

FIZIKALNA OPTIKA



INTERFERENCIJA

- kada se dva vala istovremeno nude u istoj točki prostora
- KOHERENTNI IZVORI - različita fazu dva ista vala je vremenki menjavaju

Rezultati vali: $E(t_0, x) = \vec{E}_1(t_0, x_1) + \vec{E}_2(t_0, x_2)$

$$= E_0 \cos\left(\omega t_0 - \frac{m_1 x_1}{c}\right) + E_0 \cos\left(\omega t_0 - \frac{m_2 x_2}{c}\right)$$

$$= \left\{ 2 E_0 \cos\left[\frac{\omega}{2c} (m_1 x_1 - m_2 x_2)\right] \right\} \stackrel{?}{\sim} \left[\omega t_0 - \frac{\omega}{2c} (m_1 x_1 + m_2 x_2) \right]$$

$$\Delta\phi = \frac{\omega}{c} (m_1 x_1 - m_2 x_2) = \frac{2\pi}{\lambda} (m_1 x_1 - m_2 x_2) = \frac{2\pi}{\lambda} \delta$$

$$E_{0, \text{rez}} = \left\{ 2 E_0 \cos\left[\frac{\omega}{2c} (m_1 x_1 - m_2 x_2)\right] \right\} = 2 E_0 \cos\left(\frac{\Delta\phi}{2}\right) \quad \hookrightarrow \text{razlika optičkih puteva}$$

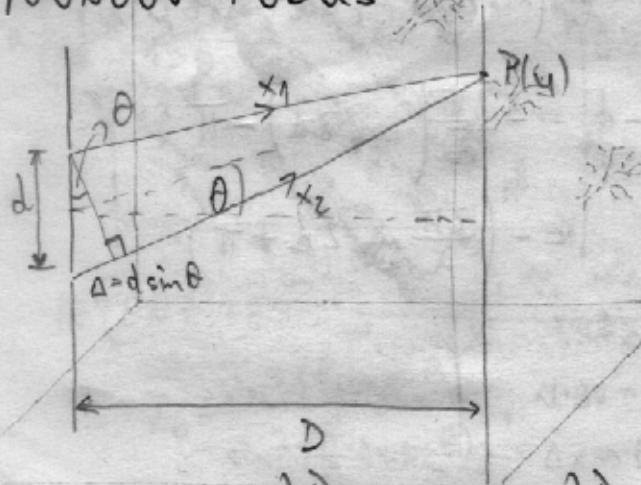
KONSTRUKTIVNA

$$\cos\left(\frac{\Delta\phi}{2}\right) = \pm 1 \Rightarrow \frac{\Delta\phi}{2} = m \cdot \pi$$

$$S_{\text{max}} = m \cdot \lambda, m_1 \neq m_2 \Rightarrow (m_2 x_2 - m_1 x_1) = \lambda_m$$

$m_1 = m_2 \Rightarrow \Delta\phi = 0$ - INTEZITET

YOUNGOV POKUS



$$y_{\text{max}} = m \frac{\lambda D}{d} \Rightarrow \Delta y_{\text{max}} = \frac{\lambda D}{d}$$

$$y_{\text{min}} = (2m+1) \frac{\lambda D}{2d}$$

DESTRUKTIVNA

$$\cos\left(\frac{\Delta\phi}{2}\right) = 0 \Rightarrow \frac{\Delta\phi}{2} = (2m+1) \frac{\pi}{2}$$

$$S_{\text{min}} = (2m+1) \frac{\lambda}{2}$$

$$\vec{E}_P = \left\{ 2 E_0 \cos\left[\frac{\omega}{2c} (m_1 x_1 - m_2 x_2)\right] \right\} \cos\left[\omega t_0 - \frac{\omega}{2c} (m_1 x_1 + m_2 x_2)\right]$$

$$m_1 = m_2 \approx 1$$

$$\delta = m \cdot |x_1 - x_2| = \Delta = d \sin \theta$$

$$d \sin \theta_{\text{max}} = m \cdot \lambda$$

$$d \sin \theta_{\text{min}} = (2m+1) \frac{\lambda}{2}$$

$$\tan \theta = \frac{y}{D} \quad \sin \theta \approx \theta$$

$$\frac{\Delta}{d} \approx \frac{y}{D}$$

INTENZITET $\left[\frac{W}{m^2}\right]$ ASITGO AUJASIGIT

$$I_0 \propto \frac{1}{2} \cdot E_0^2$$

$$I \propto (\bar{t}_{0+\epsilon})^2 = [2\bar{t}_0]^2 \left[\cos\left(\frac{\Delta\phi}{2}\right) \right]^2$$

