MASTER THESIS

me

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Abstract

We consider $\mathbb R$ as well as $\dim_{\mathcal S}, \dim_{\mathcal H} \subseteq \dim_{\mathcal W}.$

Chapter 1

Fractal Dimensions and the Einstein Relation

In this introductory chapter, we wish to briefly expose the ingredients of the Einstein relation - the Hausdorff dimension $\dim_{\mathcal{H}}$, the spectral dimension $\dim_{\mathcal{S}}$, and the walk dimension $\dim_{\mathcal{W}}$ - and state some of their properties.

1.1 Hausdorff measure and Hausdorff dimension

Although the concepts of Hausdorff measure and dimension are well-known, we give the definitions in the interest of completeness. In what follows, let (E, d) be a metric space.

Definition 1.1 (Hausdorff outer measure). For fixed $s \geq 0$, any subset $S \subseteq E$ and any $\delta > 0$, let

$$\mathcal{H}^s_{\delta}(S) := \inf \left\{ \sum_{i \in I} (\operatorname{diam} U_i)^s : |I| \le \aleph_0, S \subseteq \bigcup_{i \in I} U_i \subseteq E, \operatorname{diam} U_i \le \delta \right\},\,$$

i.e. the infimum is taken over all countable coverings of S with diameter at most δ . The s-dimensional Hausdorff outer measure of S is now defined to be

$$\mathcal{H}^s(S) := \lim_{\delta \searrow 0} \mathcal{H}^s_{\delta}(S). \tag{1.1}$$

Observe that the limit in (1.1) exists or equals ∞ , since $\mathcal{H}^s_{\delta}(S)$ is monotonically nonincreasing in δ , yet bounded from below by 0. Furthermore, it can be shown that \mathcal{H}^s defines a metric outer measure on E, and thus restricts to a measure on a σ -algebra containing the Borel σ -algera $\mathscr{B}(E)$ (cf. [Mat99, p.54ff]). By definition, the obtained measure then is the s-dimensional Hausdorff measure which we will denote by \mathcal{H}^s as well. Note that for \mathcal{H}^s to be a Radon measure, i.e. locally finite and inner regular, $\mathcal{H}^s(E) < \infty$ is sufficient.

In the special case of (E, d) being an Euclidean space, Hausdorff measures interpolate between the usual Lebesgue measures λ^n : For s = 0, we have simply $\mathcal{H}^0(S) = \#S$, whereas for any integer n > 0, it can be shown that there exists a constant $c_n > 0$ depending only on n such that $\mathcal{H}^n = c_n \lambda^n$, where the constant evaluates to the volume of the unit ball.

Since exponential functions are monotonically increasing, the Hausdorff measures' dependence on s exhibits the same behaviour for a fixed set S. At the same time, simple estimates yield that if $\mathcal{H}^s(S)$ is finite for some s, it vanishes for all s' < s, and conversely, if $\mathcal{H}^s(S) > 0$, then $\mathcal{H}^{s'}(S) = \infty$ for all s' > s. Therefore, there exists precisely one real number s where $\mathcal{H}^s(S)$ jumps from 0 to ∞ (by possibly attaining any value of $[0,\infty]$ there). This motivates the following definition of Hausdorff dimension:

Definition 1.2. The Hausdorff dimension $\dim_{\mathcal{H}}(S)$ of $S \subseteq E$ is defined as

$$\dim_{\mathcal{H}}(S) := \inf\{s \ge 0 : \mathcal{H}^s(S) > 0\}.$$

Due to the above discussion, we have the following equalities:

$$\dim_{\mathcal{H}}(S) = \inf\{s \ge 0 : \mathcal{H}^{s}(S) > 0\} = \inf\{s \ge 0 : \mathcal{H}^{s}(S) = \infty\}$$
$$= \sup\{s \ge 0 : \mathcal{H}^{s}(S) = 0\} = \sup\{s \ge 0 : \mathcal{H}^{s}(S) < \infty\},$$

providing some alternative characterisations of the Hausdorff dimension.

We further collect some important facts. To this end, let S, S' and $S_1, S_2, ...$ be subsets of E as before. Then, the following properties hold (cf. [Fal07, p.32f] for a discussion in the Euclidean setting; however all arguments adapt to our more general situation without complication):

Monotonicity. If $S \subseteq S'$ then $\dim_{\mathcal{H}}(S) \leq \dim_{\mathcal{H}}(S')$.

