# Even Shorter Quantum Circuit for Phase Estimation on Early Fault-Tolerant Quantum Computers

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2024

**Problem:** Quantum Phase Estimation.

Method: Quantum Complex Exponential Least Squares (QCELS).

Device: (Constraint) Early Fault-Tolerant Quantum Computer.

#### Introduction:

Main problem, input, complexity.

#### **Quantum Circuit and Data Generation Process**

• Data Set Preparation and Post-processing

## Method: QCELS

• Solve the optimization problem, Heuristic theorem, Intuitive analysis.

Theorem for large p<sub>0</sub> Multi QCELS

Algorithm for small  $p_0$ 

Conclusion

## Main Problem

Given a Hamiltonian  $H \in \mathbb{C}^{d \times d}$ ,  $d \gg 1$ , we estimate the smallest eigenvalue  $\lambda_0$  (Ground state energy).

Assume: 
$$\lambda_0 \in [-\pi, \pi)$$

corresponding eigenvector of  $|\psi_0\rangle\in\mathbb{C}^d$ 

## Input:

The problem is QMA-hard.

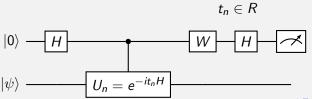
Common additional assumption: good initial state

$$|\psi\rangle = U_I |0^n\rangle$$

• Initial State:

$$|\psi\rangle \in \mathbb{C}^d$$
 
$$0 < p_0 = |\langle \psi | \psi_0 \rangle|^2 < 1$$

• Quantum Oracle: Given any



## Quantum Circuit and Hadamard Test

To achieve this, our quantum circuit above is the same as the circuit used in the Hadamard test but replaces  $U=e^{-i\tau H}$  by  $U_n=e^{-in\tau H}$  for a sequence of integers n.

#### This simple circuit:

- Uses only one ancilla qubit.
- Is suitable for early fault-tolerant quantum computers.

## Quantum Circuit Details

#### **Definitions:**

- $\bullet$  au is the time step used in the controlled time evolution.
- $t_n = n\tau$ , where  $n = 0, \dots, N-1$ , defines the discrete time steps.
- $Z_n$  is a complex-valued random variable such that  $E(Z_n) = \langle \psi | e^{-it_n H} | \psi \rangle$ .

#### **Circuit Operation:**

The circuit provides an estimate of  $\langle \psi | e^{-it_n H} | \psi \rangle$  by measuring the success probability of the first qubit.

Repeated measurements at different n provide a complex time series:

$$\{(t_n, Z_n)\}_{n=0}^{N-1}$$

where  $Z_n$  approximates  $\langle \psi | e^{-it_n H} | \psi \rangle$ .



## Complexity

$$\{t_n\}_{n=1}^N \in \mathbb{R} \longrightarrow \dots$$

#### **Maximum Running Time:**

$$T_{\mathsf{max}} = (N-1)\tau$$
 for simplicity  $T_{\mathsf{max}} = N\tau$ 

How large does our quantum computer need to be?

#### **Total Running Time:**

$$T_{\mathsf{total}} = \sum_{n=1}^{N} |t_n|$$

How long will we need to run it?



# Complexity

**Classical Result:** To ensure  $|\lambda - \lambda_0| \le \epsilon$ :

When 
$$P_0=1$$
,  $T_{\sf max}={\it O}(1)$ ,  $T_{\sf total}={\it O}\left(\frac{1}{\epsilon^2}\right)$ 

(Hadamard Test).

Any 
$$P_0$$
,  $T_{\mathsf{max}} > \frac{\pi}{\epsilon}$ ,  $T_{\mathsf{total}} = O\left(\frac{1}{\epsilon}\right)$ 

(Heisenberg-limited QPE).

**Goal:** For any  $P_0$ ,

$$T_{\mathsf{max}} > rac{\pi}{\epsilon} \quad (\mathsf{Reduce}), \quad \mathsf{such that} \quad T_{\mathsf{total}} = O\left(rac{1}{\epsilon}
ight) \quad (\mathsf{QPE}).$$



## Complexity

Our Method, for  $P_0 \ge 0.71$ ,

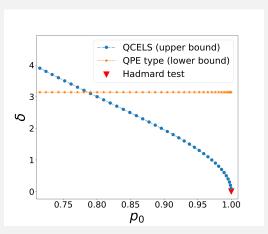
$$T_{\mathsf{max}} \leq \frac{\pi}{\epsilon}, \quad T_{\mathsf{total}} = O\left(\frac{1}{\epsilon}\right)$$