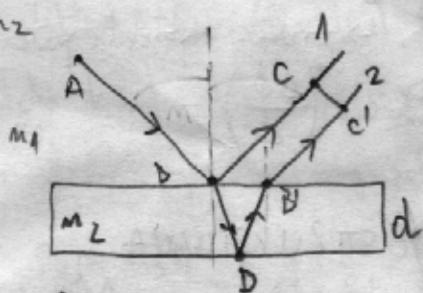
$$I_{\text{coh}} = I_0 \cos^2 \left(\frac{\Delta \phi}{2} \right) = I_0 \cos^2 \left(\frac{1}{2} \frac{2\pi}{\lambda} \delta \right)$$

$$= I_0 \cos^2 \left[\frac{\pi}{\lambda} (m_1 x_1 - m_2 x_2) \right] = I_0 \cos^2 \left(\frac{\pi}{\lambda} \Delta \right)$$

$$= I_0 \cos^2 \left(\frac{\pi}{\lambda} d \sin \theta \right)$$

TANKI LISTICÍ

$$m_1 < m_2$$



$$\bar{E}_{\text{res}} = 2E_0 \cos\left(\frac{\Delta\phi}{2}\right)$$

$$\text{MAX: } \frac{1}{2} \left(\frac{2\pi}{\lambda} m 2d + \pi \right) = m\pi$$

$$\frac{2\pi}{\lambda} m_2 d = 2m\pi - \pi$$

$$d = \frac{2m-1}{2m} \frac{\lambda}{2}$$

$$M \cap N = \frac{1}{2} \left(\frac{2\pi}{\lambda} m 2d + \pi \right) = (2m-1) \frac{\pi}{2}$$

$$\frac{2\pi}{\lambda} m 2d = (2m-2)\pi$$

$$d = \frac{m-1}{m_2} \frac{\lambda}{2}$$

$$L_1 = m_1 \overline{AB} + m_1 \overline{BC}$$

$$L_2 = m_1 \overline{AB} + m_2 \overline{BD} + m_2 \overline{DB} + m_1 \overline{BC}$$

$$\overline{BD} = \overline{DB} = d$$

$$\overline{BC} \simeq \overline{B'C'}$$

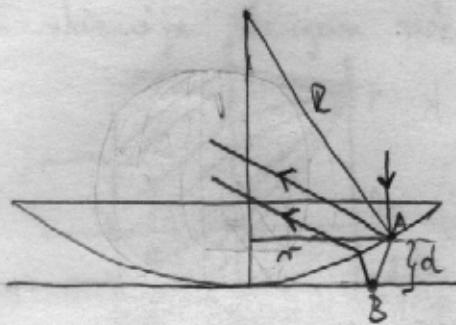
$$\phi_1 = wt - \frac{2\pi}{\lambda} (m_1 \cdot \overline{AB} + n_1 \cdot \overline{BC}) + \pi$$

$$\phi_2 = wt - \frac{2\pi i}{\lambda} (m_1 \overline{AB} + m_2 \overline{CD} + m_3 \overline{B'C'})$$

$$\Delta \phi = \phi_2 - \phi_1 = -\frac{2\pi}{\lambda} (n_2 z d) - \pi$$

$$= - \left(\frac{2\pi}{\lambda} m_2 2d + \pi \right)$$

NEWTONOVU KOLOBARU



$$|\Delta\phi| = \phi_2 - \phi_1 = \frac{2\pi}{\lambda} \cdot 2d + \pi$$

$$= \left(\frac{2\pi}{\lambda} \cdot 2d + \pi \right)$$

$$\text{MAX: } \frac{\Delta\phi}{2} = m\pi$$

$$\text{MIN: } \frac{\Delta\phi}{2} = (2m+1) \frac{\pi}{2}$$

Dva puta se reflektira
pa nema prouzine faze.