Countable Stability. For a sequence $(S_n)_{n\geq 1}$, we have the equality

$$\dim_{\mathcal{H}} \left(\bigcup_{n \geq 1} S_n \right) = \sup_{n \geq 1} \dim_{\mathcal{H}} (S_n).$$

Countable Sets. If $|S| \leq \aleph_0$ then $\dim_{\mathcal{H}}(S) = 0$.

Hölder continuous maps. If (E', d') is another metric space and $f: E \to E'$ is α -Hölder continuous for some $\alpha \in (0,1]$ then $\dim_{\mathcal{H}}(f(S)) \leq \alpha^{-1} \dim_{\mathcal{H}}(S)$. In particular, the Hausdorff dimension is invariant under a bi-Lipschitz transformation (i.e. an invertible map f with Hölder exponent $\alpha = 1$ for both f and f^{-1}).

Euclidean Case. If (E, d) happens to be an Euclidean space (or more generally a continuously differentiable manifold) of dimension n and S is an open subset then $\dim_{\mathcal{H}}(S) = n$.

We conclude this section by citing Hutchinson's theorem about the Hausdorff dimension of self-similar sets which will provide us with a plethora of interesting examples. For this, we recall that a map $F: E \to E$ on a metric space (E,d) is a strict contraction if its Lipschitz constant satisfies

$$\operatorname{Lip}_F := \sup_{\substack{x,y \in E \\ x \neq y}} \frac{d(F(x), F(y))}{d(x, y)} < 1.$$

If the stronger condition $d(F(x), F(y)) = \operatorname{Lip}_F d(x, y)$ holds for all $x, y \in E$, we call F a similitude with contraction factor Lip_F .

Theorem 1.3 (Hutchinson, [Hut81]). Let $S = \{S_1, ..., S_N\}$ be a finite set of strict contractions on the Euclidean space \mathbb{R}^n . Then there exists a unique nonempty compact set denoted

by |S| invariant under S, i.e.

$$|\mathcal{S}| = \bigcup_{i=1}^{N} S_i(|\mathcal{S}|).$$

Furthermore, assume that |S| satisfies the open set condition (OSC) which means that there exists a nonempty open set $O \subseteq E$ with the properties $S_i(O) \subseteq O$ and $S_i(O) \cap S_j(O) = \emptyset$ for all i, j = 1, ..., N with $i \neq j$. Also assume that the maps S_i are similar with contraction factor $r_i \in (0,1)$. Then, $s = \dim_{\mathcal{H}}(|S|)$ is the unique solution to the equation

$$\sum_{i=1}^{N} r_i^s = 1$$

and we have $0 < \mathcal{H}^s(|\mathcal{S}|) < \infty$.

While uniqueness and existence of |S| are still ensured for maps on a complete metric space, the open set condition is not sufficient for statements about the Hausdorff dimension, see [Sch96] for further discussion.

1.2 Weyl asymptotics and spectral dimension

The idea of introducing spectral dimension is inspired by Weyl's law for the eigenvalues of the Dirichlet-Laplace operator which we will discuss here shortly.

1.2.1 The classical case

Given a bounded open domain $E \subseteq \mathbb{R}^n$, consider the Laplace operator Δ on E acting on functions satisfying the Dirichlet boundary condition $u \equiv 0$ on ∂E . Then, the spectrum of $-\Delta$ consists of non-negative eigenvalues with a single limit point at ∞ . Hence we can order them in a non-increasing way, counting the geometric multiplicities, as

$$0 \le \lambda_1 \le \lambda_2 \le \dots \le \lambda_n \le \dots \text{ with } \lambda_n \nearrow \infty.$$
 (1.2)

In this setting, it makes sense to define the eigenvalue counting function via

$$N_{-\Delta}(x) := \max\{n \in \mathbb{N} : \lambda_n \le x\}, \ x \in \mathbb{R}_{>0}. \tag{1.3}$$

Weyl's law now states that there is the asymptotic equivalence¹

$$N_{-\Delta}(x) \sim C_n \mathcal{H}^n(E) x^{n/2}, \quad x \nearrow \infty,$$
 (1.4)

where the constant C_n is independent of the domain E (see [Wey11] and [Wey12] for the original publications). Motivated by (1.4), we define the spectral dimension of $-\Delta$ on E by

$$\dim_{\mathcal{S}}(E, -\Delta) := \lim_{x \to \infty} \frac{\log N_{-\Delta}(x)}{\log x} \tag{1.5}$$

¹We adopt the notation $f \sim g$ for the equivalence relation given by $\lim \frac{f}{g} = 1$.

which yields n/2 in the situation examined by Weyl's law. Note that the usual definition of $\dim_{\mathcal{S}}$ differs by a factor of 2 (cf. [KL93],[HKK02]) so that $\dim_{\mathcal{S}}(E,\Delta)$ normally coincides with $\dim_{\mathcal{H}}(E) = n$, however coming at the cost of an additional factor in the Einstein relation. Moreover, it can be argued that the spectral dimension is rather a property of the operator $-\Delta$ than of the underlying space E. Therefore, we take the liberty to deviate from the established convention in this minor aspect.

