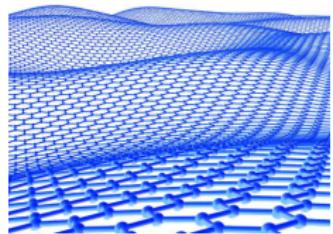
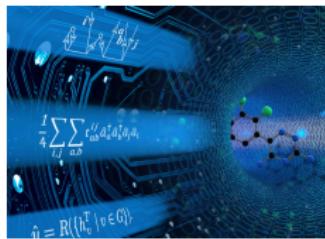


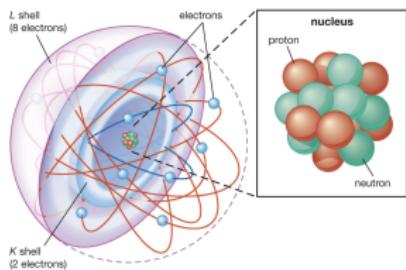
# Mathematical foundations of modern quantum simulations



Fabian M. Faulstich

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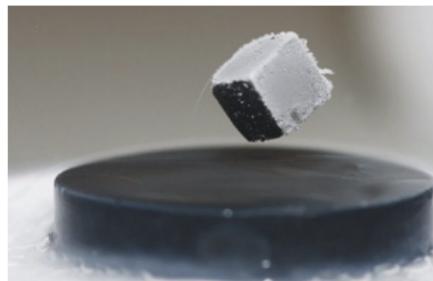
# What are quantum simulations?