$$d_{\max} = \frac{1}{2}(2m-1) \frac{\lambda}{2}$$

$$d_{\min} = \frac{1}{2} m \lambda$$

$$r^2 + (R-d)^2 = R^2$$

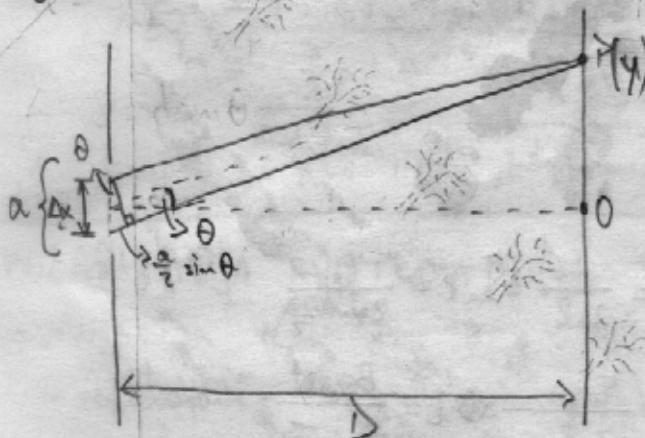
$$d = \frac{r^2}{2R}$$

$$r_{\max} = \sqrt{\frac{(2m-1)\lambda}{2}} \cdot R$$

$$r_{\min} = \sqrt{R \lambda}$$

PRIMJER Difrakcije

OGIB NA PUKOTINI



$$\Delta = \Delta x \sin \theta = \delta$$

$$\frac{\Delta\phi}{2} = \frac{\pi}{\lambda} \cdot \Delta x \sin \theta$$

$$\text{MIN: } \frac{\Delta\phi}{2} = \pm \frac{\pi}{2}$$

$$a \sin \theta_{\min} = \pm 1. \lambda$$

$$a \sin \theta_{\min} = m \lambda$$

INTENZITET:

$$I = I_0 \frac{\sin^2 \left(\frac{\pi}{\lambda} \Delta x \sin \theta \right)}{\sin^2 \left(\frac{\pi}{\lambda} \Delta x \sin \theta \right)}, \quad N \cdot \Delta x = a$$

$$\Rightarrow \sin \left(\frac{\pi}{\lambda} \Delta x \sin \theta \right) \approx \frac{\pi}{\lambda} \Delta x \sin \theta = \frac{1}{N} \frac{\pi}{\lambda} a \sin \theta$$

$$I \sim \frac{\sin^2 \left(\frac{\pi}{\lambda} a \sin \theta \right)}{\left(\frac{\pi}{\lambda} a \sin \theta \right)^2} \Rightarrow I = I_0 \frac{\sin^2 z}{z^2}$$

$$z = \frac{\pi a \sin \theta}{\lambda}$$

$$I = I_0 \left(\frac{\sin z}{z} \right)^2, \quad z = \frac{\Delta\phi}{2} = \frac{\pi a \sin \theta}{\lambda}$$

$$I = 0, \sin z = 0, z = m\pi$$

$$\Rightarrow \frac{\pi a}{\lambda} \sin \theta = m\pi$$

$$a \sin \theta = m \lambda$$

$$a = \lambda, \theta_{\min} = 30^\circ$$

INTERFERENCIJA NA DVije PUKOTINE

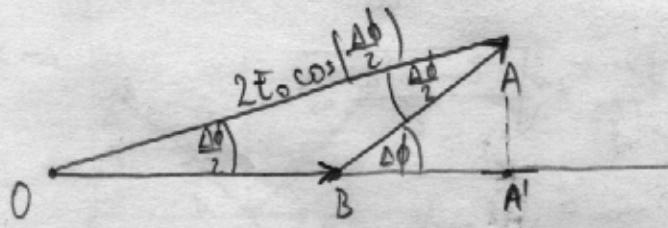
$$E_1 = E_0 \cos(\omega t - k_1 x_1) \\ = E_0 \cos \phi_1$$

$$\bar{E}_2 = E_0 \cos \phi_2$$

$$\Delta \phi = \phi_2 - \phi_1 = (\omega t - k_2 x_2) - (\omega t - k_1 x_1) = k_1 x_1 - k_2 x_2$$

$$k_i = \frac{2\pi}{\lambda_i} = \frac{2\pi}{(\frac{\lambda}{n_i})} = 2\pi n_i \frac{1}{\lambda}$$

$$\bar{E}_2 = E_0 \cos(\phi_1 + \Delta \phi)$$



$$\triangle OAB : \alpha = \frac{\Delta \phi}{2}$$

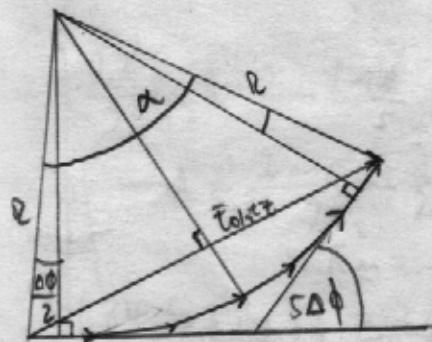
$$\sin \alpha = \frac{\overline{AA'}}{E_{0/\text{ref}} z} = \sin \frac{\Delta \phi}{2}$$

$$\triangle AA'B \sin \Delta \phi = \frac{\overline{AA'}}{E_0}$$

$$E_0 \sin \Delta \phi = E_{0/\text{ref}} \sin \frac{\Delta \phi}{2}$$

$$E_{0/\text{ref}} = 2 E_0 \cos \frac{\Delta \phi}{2}$$

Pojašnjenje na 6 izvora:



$$\alpha = 6 \Delta \phi$$

$$\sin \frac{\Delta \phi}{2} = \frac{E_0}{2 R} \Rightarrow R = \frac{E_0}{2 \sin \frac{\Delta \phi}{2}}$$

$$\sin \frac{\alpha}{2} = \frac{E_{0/\text{ref}}}{2 R} \Rightarrow R = \frac{E_{0/\text{ref}}}{2 \sin \frac{\alpha}{2}}$$

$$E_{0/\text{ref}} = E_0 \frac{\sin \frac{\alpha}{2}}{\sin \frac{\Delta \phi}{2}} = E_0 \frac{\sin \frac{6 \Delta \phi}{2}}{\sin \frac{\Delta \phi}{2}}$$

$$I(\theta) = \frac{I_0}{36} \cdot \frac{\sin^2 \frac{6 \Delta \phi}{2}}{\sin^2 \frac{\Delta \phi}{2}}$$

$$\text{MAX: } d \sin \theta_m = m \lambda$$

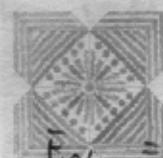
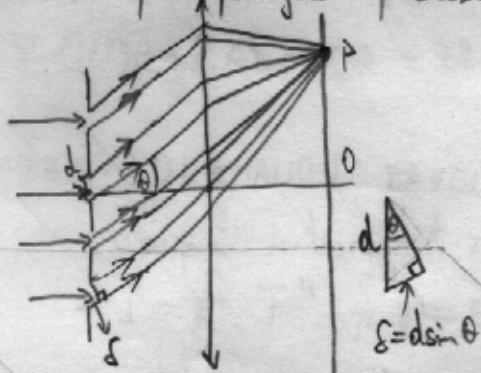
$$\text{MIN: } d \sin \theta = m \frac{\lambda}{N} \quad (m \neq kN)$$

$$\text{SPOL. MAX: } \frac{\pi}{\lambda} d \sin \theta = \pm \frac{3\pi}{N}, \pm \frac{5\pi}{N}, \dots$$

$$\text{Niz vrata } I(\theta) = \frac{I_0}{N^2} \frac{\sin^2 \frac{m \Delta \phi}{2}}{\sin^2 \frac{\Delta \phi}{2}}, \Delta \phi = \frac{2\pi}{\lambda} d \sin \theta$$

OPTIČKA REŠETKA

- generalizacija Youngova pokusa



$$\bar{E}_{\text{Omega}} = E_0 \frac{\sin \frac{N\Delta\phi}{2}}{\sin \frac{\Delta\phi}{2}}$$

$$I = I_0 \frac{\sin^2 \frac{N\Delta\phi}{2}}{\sin^2 \frac{\Delta\phi}{2}}$$

$$\text{MAX: } \frac{\Delta\phi}{2} = 0, \pm\pi, \pm 2\pi, \dots = m\pi$$

$$d \sin \theta_m = m\lambda$$

$$\text{MIN: } d \sin \theta_{\min} = m \frac{\lambda}{N}$$

INTERFERENCIJA NA ŠIROKOJ PUKOTINI

$$E = 2E_0 \cos\left(\frac{\Delta\phi}{2}\right) = 2E_0 \cos\left(\frac{\pi}{\lambda} d \sin \theta\right)$$

$$\bar{E}_0 \rightarrow E_0 \frac{\sin y}{4} ; y = \frac{1}{2} \cdot \frac{2\pi}{\lambda} d \sin \theta$$

Za dva izvora:

$$\cos\left(\frac{\Delta\phi}{2}\right) = \frac{\sin\left(\frac{2\Delta\phi}{2}\right)}{\sin\left(\frac{\Delta\phi}{2}\right)}$$

$$E = E_0 \frac{\sin y}{4} \cdot \frac{\sin\left(\frac{2\Delta\phi}{2}\right)}{\sin\left(\frac{\Delta\phi}{2}\right)}$$

$$E = E_0 \underbrace{\frac{\sin\left(\frac{\pi}{\lambda} d \sin \theta\right)}{\frac{\pi}{\lambda} d \sin \theta}}_{\text{difrakcija}} \cdot \underbrace{\frac{\sin\left(2\frac{\pi}{\lambda} d \sin \theta\right)}{\sin\left(\frac{\pi}{\lambda} d \sin \theta\right)}}_{\text{interferencija}}$$

$$I = I_0 \left(\frac{\sin\left(\frac{\pi}{\lambda} d \sin \theta\right)}{\frac{\pi}{\lambda} d \sin \theta} \right)^2 \cdot \left(\frac{\sin\left(N\frac{\pi}{\lambda} d \sin \theta\right)}{\sin\left(\frac{\pi}{\lambda} d \sin \theta\right)} \right)^2$$

d - razmak među polukotinama
a - širina pukotine

POLARIZACIJA SVIJETLOSTI

DOBIVANJE:

- refleksijom - dobita na ravni refleksije
- raspršenjem
- dvodom - daje zrake dobito polarizirane
- dikroizam -

POLARIZACIJA REFLEKSIJOM

- Brewsterov kut
- $\tan \alpha_B = n$
- lomlja i reflektiranje

DIKROIZAM - "polaroidi"

- Malusov zakon

$$I(\theta) = I(0) \cos^2 \theta$$

ZRAČENJE : ATOMI

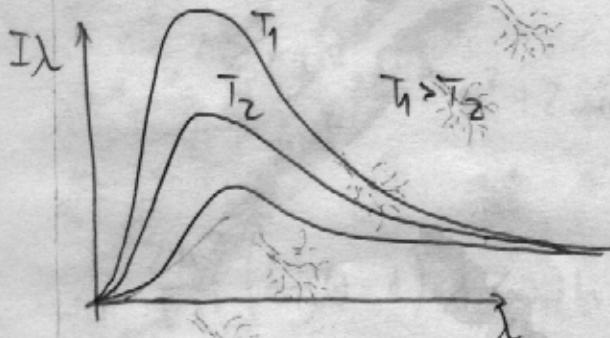
ZAKONI ZRAČENJA - ZRAČENJE ABSOLUTNO ČENOG TIJERA

STEFAN-BOLTZMANOV ZAKON

- ukupni intenzitet zračenja čenog tijela ovisan je o temp.

$$- I = \sigma \cdot T^4, \quad \sigma = 5,67 \cdot 10^{-8} \frac{W}{m^2 K^4}$$

- spektralna gustoća zračenja: $I_\lambda = \frac{dI}{d\lambda}$



$$I_\lambda = f(\lambda, T)$$

Wienov zakon ponaša: $\lambda_{max} \cdot T = k_w$, $k_w = 0,0029 \text{ m} \cdot \text{K}$;

PLANCKOV ZAKON

Rayleigh-Jeans-ova formula

- ultrafijubičasta katastrofa

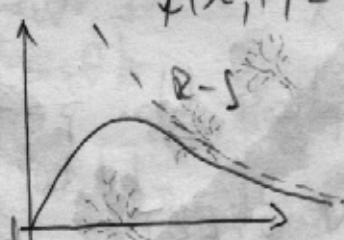
- dobije sebeckomačno velik intenzitet

$$- f(\lambda, T) = \frac{2\pi}{c} \nu^2 h T$$

$$- f(\lambda, T) = \frac{2\pi c h T}{\lambda^3}$$

$$f(\lambda, T) = \frac{A}{\lambda^5} e^{-\frac{b}{\lambda T}}$$

- loše za velike λ



PLANCK:

- energija u λ je paketičarna

$$- E_0 = h\nu = \frac{hc}{\lambda} - minimum energija, kvant energije$$

$$f(\lambda, T) = \frac{2\pi h c^2}{\lambda^5} \cdot \frac{1}{e^{hc/\lambda kT} - 1}$$

$$I = \int_0^\infty f(\lambda, T) d\lambda = \frac{2\pi^5 k^4}{15 c^2 h^3} \cdot T^4 - STEFAN-BOLTZMAN$$

PLANCKOVÝ ZVOD

$$\bar{E} = \sum_n E_n p_n$$

$$E_n = n h \nu$$

$$p_n = P(E_n) - \text{významnost} \Rightarrow p_n = c^1 \cdot e^{-\frac{E_n}{kT}}$$

je na tvar. T podstatné stojí E_n

$$\sum_n p_n = 1 \quad \sum_n p_n = c^1 \sum_n e^{-\frac{n h \nu}{kT}}, \quad \frac{h \nu}{kT} \rightarrow x$$