1.2.2 The general case

How can we generalise the concepts just introduced to sets E which are not bounded open subsets of \mathbb{R}^n ? For this, suppose we are given a metric measure space (E, d, μ) , where (E, d) is a locally compact separable metric space and μ is a Radon measure on E.

Of course, the notion of an eigenvalue counting function as outlined above works for any operator A whose set of eigenvalues possesses only one limit point at $+\infty$. However, as we will explain in the next section, we also wish to associate a reasonably well-behaved Markov process with state space E to A. Therefore, we choose to impose the following conditions on A:

Assumptions 1.4. For an operator $A: L^2(E,d) \supseteq \mathcal{D}(A) \to L^2(E,d)$, we assume the following holds:

Self-adjointness. A is a densely defined, self-adjoint operator on the Hilbert space $L^2(E,\mu)$. **Eigenvalues.** The set of eigenvalues is a subset of $\mathbb{R}_{\geq 0}$ and can be enumerated as in (1.2). **Dissipativeness.** -A is dissipative. In other words, for all $f \in \mathcal{D}(A)$ and all $\lambda > 0$, we have $\|(\lambda + A)f\| \geq \lambda \|f\|$.

The first of these assumptions guarantees that A is a closed operator, whereas the second ensures that $\lambda + A$ is surjective for at least one $\lambda > 0$. Thus, the Hille-Yosida theorem states that there is a strongly continuous semigroup of contractive linear operators T_t on H such that -A is its infinitesimal generator. That is to say:

Definition 1.5. A strongly continuous semigroup $(T_t)_{t\geq 0}$ on a Hilbert space H is a monoid homomorphism $t\mapsto T_t$ from $(\mathbb{R}_{\geq 0},+)$ to the space of bounded linear operators $(\mathbb{B}(H),\cdot)$ on H (equipped with composition) satisfying for all $f\in H$ the additional property

$$\lim_{t \searrow 0} ||T_t f - f|| = 0.$$

The infinitesimal generator $(-A, \mathcal{D}(A))$ of $(T_t)_{t\geq 0}$ is defined via

$$(-A)f = \lim_{t \searrow 0} \frac{1}{t} (T_t f - f), \ f \in \mathcal{D}(A),$$

where $\mathcal{D}(A)$ is the set of elements in H for which this limit exists.

Theorem 1.6 (Hille–Yosida). An operator $(-A, \mathcal{D}(A))$ is the generator of a strongly continuous semigroup $(T_t)_{t\geq 0}$ with $||T_t|| \leq 1$ for all $t\geq 0$ if and only if -A is a densely defined, closed, dissipative operator such that for some $\lambda > 0$, the map $\lambda + A$ is surjective.

It can be shown that there is a one-to-one correspondence between contractive semigroups and operators that satisfy the Hille-Yosida theorem, that is, the semigroup in the above theorem is uniquely determined by A.

Having discussed the motivation for the assumptions 1.4, we now proceed to adapt the definitions made in (1.3) and (1.5) in a rather straightforward way:

Definition 1.7. Given an operator $(A, \mathcal{D}(A))$ on $L^2(E, \mu)$ satisfying the assumptions 1.4, its eigenvalue counting function is defined by

$$N_A(x) := \max\{n \in \mathbb{N} : \lambda_n \le x\}, \ x \in \mathbb{R}_{>0}, \tag{1.6}$$

and the spectral dimension of A by

$$\dim_{\mathcal{S}}(E, A) := \lim_{x \to \infty} \frac{\log N_A(x)}{\log x}.$$
 (1.7)

1.3 Markov processes and walk dimension

1.3.1 From Dirichlet forms to Markov processes

The theory presented here is mostly taken from [FOT11] and [MR12, ch. 4]. Set $H = L^2(E, \mu)$ where μ is a σ -finite Borel-measure on E.

