

# FEMTOSECOND-LASER INDUCED DYNAMICS OF CO ON Ru(0001):

## NEW INSIGHTS FROM A HOT-ELECTRON, ELECTRONIC FRICTION MODEL INCLUDING SURFACE MOTION

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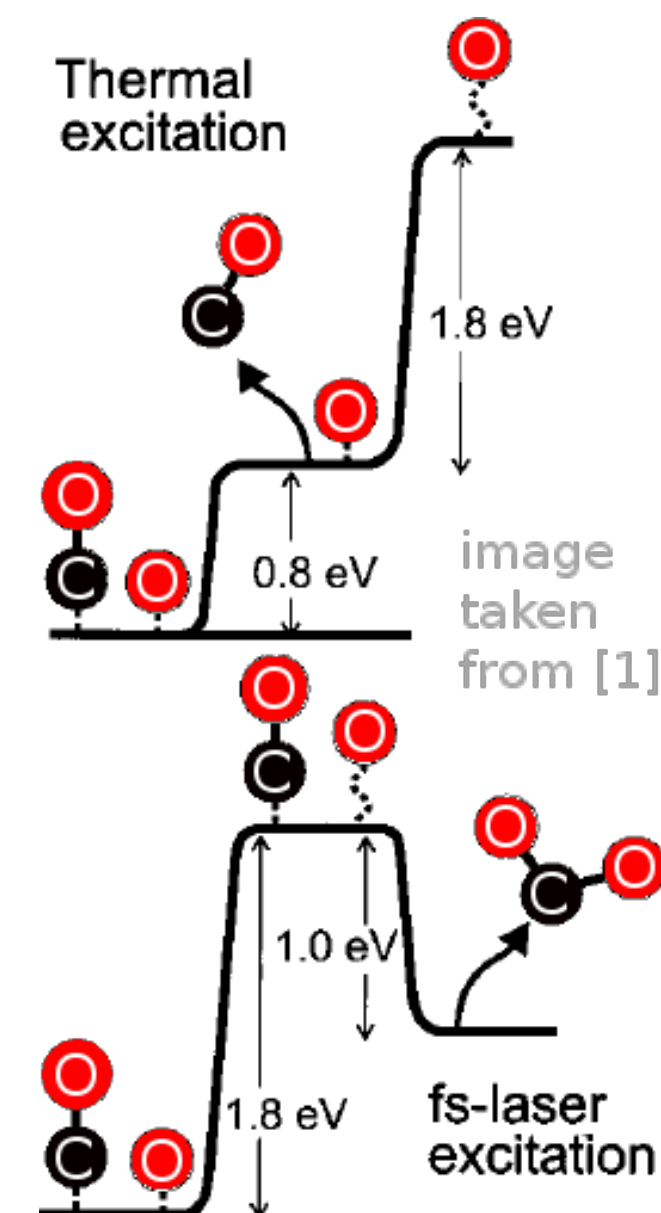
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## Introduction