$$= c^1 \sum_n e^{-nx}$$

$$= c^1 (1 + e^{-x} + e^{-2x} + e^{-3x} + \dots) =$$

$$= c^1 \underbrace{(1 + e^{-x} + (e^{-x})^2 + (e^{-x})^3)}$$

$$c^1 \frac{1}{1-e^{-x}} = 1 \quad \frac{1}{1-e^{-x}}$$

$$c^1 = 1 - e^{-x}$$

$$= 1 - e^{-\frac{h \nu}{kT}}$$

$$\begin{aligned} \bar{E} &= \sum_n E_n p_n = h \nu \left(1 - e^{-\frac{h \nu}{kT}}\right) \cdot \sum_n n \cdot e^{-\frac{nh\nu}{kT}} \\ &= h \nu \left(1 - e^{-x}\right) \sum_n n e^{-nx} \\ &= h \nu \left(1 - e^{-x}\right) \left[e^{-x} + 2e^{-2x} + 3e^{-3x} + \dots \right] \\ &= h \nu \left(1 - e^{-x}\right) (-1) \frac{d}{dx} \left[e^{-x} + (e^{-x})^2 + (e^{-x})^3 + \dots \right] \\ &= -h \nu \left(1 - e^{-x}\right) \frac{d}{dx} (e^{-x}) \underbrace{\left[1 + e^{-x} + (e^{-x})^2 + (e^{-x})^3 + \dots\right]}_{\frac{1}{1-e^{-x}}} \\ &= -h \nu \left(1 - e^{-x}\right) \frac{d}{dx} \left(\frac{1}{e^{-x}-1}\right) \Rightarrow h \nu \frac{1}{e^{-x}-1} = h \nu \frac{1}{e^{x}-1} \end{aligned}$$

$$\frac{8\pi}{c^3} \cdot j^2 \cdot \left(\frac{h\nu}{e^{\frac{h\nu}{kT}} - 1} \right) = f(j, T)$$

Laijhi

$$T_{\text{max}} \quad h\nu = \frac{hc}{\lambda} \ll kT$$

Wien

$$h\nu = \frac{hc}{\lambda} \gg kT$$

$$f(\lambda, T) = \frac{2\pi c}{\lambda^5} \cdot kT$$

$$f(\lambda, T) = \frac{2\pi hc^2}{\lambda^5} \cdot e^{-\frac{hc}{\lambda kT}}$$

$$I = \int_0^\infty f(\lambda, T) d\lambda = \sigma \cdot T^4 \quad \text{STEFAN-BOLTZMANN}$$

FOTOEFEKT

(1) $\nu > \nu_g$ $I \sim \text{Intensita}$

(2) $\nu < \nu_g$, $I = 0$, $I_{\text{int.}} \rightarrow \infty$

(3) V_2 - zavistorni napon

$$\bar{E}_{\text{kinet.}} = \frac{mv_{\text{zak.}}^2}{2} = e \cdot V_2$$

(4) velika priroda ne može dajemiti
veliku izduze u zvisi intenzitet

EINSTEIN

- kvantna priroda svjetlosti

$$E_f = h\nu = \frac{hc}{\lambda} = W + E_k, W = h \cdot \nu_g, \bar{E}_{\text{kinet.}} = eV_2$$

- čestica svjetlosti foton $p = \frac{E}{c} = \frac{h\nu}{c} = \frac{h}{\lambda}$

$$p^2mc^2 = mc^2 + \bar{E}_k$$

COMPTONOV EFEKT

- raspršuje foton

$$\Delta\lambda = \lambda' - \lambda = \frac{h}{mc}(1 - \cos\theta)$$

IZVOD

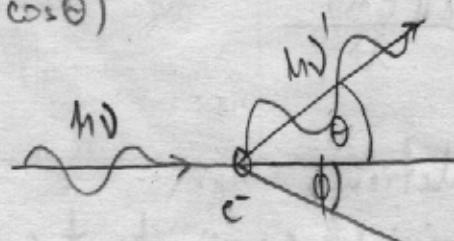
$$E_f = h\nu = \frac{hc}{\lambda}$$

$$p_f = \frac{E}{c} = \frac{h}{\lambda}$$

$$p'_f = \frac{h}{\lambda'} - \text{iduzi foton}$$

$$p'_{fM} = \frac{h}{\lambda'} \cos\theta$$

$$p'_{fN} = \frac{h}{\lambda'} \sin\theta$$



elektron mireje

$$p'_e = \gamma m_0 v$$

$$p'_{eM} = \gamma m_0 v \cos\phi$$

$$p'_{eN} = \gamma m_0 v \sin\phi$$

$$20E \frac{hc}{\lambda} + mc^2 = \frac{hc}{\lambda'} + p'mc^2$$

ZOKG:

$$\text{horizont: } \frac{h}{\lambda} = \frac{h}{\lambda'} \cos\theta + p'm \frac{\cos\theta}{m}$$

$$\text{ver: } 0 = \frac{h}{\lambda'} \sin\theta + p'm \frac{\sin\theta}{m}$$

$$\frac{h}{\lambda} - \frac{h}{\lambda'} \cos\theta = p'm \cos\phi \quad \left. \right\}$$

$$-\frac{h}{\lambda'} \sin\theta = p'm \sin\phi \quad \left. \right\}$$

$$\frac{h^2}{\lambda^2} - 2 \frac{h^2}{\lambda\lambda'} \cos\theta + \frac{h^2}{\lambda'^2} \cos^2\theta + \frac{h^2}{\lambda'^2} \sin^2\theta \\ = p'm^2 v^2 (\cos^2\phi + \sin^2\phi)$$

