Multi-dimensional Femtosecond-laser induced dynamics of CO on metals:

ACCOUNTING FOR ELECTRONIC FRICTION AND SURFACE MOTION WITH COMBINED MODELS

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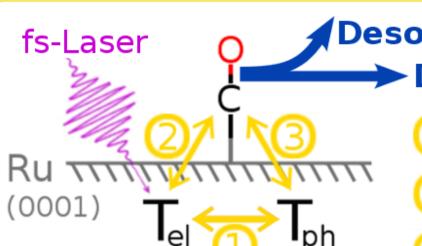
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Introduction

Motivation

- Aim: gain precise understanding of adsorbate bonding on metals ⇒ Important for Catalysis
- Why femtosecond(fs)-lasers?
- -produce non-equilibrium 2-T-states ⇒ different than normal heating
- -further **tool** besides STM and scattering -direct future applications possible \Rightarrow "femtochemistry"[1]
- Why CO/Ru(001) and CO/Cu(100)?
- -both are well studied **model systems**
- -recently, interesting **fs-laser experiments**[2][3]
- -also, **ab-initio** based **6-dim. potentials** available [4][5]

How do fs-lasers affect adsorbate-metal systems?



Desorption

- Diffusion (and possibly Reactions)
- Electron-phonon coupling
- Electronic friction
- Phonon-adsorbate interaction

high

Fermi-Dirac-Distribution

t/ps

 T_{el} (Ru)

 $_{---}$ T_{ph} $(Ru)|_{I}$

- only electrons of metal absorb laser $\wedge f(E)$ low T_{el}
- electron-hole pairs thermalize fast ⇒ "hot" Fermi-Dirac-distribution
- electrons transfer energy to ion lattice,
- via 1 electron-phonon coupling • equilibration within ps-timescale
- \Rightarrow Thus, for few ps **two temperatures**:
 - $-T_{\rm el}$ electron temperature
 - $-T_{
 m ph}$ phonon temperature
- both can **couple** to adsorbed **molecule**
- low electron heat **capacity** \Rightarrow $T_{\rm el}$ higher

Two-Temperature Model (2TM)[6] • describes interaction of electrons with phonons and laser Femtochemistry example

$$C_{\rm el} \frac{\partial T_{\rm el}}{\partial t} = \frac{\partial}{\partial z} \kappa \frac{\partial}{\partial z} T_{\rm el} - g(T_{\rm el} - T_{\rm ph}) + S(z, t),$$

$$C_{\rm ph} \frac{\partial T_{\rm ph}}{\partial t} = g(T_{\rm el} - T_{\rm ph}).$$
Ru Type (00001)
$$C_{\rm ph} \frac{\partial T_{\rm ph}}{\partial t} = g(T_{\rm el} - T_{\rm ph}).$$
Ru Type (depth z Tel

• Basis for dynamics: **precomputed PES** from DFT (rPBE + D2)

• all 6 dimensions of the adsorbate

• analytical PES and gradients ⇒ very fast

⇒ number, length of trajectories can be large

• downsides: - surface frozen \Rightarrow no phonons

- had to be constructed first

 \Rightarrow get $T_{\rm el}$ and $T_{\rm ph}$ as f(z,t) from laser parameters and material properties

fs-Laser

Local Density Friction Approximation (LDFA)[8]

Electronic Friction: Langevin Dynamics[7] and

Models and Methods

Six-dimensional Potential Energy Surface (6D PES)[4]

• Langevin equation of motion, a stochastical differential equation: $m_k \frac{d^2 \underline{r}_k}{dt^2} = -\underline{\nabla}_k V(\underline{r}_1, \underline{r}_2) - \eta_{\text{el},k}(\underline{r}_k) \frac{d\underline{r}_k}{dt} + \underline{R}_{\text{el},k}(t).$ Friction force Force due Force on to PES Atom *k* slows movement from e-h pairs

-1.60

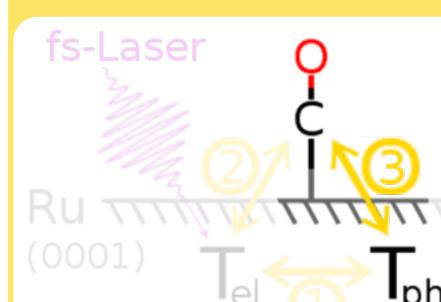
Potential for CO/Ru(001)

fs-Laser



- Local Density Friction Approx. (LDFA): simple model to get friction coefficients $\eta_{el,k}$ $-\mathbf{Atom}\ k\ \mathbf{embedded}\ \mathbf{in}\ \mathbf{free}\ \mathbf{electron}\ \mathbf{gas}$ with density of bare surface at current position \underline{r}_k
- Random forces $\underline{R}_{el,k}$: white noise, dependent on both $\eta_{el,k}$ (from LDFA) and T_{el} (from 2TM) -justified by 2. fluctuation dissipation theorem[9] (relating friction and thermal movement)