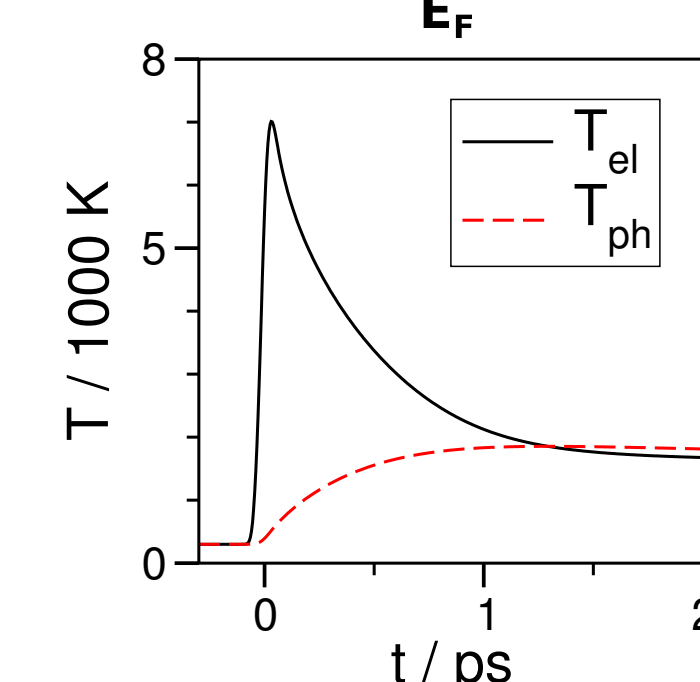
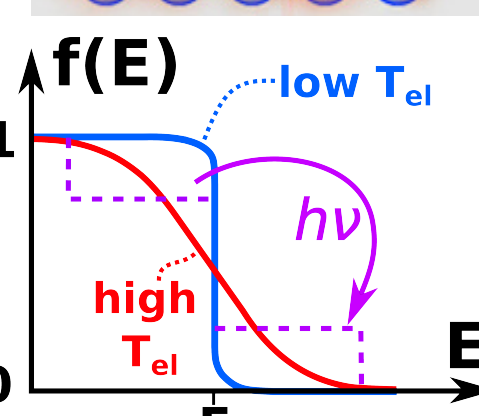
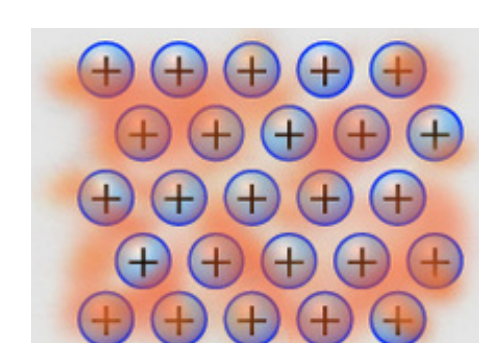
### Motivation

- research on small molecules adsorbed to metals is important for:
  - catalytic applications
  - fundamental understanding of bonding
- femtosecond(fs)-lasers  $\Rightarrow$  very valuable tool
  - allow for investigations on small timescales
  - open up new processes compared to heating or “normal” lasers  $\Rightarrow$  femtochemistry, e.g. [1]:
  - may enable specific control over catalytic reactions  $\Rightarrow$  photocatalysis
- specific motivation for system CO/Ru(0001)
  - experimentally well studied regarding fs-laser irradiation, e.g. [2, 3]
  - fulldimensional *ab-initio* potential recently developed in our group[4]
  - details of this indicate interpretation of experiment [3] may be wrong