-OMPTON-

$$\frac{h^2}{\lambda^2} - 2 \frac{h^2}{\lambda \lambda'} \cos \theta + \frac{h^2}{\lambda'^2} = \gamma^2 m^2 c^2 \quad (*)$$

$$ZOE: \frac{hc}{\lambda} + mc^2 - \frac{hc}{\lambda'} = \gamma mc^2 \quad |^2$$

$$\frac{m^2 c^2}{\lambda^2} + \frac{h^2 c^2}{\lambda'^2} + m^2 c^4 - \frac{2 h^2 c^2}{\lambda \lambda'} + 2 m c^2 h \left(\frac{1}{\lambda} - \frac{1}{\lambda'} \right) = \gamma^2 m^2 c^4 \quad | : c^2$$

$$\frac{h^2}{\lambda^2} + \frac{h^2}{\lambda'^2} - \frac{2 h^2}{\lambda \lambda'} + m^2 c^2 + 2 m c h \left(\frac{1}{\lambda} - \frac{1}{\lambda'} \right) = \gamma^2 m^2 c^2 \quad - (*)$$

$$\frac{2 h^2}{\lambda \lambda'} (\cos \theta - 1) + 2 m c h \left(\frac{1}{\lambda} - \frac{1}{\lambda'} \right) = 0 \quad |, \lambda \lambda'$$

$$2 h^2 (\cos \theta - 1) = - 2 m c h (\lambda' - \lambda)$$

$$\lambda' - \lambda = \Delta \lambda = \frac{h}{mc} (1 - \cos \theta)$$

$$\lambda_c = \frac{h}{mc} \quad \text{Compton } \lambda \text{ electrons (m-electrons)}$$

THOMPSONOV model atoma

- puding sa šljivom
- 10^{-10} nm
- titrage "šljiva" do ravnotečnog položaja emitiraju EM zračenje

RUTHERFORD

- α -čestica bombardirali litičku zlata; gledali odijanje
- 1911. - jezgra 10^{-14} nm
- elektroni 10^{-10} nm
- planetarni model atoma

SPETROSKOPIJA vodikovog atoma, BOHR

$$\text{BALMERova formula: } - \lambda = (364,56 \text{ nm}) \cdot \frac{m^2}{m^2 - 4}, \quad m = 3, 4, 5, \dots$$

$$- \frac{1}{\lambda} = R_H \left(\frac{1}{2^2} - \frac{1}{m^2} \right) \quad m = 3, 4, 5, \dots$$

- R_H - Rydbergova konstanta

BOHR

- elektron gubi energiju
- $a_{cp} = \frac{v^2}{r} \Rightarrow P = \frac{dE}{dt} = \frac{e^2 a_{cp}^2}{6\pi \epsilon_0 c^3}$
- $v \sim 10^{11} \rightarrow$ elektron bi pas na jezgru

POSTULATI:

P-1. el. se giba do jezgre samo na određenim stazama, tada ne zravi i u stacionarnom je stanju.

P-2. Dovedjimo ustajajuće koja krit. kol. gib. jedinaka $\frac{h}{2\pi} \Rightarrow L = \tau_m m v_m = m \frac{h}{2\pi}$
 $L = \underbrace{m \tau_m}_\text{glavni broj} \rightarrow$ broj

$$\int_0^{2\pi} P_\theta d\theta = mh$$

$$P_\theta \int_0^{2\pi} d\theta = mh$$

$$P_\theta = \frac{mh}{2\pi}$$

BOMR - nast.

P-3. Kad je el. skoci iz višje u nižu stanju emitira visok energije:

$$h\nu = E_i - E_f$$

KVANTIZACIJA ENERGIJE

P-1

$$F_{\text{atom}} = \bar{F}_C \Rightarrow \frac{1}{4\pi\epsilon_0} \frac{e^2}{r_m^2} = \frac{m v_m^2}{r_m}$$