Any  $P_0$ ,

$$T_{\mathsf{max}} > rac{\pi}{\epsilon}, \quad T_{\mathsf{total}} = O\left(rac{1}{\epsilon}
ight) \quad \mathsf{(QPE)}.$$

# Informal Theory

$$T_{\mathsf{max}} = O\left(\frac{\sqrt{1-p_0}}{\epsilon}\right)$$



## Quantum Circuit and Data Generation Process

- Quantum Circuit: Fig. 1 shows a quantum circuit designed to generate data points for the Hamiltonian *H*.
- Random Variables:
  - Set W = I, measure ancilla qubit:

$$X_n = \begin{cases} 1 & \text{if outcome is } 0 \\ -1 & \text{if outcome is } 1 \end{cases}$$

$$E(X_n) = \operatorname{Re}\langle \psi | \exp(-in\tau H) | \psi \rangle$$

• Set  $W = S^{\dagger}$ , measure ancilla qubit:

$$Y_n = egin{cases} 1 & ext{if outcome is 0} \\ -1 & ext{if outcome is 1} \end{cases}$$

$$E(Y_n) = \operatorname{Im}\langle \psi | \exp(-in\tau H) | \psi \rangle$$

# Data Set Preparation and Postprocessing

- Input Parameters: Given two preset parameters  $N, N_s > 0$  and time step  $\tau > 0$ , the data set is generated.
- Data Set  $D_H$ :

$$D_H = \{(n\tau, Z_n)\}_{n=0}^{N-1}$$

By running the quantum circuit  $N_s$  times  $Z_n$  is calculated as:

$$Z_n = \frac{1}{N_s} \sum_{k=1}^{N_s} (X_{k,n} + iY_{k,n})$$

• **Key Result:** In the limit  $N_s \to \infty$ ,

$$Z_n = \langle \psi | \exp(-in\tau H) | \psi \rangle$$

- Simulation Time:
  - Maximal simulation time:  $T_{\sf max} = (N-1) au$
  - Total simulation time:  $\frac{NN_sT_{\text{max}}}{2}$



## Quantum Complex Exponential Least Squares

- Time t
- Sampled real and imaginary parts over time

## Objective:

$$(r^*, \theta^*) = \arg\min_{r \in \mathbb{C}, \theta \in \mathbb{R}} L(r, \theta)$$

where the loss function is defined as:

$$L(r,\theta) = \frac{1}{N} \sum_{n=0}^{N-1} |Z_n - re^{-it_n\theta}|^2$$

**Output:**  $\theta^*$  is the estimate for the phase  $\lambda_0$ .

**Note:** Fitting may not be exact if  $p_0 < 1$ , but it can still estimate  $\lambda_0$  accurately.

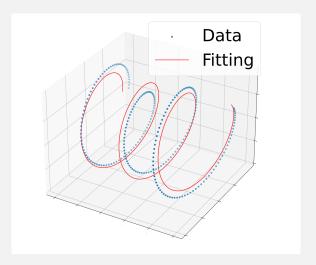


Figure: Data fitting

# Solve the Optimization Problem

## Step 1: Fix $\theta$ , optimize r:

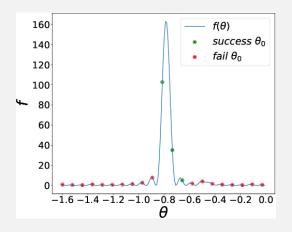
$$\min_{r \in \mathbb{C}} L(r, \theta) = \frac{1}{N} \sum_{n=0}^{N-1} |Z_n|^2 - \frac{1}{N} \left| \sum_{n=0}^{N-1} Z_n e^{i\theta n\tau} \right|^2$$

## Step 2: Optimize with respect to $\theta$ :

$$\theta^* = \arg\max_{\theta \in \mathbb{R}} f(\theta), \quad f(\theta) = \frac{1}{N} \left| \sum_{n=0}^{N-1} Z_n e^{i\theta n\tau} \right|^2$$

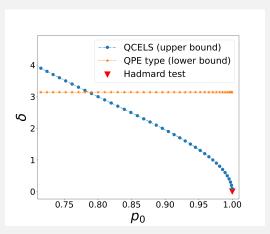
**Energy Landscape:** The landscape is rugged but can be handled classically since  $\theta$  is a scalar.