### How does fs-laser-irradiation affect metal surfaces?

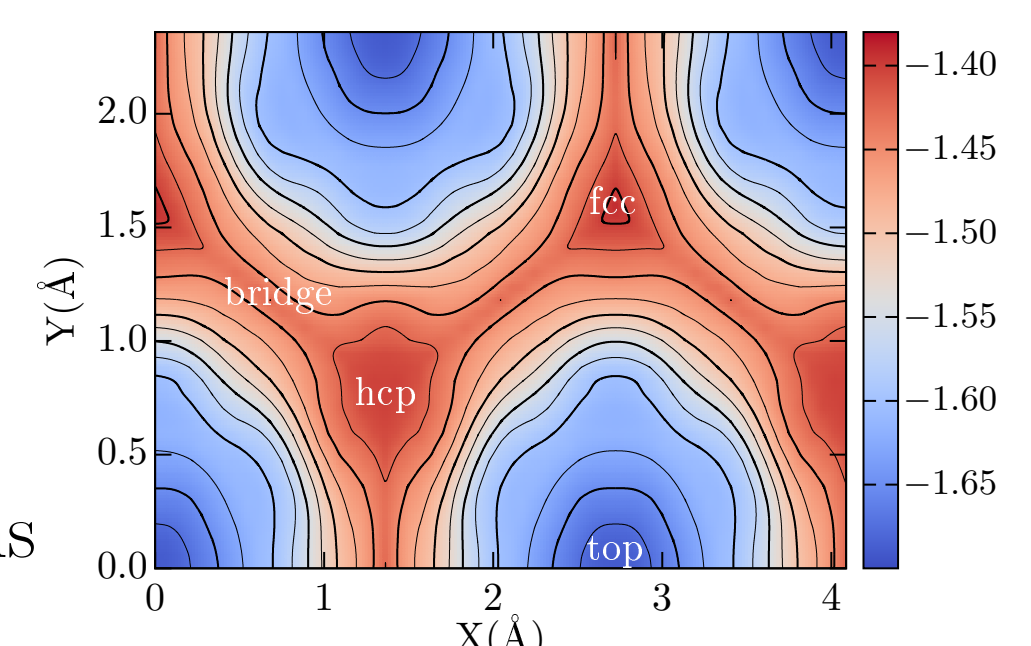
- metals: ion lattice plus quasi-free electron gas
  - visible light is absorbed only by the electrons
  - produced electron-hole pairs thermalize quickly  $\Rightarrow$  “hot” Fermi-Dirac-distribution (after  $\sim 10$  fs)
  - electrons transfer part of energy to ion lattice, via ① **electron-phonon coupling** (phonons = lattice vibrations; quasi-particles)
    - electrons couple to phonons as their fast movement causes “shockwaves” in ion lattice
    - equilibration process completes after  $\sim 1$  ps
- $\Rightarrow$  Thus, with fs-lasers, two different temperatures:
- $T_{el}$  - electron temperature
  - $T_{ph}$  - phonon temperature
- can be simulated using a Two-Temperature Model (2TM)[5] (see right)



## Models and Methods

### Six-dimensional Potential Energy Surface (6D PES)[4]

- Basis for dynamics: precomputed PES from DFT (rPBE + D2)
    - all 6 dimensions of the adsorbate
    - analytical PES and gradients  $\Rightarrow$  very fast
- $\Rightarrow$  number and length of trajectories can be large
- downsides:
    - surface atoms frozen  $\Rightarrow$  no phonons
    - had to be constructed first



### Two-Temperature Model (2TM)[5]

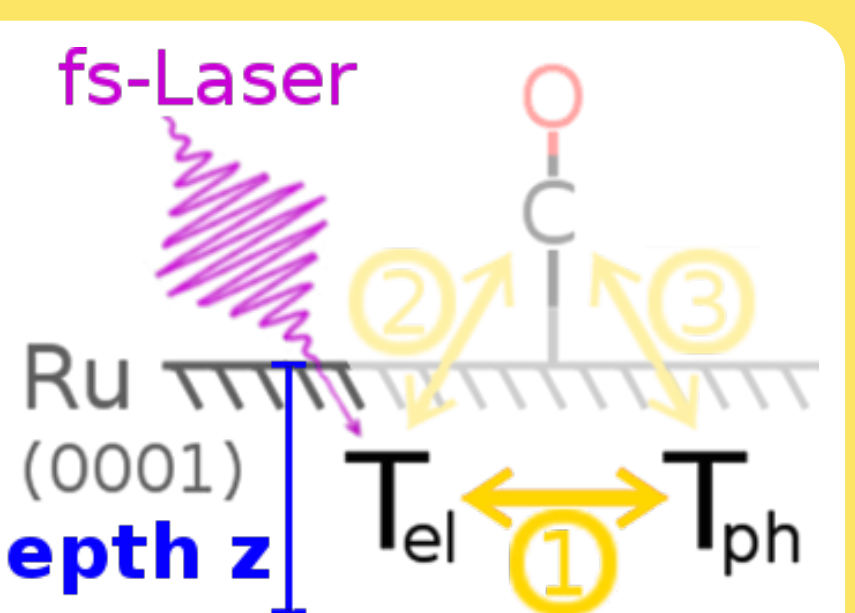
- describes interaction of metal with laser, using two differential equations:

$$C_{el} \frac{\partial T_{el}}{\partial t} = \frac{\partial}{\partial z} \kappa \frac{\partial T_{el}}{\partial z} - g(T_{el} - T_{ph}) + S(z, t),$$