P-2

$$v_m = \frac{m\bar{v}}{mr_m} \Rightarrow r_m = \frac{mv}{m\bar{v}}$$

P-1; P-2

$$\frac{1}{4\pi\epsilon_0} \frac{e^2}{r_m^2} = \frac{mv_m^2}{m^2 \frac{v^2}{r_m^2}}$$

$$r_m = \frac{4\pi\epsilon_0 m^2 \bar{v}^2}{me^2} \quad m = 1, 2, 3, \dots$$

$$v_m = \frac{1}{m} \frac{e^2}{4\pi\epsilon_0 \bar{v}} = \frac{v_1}{m}$$

$$E_m = \frac{1}{2} m v_m^2 = \frac{1}{m} \frac{e^2}{4\pi\epsilon_0 \bar{v}}$$

$$E_m = -\frac{1}{m^2} \cdot \frac{me^4}{32\pi^2 \epsilon_0^2 \bar{v}^2}$$

$$E_1 = -13,58 \text{ eV}$$

energija ionizacije/vezanja

$$P-3 \quad h\nu = E_i - E_f$$

$E_i > E_f$ emisija fotona

$$E_i = \frac{-E_i}{m_i^2} \quad E_f = \frac{-E_f}{m_f^2}$$

$$\nu = \frac{E_i}{h} - \frac{E_f}{h} = \frac{E_i}{h} \left(\frac{1}{m_f^2} - \frac{1}{m_i^2} \right)$$

$$\frac{1}{\lambda} = \frac{E_i}{hc} \left(\frac{1}{m_f^2} - \frac{1}{m_i^2} \right)$$

$$\frac{1}{\lambda} = \frac{me^4}{(4\pi\epsilon_0)^2 \cdot 4\pi c^3} \left(\frac{1}{m_f^2} - \frac{1}{m_i^2} \right)$$

$$\frac{1}{\lambda} = R_H \left(\frac{1}{m_f^2} - \frac{1}{m_i^2} \right)$$

Balmer $m_f=2$

Lymar $m_f=1$

Packen $m_f=3$

DE BROGLIEVA HIPOTEZA

- pretpostavlja da i obične čestice imaju vlastin karakter

$$E = h\nu = \frac{h}{2\pi} 2\pi\nu = \hbar w = \frac{hc}{\lambda} \Rightarrow$$

$$P = \frac{E}{c} = \frac{h}{\lambda} = \frac{h}{2\pi(\frac{\lambda}{2\pi})} = \frac{h}{\lambda/2\pi} = \hbar k$$

$$P = \frac{h}{\lambda_B} = \hbar k_B = mv$$

DE BROGLIEEV VALI

$$\lambda_m = 2r_m \pi = m\lambda = \frac{mh}{P_m} = \frac{mh}{mv_m}$$

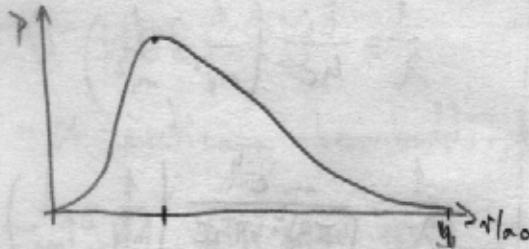
$$2\pi r_m m v_m = mh$$

$$r_m m v_m = m \frac{h}{2\pi} = m\bar{v}$$

KVANTNA MEHANIKA

SCHRÖDINGEROVA JDŽB.

- vjerojatnost nalazeња elektrona na nekoј energiji i vjerojatnost prijelaza



- glavni kvantni broj n ($1, 2, 3, \dots$)
- orbitalni kvantni broj l (L) ($0, 1, 2, \dots, n-1$)
- magnetski kvantni broj m_l (l_z) ($-l, -l+1, \dots, l-1, l$)
 - međudjelovanje el. i mag. polja
- spin $S_z = \text{magnit} \quad (-\frac{1}{2}, +\frac{1}{2})$

$l=0 \quad S \quad l=1 \quad P \quad l=2 \quad D \quad l=3 \quad F$

RELACIJE NEODREĐENOSTI

Heisenberg - 1927.

- ne možemo istovremeno precizno odrediti položaj i brzinu elektrona
- gledaju fotonom koji ima kol. gib $\tau = \frac{\hbar}{\lambda}$

$$\Delta x \cdot \Delta p_x \geq \frac{\hbar}{2}$$

$$\Delta p_x \geq \frac{\hbar}{2\Delta x}$$

PAULI - dva elektrona u atomu ne mogu imati sva četiri kvantna broja ista.

LASERI

$$E_1 + h\nu = E_2$$

$$\text{makro } 10^{-3} - 10^{-8} \text{ s}$$

$$E_2 = h\nu + E_1 \text{ ili stimulirano}$$

$$E_2 (+h\nu) = 2h\nu + E_1$$

- foton iste energije nestručnojelyje se elektronom i potice ga da "padne" - emitira isti blizu foton
- lančana reakcija