## MASTER THESIS

me

 $March\ 27,\ 2018$ 

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### Abstract

We consider  $\mathbb R$  as well as  $\dim_{\mathcal S}, \dim_{\mathcal H} \subseteq \dim_{\mathcal W}.$ 

## Chapter 1

# Fractal Dimensions and the Einstein Relation

In this introductory chapter, we wish to briefly expose the ingredients of the Einstein relation - the Hausdorff dimension  $\dim_{\mathcal{H}}$ , the spectral dimension  $\dim_{\mathcal{S}}$ , and the walk dimension  $\dim_{\mathcal{W}}$  - and state some of their properties.

#### 1.1 Hausdorff measure and Hausdorff dimension

Although the concepts of Hausdorff measure and dimension are well-known, we give the definitions in the interest of completeness. In what follows, let (E, d) be a metric space.

**Definition 1.1** (Hausdorff outer measure). For fixed  $s \geq 0$ , any subset  $S \subseteq E$  and any  $\delta > 0$ , let

$$\mathcal{H}^s_{\delta}(S) := \inf \left\{ \sum_{i \in I} (\operatorname{diam} U_i)^s : |I| \le \aleph_0, S \subseteq \bigcup_{i \in I} U_i \subseteq E, \operatorname{diam} U_i \le \delta \right\},\,$$

i.e. the infimum is taken over all countable coverings of S with diameter at most  $\delta$ . The s-dimensional Hausdorff outer measure of S is now defined to be

$$\mathcal{H}^s(S) := \lim_{\delta \searrow 0} \mathcal{H}^s_{\delta}(S). \tag{1.1}$$

Observe that the limit in (1.1) exists or equals  $\infty$ , since  $\mathcal{H}^s_{\delta}(S)$  is monotonically nonincreasing in  $\delta$ , yet bounded from below by 0. Furthermore, it can be shown that  $\mathcal{H}^s$  defines a metric outer measure on E, and thus restricts to a measure on a  $\sigma$ -algebra containing the Borel  $\sigma$ -algera  $\mathscr{B}(E)$  (cf. [Mat99, p.54ff]). By definition, the obtained measure then is the s-dimensional Hausdorff measure which we will denote by  $\mathcal{H}^s$  as well. Note that for  $\mathcal{H}^s$  to be a Radon measure, i.e. locally finite and inner regular,  $\mathcal{H}^s(E) < \infty$  is sufficient.

In the special case of (E, d) being an Euclidean space, Hausdorff measures interpolate between the usual Lebesgue measures  $\lambda^n$ : For s = 0, we have simply  $\mathcal{H}^0(S) = \#S$ , whereas for any integer n > 0, it can be shown that there exists a constant  $c_n > 0$  depending only on n such that  $\mathcal{H}^n = c_n \lambda^n$ , where the constant evaluates to the volume of the unit ball.

Since exponential functions are monotonically increasing, the Hausdorff measures' dependence on s exhibits the same behaviour for a fixed set S. At the same time, simple estimates yield that if  $\mathcal{H}^s(S)$  is finite for some s, it vanishes for all s' < s, and conversely, if  $\mathcal{H}^s(S) > 0$ , then  $\mathcal{H}^{s'}(S) = \infty$  for all s' > s. Therefore, there exists precisely one real number s where  $\mathcal{H}^s(S)$  jumps from 0 to  $\infty$  (by possibly attaining any value of  $[0,\infty]$  there). This motivates the following definition of Hausdorff dimension:

**Definition 1.2.** The Hausdorff dimension  $\dim_{\mathcal{H}}(S)$  of  $S \subseteq E$  is defined as

$$\dim_{\mathcal{H}}(S) := \inf\{s \ge 0 : \mathcal{H}^s(S) > 0\}.$$

Due to the above discussion, we have the following equalities:

$$\dim_{\mathcal{H}}(S) = \inf\{s \ge 0 : \mathcal{H}^{s}(S) > 0\} = \inf\{s \ge 0 : \mathcal{H}^{s}(S) = \infty\}$$
$$= \sup\{s \ge 0 : \mathcal{H}^{s}(S) = 0\} = \sup\{s \ge 0 : \mathcal{H}^{s}(S) < \infty\},$$

providing some alternative characterisations of the Hausdorff dimension.

We further collect some important facts. To this end, let S, S' and  $S_1, S_2, ...$  be subsets of E as before. Then, the following properties hold (cf. [Fal07, p.32f] for a discussion in the Euclidean setting; however all arguments adapt to our more general situation without complication):

Monotonicity. If  $S \subseteq S'$  then  $\dim_{\mathcal{H}}(S) \leq \dim_{\mathcal{H}}(S')$ .