$$C_{ph} \frac{\partial T_{ph}}{\partial t} = g(T_{el} - T_{ph}).$$

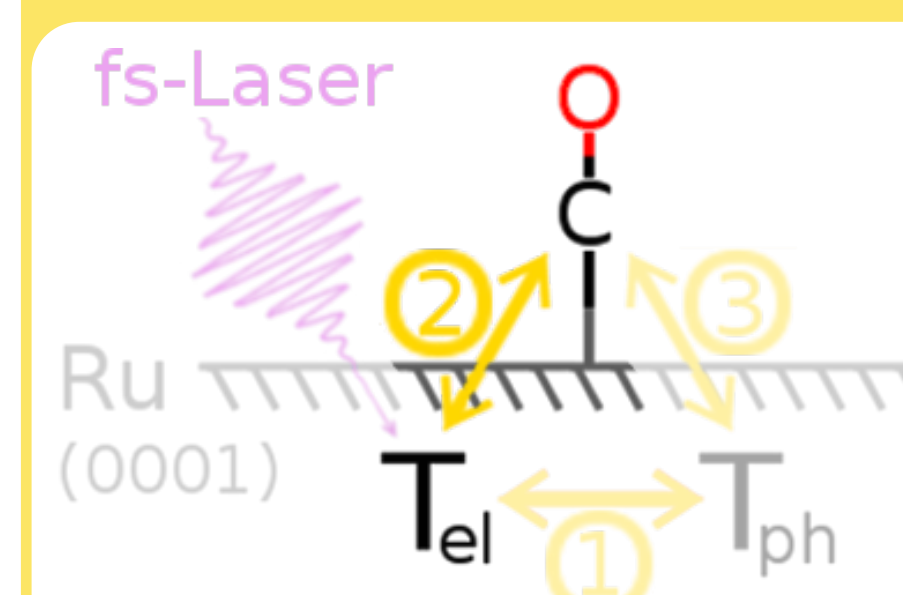
$\Rightarrow$  get  $T_{el}$  and  $T_{ph}$  as  $f(z, t)$  from laser parameters and material properties:

- laser wavelength  $\lambda$  (affects penetration depth into material)
- electron and phonon heat capacities  $C_{el}$  and  $C_{ph}$
- (effective) absorbed fluence  $F$  (energy/area)
- electron heat conductivity  $\kappa$
- pulse duration  $\tau$  (all three appear in the “source term”  $S(z, t)$ )
- electron-phonon coupling constant  $g$



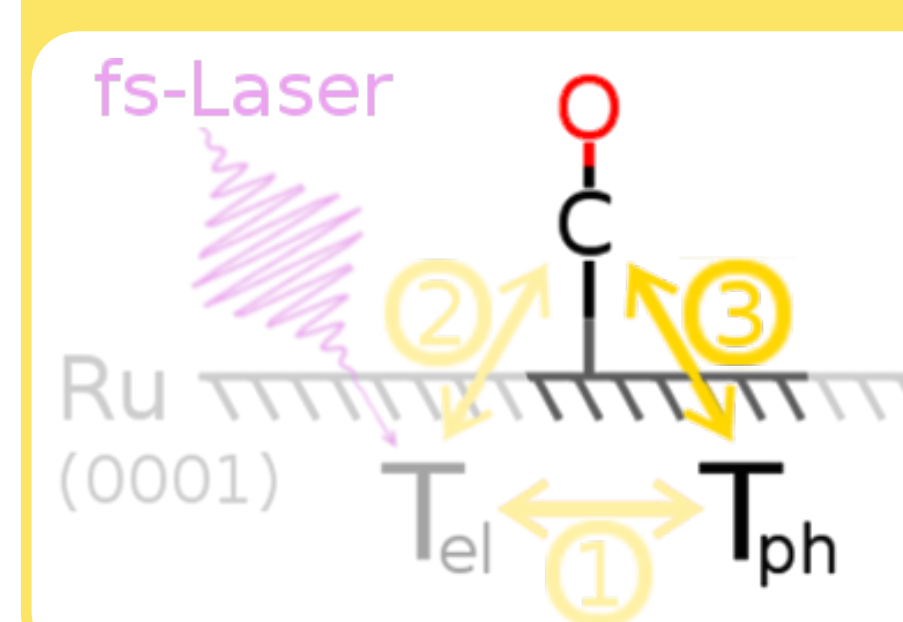
### Electronic Friction: Langevin Dynamics[6] and Local Density Friction Approximation (LDFA)[7]

