} \dim_{\mathcal{H}}(S_n).$$

**Countable Sets.** If  $|S| \leq \aleph_0$  then  $\dim_{\mathcal{H}}(S) = 0$ .

**Hölder continuous maps.** If  $(E', d')$  is another metric space and  $f : E \rightarrow E'$  is  $\alpha$ -Hölder continuous for some  $\alpha \in (0, 1]$  then  $\dim_{\mathcal{H}}(f(S)) \leq \alpha^{-1} \dim_{\mathcal{H}}(S)$ . In particular, the Hausdorff dimension is invariant under a bi-Lipschitz transformation (i.e. an invertible map  $f$  with Hölder exponent  $\alpha = 1$  for both  $f$  and  $f^{-1}$ ).

**Euclidean Case.** If  $(E, d)$  happens to be an Euclidean space (or more generally a continuously differentiable manifold) of dimension  $n$  and  $S$  is an open subset then  $\dim_{\mathcal{H}}(S) = n$ .

We conclude this section by citing Hutchinson's theorem about the Hausdorff dimension of self-similar sets which will provide us with a plethora of interesting examples. For this, we recall that a map  $F : E \rightarrow E$  on a metric space  $(E, d)$  is a strict contraction if its Lipschitz constant satisfies

$$\text{Lip}_F := \sup_{\substack{x, y \in E \\ x \neq y}} \frac{d(F(x), F(y))}{d(x, y)} < 1.$$

If the stronger condition  $d(F(x), F(y)) = \text{Lip}_F d(x, y)$  holds for all  $x, y \in E$ , we call  $F$  a similitude with contraction factor  $\text{Lip}_F$ .

**Theorem 1.3** (Hutchinson, [Hut81]). *Let  $\mathcal{S} = \{S_1, \dots, S_N\}$  be a finite set of strict contractions on the Euclidean space  $\mathbb{R}^n$ . Then there exists a unique nonempty compact set denoted by  $|\mathcal{S}|$  invariant under  $\mathcal{S}$ , i.e.*

$$|\mathcal{S}| = \bigcup_{i=1}^N S_i(|\mathcal{S}|).$$

*Furthermore, assume that  $|\mathcal{S}|$  satisfies the open set condition (OSC) which means that there exists a nonempty open set  $O \subseteq E$  with the properties  $S_i(O) \subseteq O$  and  $S_i(O) \cap S_j(O) = \emptyset$  for all  $i, j = 1, \dots, N$  with  $i \neq j$ . Also assume that the maps  $S_i$  are similitudes with contraction factor  $r_i \in (0, 1)$ . Then,  $s = \dim_{\mathcal{H}}(|\mathcal{S}|)$  is the unique solution to the equation*

$$\sum_{i=1}^N r_i^s = 1$$

*and we have  $0 < \mathcal{H}^s(|\mathcal{S}|) < \infty$ .*

While uniqueness and existence of  $|\mathcal{S}|$  are still ensured for maps on a complete metric space, the open set condition is not sufficient for statements about the Hausdorff dimension, see [Sch96] for further discussion.

## 1.2 Weyl asymptotics and spectral dimension

The idea of introducing spectral dimension is inspired by Weyl's law for the eigenvalues of the Dirichlet-Laplace operator which we will discuss here shortly:

### 1.2.1 The classical case

Given a bounded open domain  $E \subseteq \mathbb{R}^n$ , consider the Laplace operator  $\Delta$  on  $E$  acting on functions satisfying the Dirichlet boundary condition  $u \equiv 0$  on  $\partial E$ . Then, the spectrum of  $-\Delta$  consists of non-negative eigenvalues with a single limit point at  $\infty$ . Hence we can order them in a non-increasing way, counting the geometric multiplicities, as

$$0 \leq \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_n \leq \dots \text{ with } \lambda_n \nearrow \infty.$$

In this setting, it makes sense to define the eigenvalue counting function via

$$N_{-\Delta}(x) := \max\{n \in \mathbb{N} : \lambda_n \leq x\}. \quad (1.2)$$

Weyl's law now states that there is the asymptotic equivalence<sup>1</sup>

$$N_{-\Delta}(x) \sim C_n \mathcal{H}^n(E) x^{n/2}, \quad x \nearrow \infty, \quad (1.3)$$

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<sup>1</sup>We adopt the notation  $f \sim g$  for the equivalence relation given by  $\lim \frac{f}{g} = 1$ .

where the constant  $C_n$  is independent of the domain  $E$  (see [Wey11] and [Wey12] for the original publications). Motivated by (1.3), we define the spectral dimension of  $\Delta$  on  $E$  by

$$\dim_{\mathcal{S}}(E, \Delta) := \lim_{x \rightarrow \infty} \frac{\log N_{-\Delta}(x)}{\log x} \quad (1.4)$$

which yields  $n/2$  in the situation examined by Weyl's law. Note that the usual definition of  $\dim_{\mathcal{S}}$  differs by a factor of 2 (cf. [QUOTE!!]) so that  $\dim_{\mathcal{S}}(E, \Delta)$  coincides with  $\dim_{\mathcal{H}}(E) = n$ , however coming at the cost of an additional factor in the Einstein relation. Therefore, we take the liberty to deviate from the established convention in this minor aspect.