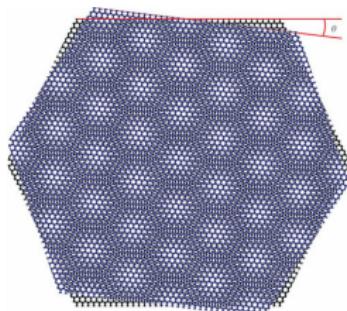
Stability of matter



Color of gold

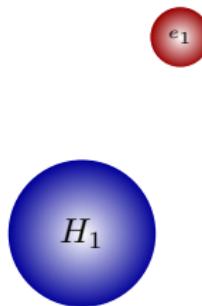


Superconductivity

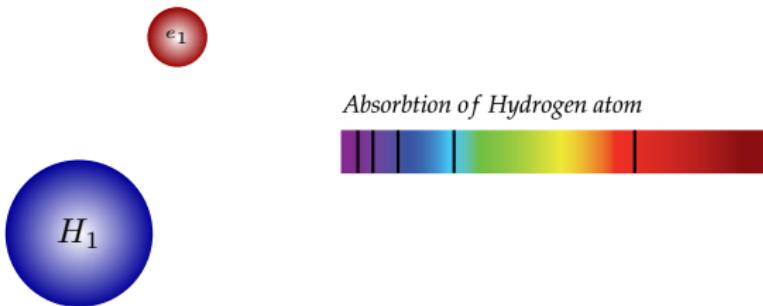


Solid state physics

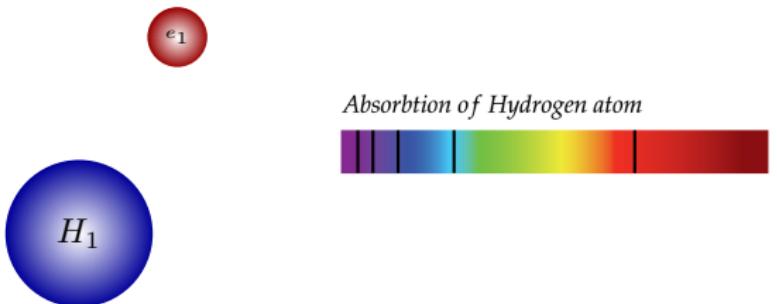
## Towards a mathematical formulation



# Towards a mathematical formulation



# Towards a mathematical formulation

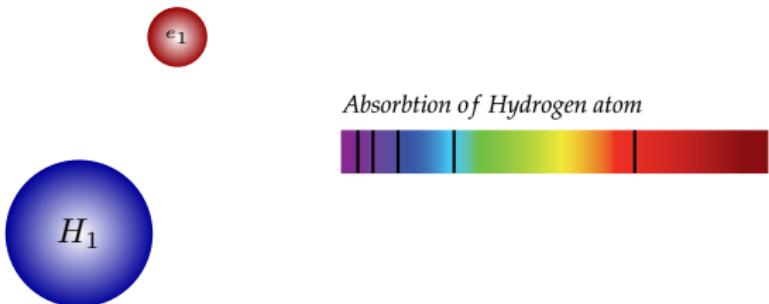


Two energetic contributions:

1. Kinetic energy:  $h_{\text{kin}} = -\frac{1}{2}\Delta$
2. Coulomb energy:  $V_{\text{coul}} = -\frac{1}{|r - R|}$



# Towards a mathematical formulation



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2. Coulomb energy:  $V_{\text{coul}} = -\frac{1}{|r - R|}$



Birth of quantum chemistry<sup>a</sup>:

$$\Delta \psi + \frac{2m}{K^2} \left( E + \frac{e^2}{r} \right) \psi = 0$$

<sup>a</sup>Schrödinger, Annalen der physik (1926)

# Schrödinger equation

The goal is to solve

$$H\Psi = E\Psi,$$

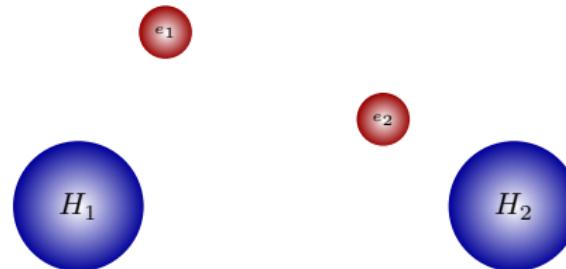
where

$$H = -\frac{1}{2} \sum_{i=1}^N \Delta_{r_i} - \sum_{i=1}^N \sum_{j=1}^{N_{\text{nuc}}} \frac{Z_j}{|r_i - R_j|} + \sum_{i=1}^N \sum_{j>i}^N \frac{1}{|r_i - r_j|}$$

and  $\Psi$  is a function defined on  $\mathbb{R}^{3N}$  for  $N$  electrons.

$r_i$  position of the  $i$ th electron

$R_j$ ,  $Z_j$  position (fixed) and charge of the  $j$ th nucleus



## What is the problem?

Discretization scales *exponentially* in the number of electrons  $N$

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<sup>1</sup>The estimated number of particles in the universe is  $\sim 10^{80}$

## What is the problem?

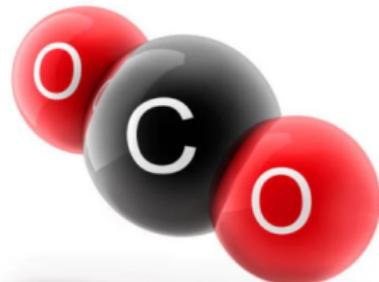
Discretization scales *exponentially* in the number of electrons  $N$

Example: carbon dioxide ( $\text{CO}_2$ )

Oxygen has 8 electrons

Carbon has 6 electrons

In total:  $N = 2 * 8 + 6 = 22$



---

<sup>1</sup>The estimated number of particles in the universe is  $\sim 10^{80}$

## What is the problem?

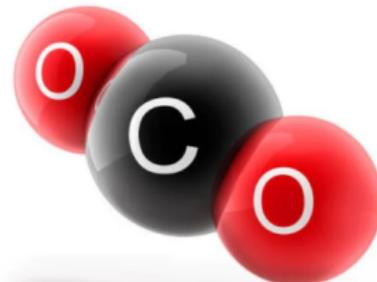
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Naïve grid with 10 points per axis yields  $10^{3 \cdot 22} = 10^{66}$  grid points<sup>1</sup>

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## What is the problem?