Definition 1.8. A map $\mathcal{E}: \mathcal{D}(\mathcal{E}) \times \mathcal{D}(\mathcal{E}) \to \mathbb{R}$ is a Dirichlet form if it satisfies the following conditions:

- i. The domain $\mathcal{D}(\mathcal{E}) \subseteq H$ of \mathcal{E} is a dense linear subspace.
- ii. \mathcal{E} is a symmetric, non-negative definite bilinear form.
- iii. This form is closed, that is, the inner product space $(\mathcal{D}(\mathcal{E}), \mathcal{E}_{\alpha})$ equipped with the scalar product

$$\mathcal{E}_{\alpha}(u,v) := \mathcal{E}(u,v) + \alpha \langle u,v \rangle \text{ for } u,v \in \mathcal{D}(\mathcal{E}_{\alpha}) = \mathcal{D}(\mathcal{E}), \ \alpha > 0,$$

is complete (and thus itself a Hilbert space).

iv. \mathcal{E} is a Markovian form, i.e. for all $u \in \mathcal{D}(\mathcal{E})$, $v := (0 \lor u) \land 1 \in \mathcal{D}(\mathcal{E})$ and we have $\mathcal{E}[v] \leq \mathcal{E}[u]$ for the quadratic form of \mathcal{E} .

We remark that the choice of $\alpha > 0$ is irrelevant for the completeness of $(\mathcal{D}(\mathcal{E}), \mathcal{E}_{\alpha})$ since all induced norms are equivalent to each other.

Definition 1.9. A Dirichlet form $(\mathcal{E}, \mathcal{D}(\mathcal{E}))$ on $L^2(E, \mu)$ is said to be

- i. regular if it possesses a core, that is, the space $\mathcal{D}(\mathcal{E}) \cap C_c(E)$ is simultaneously dense in $\mathcal{D}(\mathcal{E})$ with respect to the \mathcal{E}_1 -norm and in $C_c(E)$ with respect to the uniform norm.
- ii. local if $\mathcal{E}(u,v) = 0$ whenever $u,v \in \mathcal{D}(\mathcal{E})$ have disjoint compact support.
- iii. strongly local if $\mathcal{E}(u,v) = 0$ whenever $u,v \in \mathcal{D}(\mathcal{E})$ have compact support and v is constant on a neighbourhood of $\mathrm{supp}(u)$.

If additionally $\mu(E) < \infty$, we say that \mathcal{E} is

- iv. conservative if $1 \in \mathcal{D}(\mathcal{E})$ and $\mathcal{E}[1] = 0$.
- v. irreducible if it is conservative and $\mathcal{E}[f]$ implies that f is constant.

We can uniquely attach a positively semidefinite operator A to a Dirchlet form (and vice versa) via the relation

$$\mathcal{E}(u,v) = \langle Au, v \rangle, \ u \in \mathcal{D}(A), v \in \mathcal{D}(\mathcal{E}). \tag{1.8}$$

In particular, if A meets the requirements of 1.4, we not only have precisely one strongly continuous contraction semigroup on H as explained by theorem 1.6, but also a unique Dirichlet form thanks to (1.8). In similar style, we would also like to attach a unique Markov process to A - or, equivalently, to the semigroup or the Dirichlet form.

To define a suitable stochastic process with values in E, we first adjoin a cemetary state * in such a way that if E is non-compact, $E_* := E \sqcup \{*\}$ is the one-point compactification of E, whereas * is supposed to be an isolated point if E is compact. Let $X = (\Omega, \mathscr{A}, (X_t)_{t \geq 0}, (\mathbf{P}_x)_{x \in E_*})$ be a stochastic process on a measurable space (Ω, \mathscr{A}) with values in E_* , where we adapt the notation that $\mathbf{P}_x[X_0 = x] = 1$ for all $x \in E_*$ and $\mathbf{P}_*[X_t = *] = 1$ for all $t \geq 0$. Note that X induces a filtration $\mathscr{F} = (\mathscr{F}_t)_{t \geq 0}$ on \mathscr{A} by

$$\mathscr{F}_t = \bigcap_{\mathbf{P} \in \mathcal{M}_1^+(\Omega, \mathscr{A})} \left(\sigma \{ X_s : 0 \le s \le t \} \right)^{\mathbf{P}}.$$

Here, \mathcal{M}_1^+ denotes the set of all probability measures on (Ω, \mathscr{A}) , $\sigma\{\cdot\}$ denotes the the σ -algebra generated by $\{\cdot\}$ and $\mathscr{B}^{\mathbf{P}}$ denotes the completion of a σ -algebra \mathscr{B} with respect to the measure \mathbf{P} . Henceforth, we will only consider stochastic processes X that satisfy the strong Markov property with respect to \mathscr{F} and are time-homogenous. Such X is called Hunt process if it additionally has right-continuous trojectories and is quasi-left-continuous, i.e. any sequence $\tau_n \nearrow \tau$ of \mathscr{F} -stopping times satisfies

$$\mathbf{P}_{\alpha} \left[\lim_{n \to \infty} X_{\tau_n} = X_{\tau}, \tau < \infty \right] = \mathbf{P}_{\alpha} [\tau < \infty]$$

for any initial distribution α . We can now easily translate Markov processes to contractive semigroups by setting

$$(T_t f)(x) := \mathbf{E}_x [f(X_t)], \ t \ge 0.$$
 (1.9)