- Langevin equation of motion, a stochastic differential equation:
 
$$m_k \frac{d^2 \underline{r}_k}{dt^2} = \underbrace{-\nabla_k V(\underline{r}_1, \underline{r}_2)}_{\text{Force on Atom } k} - \underbrace{\eta_{el,k}(\underline{r}_k) \frac{d\underline{r}_k}{dt}}_{\text{Friction force slows movement}} + \underbrace{\underline{R}_{el,k}(t)}_{\text{Random force from e-h pairs}}.$$
- describes movement of CO on the PES and interaction with electron-hole pairs (friction and excitation)
- Local Density Friction Approx. (LDFA): most simple model to calculate friction coefficients  $\eta_{el,k}$ 
  - Atom  $k$  embedded in free electron gas with density of bare surface at current position  $\underline{r}_k$
- Random forces  $\underline{R}_{el,k}$ : gaussian white noise, dependent on both  $\eta_{el,k}$  (from LDFA) and  $T_{el}$  (from 2TM)
  - justified by the 2. fluctuation dissipation theorem [11], which relates friction and thermal movement



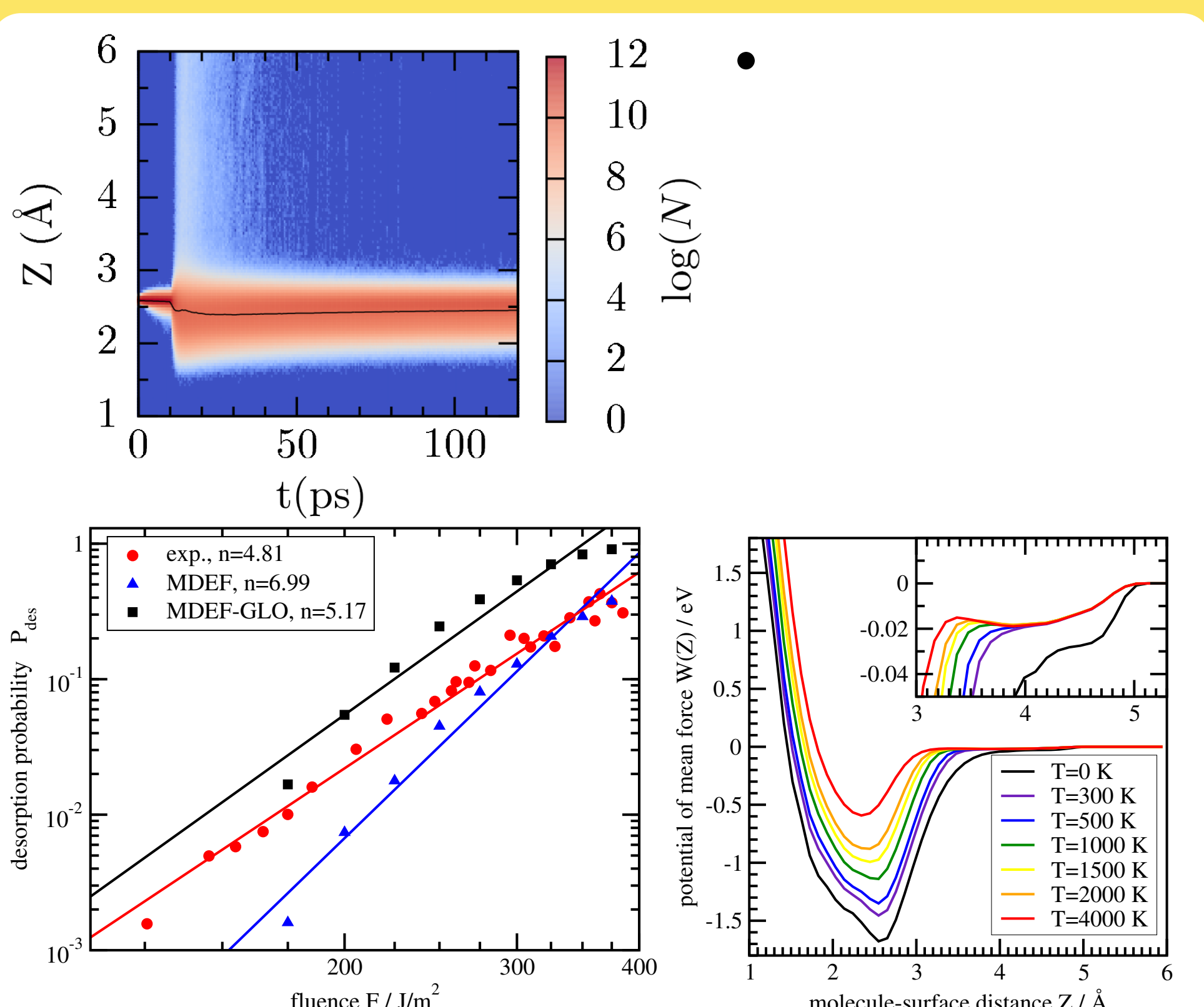
### Inclusion of Phonons: Generalized Langevin Oscillator(GLO)-model[8, 9, 10]