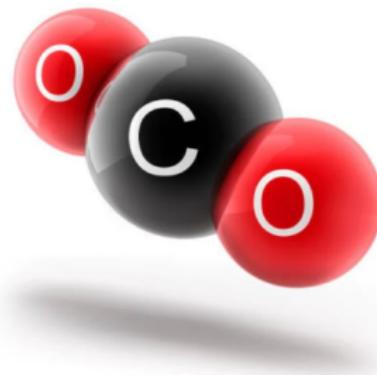
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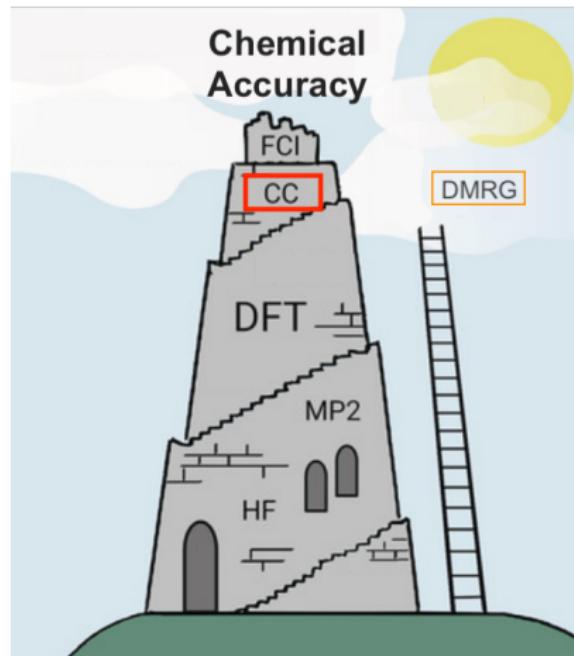
Naïve grid with 10 points per axis yields  $10^{3 \cdot 22} = 10^{66}$  grid points<sup>1</sup>

⇒ Approximation is key!

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<sup>1</sup>The estimated number of particles in the universe is  $\sim 10^{80}$

# Quantum chemical methods



## Strongly correlated quantum chemistry



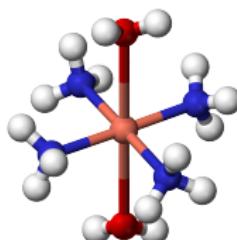
# Strongly correlated quantum chemistry



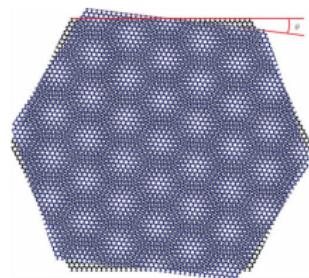
This happens at various scales



Small molecules



Transition metal  
complexes



Twisted bilayer  
graphene

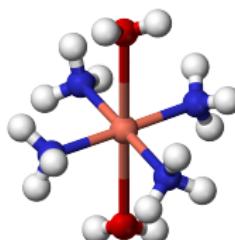
# Strongly correlated quantum chemistry



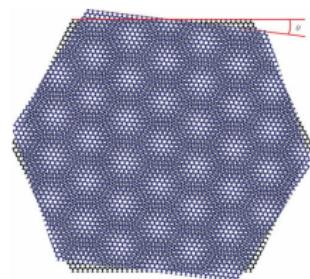
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Strongly correlated systems are very difficult to compute

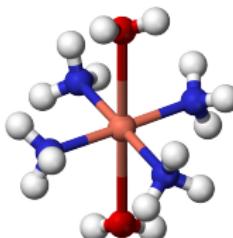
# Strongly correlated quantum chemistry



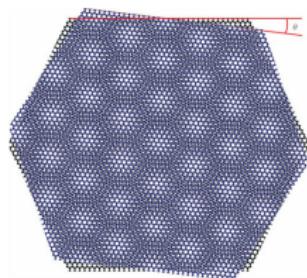
This happens at various scales



Small molecules



Transition metal  
complexes



Twisted bilayer  
graphene

Strongly correlated systems are very difficult to compute

Important failure modes for existing methods

⇒ The research frontier of computational chemistry lies in strongly correlated systems