# Optimization Landscape



# Informal Theory

$$T_{\mathsf{max}} = O\left(\frac{\sqrt{1-p_0}}{\epsilon}\right)$$



## Heuristic Theorem: Maximal Run Time

Heuristic Theorem (General Result) for large P<sub>0</sub>:

Assume the overlap  $p_0$  between the initial state and the target eigenstate is very large. In this case,

$$\delta = \Theta(\sqrt{1-p_0}), \quad T_{\mathsf{max}} = rac{\delta}{\epsilon}$$

Let  $\theta^*$  be the optimizer:

$$|(\theta^* - \lambda_0) \mod [-\pi/\tau, \pi/\tau]| < \epsilon.$$

**Key Result:** 

$$T_{\mathsf{max}} = O\left(\frac{\delta}{\epsilon}\right)$$

where  $\delta \sim \sqrt{1-p_0}$ . This holds for sufficiently large  $p_0$ .

# Intuitive Analysis of Basic Version of QCELS

Recall:

$$heta^* = rg \max_{ heta \in \mathbb{R}} f( heta), \quad f( heta) = rac{1}{N} \left| \sum_{n=0}^{N-1} Z_n e^{i heta n au} \right|^2.$$

Objective: Bound

$$R_0 = |(\lambda_0 - \theta^*)\tau \mod [-\pi, \pi]|.$$

Lower bound on  $f(\lambda_0)$ :

$$(2p_0-1)N \leq \sqrt{f(\lambda_0)}.$$

**Upper bound on**  $f(\theta^*)$ :

$$\sqrt{f(\theta^*)} \leq \left| \frac{\sin(NR_0/2)}{\sin(R_0/2)} \right| + (1-p_0)N.$$

## Intuitive Analysis of Basic Version of QCELS

#### **Optimality Condition:**

$$\sqrt{f(\theta^*)} \ge \sqrt{f(\lambda_0)} \quad \Rightarrow \quad \left| \frac{\sin(NR_0/2)}{\sin(R_0/2)} \right| \ge (3p_0 - 2)N.$$

Approximation:

$$\frac{\sin(\textit{N}(\delta/2\textit{N}))}{\sin(\delta/2\textit{N})} \approx \textit{N}\left(1 - \frac{\delta^2}{24}\right).$$

$$\delta^2 \approx 72(1-p_0) \quad \Rightarrow \quad \delta \to 0 \text{ as } p_0 \to 1.$$

Since  $\frac{\sin(Nx)}{\sin(x)}$  is decreasing on  $[0, \pi/(2N)]$ , we conclude:

$$R_0 \leq rac{\delta}{N} \quad ext{or} \quad |(\lambda_0 - heta^*) \mod [-\pi/ au, \pi/ au]| < rac{\delta}{T_{ ext{max}}} = \epsilon.$$

This gives  $T_{\max} = \frac{\delta}{\epsilon}$ : short runtime!

# Theorem for Large but specific $p_0$

#### Theorem 1:

If  $p_0 > 0.71$ , choose

$$\delta = \Theta\left(\sqrt{1-p_0}\right)$$
 .

There exists an algorithm that uses 1 ancilla qubit to estimate  $\lambda_0$  to precision  $\epsilon$  with:

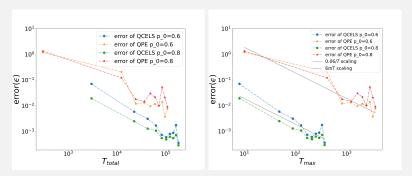
$$T_{\mathsf{max}} = rac{\delta}{\epsilon}, \quad T_{\mathsf{total}} = ilde{\Theta}\left(\delta^{-(1+o(1))}\epsilon^{-1}
ight).$$

#### **Distinct Feature:**

• The preconstant  $\delta$  can be arbitrarily small as  $p_0 \to 1$ .

#### Numerical evidence

## Transverse Field Ising Model (TFIM).



- Numerical performance is much better than theoretical prediction, and the bound 0.71 can be pushed downward.
- Two order of magnitude reduction of maximal runtime!

## Convergence - Basic Version of QCELS

#### Theorem (Basic version of QCELS):

Given  $p_0 > 0.71$ , we can choose:

$$\delta = \Theta(\sqrt{1-p_0}), \quad T_{\mathsf{max}} = rac{\delta}{\epsilon}, \quad \mathsf{NN_s} = \Omega(\delta^{-(2+o(1))})$$

Let  $\theta^*$  be the optimizer. Then with high probability:

$$|(\theta^* - \lambda_0) \mod [-\pi/\tau, \pi/\tau]| < \epsilon.$$

#### **Key Results:**

- Short maximal runtime (circuit depth).
- Does not achieve Heisenberg-limited scaling.