Countable Stability. For a sequence  $(S_n)_{n\geq 1}$ , we have the equality

$$\dim_{\mathcal{H}} \left( \bigcup_{n \geq 1} S_n \right) = \sup_{n \geq 1} \dim_{\mathcal{H}} (S_n).$$

Countable Sets. If  $|S| \leq \aleph_0$  then  $\dim_{\mathcal{H}}(S) = 0$ .

**Hölder continuous maps.** If (E', d') is another metric space and  $f: E \to E'$  is  $\alpha$ -Hölder continuous for some  $\alpha \in (0, 1]$  then  $\dim_{\mathcal{H}}(f(S)) \leq \alpha^{-1} \dim_{\mathcal{H}}(S)$ . In particular, the Hausdorff dimension is invariant under a bi-Lipschitz transformation (i.e. an invertible map f with Hölder exponent  $\alpha = 1$  for both f and  $f^{-1}$ ).

**Euclidean Case.** If (E, d) happens to be an Euclidean space (or more generally a continuously differentiable manifold) of dimension n and S is an open subset then  $\dim_{\mathcal{H}}(S) = n$ .

We conclude this section by citing Hutchinson's theorem about the Hausdorff dimension of self-similar sets which will provide us with a plethora of interesting examples. For this, we recall that a map  $F: E \to E$  on a metric space (E,d) is a strict contraction if its Lipschitz constant satisfies

$$\operatorname{Lip}_F := \sup_{\substack{x,y \in E \\ x \neq y}} \frac{d(F(x), F(y))}{d(x, y)} < 1.$$

If the stronger condition  $d(F(x), F(y)) = \operatorname{Lip}_F d(x, y)$  holds for all  $x, y \in E$ , we call F a similitude with contraction factor  $\operatorname{Lip}_F$ .

**Theorem 1.3** (Hutchinson, [Hut81]). Let  $S = \{S_1, ..., S_N\}$  be a finite set of strict contractions on the Euclidean space  $\mathbb{R}^n$ . Then there exists a unique nonempty compact set denoted by |S| invariant under S, i.e.

$$|\mathcal{S}| = \bigcup_{i=1}^{N} S_i(|\mathcal{S}|).$$

Furthermore, assume that |S| satisfies the open set condition (OSC) which means that there exists a nonempty open set  $O \subseteq E$  with the properties  $S_i(O) \subseteq O$  and  $S_i(O) \cap S_j(O) = \emptyset$  for all i, j = 1, ..., N with  $i \neq j$ . Also assume that the maps  $S_i$  are similar with contraction factor  $r_i \in (0,1)$ . Then,  $s = \dim_{\mathcal{H}}(|S|)$  is the unique solution to the equation

$$\sum_{i=1}^{N} r_i^s = 1$$

and we have  $0 < \mathcal{H}^s(|\mathcal{S}|) < \infty$ .

While uniqueness and existence of |S| are still ensured for maps on a complete metric space, the open set condition is not sufficient for statements about the Hausdorff dimension, see [Sch96] for further discussion.

## 1.2 Weyl asymptotics and spectral dimension

The idea of introducing spectral dimension is inspired by Weyl's law for the eigenvalues of the Dirichlet-Laplace operator which we will discuss here shortly.

#### 1.2.1 The classical case

Given a bounded open domain  $E \subseteq \mathbb{R}^n$ , consider the Laplace operator  $\Delta$  on E acting on functions satisfying the Dirichlet boundary condition  $u \equiv 0$  on  $\partial E$ . Then, the spectrum of  $-\Delta$  consists of non-negative eigenvalues with a single limit point at  $\infty$ . Hence we can order them in a non-increasing way, counting the geometric multiplicities, as

$$0 \le \lambda_1 \le \lambda_2 \le \dots \le \lambda_n \le \dots \text{ with } \lambda_n \nearrow \infty.$$
 (1.2)

In this setting, it makes sense to define the eigenvalue counting function via

$$N_{-\Delta}(x) := \max\{n \in \mathbb{N} : \lambda_n \le x\}. \tag{1.3}$$

Weyl's law now states that there is the asymptotic equivalence<sup>1</sup>

$$N_{-\Delta}(x) \sim C_n \mathcal{H}^n(E) x^{n/2}, \quad x \nearrow \infty,$$
 (1.4)

<sup>&</sup>lt;sup>1</sup>We adopt the notation  $f \sim g$  for the equivalence relation given by  $\lim \frac{f}{g} = 1$ .

where the constant  $C_n$  is independent of the domain E (see [Wey11] and [Wey12] for the original publications). Motivated by (1.4), we define the spectral dimension of  $\Delta$  on E by

$$\dim_{\mathcal{S}}(E, \Delta) := \lim_{x \to \infty} \frac{\log N_{-\Delta}(x)}{\log x} \tag{1.5}$$

which yields n/2 in the situation examined by Weyl's law. Note that the usual definition of  $\dim_{\mathcal{S}}$  differs by a factor of 2 (cf. [KL93],[HKK02]) so that  $\dim_{\mathcal{S}}(E,\Delta)$  normally coincides with  $\dim_{\mathcal{H}}(E) = n$ , however coming at the cost of an additional factor in the Einstein relation. Therefore, we take the liberty to deviate from the established convention in this minor aspect.