- influence of phonons modeled in an effective way (augments frozen surface)
- entire surface understood as 3D oscillator (coordinates  $\underline{r}_s$ , mass  $m_s = m_{Ru}$ )
- coupling to molecule via shifting:  $V_{GLO}(\underline{r}_C, \underline{r}_O; \underline{r}_s) = V(\underline{r}_C - \underline{r}_s, \underline{r}_O - \underline{r}_s)$
- additionally coupled to ghost oscillator  $\underline{r}_g$  to model influence of the bulk
  - ghost oscillator is subject to friction  $\eta_{ph}$  and random forces  $\underline{R}_{ph}(T_{ph})$

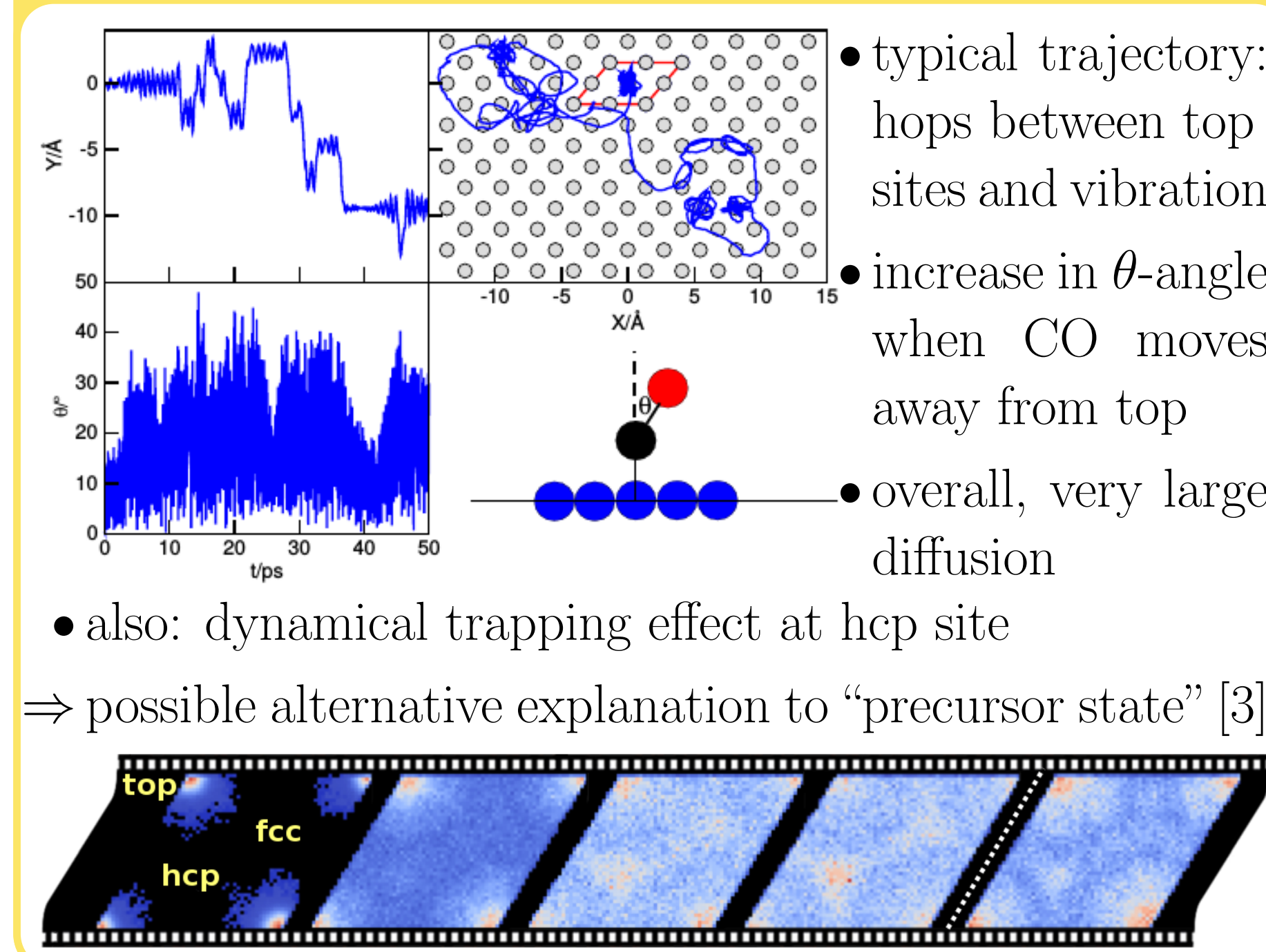


## Results

### Desorption

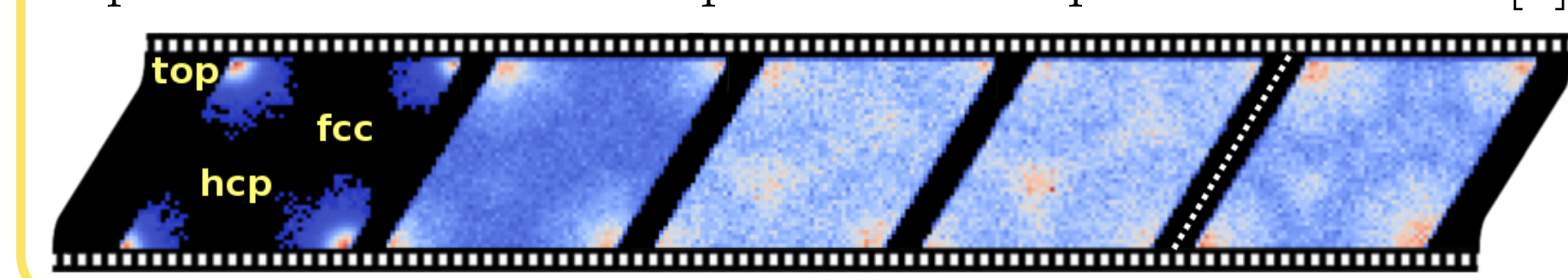


### Diffusion



- typical trajectory: hops between top sites and vibration
- increase in  $\theta$ -angle when CO moves away from top
- overall, very large diffusion
- also: dynamical trapping effect at hcp site

$\Rightarrow$  possible alternative explanation to “precursor state” [3]



## References

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