# Roadmap



1. Quantum mechanics of the hydrogen atom
2. The hydrogen molecule and the Hartree-Fock method
3. Second quantization I
4. Second quantization II
5. Second quantization III
6. The linear combination of atomic orbitals ansatz
7. Gaussian-type atomic orbitals and high-dimensional integration
8. Size consistency and the coupled cluster ansatz
9. Truncating the coupled cluster ansatz and its working equations
10. Finding one root to the coupled cluster equations
11. Finding all roots to the coupled cluster equations

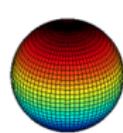
# Quantum mechanics of the hydrogen atom

Analytic solution to the eigenvalue problem

$$\left( \frac{1}{2} \Delta_r + \frac{1}{|r - R|} \right) \Psi(r) = -E\Psi(r)$$

3D partial differential equation!

$$Y_0^0 = 1$$



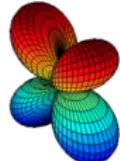
$$Y_1^0 = \cos\theta$$



$$Y_2^0 = 3\cos^2\theta - 1$$



$$Y_2^1 = \cos\theta \sin\theta \sin\phi$$



$$Y_3^0 = 5\cos^3\theta - 3\cos\theta$$



$$Y_3^1 = (5\cos^2\theta - 1)\sin\theta \cos\phi$$



# The hydrogen molecule and the Hartree-Fock method

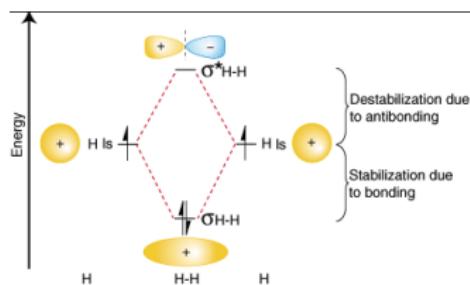
The Hamiltonian reads

$$H = -\frac{1}{2}\Delta_{r_1} - \frac{1}{|r_1 - R_1|} - \frac{1}{2}\Delta_{r_2} - \frac{1}{|r_2 - R_2|} \\ - \frac{1}{|r_2 - R_1|} - \frac{1}{|r_1 - R_2|} + \frac{1}{|r_1 - r_2|}$$

we want to solve

$$H\Psi(r_1, r_2) = E\Psi(r_1, r_2)$$

⇒ Not possible!



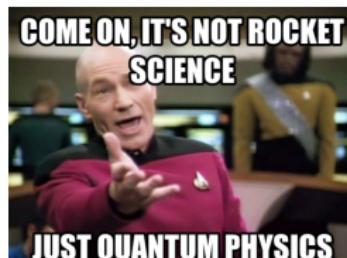
## Second quantization I:

Fock space, creation and annihilation operators, CAR algebra

High-dimensional PDE's are hard!

- Can we think of matrices instead?
- Can we manipulate these matrices fast?

# Second Quantization



## Second quantization II:

Excitation operators, CCR algebra, nilpotent Lie algebra

The second quantization shows its teeth



## Second quantization III:

Slater-Condon rules, Wick's theorem, Hartree-Fock revisited

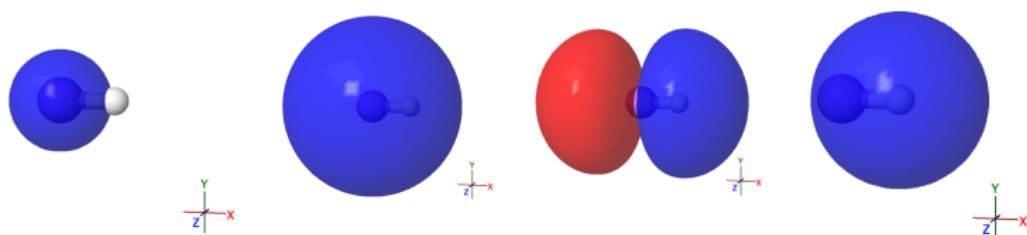
Let's bring out the big guns...



