

$$\text{已求得: } \begin{cases} G_r = \frac{A}{r^2} + B(1+2\ln r) + 2C \\ G_\theta = -\frac{A}{r^2} + B(3+2\ln r) + 2C \\ \tau_{r\theta} = \tau_{\theta r} = 0 \end{cases}$$

法一:
若天书: G_r, G_θ 当 $r \rightarrow 0$ 时
→ ∞ , 不符合事实
故 $B=0$.

对于平面应力问题, 有物理方程: (左侧为方程)

也可通过下列
积分说明.

$$\begin{aligned} \frac{\partial u_r}{\partial r} &= \varepsilon_r = \frac{1}{E} (G_r - \mu G_\theta) \\ &= \frac{1}{E} \left(\frac{A}{r^2} + B(1+2\ln r) + 2C + \mu \frac{A}{r^2} - \mu B(3+2\ln r) - 2C\mu \right) \\ &= \frac{1}{E} \left[(1+\mu) \frac{A}{r^2} + (1-3\mu) B + 2(1-\mu) B \ln r + 2C(1-\mu) \right] \quad (1) \end{aligned}$$

$$\begin{aligned} \frac{1}{r} \frac{\partial u_\theta}{\partial \theta} + \frac{u_r}{r} &= \varepsilon_\theta = \frac{1}{E} [G_\theta - \mu G_r] \\ &= \frac{1}{E} \left[-\frac{A}{r^2} + B(3+2\ln r) + 2C - \mu \frac{A}{r^2} - \mu B(1+2\ln r) - 2C\mu \right] \\ &= \frac{1}{E} \left[-(1+\mu) \frac{A}{r^2} + B(3-\mu) + 2(1-\mu) B \ln r + 2C(1-\mu) \right] \quad (2) \end{aligned}$$

$$\gamma_{r\theta} = \frac{1}{r} \frac{\partial u_r}{\partial \theta} + \frac{\partial u_\theta}{\partial r} - \frac{u_\theta}{r} = 0 \quad (3) \quad \int \ln x = x \ln x - x + C$$

则由①积分得

$$\begin{aligned} u_r &= \frac{1}{E} \left[(1+\mu) \cdot \frac{A}{r} + (1-3\mu) Br + 2(1-\mu) B(r \ln r - r) + 2C(1-\mu)r \right] \\ &= \frac{1}{E} \left[(1+\mu) \cdot \frac{A}{r} + 2(1-\mu) Br (\ln r - 1) + (1-3\mu) Br + 2(1-\mu) Cr \right] + f(\theta) \end{aligned}$$

代入②:

$$\begin{aligned} \frac{\partial u_\theta}{\partial \theta} &= \frac{r}{E} \left[-\cancel{(1+\mu)} \frac{A}{r^2} + B(3-\mu) + 2\cancel{(1-\mu)} B \ln r + 2C\cancel{(1-\mu)} \right] \\ &\quad - \frac{r}{E} \left[-\cancel{(1+\mu)} \frac{A}{r^2} + 2\cancel{(1-\mu)} B (\ln r - 1) + (1-3\mu) B + 2\cancel{(1-\mu)} C \right] - f(\theta) \\ &= \frac{r}{E} \left[\cancel{(3-\mu)} B + 2\cancel{(1-\mu)} B - \cancel{(1-3\mu)} B \right] - f(\theta) \\ &= \frac{4Br}{E} - f(\theta) \end{aligned}$$

$3-\mu+2-\mu-1+3\mu=4$

故有: $u_\theta = \frac{4Br\theta}{E} - \int f(\theta) d\theta + f_1(r)$ (偏积分)

代入: $\gamma_{r\theta} = \frac{1}{r} \frac{\partial u_r}{\partial \theta} + \frac{\partial u_\theta}{\partial r} - \frac{u_\theta}{r} = 0$

$$\text{则: } \frac{1}{r} f'(\theta) + f_1'(r) - \frac{4B\theta}{E} + \frac{1}{r} \int f(\theta) d\theta - \frac{f_1(r)}{r} = 0$$

$$\sigma_{\theta\theta} = \frac{E}{1-\nu} \left(\frac{u}{r} + \nu \frac{1}{r} \frac{du}{dr} \right)$$

得:
$$-rf_1'(r) + f_1(r) = \int f(\theta) d\theta + f'(\theta) - \frac{4Br\theta}{E}$$

此时: 左边仅 r 函数, 右边仅 θ 函数 故仅能等于常数 F

$$\begin{cases} f_1(r) - rf_1'(r) = F & (1) \end{cases}$$

$$\frac{df(\theta)}{d\theta} + \int f(\theta) d\theta - \frac{4Br\theta}{E} = F \quad (2)$$

我们得 $B=0$

对 (1) 式积分: 有:

$$\frac{df_1}{f_1(r) - F} = \frac{dr}{r} \quad \text{即: } \ln(f_1(r) - F) = \ln Cr$$

~~此时若 $B \neq 0$, 则 θ 是变动的, 故该式非常数, 不满足位移单值条件, 因此有: $B=0$~~

$$\therefore f_1(r) = F + Cr = Hr + F, \text{ 其中 } H \text{ 为常数}$$

对于 (2) 式, 代入 $B=0$, 有:

$$\frac{df(\theta)}{d\theta} + \int f(\theta) d\theta = F$$

$$\text{则: } \frac{d^2 f(\theta)}{d\theta^2} + f(\theta) = 0 \xrightarrow{\text{通解}} f(\theta) = I \cos \theta + K \sin \theta$$

得到位移表达式:

$$\begin{cases} u_\theta = \frac{4Br\theta}{E} - \int (I \cos \theta + K \sin \theta) d\theta + Hr + F \end{cases}$$

$$u_r = \frac{1}{E} \left[(H/\mu) \cdot \frac{A}{r} + 2(H/\mu)Br \ln r - (1-3\mu)Br + (1-2\mu)Cr \right] + I \cos \theta + K \sin \theta$$

其中: 代入 $B=0$, 而有轴对称应力情况下:

有 $u_\theta = 0$, 则此时 $\Rightarrow I = K = H = F = 0$,

则: 有位移表达式:

例: 有位移表达式:

$$\begin{cases} u_r = \frac{1}{E} \left[-(1+\mu) \frac{A}{r} + 2(1-\mu) Cr \right] \\ u_\theta = 0 \end{cases}$$