#### 1.2.2 The general case

How can we generalise the concepts just introduced to sets E which are not bounded open subsets of  $\mathbb{R}^n$ ? For this, suppose we are given a metric measure space  $(E, d, \mu)$ , where (E, d) is a separable metric space and  $\mu$  is a Radon measure on E.

Of course, the notion of an eigenvalue counting function as outlined above works for any operator A whose set of eigenvalues possesses only one limit point at  $+\infty$ . However, as we will explain in the next section, we also wish to associate a reasonably well-behaved Markov process to A. Therefore, we choose to impose the following conditions on A:

**Self-adjointness.** A is a densely defined, self-adjoint operator on the Hilbert space  $L^2(E,\mu)$ .

**Eigenvalues.** The set of eigenvalues is a subset of  $\mathbb{R}_{>0}$  and can be enumerated as in (1.2).

**Dissipativeness.** -A is dissipative. In other words, for all  $f \in \mathcal{D}(A)$  and all  $\lambda > 0$ , we have  $\|(\lambda + A)f\| \ge \lambda \|f\|$ .

The first of these assumptions guarantees that A is a closed operator, whereas the second ensures that  $\lambda + A$  is surjective for at least one  $\lambda > 0$ . Thus, the Hille-Yosida theorem states that there is a strongly continuous semigroup of contractive linear operators  $T_t$  on H such that -A is its infinitesimal generator. That is to say:

**Theorem 1.4** (Hille-Yosida). There exists a monoid morphism  $t \mapsto T_t$  from  $\mathbb{R}_{\geq 0}$  to the space of bounded linear operators on H,  $\mathbb{B}(H)$ , with  $||T_t|| \leq 1$  for all  $t \geq 0$  and  $\lim_{t \searrow 0} ||T_t f - f|| = 0$  for all  $f \in H$ , such that  $(-A, \mathcal{D}(A))$  is the operator defined via

$$(-A)f = \lim_{t \searrow 0} \frac{T_t f - f}{t},$$

where  $\mathcal{D}(A)$  is the set of elements in H for which this limit exists, if and only if -A is a densely defined, closed, dissipative operator such that for some  $\lambda > 0$ , the map  $\lambda + A$  is surjective.

## 1.3 Diffusion processes and walk dimension

#### 1.3.1 From Dirichlet forms to Markov processes

For the following definition taken from [FOT11, p.3ff], set  $H = L^2(E, \mu)$  where  $\mu$  is a  $\sigma$ -finite Borel-measure on E.

**Definition 1.5.** A map  $\mathcal{E}: \mathcal{D}(\mathcal{E}) \times \mathcal{D}(\mathcal{E}) \to \mathbb{R}$  is a Dirichlet form if it satisfies the following conditions:

- i. The domain  $\mathcal{D}(\mathcal{E}) \subseteq H$  of  $\mathcal{E}$  is a dense linear subspace.
- ii.  $\mathcal{E}$  is a symmetric, non-negative definite bilinear form.
- iii. This form is closed, that is, the inner product space  $(\mathcal{D}(\mathcal{E}), \mathcal{E}_{\alpha})$  equipped with the scalar product

$$\mathcal{E}_{\alpha}(u,v) := \mathcal{E}(u,v) + \alpha \langle u,v \rangle \text{ for } u,v \in \mathcal{D}(\mathcal{E}_{\alpha}) = \mathcal{D}(\mathcal{E}), \ \alpha > 0,$$

is complete (and thus itself a Hilbert space).

iv.  $\mathcal{E}$  is a Markovian form, i.e. for all  $u \in \mathcal{D}(\mathcal{E})$ ,  $v := (0 \lor u) \land 1 \in \mathcal{D}(\mathcal{E})$  and we have  $\mathcal{E}[v] \leq \mathcal{E}[u]$  for the quadratic form of  $\mathcal{E}$ .

We remark that the choice of  $\alpha > 0$  is irrelevant for the completeness of  $(\mathcal{D}(\mathcal{E}), \mathcal{E}_{\alpha})$  since all induced norms are equivalent to each other.