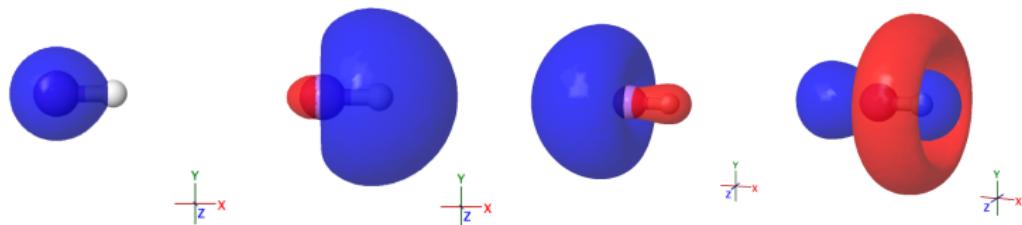
(Wick'ed guys)

# The linear combination of atomic orbitals ansatz

## Atomic orbitals



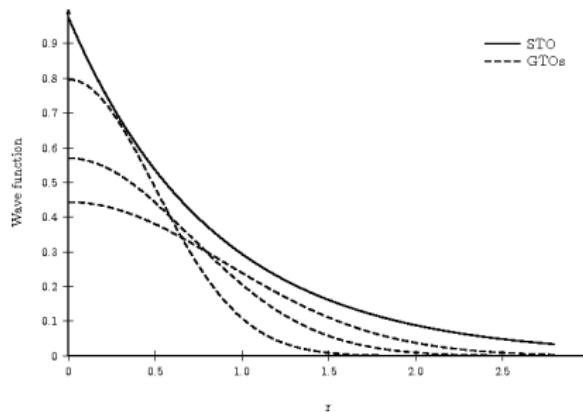
## Molecular orbitals



# Gaussian-type atomic orbitals and high-dimensional integration

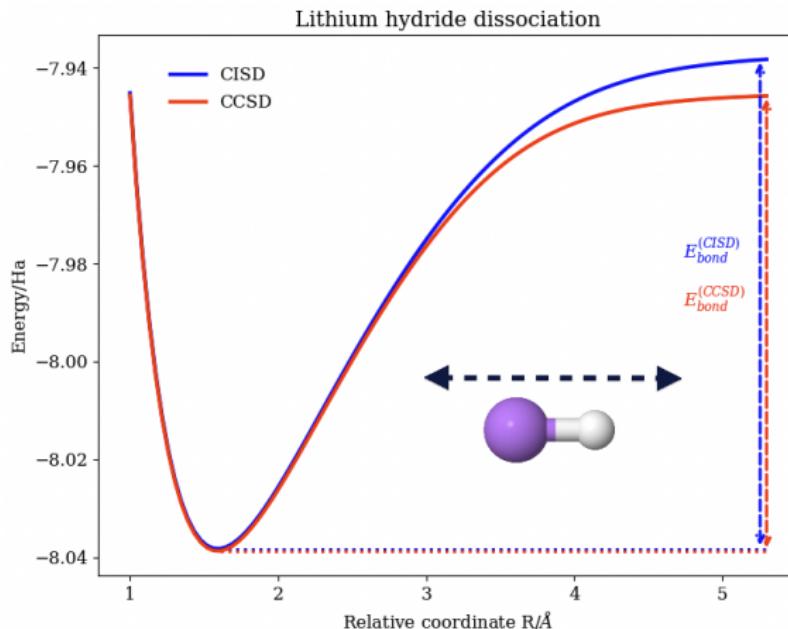
We will face the following integral evaluation:

$$v_{p,q,r,s} = \int_{X \times X} \frac{\xi_p^*(x_1)\xi_q(x_1)\xi_r^*(x_2)\xi_s(x_2)}{|r_1 - r_2|} d\lambda(x)d\lambda(x),$$