## 1.3 Diffusion processes and walk dimension

### 1.3.1 From Dirichlet forms to Markov processes

For the following definition taken from [FOT11, p.3ff], set  $H = L^2(E, \mu)$  where  $\mu$  is a  $\sigma$ -finite Borel-measure on  $E$ .

**Definition 1.4.** A map  $\mathcal{E} : \mathcal{D}(\mathcal{E}) \times \mathcal{D}(\mathcal{E}) \rightarrow \mathbb{R}$  is a Dirichlet form if it satisfies the following conditions:

- i. The domain  $\mathcal{D}(\mathcal{E}) \subseteq H$  of  $\mathcal{E}$  is a dense linear subspace.
- ii.  $\mathcal{E}$  is a symmetric, non-negative definite bilinear form.
- iii. This form is closed, that is, the inner product space  $(\mathcal{D}(\mathcal{E}), \mathcal{E}_\alpha)$  equipped with the inner product

$$\mathcal{E}_\alpha(u, v) := \mathcal{E}(u, v) + \alpha \langle u, v \rangle \quad \text{for } u, v \in \mathcal{D}(\mathcal{E}_\alpha) = \mathcal{D}(\mathcal{E}), \quad \alpha > 0,$$

is complete (and thus itself a Hilbert space).

- iv.  $\mathcal{E}$  is a Markovian form, i.e. for all  $u \in \mathcal{D}(\mathcal{E})$ ,  $v := (0 \vee u) \wedge 1 \in \mathcal{D}(\mathcal{E})$  and we have  $\mathcal{E}[v] \leq \mathcal{E}[u]$  for the quadratic form of  $\mathcal{E}$ .

We remark that the choice of  $\alpha > 0$  is irrelevant for the completeness of  $(\mathcal{D}(\mathcal{E}), \mathcal{E}_\alpha)$  since all induced norms are equivalent to each other.

We start with the following definition (cf. [MR12, Def. IV.1.13]):

**Definition 1.5.** Given a filtered probability space  $(\Omega, \mathcal{A}, \mathcal{F} = (\mathcal{F}_t)_{t \geq 0}, \mathbf{P})$  satisfying the usual conditions, an  $\mathcal{F}$ -adapted time-homogenous Markov process  $X = (X_t)_{t \geq 0}$  with state space  $E_\Delta$  is called a right process if it satisfies the strong Markov property for all  $\mathcal{F}$ -stopping times and all its trajectories are right continuous.

## 1.4 Other versions of the Einstein relation

Telcs, Mandelbrot, etc...



## Chapter 2

# Examples and Non-examples

In this chapter, we will explore the Einstein relation by examining some examples and will thusly motivate some of the more general results of the next chapter.

### 2.1 Euclidean Space

### 2.2 Sierpinski Gasket

The Sierpinski Gasket is a simple example of an iterated function fractal and can be described according to theorem 1.3 as the unique non-empty compact set  $SG \subseteq \mathbb{R}^2$  which is invariant under the three similitudes

$$S_1(x, y) = \left(\frac{x}{2}, \frac{y}{2}\right), S_2(x, y) = \left(\frac{x+1}{2}, \frac{y}{2}\right), S_3(x, y) = \left(\frac{2x+1}{4}, \frac{2y+\sqrt{3}}{4}\right),$$

see PICTURE! Since SG satisfies the (OSC), e.g. by taking the open equilateral triangle with corners  $(0, 0)$ ,  $(0, 1)$  and  $(1/2, \sqrt{3}/2)$ , we obtain both  $s = \dim_{\mathcal{H}}(SG) = \frac{\ln 3}{\ln 2}$  and  $\mathcal{H}^s(SG) \in (0, \infty)$  by a second appeal to Hutchinson's theorem.

### 2.3 Combs and inhomogenous graphs

### 2.4 Bounded metric spaces

## Chapter 3

# The Einstein Relation under Change of Measure and Metric??

Vague ideas only, no plans (yet) :/

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