

### 3.1.2 Poisson 公式

$$\text{对 } \begin{cases} \frac{\partial u}{\partial t} - \alpha^2 \frac{\partial^2 u}{\partial x^2} = f(x, t) & ① \\ u(x, 0) = \varphi(x) & ② \end{cases} \text{ , 未提供 } \psi(x) \text{ 边界}$$

对 ① Fourier 变换, 有  $\left(\frac{d^m u}{dx^m}\right)^{\wedge} = (j\omega)^m \hat{u}$

$$\frac{\partial \hat{u}}{\partial t} - \alpha^2 (j\omega)^2 \hat{u} = \hat{f}(\omega, t)$$

$$\begin{cases} ① \frac{\partial \hat{u}}{\partial t} + \alpha^2 \lambda^2 \hat{u} = \hat{f}(\lambda, t) \\ ② \hat{u}(\lambda, 0) = \hat{\varphi}(\lambda) \end{cases}$$

$$\text{其中 } \hat{u}(\lambda, t) \text{ 为解 } u(x, t) \text{ 关于 } x \text{ 的 Fourier 变式}$$

对一阶偏微分方程, 特征方程为  $\lambda + \alpha^2 \lambda^2 = 0$

通解取  $\gamma = -\alpha^2 \lambda^2$ , 则  $\hat{u} = C \cdot e^{-\alpha^2 \lambda^2 t}$ , 有

$t=0$  时, 得:  $C = \hat{\varphi}(\lambda)$  代入有:

$$-\alpha^2 \lambda^2 \hat{\varphi}(\lambda) e^{-\alpha^2 \lambda^2 t} + \alpha^2 \lambda^2 \hat{\varphi}(\lambda) e^{-\alpha^2 \lambda^2 t} = \hat{f}(\lambda, t)$$

$$\alpha^2 \lambda^2 e^{-\alpha^2 \lambda^2 t} (\hat{\varphi}(\lambda) - \hat{\varphi}_t(\lambda)) = \int_{-\infty}^{+\infty} \hat{f}(\lambda, t) e^{-j\lambda t} dt$$

取一特解为  $u(x, t)$  关于  $x$  的 Fourier 变式, 为:

$$u = \int_0^t \hat{f}(\lambda, \tau) e^{-\alpha^2 \lambda^2 (t-\tau)} d\tau, \text{ 则 } \frac{\partial u}{\partial t} + \alpha^2 \lambda^2 u$$

$$\text{由 } \frac{du}{dt} + p(t) \cdot u = Q(t) = \hat{f}(\lambda, t) + \alpha^2 \lambda^2 \int_0^t \hat{f}(\lambda, \tau) e^{-\alpha^2 \lambda^2 (t-\tau)} d\tau$$

计算积分因子:  $\mu(t) = e^{\int p(t) dt} = e^{\int \alpha^2 \lambda^2 dt}$  两边同乘  $\mu(t)$  有

同乘有:  $e^{\alpha^2 \lambda^2 t} \frac{d\hat{u}}{dt} + e^{\alpha^2 \lambda^2 t} \cdot \alpha^2 \lambda^2 \hat{u} = e^{\alpha^2 \lambda^2 t} \hat{f}(\lambda, t)$

一阶偏微分方程特解求法

乘因子后, 可合为一项

$$\frac{d}{dt} (\hat{u} e^{\alpha^2 \lambda^2 t}) = e^{\alpha^2 \lambda^2 t} \hat{f}(\lambda, t), \text{ 两边积分有}$$

$$\hat{u} e^{\alpha^2 \lambda^2 t} = \int_0^t e^{\alpha^2 \lambda^2 \tau} \hat{f}(\lambda, \tau) d\tau$$

$$\text{故 } \hat{u} = \int_0^t \hat{f}(\lambda, \tau) e^{-\alpha^2 \lambda^2 (t-\tau)} d\tau \text{ 为一特解} \quad \star$$

故解为

$$\hat{u} = \hat{\varphi}(\lambda) e^{-\alpha^2 \lambda^2 t} + \int_0^t \hat{f}(\lambda, \tau) e^{-\alpha^2 \lambda^2 (t-\tau)} d\tau$$



$$\text{由于 } \hat{u}(x, t) = \hat{\varphi} e^{-a^2 \lambda^2 t} + \int_0^t (f(\lambda, \tau) e^{-a^2 \lambda^2 (t-\tau)}) d\tau$$

则反变换有:

$$u(x, t) = (\hat{\varphi} e^{-a^2 \lambda^2 t})^V + \int_0^t (f(\lambda, \tau) e^{-a^2 \lambda^2 (t-\tau)})^V d\tau$$

$$\text{考虑到 } (e^{-A x^2})^\wedge = \frac{1}{\sqrt{2A}} e^{-\frac{\lambda^2}{4A}}, \text{ 取 } \frac{1}{4A} = a^2 t \rightarrow A = \frac{1}{4a^2 t}$$

$$\text{则 } (e^{-\frac{x^2}{4a^2 t}})^\wedge = \sqrt{2a^2 t} e^{-a^2 \lambda^2 t} \rightarrow (e^{-a^2 \lambda^2 t})^V = \frac{1}{a \sqrt{2t}} e^{-\frac{x^2}{4a^2 t}}$$

取为  $g(x, t)$

则有:

$$(g e^{-a^2 \lambda^2 t})^V = (\hat{\varphi} \hat{g})^V \stackrel{\text{卷积性质}}{=} ((\varphi * g)^\wedge)^V = \varphi * g$$

$$= \frac{1}{a \sqrt{2t}} \int_{-\infty}^{+\infty} \varphi(\xi) e^{-\frac{(x-\xi)^2}{4a^2 t}} d\xi \quad \text{①} \quad \frac{1}{a \sqrt{2t}} \leftarrow K(x-\xi, t)$$

同理有:

$$(f(\lambda, t) e^{-a^2 \lambda^2 (t-\tau)})^V, \text{ 取 } A = \frac{1}{4a^2 (t-\tau)} \text{ 有:}$$

$$= \frac{1}{a \sqrt{2(t-\tau)}} \int_{-\infty}^{+\infty} f(\xi, \tau) e^{-\frac{(x-\xi)^2}{4a^2 (t-\tau)}} d\xi \quad \text{②} \quad \Rightarrow \text{令其} = K(x-\xi, t-\tau)$$

热核函数

代入  $u(x, t)$  表达式中, 则得:

$$u(x, t) = \frac{1}{a \sqrt{2t}} \int_{-\infty}^{+\infty} K(x-\xi, t) \varphi(\xi) d\xi + \int_0^t \int_{-\infty}^{+\infty} K(x-\xi, t-\tau) f(\xi, \tau) d\xi d\tau$$