# Size consistency and the coupled cluster ansatz

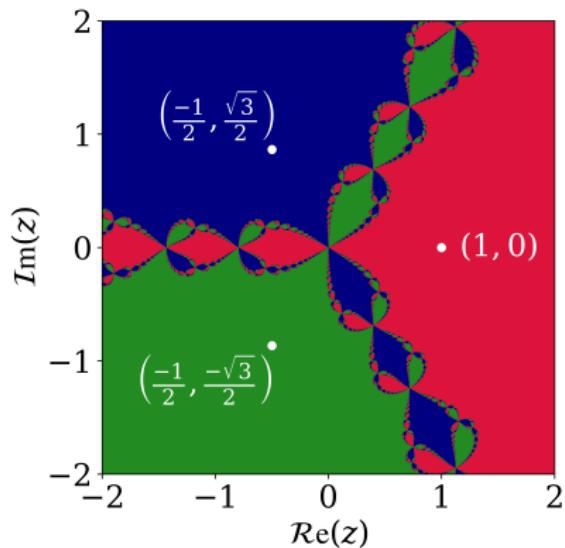
What do we really care about?



# Truncating the coupled cluster ansatz and its working equations

$$\begin{aligned}
E(t) &= \langle \Psi_0, H\Psi_0 \rangle + \sum_{IA} f_{IA} t_I^A + \frac{1}{4} \sum_{IJAB} \langle IJ \parallel AB \rangle t_I^A t_J^B + \frac{1}{2} \sum_{IJAB} \langle IJ \parallel AB \rangle t_I^A t_J^B, \\
f(t)_I^A &= f_{IA} + \sum_C f_{AC} t_I^C - \sum_K f_{KI} t_K^A + \sum_{KC} \langle KA \parallel CI \rangle t_C^K + \sum_{KC} f_{KC} t_K^{AC} \\
&\quad + \frac{1}{2} \sum_{KCD} \langle KA \parallel CD \rangle t_{KI}^{CD} - \frac{1}{2} \sum_{KLC} \langle KL \parallel CI \rangle t_{KL}^{CA} - \sum_{KC} f_{KC} t_I^C t_K^A - \sum_{KLC} \langle KL \parallel CI \rangle t_K^C t_L^A \\
&\quad + \sum_{KCD} \langle KA \parallel CD \rangle t_K^C t_D^P - \sum_{KLCD} \langle KL \parallel CD \rangle t_K^C t_I^D t_L^A + \sum_{KLCD} \langle KL \parallel CD \rangle t_C^K t_{LI}^{DA} \\
&\quad - \frac{1}{2} \sum_{KLCD} \langle KL \parallel CD \rangle t_{KI}^{CD} t_L^A - \frac{1}{2} \sum_{KLCD} \langle KL \parallel CD \rangle t_{KL}^{CA} t_I^B \\
f(t)_{IJ}^{AB} &= \langle IJ \parallel AB \rangle + \sum_C (f_{BC} t_{IJ}^{AC} - f_{AC} t_{IJ}^{BC}) - \sum_K (f_{KJ} t_{IK}^{AB} - f_{KI} t_{JK}^{AB}) \\
&\quad + \frac{1}{2} \sum_{KL} \langle KL \parallel IJ \rangle t_{KL}^{AB} + \frac{1}{2} \sum_{CD} \langle AB \parallel CD \rangle t_{IJ}^{CD} + P(IJ)P(AB) \sum_{KC} \langle KB \parallel CJ \rangle t_{IK}^{AC} \\
&\quad + P(IJ) \sum_C \langle AB \parallel CJ \rangle t_I^C - P(AB) \sum_K \langle KB \parallel IJ \rangle t_A^K \\
&\quad + \frac{1}{2} P(IJ)P(AB) \sum_{KLCD} \langle KL \parallel CD \rangle t_{IK}^C t_{LJ}^{DB} + \frac{1}{4} \sum_{KLCD} \langle KL \parallel CD \rangle t_{IJ}^{CD} t_{KL}^{AB} \\
&\quad + \frac{1}{2} P(AB) \sum_{KLCD} \langle KL \parallel CD \rangle t_{IJ}^C t_{KL}^{BD} - \frac{1}{2} P(IJ) \sum_{KLCD} \langle KL \parallel CD \rangle t_{IK}^{AB} t_{JL}^{CD} \\
&\quad + \frac{1}{2} P(AB) \sum_{KL} \langle KL \parallel IJ \rangle t_K^A t_L^B + \frac{1}{2} P(IJ) \sum_{CD} \langle AB \parallel CD \rangle t_I^C t_J^D \\