## **Scalings:**

$$T_{\text{max}} = N\tau \quad \Rightarrow \quad N = O(\epsilon^{-1}) \text{ if } \tau \text{ is small.}$$

$$T_{\mathsf{total}} = \tau N_{\mathsf{s}} N(N-1)/2 = O(\epsilon^{-2}).$$

## Multi-level QCELS

#### Multi-level QCELS:

The result  $T_{\max} = \frac{\delta}{\epsilon}$ ,  $NN_s = \Omega(\delta^{-(2+o(1))})$  is independent of  $\tau$ .

$$|(\theta^* - \lambda_0) \mod [-\pi/\tau, \pi/\tau]| < \epsilon.$$

**Algorithm:** For  $j = 1, \ldots, J$ :

- Generate data set  $D_{H,j} = \{(n\tau_j, Z_{n,j})\}_{n=0}^{N-1}$ .
  - $2 \ \, \mathsf{Solve} \, \left( r_j^*, \theta_j^* \right) = \mathsf{arg} \, \mathsf{min}_{r \in \mathbb{C}, \theta \in [-\lambda_{\mathsf{min}}, \lambda_{\mathsf{max}}]} \, L(r, \theta).$
  - Shrink search interval:

$$\lambda_{\min} = \theta_j^* - \frac{\pi}{2\tau_j}, \quad \lambda_{\max} = \theta_j^* + \frac{\pi}{2\tau_j}.$$

## Convergence - Multi-level QCELS

## Theorem (Multi-level QCELS):

If  $p_0 > 0.71$ , choose:

$$\delta = \Theta(\sqrt{1-p_0}), \quad T_{\mathsf{max}} = rac{\delta}{\epsilon}, \quad T_{\mathsf{total}} = \Theta\left(\delta^{-(1+o(1))}\epsilon^{-1}
ight).$$

Let  $\theta^*$  be the output of multi-level QCELS. Then with high probability:

$$|(\theta^* - \lambda_0) \mod [-\pi, \pi]| < \epsilon.$$

#### **Key Results:**

- Short maximal runtime (circuit depth).
- Achieves Heisenberg-limited scaling.

# Algorithm for Small $p_0$

## **Proposed Algorithm:**

• Combines the multilevel QCELS algorithm with the Fourier-filtering technique from Ref. [15] to estimate  $\lambda_0$ .

Relative Overlap  $p_r(I, I')$ :

$$p_r(I,I') = \frac{|\langle \psi | \psi_0 \rangle|^2 1_I(\lambda_0)}{\sum_{\lambda_k \in I'} |\langle \psi | \psi_k \rangle|^2},$$

where  $1_I(\cdot)$  is the indicator function on I, and  $\lambda_0 \in I$ .

Scenario with Spectral Gap  $\Delta = \lambda_1 - \lambda_0$ :

- Choose intervals  $I = [-\pi, \lambda_{prior} + \Delta/4]$  and  $I' = [-\pi, \lambda_{prior} + 3\Delta/4]$ .
- The distance between I and  $(I')^c$  is  $D = \Delta/2$ .
- Relative overlap  $p_r(I, I')$  is 1 when  $|\lambda_{prior} \lambda_0| \leq \Delta/4$ .

**Note:** Even with a small spectral gap, appropriate intervals I and I' can result in a larger distance D and still achieve a large relative overlap.

# Small p<sub>0</sub>

#### Theorem

Given relative overlap  $p_r(I, I') \ge 0.71$  and  $D = \min_{x_1 \notin I', x_2 \in I} |x_1 - x_2|$ , choose

$$\delta = \Theta\left(\sqrt{1 - p_r(I, I')}\right).$$

There exists an algorithm that uses 1 ancilla qubit to estimate  $\lambda_0$  to precision  $\epsilon$  with:

$$T_{\mathsf{max}} = \tilde{\Theta}(D^{-1}) + \frac{\delta}{\epsilon}, \quad T_{\mathsf{total}} = \tilde{\Theta}\left(p_0^{-2}\delta^{-(2+o(1))}(D^{-1} + \frac{\delta}{\epsilon})\right).$$

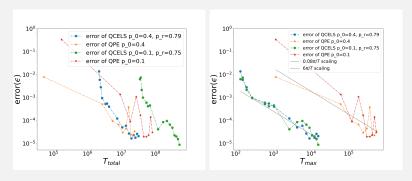
#### **Distinct Features:**

- Uses information of relative overlap (previous algorithms are agnostic to this).
- Reduces circuit depth when  $D \ll \epsilon$  and  $p_r(I, I')$  is large.



## Numerical evidence

#### Hubbard.



• Two order of magnitude reduction of maximal runtime!

## Conclusion

#### Conclusion:

- QCELS is efficient for estimating the ground state energy on early quantum computers.
- It reduces the maximal and total runtime significantly, making it a powerful method for quantum phase estimation.

Thank you!