We start with the following definition (cf. [MR12, Def. IV.1.13]):

**Definition 1.6.** Given a filtered probability space  $(\Omega, \mathscr{A}, \mathscr{F} = (\mathscr{F}_t)_{t\geq 0}, \mathbf{P})$  satisfying the usual conditions, an  $\mathscr{F}$ -adapted time-homogenous Markov process  $X = (X_t)_{t\geq 0}$  with state space  $E_{\Delta}$  is called a right process if it satisfies the strong Markov property for all  $\mathscr{F}$ -stopping times and all its trajectories are right continuous.

### 1.4 Other versions of the Einstein relation

Some references to further literature, such as Telcs, Mandelbrot, etc... is planned to be included here.

## Chapter 2

## Examples and Non-examples

In this chapter, we will explore the Einstein relation by examining some examples and will thusly motivate some of the more general results of the next chapter.

## 2.1 Euclidean Space

## 2.2 Sierpinski Gasket

The Sierpinski Gasket is a simple example of an iterated function fractal and can be described according to theorem 1.3 as the unique non-empty compact set  $SG \subseteq \mathbb{R}^2$  which is invariant under the three similitudes

$$S_1(x,y) = \left(\frac{x}{2}, \frac{y}{2}\right), \ S_2(x,y) = \left(\frac{x+1}{2}, \frac{y}{2}\right), \ S_3(x,y) = \left(\frac{2x+1}{4}, \frac{2y+\sqrt{3}}{4}\right),$$

see PICTURE! Since SG satisfies the (OSC), e.g. by taking the open equilateral triangle with corners (0,0),(0,1) and  $(1/2,\sqrt{3}/2)$ , we obtain both  $s=\dim_{\mathcal{H}}(SG)=\frac{\ln 3}{\ln 2}$  and  $\mathcal{H}^s(SG)\in(0,\infty)$  by a second appeal to Hutchinson's theorem.

## 2.3 Combs and inhomogenous graphs

## 2.4 Bounded metric spaces

## Chapter 3

# The Einstein Relation under Change of Measure and Metric??

Vague ideas only, no plans (yet) :/

# Bibliography

- [Ber06] D. Bertacchi. "Asymptotic Behaviour of the Simple Random Walk on the 2-dimensional Comb". In: *Electron. J. Probab.* 11 (2006), pp. 1184–1203. DOI: 10.1214/EJP.v11-377. URL: https://doi.org/10.1214/EJP.v11-377.
- [Fal07] K. Falconer. Fractal Geometry: Mathematical Foundations and Applications. 2nd ed. Wiley, Hoboken, 2007.
- [FOT11] M. Fukushima, Y. Oshima, and M. Takeda. Dirichlet Forms and Symmetric Markov Processes. 2nd ed. De Gruyter studies in mathematics. De Gruyter, Berlin Boston, 2011.
- [HKK02] B. M. Hambly, J. Kigami, and T. Kumagai. "Multifractal formalisms for the local spectral and walk dimensions". In: Mathematical Proceedings of the Cambridge Philosophical Society 132.3 (2002), pp. 555–571.
- [Hut81] J.E. Hutchinson. "Fractals and self similarity". In: *Indiana University Mathematics Journal* 30.5 (1981), pp. 713–747.
- [Kal02] O. Kallenberg. Foundations of Modern Probability. 2nd ed. Probability and Its Applications. Springer, New York, 2002.
- [KL93] J. Kigami and M.L. Lapidus. "Weyl's problem for the spectral distribution of Laplacians on pcf self-similar fractals". In: *Communications in mathematical physics* 158.1 (1993), pp. 93–125.
- [Mat99] P. Mattila. Geometry of Sets and Measures in Euclidean Spaces: Fractals and Rectifiability. Cambridge Studies in Advanced Mathematics. Cambridge University Press, Cambridge, 1999.
- [MR12] Z.M. Ma and M. Röckner. Introduction to the Theory of (Non-Symmetric) Dirichlet Forms. Universitext. Springer, Berlin Heidelberg, 2012.
- [Sch96] A. Schief. "Self-similar sets in complete metric spaces". In: *Proceedings of the American Mathematical Society* 124.2 (1996), pp. 481–490.
- [Tel06] A. Telcs. *The Art of Random Walks*. Lecture Notes in Mathematics. Springer, Berlin Heidelberg, 2006.
- [Wey11] H. Weyl. "Ueber die asymptotische Verteilung der Eigenwerte". In: Nachrichten von der Gesellschaft der Wissenschaften zu Göttingen, Mathematisch-Physikalische Klasse 1911 (1911), pp. 110–117.

### BIBLIOGRAPHY

[Wey12] H. Weyl. "Das asymptotische Verteilungsgesetz der Eigenwerte linearer partieller Differentialgleichungen (mit einer Anwendung auf die Theorie der Hohlraumstrahlung)".
In: Mathematische Annalen 71 (1912), pp. 441–479.