Inclusion of Phonons: Generalized Langevin Oscillator(GLO)-model[10]

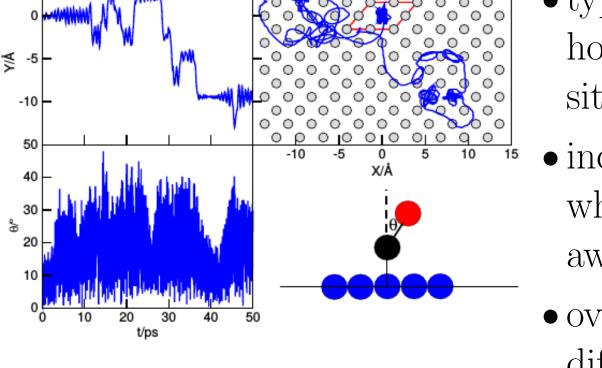


- influence of phonons effectifely modeled (augments frozen surface)
- entire surface understood as 3D oscillator (coords. \underline{r}_s , mass 1 atom)
- coupling to molecule via shifting: $V_{\text{GLO}}(\underline{r}_{\text{C}},\underline{r}_{\text{O}};\underline{r}_{s}) = V(\underline{r}_{\text{C}} \underline{r}_{s},\underline{r}_{\text{O}} \underline{r}_{s})$
- additionally coupled to **ghost oscillator** \underline{r}_q , **models** influence of **bulk**
- -ghost oscillator is subject to friction $\eta_{\rm ph}$ and random forces $\underline{R}_{\rm ph}(T_{\rm ph})$

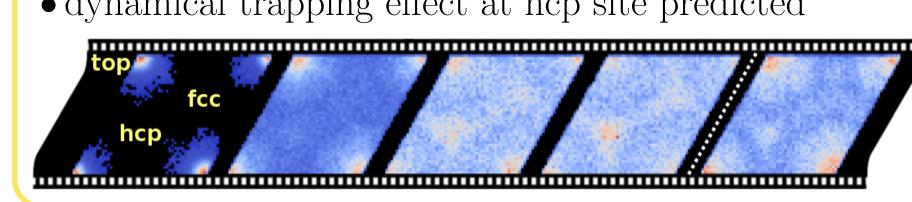
Results

Desorption (Data for Ru) • desorption mainly dur-5 ing first 50 ps (\mathring{A}) • fluence dependence of N = 3desorption yield close to experiment • no barrier in PMF 100 t(ps)exp., n=4.81 MDEF, n=6.99 ■ MDEF-GLO, n=5.17 - T=300 K T=500 K T=1500 KT=2000 KT=4000 K 300 fluence F / J/m² molecule-surface distance Z / Å

Diffusion (Data for Ru, but Cu similar)

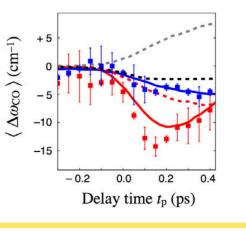


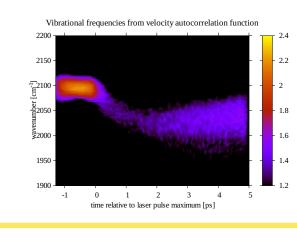
- typical trajectory: hops between top sites and vibration
- increase in θ -angle when CO moves away from top
- overall, very large diffusion
- also, nonisotropic diffusion behaviour observed:
- dynamical trapping effect at hcp site predicted



Vibrations (Data for Cu)

• Frequency-shift from time-resolved SFG





• Preliminary results from our dynamics

Conclusions

- 6D Langevin dynamics of CO on Ru and Cu
- based on first principles, no "free" parameters
- accounting for (via LDFA) electronic friction, hot electron excitation and (via GLO) substrate motion
- allows for detailed time- and space-resolved insights
- no physisorbed state, molecules desorp directly

Outlook

- better electronic friction ($\eta(T_{\rm el})$) and beyond LDFA) \Rightarrow Long term goal: use tensorial friction (exact)
- non-equilibrium lattice model (NLM) instead 2TM
- simulate other coveages (bigger super cells)
- simulate bigger systems (CO + H; hydrocarbons)
- include interaction between adsorbate molecules

[1] M. Bonn, S. Funk, Ch. Hess, D.N. Denzler et al., Science 285, 1042 (1999).