The other direction is more involved, and the process attached to a Dirichlet form is generally non-unique. We have, however, (cf. [FOT11, theorems 7.2.1 and 7.2.2])

Theorem 1.10. Let \mathcal{E} be a regular Dirichlet form on $L^2(X,\mu)$. Then, there exists a Hunt process X on (E,d) such that the operators $T_t, t \geq 0$, from (1.9) are symmetric and \mathcal{E} is the Dirichlet form belonging to this semigroup.

Moreover, if \mathcal{E} is local, X is a diffusion process.

As hinted above, those processes are not unique: One can modify X to \tilde{X} by killing the process on a polar set and obtain the same semigroup for both. See section 7.2.2. in [FOT11] for further discussion.

1.3.2 Local walk dimension and Einstein relation

For the next definition to make sense, we need to impose some additional assumption on the metric space (E, d). We choose to demand for now that E is path connected, but will discuss other scenarios in the next chapter.

Definition 1.11. We define the quantity

$$\dim_{\mathcal{W}}(E, X, x) = \lim_{r \searrow 0} \frac{\log \mathbf{E}_x [\tau_r]}{\log r}$$

and call it the (local) walk dimension of (E, d) at $x \in E$ with respect to the Markov process $(X_t)_{t\geq 0}$. If $\dim_{\mathcal{W}}(E, X, x)$ is μ -a.e. constant on E, we shorten our notation to $\dim_{\mathcal{W}}(E, X)$.

We are now finally able to state the Einstein relation:

Definition 1.12. Let (E, d, μ) be a locally compact separable metric measure space and let $(A, \mathcal{D}(A))$ be an operator on $L^2(E, \mu)$ satisfying assumptions 1.4. Suppose $X = ((X_t)_{t\geq 0}, (\mathbf{P}_x)_{x\in E_*})$ is a Markov process associated to A via the Dirichlet form $\mathcal{E}(\cdot, \cdot) = \langle A \cdot, \cdot \rangle$. We then say that the Einstein relation with constant c holds on E with respect to A if

$$\dim_{\mathcal{H}}(E) = c \dim_{\mathcal{S}}(E, A) \dim_{\mathcal{W}}(E, X). \tag{1.10}$$

We omit mentioning the constant if c = 1.

1.4 Other versions of the Einstein relation

TO DO: A discussion of further literature, such as Telcs, Mandelbrot, etc... are planned to be included here.

Chapter 2

Examples and Non-examples

In this chapter, we will explore the Einstein relation by examining some examples and will thusly motivate some of the more general results of the next chapter.

2.1 Euclidean Space

2.2 Sierpinski Gasket

The Sierpinski Gasket is a simple example of an iterated function fractal and can be described according to theorem 1.3 as the unique non-empty compact set $SG \subseteq \mathbb{R}^2$ which is invariant under the three similitudes

$$S_1(x,y) = \left(\frac{x}{2}, \frac{y}{2}\right), \ S_2(x,y) = \left(\frac{x+1}{2}, \frac{y}{2}\right), \ S_3(x,y) = \left(\frac{2x+1}{4}, \frac{2y+\sqrt{3}}{4}\right),$$

see PICTURE! Since SG satisfies the (OSC), e.g. by taking the open equilateral triangle with corners (0,0),(0,1) and $(1/2,\sqrt{3}/2)$, we obtain both

$$s = \dim_{\mathcal{H}}(SG) = \frac{\ln 3}{\ln 2} \tag{2.1}$$

and $\mathcal{H}^s(SG) \in (0, \infty)$ by a second appeal to Hutchinson's theorem.

2.3 Combs and inhomogenous graphs

In this section, we start by considering the graph C_2 , called the two-dimensional integer comb, with vertex set \mathbb{Z}^2 and edge set given by

$$\{\{(n_1, n_2), (m_1, m_2)\} \in \mathbb{Z}^2 \times \mathbb{Z}^2 : |n_1 - m_1| = 1, n_2 = m_2 = 0 \text{ or } |n_2 - m_2| = 1, n_1 = m_1\}$$

2.4 Bounded metric spaces

Chapter 3

The Einstein Relation under Change of Measure and Metric??

Vague ideas only, no plans (yet) :/

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