&\quad - P(IJ)P(AB) \sum_{KC} \langle KB \parallel IC \rangle t_K^A t_J^C + P(AB) \sum_{KC} f_{KC} t_K^A t_J^{BC} \\
&\quad + P(IJ) \sum_{KC} f_{KC} t_I^C t_J^{AB} - P(IJ) \sum_{KLC} \langle KL \parallel CI \rangle t_C^K t_L^{AB} \\
&\quad + P(AB) \sum_{KCD} \langle KA \parallel CD \rangle t_K^C t_{IJ}^{DB} + P(IJ)P(AB) \sum_{KCD} \langle AK \parallel DC \rangle t_I^D t_{JK}^{BC} \\
&\quad + P(IJ)P(AB) \sum_{KLC} \langle KL \parallel IC \rangle t_L^A t_{JK}^{BC} + \frac{1}{2} P(IJ) \sum_{KLC} \langle KL \parallel CJ \rangle t_I^C t_{KL}^{AB} \\
&\quad - \frac{1}{2} P(AB) \sum_{KCD} \langle KB \parallel CD \rangle t_K^A t_{IJ}^{CD} + \frac{1}{2} P(IJ)P(AB) \sum_{KLC} \langle KB \parallel CD \rangle t_I^C t_K^A t_J^D \\
&\quad + \frac{1}{2} P(IJ)P(AB) \sum_{KLC} \langle KL \parallel CJ \rangle t_I^C t_K^A t_L^B - P(IJ) \sum_{KLCD} \langle KL \parallel CD \rangle t_K^C t_I^D t_{LJ}^{AB} \\
&\quad - P(AB) \sum_{KLCD} \langle KL \parallel CD \rangle t_K^C t_L^A t_{IJ}^{DB} - \frac{1}{4} P(IJ) \sum_{KLCD} \langle KL \parallel CD \rangle t_I^C t_J^D t_{KL}^{AB} \\
&\quad + \frac{1}{4} P(AB) \sum_{KLCD} \langle KL \parallel CD \rangle t_K^A t_L^B t_{IJ}^{CD} + P(IJ)P(AB) \sum_{KLCD} \langle KL \parallel CD \rangle t_I^C t_L^B t_{KJ}^{AD} \\
&\quad + \frac{1}{4} P(IJ)P(AB) \sum_{KLCD} \langle KL \parallel CD \rangle t_I^C t_K^A t_J^D t_L^B
\end{aligned}$$

# Finding one root to the coupled cluster equations



## Finding all roots to the coupled cluster equations

Seeking:

$$f(t) = \begin{bmatrix} f_1(t) \\ \vdots \\ f_m(t) \end{bmatrix} = \begin{bmatrix} f_1(t_1, \dots, t_m) \\ \vdots \\ f_m(t_1, \dots, t_m) \end{bmatrix} = 0.$$

Find a start system  $g$ , and then continuously deform  $g$  into  $f$

Davidenko differential equation:

$$\frac{\partial}{\partial t} H(t, \lambda) \left( \frac{d}{d\lambda} t(\lambda) \right) + \frac{\partial}{\partial \lambda} H(t, \lambda) = 0, \quad t(1) = s_0,$$

# Finding all roots to the coupled cluster equations

Seeking:

$$f(t) = \begin{bmatrix} f_1(t) \\ \vdots \\ f_m(t) \end{bmatrix} = \begin{bmatrix} f_1(t_1, \dots, t_m) \\ \vdots \\ f_m(t_1, \dots, t_m) \end{bmatrix} = 0.$$

